

Title: An Asymptotic Framework for Gravitational Scattering

Speakers: Samuel Gralla

Series: Strong Gravity

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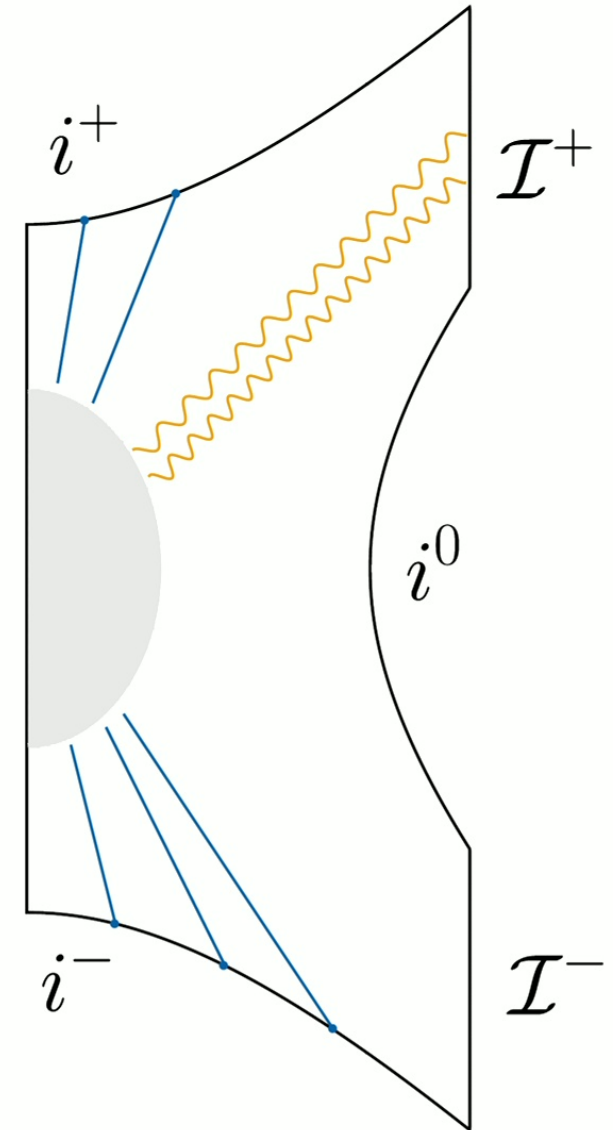
Abstract: Gravitational scattering provides valuable insight into classical dynamics and is likely of fundamental importance in quantum gravity. However, a complete framework for gravitational scattering does not yet exist. The lack of a clear framework hinders progress; for example, different groups have reported different results for the loss of angular momentum in post-Minkowskian scattering. In this talk I will report on recent work 2303.17124 with Compere and Wei constructing a general framework for classical gravitational scattering of finite-sized massive bodies in four spacetime dimensions. We formulate assumptions and definitions such that the five asymptotic regions (past/future timelike/null infinity and spatial infinity) share a single Bondi-Metzner-Sachs (BMS) group of symmetries and associated charges and derive global conservation laws stating that the total change in charge is balanced by the corresponding radiative flux. Our assumptions are compatible with all known properties of scattering spacetimes, including certain logarithmic corrections that invalidate common falloff assumptions. Among the new implications are rigorous definitions for quantities like initial/final spin, scattering angle, and impact parameter in multi-body spacetimes, without the use of any preferred background structure. We show that spin is supertranslation-invariant, while impact parameter is not. To complement these derivations I will emphasize a helpful "puzzle piece" diagram that faithfully represents all five asymptotic regions, illustrating their roles in the scattering problem.

Zoom link: <https://pitp.zoom.us/j/94358515412?pwd=cUFocVNjQ1pqUmp4MDN0RmRLbjE0QT09>

An asymptotic framework for gravitational scattering

Sam Gralla, University of Arizona

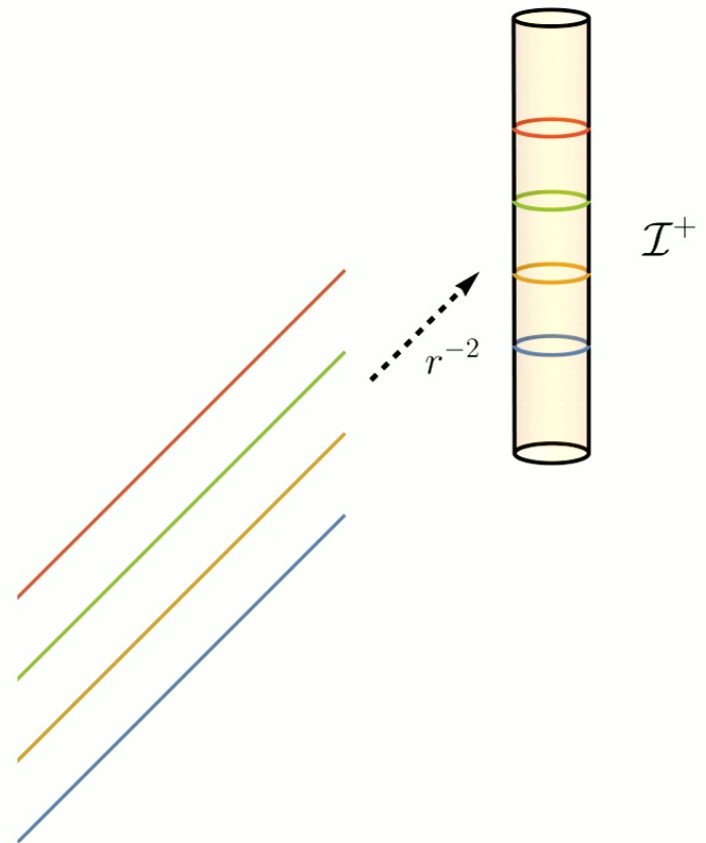
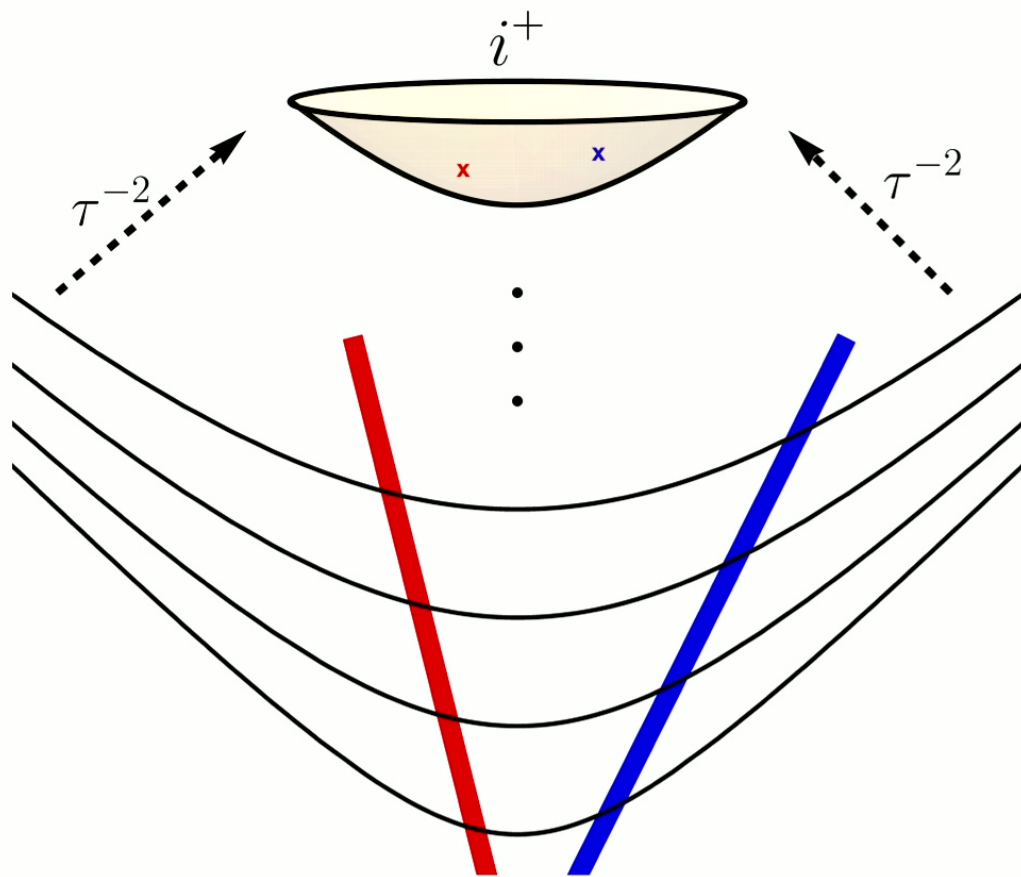
2303.17124 w/ Compere and Wei



Why a framework for scattering?

- Physical insight – conserved quantities in GR
- Quantum Gravity – symmetries of S-matrix (if it exists)
- “Practical” application – Compare different calculations
 - Need to understand “gauge-invariant observables”
 - Different groups are reporting different results for the change in angular momentum in weak-field scattering!

Need to understand “non-local observables” in gravity



Charge	name	generator
E	Energy	$T = 1$
P^i	Momentum	$T = n^i(\theta, \phi)$
$P_{\ell m}$	Supermomentum	$T = Y_{\ell m}(\theta, \phi)$
L^i	Angular Momentum	$Y^A = -\epsilon^{AB} \partial_B n^i(\theta, \phi)$
N^i	Mass Moment	$Y^A = \partial^A n^i(\theta, \phi)$

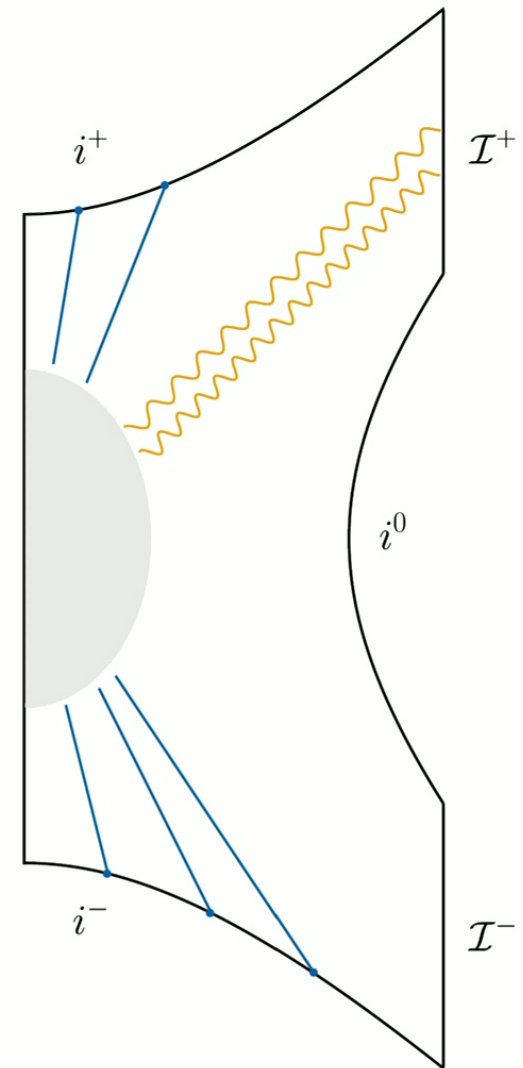
For each BMS charge (choice of T or Y):

$$\sum_{n=1}^{N^+} Q_n^{i^+} + \Delta Q^{\mathcal{I}^+} = Q^{i^0} = \sum_{n=1}^{N^-} Q_n^{i^-} + \Delta Q^{\mathcal{I}^-}$$

Charge	name	generator
E	Energy	$T = 1$
P^i	Momentum	$T = n^i(\theta, \phi)$
P_{lm}	Supermomentum	$T = Y_{lm}(\theta, \phi)$
L^i	Angular Momentum	$Y^A = -\epsilon^{AB} \partial_B n^i(\theta, \phi)$
N^i	Mass Moment	$Y^A = \partial^A n^i(\theta, \phi)$

Summary

- All 5 infinities share a single BMS group
- A charge for each BMS generator
- A conservation (flux-balance) law for each charge
- Define initial and final quantities using charges

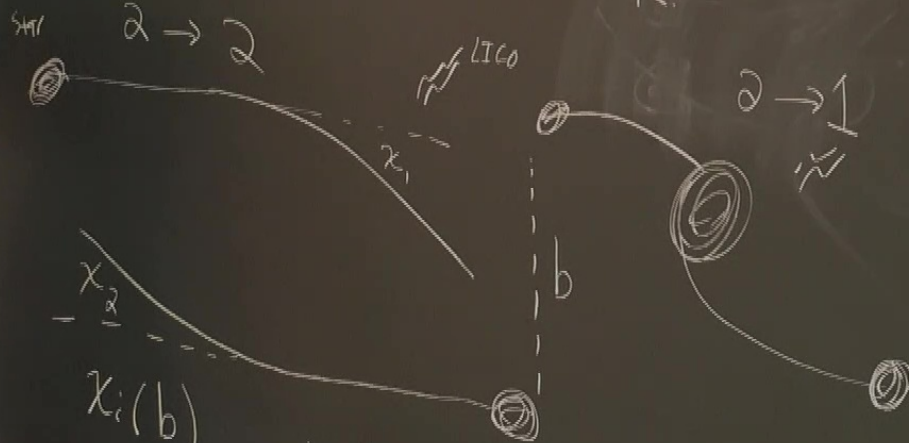


An Asymptotic Framework for Gravitational Scattering

Sam Gralla, U. of Arizona

2303.17124 w/ Compère, Wei

Classical $m \rightarrow n$ Scattering of finite-sized bodies in 4D GR.

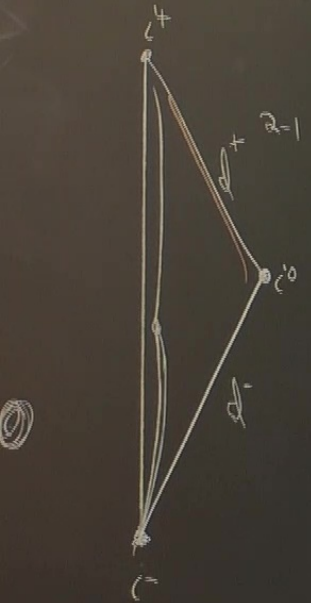


$$b \sim \frac{L_{orb} b}{P} \text{ in CM frame}$$

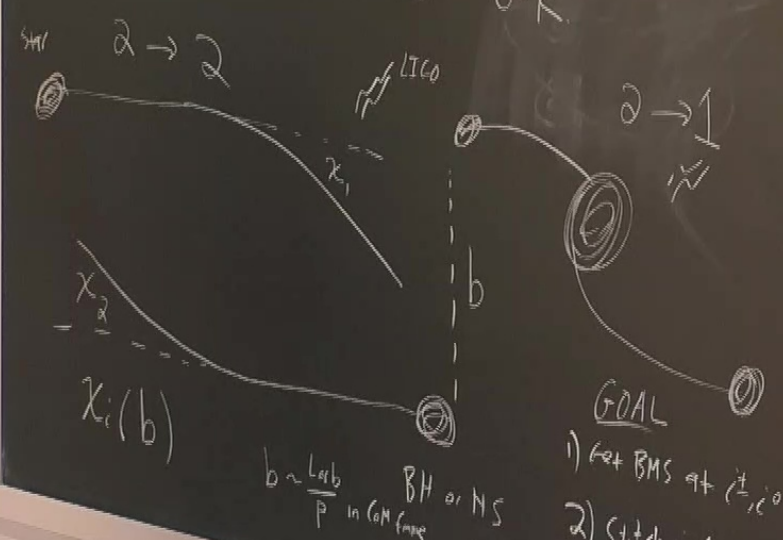
BH or NS

2305 17124 w/ Compère, Wei

of finite-sized



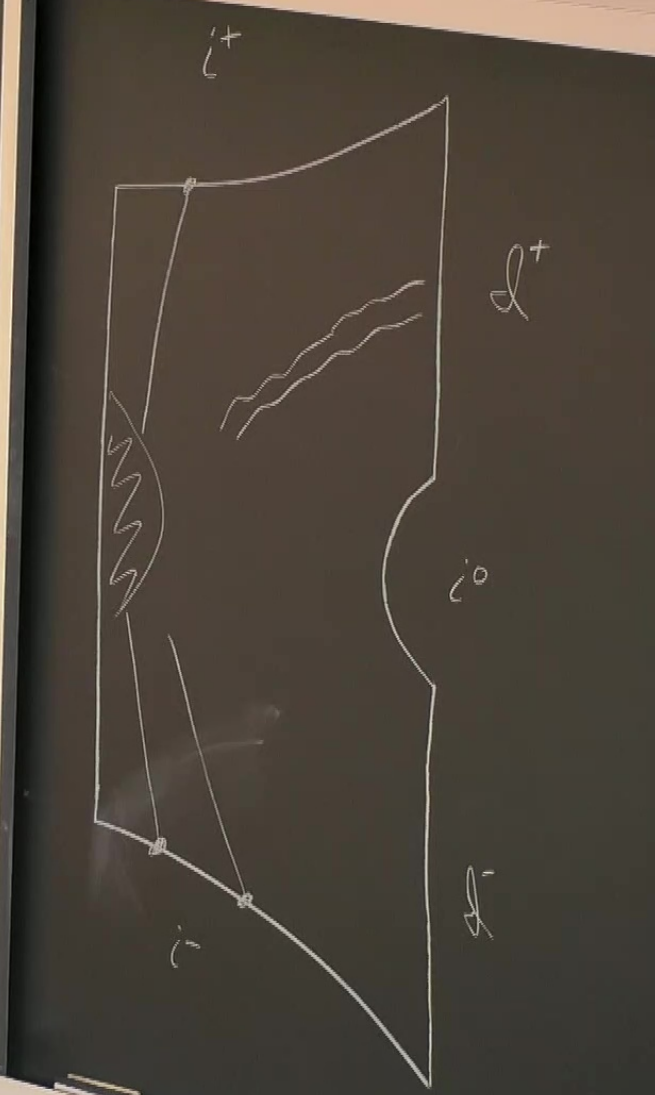
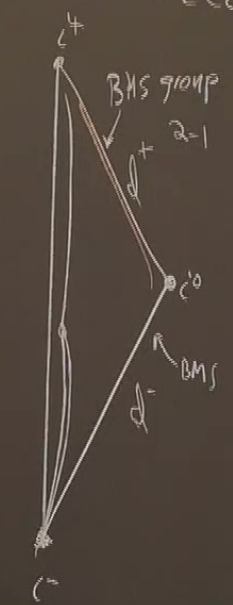
Classical $m \rightarrow n$ Scattering of finite-sized bodies in 4D GR.



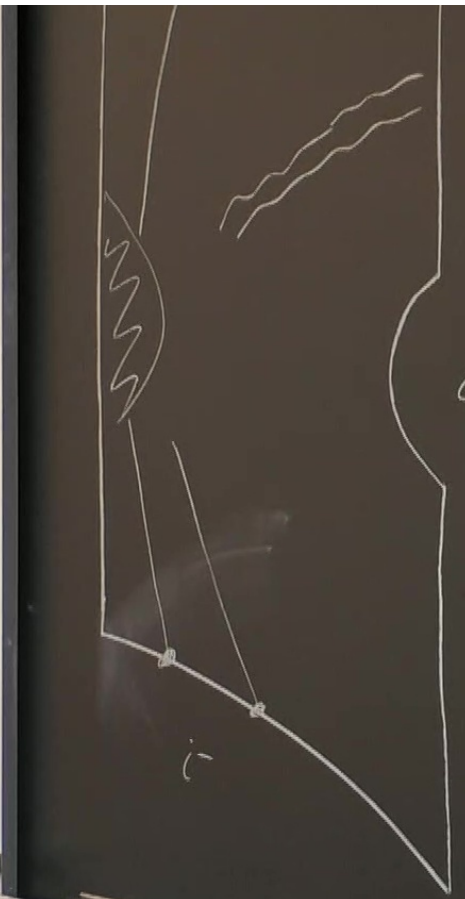
$x_i(b)$

$b \sim \frac{L_{\text{orb}}}{P}$ in CM frame
 BH or NS

- GOAL
- 1) Get BMS at (i^\pm, i^0)
 - 2) Stitch infinities together
 One BMS group.



2-sized
 BMS group
 d^+_{2-1}
 i^0
 BMS
 d^-



d^+

i^0

d^-

flat:

$x = Vxt$ Body has no limit in this approach.
 Introduce new coords that are const on outgoing bodies



$(\rho, \tau, \theta, \phi)$

rapidity = $\text{arctanh } v$

$\rho = \text{arctanh } \frac{r}{t}$
 proper time

$\tau = \sqrt{t^2 - r^2}$

$ds^2 = -d\tau^2 + \rho^2 (d\rho^2 + \sinh^2 \rho d\Omega^2)$

new coolds that are const on outgoing bodies

flat:



$$(r, \tau, \theta, \phi)$$

$$\text{rapidity} = \text{arctanh } v$$

$$p = \text{arctanh } \frac{r}{t}$$

proper time

$$\tau = \sqrt{t^2 - r^2}$$

$$ds^2 = -d\tau^2 + \tau^2 (dp^2 + \sinh^2 p d\Omega^2)$$

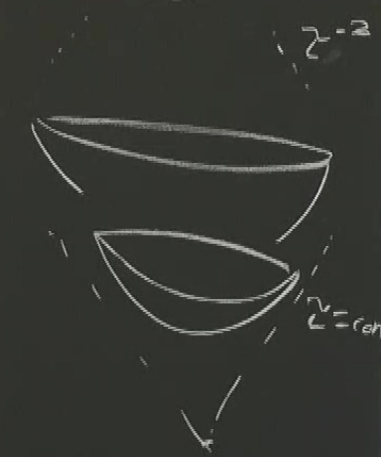
Spacetime is asymp. flat if } coolds sit. (r, θ, ϕ)

\mathcal{H} , Euclidean AdS_3 .

$$ds^2 \rightarrow$$

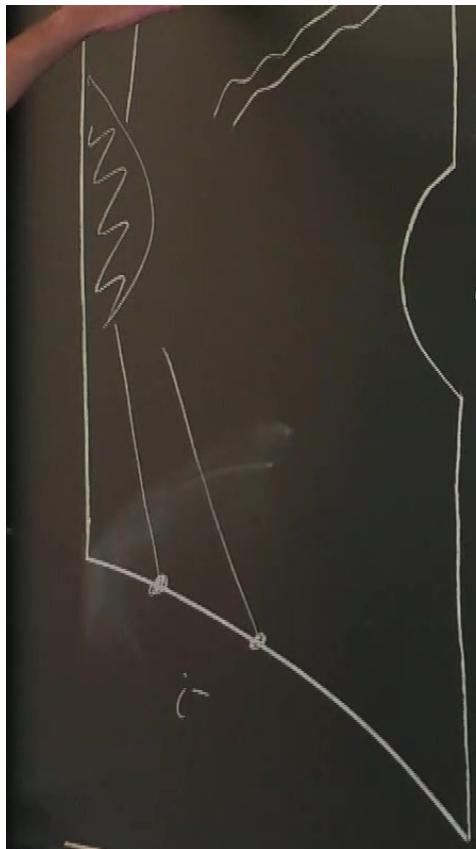
as $\tau \rightarrow \infty$.

\mathcal{I}^+ , $(\mathcal{H}(t))$



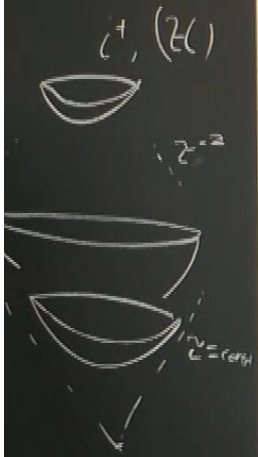
τ^{-2}

$\tau = \text{const}$



this approach.

outgoing bodies



Reig-Schmid

$$ds^2 = -\left(1 + \frac{2V}{r} + \frac{V^2}{r^2} + \dots\right) dt^2 + r^2 \left(h_{ab} + \frac{1}{r} (K_{ab} - \nabla h_{ab}) + \frac{\log r}{r^2} L_{ab} + \frac{1}{r^2} J_{ab} + \dots \right) d\phi^a d\phi^b$$

$\phi^a = (p, \theta, \phi)$ $r \rightarrow \infty$

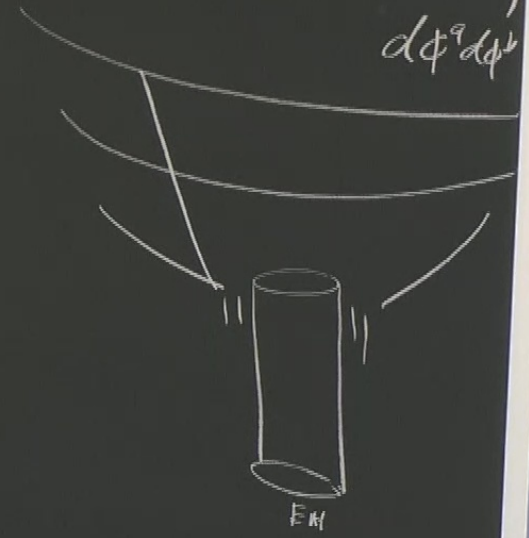
\mathcal{I}^+ : space spanned by $(p, \theta, \phi) = \phi^a$

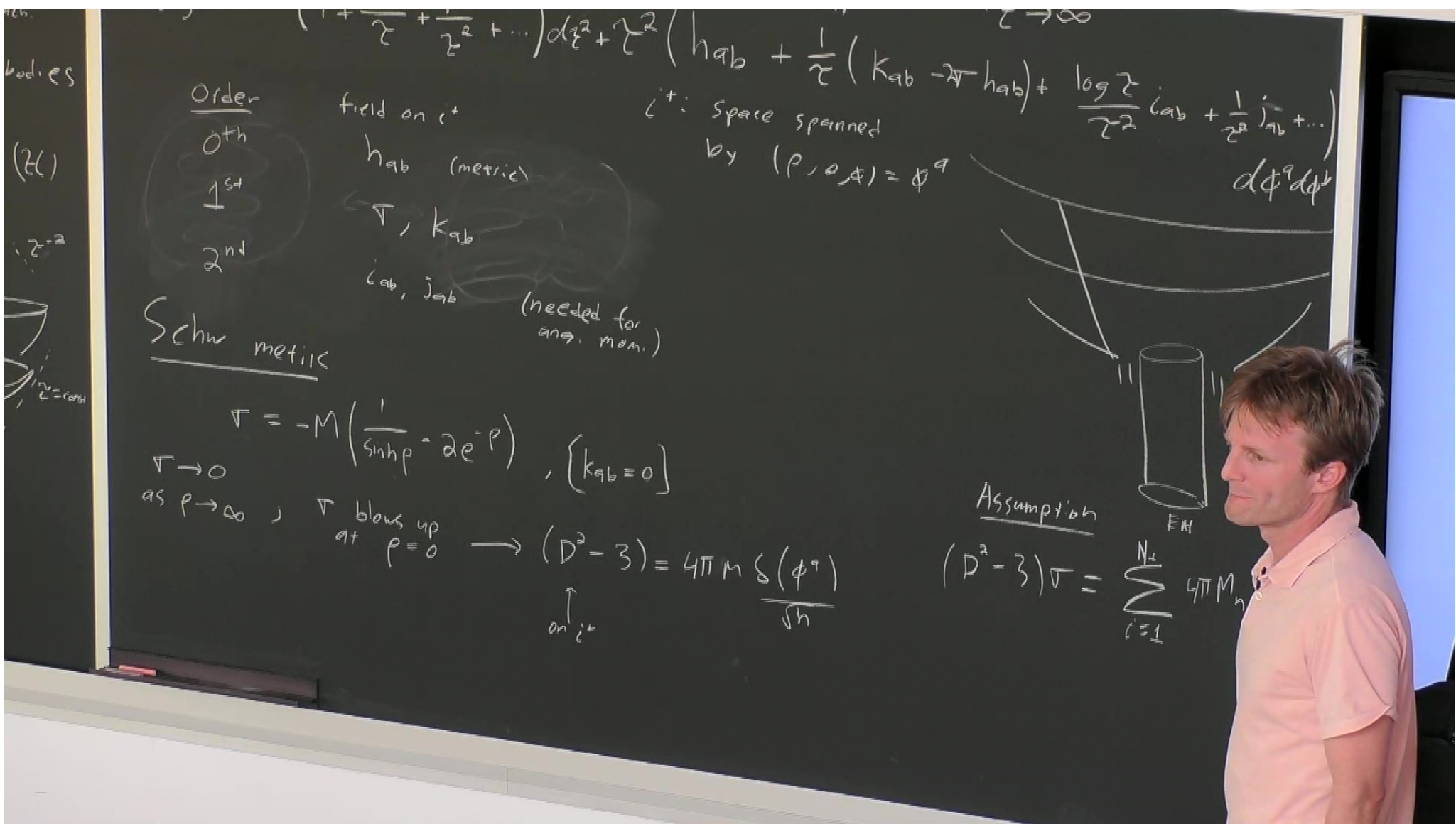
- Order
- 0th h_{ab} (metric)
 - 1st ∇, K_{ab}
 - 2nd L_{ab}, J_{ab} (needed for ang. mom.)

Schw metric

$$V = -M \left(\frac{1}{\sinh p} - a e^{-p} \right), [K_{ab} = 0]$$

$V \rightarrow 0$ as $p \rightarrow \infty$, V blows up at $p = 0$

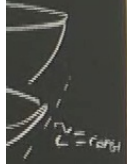




bodies

(\mathbb{Z})

\mathbb{Z}^2



Order

0th

1st

2nd

field on i^+

h_{ab} (metric)

∇, K_{ab}

L_{ab}, J_{ab}

(needed for ang. mom.)

Schw metric

$$\nabla = -M \left(\frac{1}{\sinh \rho} - 2e^{-\rho} \right), [K_{ab} = 0]$$

$\nabla \rightarrow 0$

as $\rho \rightarrow \infty$

∇ blows up at $\rho = 0$

$$\rightarrow (D^2 - 3) = 4\pi M \frac{\delta(\phi^a)}{\sqrt{h}}$$

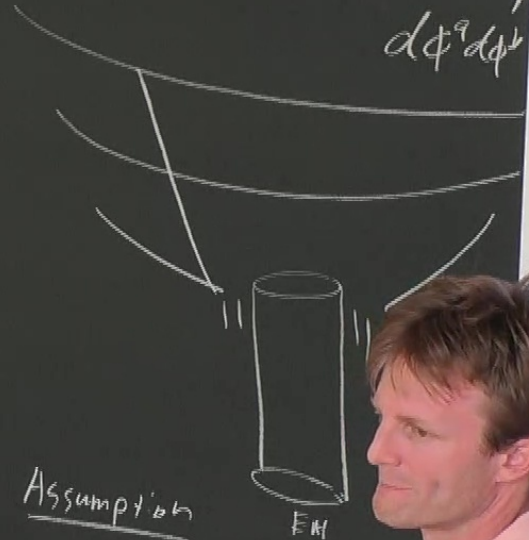
↑
on i^+

i^+ : space spanned

by $(\rho, \theta, \phi) = \phi^a$

$\mathbb{Z} \rightarrow \infty$

$$\left(h_{ab} + \frac{1}{\mathbb{Z}} (K_{ab} - 2\nabla h_{ab}) + \frac{\log \mathbb{Z}}{\mathbb{Z}^2} (L_{ab} + \frac{1}{\mathbb{Z}^2} J_{ab} + \dots) \right) d\phi^a d\phi^b$$



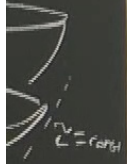
Assumption

$$(D^2 - 3)\nabla = \sum_{i=1}^{N_c} 4\pi M_i \frac{\delta(\phi^a)}{\sqrt{h}}$$

bodies

(Z)

z^{-2}



- Order
- 0th
 - 1st
 - 2nd

field on i^+

h_{ab} (metric)

∇, K_{ab}

L_{ab}, J_{ab}

(needed for ang. mom.)

Schw metric

$$\nabla = -M \left(\frac{1}{\sinh \rho} - 2e^{-\rho} \right), \quad [K_{ab} = 0]$$

$\nabla \rightarrow 0$

as $\rho \rightarrow \infty$

∇ blows up at $\rho = 0$

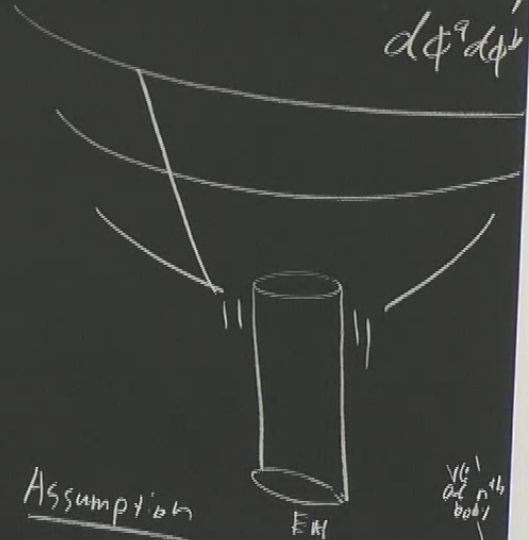
$$\rightarrow (D^2 - 3) = 4\pi M \frac{\delta(\phi^a)}{\sqrt{h}}$$

↑
on i^+

i^+ : space spanned by $(\rho, \theta, \phi) = \phi^a$

$z \rightarrow \infty$

$$\left(h_{ab} + \frac{1}{z} (K_{ab} - 2\nabla h_{ab}) + \frac{\log z}{z^2} L_{ab} + \frac{1}{z^2} J_{ab} + \dots \right) d\phi^a d\phi^b$$



Assumption

$$(D^2 - 3)\nabla = \sum_{i=1}^{N_b} 4\pi M_i \frac{\delta(\phi^a - \phi_i^a)}{\sqrt{h}}$$