

Title: Quantum Gravity Lecture (230502)

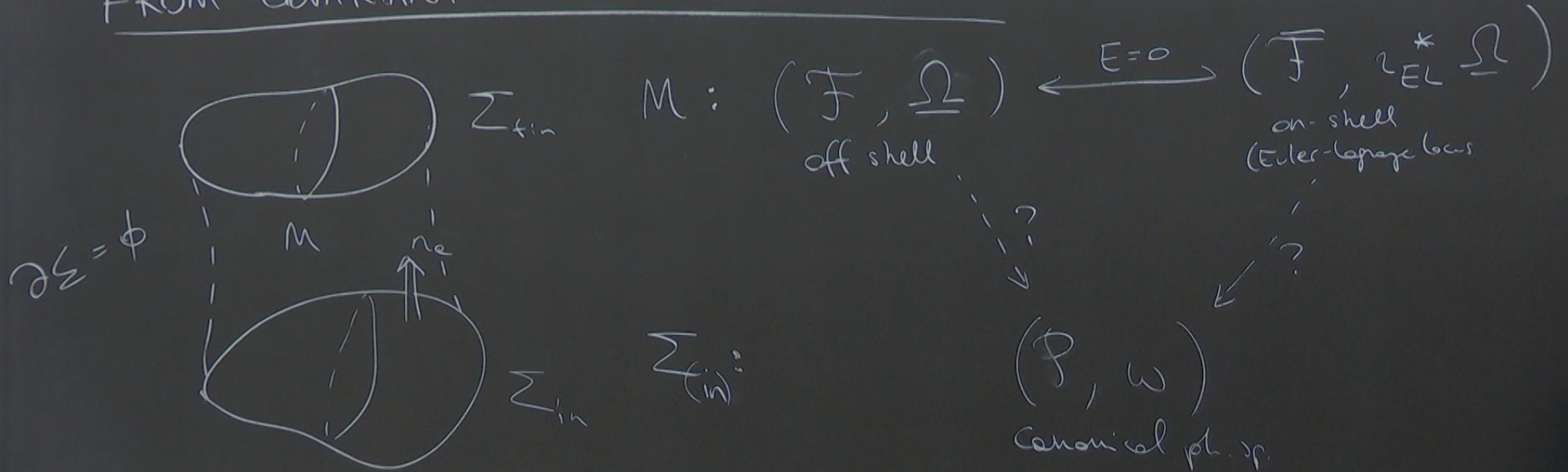
Speakers: Aldo Riello

Collection: Quantum Gravity (2022/2023)

Date: May 02, 2023 - 9:00 AM

URL: <https://pirsa.org/23050004>

# FROM COVARIANT TO CANONICAL PH. SPACE



$\rightarrow (\overline{\mathcal{F}}, \overset{*}{\mathcal{L}}_{EL} \underline{\Omega})$   
on-shell  
(Euler-Lagrange laws)

?

ph. sp.

## Example 1

Particle (0+1d field theory)

$$S[\gamma(t)] = \int \left( \frac{m}{2} \dot{\gamma}^2 - V(\gamma(t)) \right) dt$$

$$\mathcal{F} = \text{Maps}(\mathbb{R}, \mathbb{R}^3)$$

$$\underline{\Omega} = (\partial_t d\gamma(t))_1 d\gamma(t)$$

How to describe  $(\overline{\mathcal{F}}, \overset{*}{\mathcal{L}}_{EL} \underline{\Omega})$  explicitly?!

If eom. are nice,  $\exists$  map btw on-shell histories & initial conditions @  $\Sigma_{in}$ .

$M : (F, \underline{\Omega})$   
off shell

on-shell  
(Euler-Lagrange laws)

$(\mathcal{P}, \omega)$

canonical ph. sp.

$\underline{\Omega} \in \Omega^{top-1, 2}(M \times F)$

wants to be  
integrated over  
 $\Sigma$

$\Sigma$

## Example 1

Particle (0+1d field theory)

$$S[\gamma(t)] = \int \left( \frac{m}{2} \dot{\gamma}^2 - V(\gamma(t)) \right) dt$$

$$\mathcal{F} = \text{Maps}(\mathbb{R}, \mathbb{R}^3)$$

$$\Omega = (\partial_t d\gamma(t)) \lrcorner d\gamma(t)$$

How to describe  $(\overline{\mathcal{F}}, \overline{\Omega})$  explicitly?!

If con. or nice,  $\exists$  map btw on-shell histories & initial conditions @  $\Sigma_{in}$  -

$$\Sigma_{in} = \{t = t_0\}$$

$$\Omega_\Sigma = \Omega|_{t=t_0} = d\dot{\gamma}(t_0) \lrcorner d\gamma(t_0)$$

Rmk  $\Omega_\Sigma$  has a huge kernel in  $\mathcal{F}$

$$X \in \ker(\Omega_\Sigma) \text{ iff } \overset{\circ}{i}_X \Omega_\Sigma = 0$$

$$X = \left[ dt \left( \frac{\delta X}{\delta x} \right) \frac{\delta}{\delta \gamma(t)} \right]$$

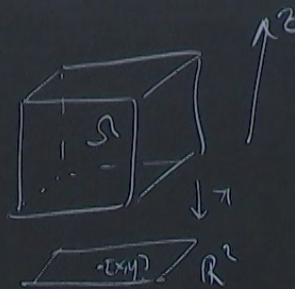
$$\text{iff } \delta_X \gamma(t_0) = (\partial_t \delta_X \gamma)(t_0) = 0$$

$$X = X^1 \frac{\partial}{\partial x} + X^2 \frac{\partial}{\partial y} + X^3 \frac{\partial}{\partial z}$$

$$\Omega = dx \wedge dy$$

$$\omega^1 = i_X \Omega = X^1 dy - X^2 dx$$

$$X^1 = X^2 = 0$$



$$\Sigma_{in} = \{t = t_0\}$$

$$\Omega_\Sigma = \Omega|_{t=t_0} = m(\partial_t dY(t_0) + dY(t_0))$$

Rmk  $\Omega_\Sigma$  has a huge kernel in  $F$

$$X \in \ker(\Omega_\Sigma) \text{ iff } \int_X \Omega_\Sigma = 0 \text{ as a 1-form}$$

$$X = \left[ \begin{array}{c} \partial_t (\delta Y) + \frac{\delta}{\delta Y(t_0)} \\ \delta Y(t_0) \end{array} \right]$$

$$\text{iff } \delta_X Y(t_0) = (\partial_t \delta_X Y)(t_0) = 0$$

Any two histories  $Y_1(t), Y_2(t)$

with the same  $Y_1(t_0) = Y_2(t_0)$

$$\dot{Y}_1(t_0) = \dot{Y}_2(t_0)$$

lie along the kernel of  $\Omega_\Sigma$

$\Rightarrow$  If I want a symplectic (Space  $P$  I need to identify

$$\rightarrow [\gamma(t)]_{q,p} = \{ \gamma \in F \}$$

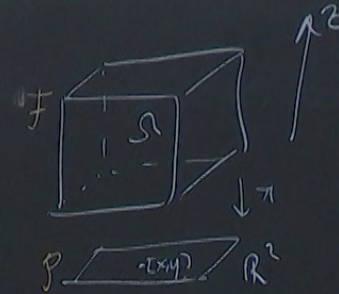
$$\rightarrow P := F / \ker(\Omega_\Sigma)$$

$$X = X^1 \frac{\partial}{\partial x} + X^2 \frac{\partial}{\partial y} + X^3 \frac{\partial}{\partial z}$$

$$\Omega = dx \wedge dy$$

$$\omega = i_X \Omega = X^1 dy - X^2 dx$$

$$X^1 = X^2 = 0$$





Any two histories  $\gamma_1(t), \gamma_2(t)$

with the same  $\gamma_1(t_0) = \gamma_2(t_0)$

$$\dot{\gamma}_1(t_0) = \dot{\gamma}_2(t_0)$$

lie along the kernel of  $\omega \in \text{Ker}(\Omega_\Sigma)$

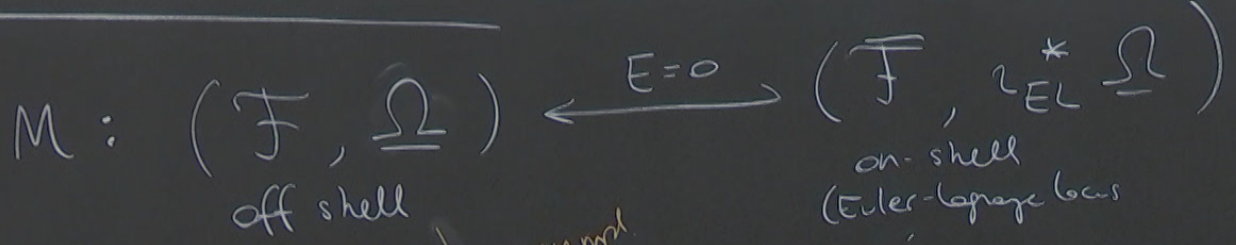
$\Rightarrow$  If I want a symplectic (non degenerate)

Space  $\mathcal{P}^{\text{off}}$  I need to identify all these config.

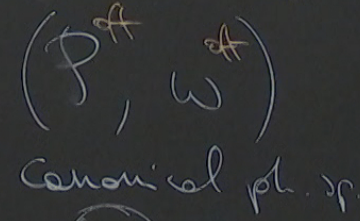
$$\rightarrow [\gamma(t)]_{q,p} = \{ \gamma \in \mathcal{F} : \gamma(t_0) = q, m\dot{\gamma}(t_0) = p \}$$

$$\rightarrow \mathcal{P}^{\text{off}} := \mathcal{F} / \text{Ker}(\Omega_\Sigma) \simeq \{ (q, p) \in \mathbb{R}^2 \} \quad \text{presymplectic reduction}$$

TO CANONICAL PH. SPACE



presympt. red.  
 $\int_{\text{Ker}(S_0)}$



$\underline{\Omega} \in \Omega^{\text{top-1}, 2}(M \times \mathcal{F})$

wants to be integrated over  $\Sigma$

Example 1

Particle (or 1d field theory)

$S[\gamma(t)] = \int \left( \frac{m}{2} \dot{\gamma}^2 - \dots \right)$

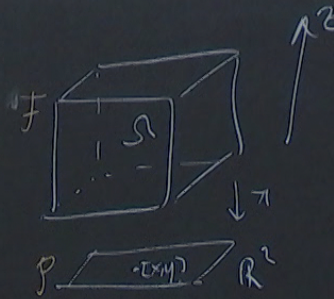
$\mathcal{F} = \text{Maps}(\mathbb{R}, \mathbb{R}^3)$

$\underline{\Omega} = (\int_t d\gamma(t)) \wedge d\gamma(t)$

How to describe  $(\overline{\mathcal{F}}, \int_{EL}^* \underline{\Omega})$

If eom. are nice,  $\exists$  m. histories & initial cond.

$$X = X^1 \frac{\partial}{\partial x} + X^2 \frac{\partial}{\partial y} + X^3 \frac{\partial}{\partial z}$$



$$\Omega = dx \wedge dy$$

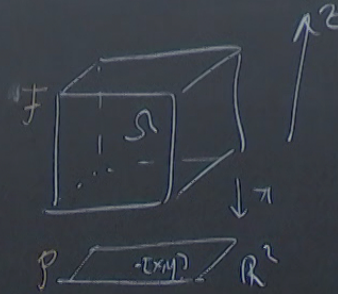
$$\omega^1 = i_X \Omega = X^1 dy - X^2 dx$$

$$X^1 = X^2 = 0$$

on  $\mathcal{P}^{\text{off}}$ ,  $\Omega_\Sigma$  induces a non-degenerate

2-form  $\omega^{\text{off}} = dp \wedge dq$  ( $d$  in this ex is one) finite dim space

$$X = X^1 \frac{\partial}{\partial x} + X^2 \frac{\partial}{\partial y} + X^3 \frac{\partial}{\partial z}$$



$$\Omega = dx \wedge dy$$

$$\omega^1 \lrcorner \Omega = X^1 dy - X^2 dx$$

$$X^1 = X^2 = 0$$

on  $\mathcal{P}^{\text{off}}$ ,  $\Omega_\Sigma$  induces a non-degenerate

2-form  $\omega^{\text{off}} = dp \wedge dq$  ( $d$  in this ex is on  $\mathcal{P}$ )  
 finite dim space

Lucky:  $(p, q)$  are also in 1-to-1 correspondence with on-shell histories.

# CANONICAL PH. SPACE

$$(\mathcal{F}, \underline{\Omega}) \xleftrightarrow{E=0} (\overline{\mathcal{F}}, \mathbb{Z}_{EL}^* \underline{\Omega})$$

off shell

on-shell  
(Euler-Lagrange laws)

presympt. red.  
%/ker(S<sub>0</sub>)

lucky case

$$(\mathcal{P}^{off}, \omega^{off}) = (\mathcal{P}, \omega)$$

Canonical ph. sp.

$$\underline{\Omega} \in \Omega^{top-1, 2} (M \times \mathcal{F})$$

wants to be  
integrated over  
 $\Sigma$

## Example 1

Particle (0+1d)

$$S[\gamma(t)] = \int (\dots)$$

$$\mathcal{F} = \text{Maps}(\mathbb{R}, \dots)$$

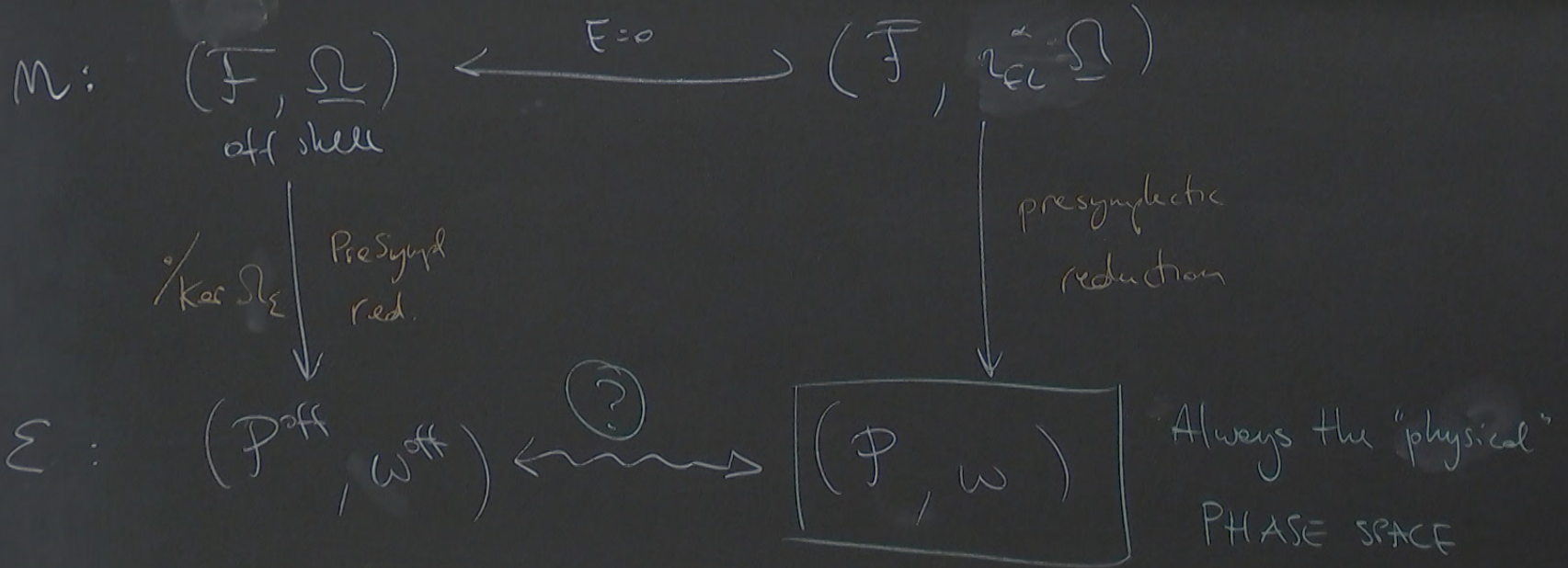
$$\underline{\Omega} = (\partial_t d \gamma(t))$$

How to describe

If eom. or nic  
histories & in

For particle & scalar field theory

(?) = an isomorphism.



Scalar field theory

$$\mathcal{F} = C^\infty(M, \mathbb{R})$$

$$\Omega_\Sigma = \int_\Sigma \sqrt{|h|} (\nabla_n d\varphi(t, x)) \wedge d\varphi(x, t)$$

presympl red:  $(\mathcal{P}^{\text{off}}, \omega^{\text{off}}) \simeq (\mathcal{P}, \omega)$

$$\mathcal{P}^{\text{off}} = \left\{ [\varphi(x, t)] : \begin{cases} \varphi(t_0, x) = \phi(x) \\ \sqrt{|h|} \nabla_n \varphi(t_0, x) = \pi(x) \end{cases} \right\}$$

function on  $\Sigma$       density on  $\Sigma$

↓                      ↓

$$= \left\{ (\phi(x), \pi(x)) \right\}$$

Maxwell (YM very similar)

$$F = \Omega'(M)$$

$$S = \frac{1}{2} \int *F \wedge F, \quad F = dA$$

$$\underline{\Omega} = *dF \wedge dA$$

$$\underline{\Omega}_\Sigma = \int_\Sigma *dF(A) \wedge dA = \int_\Sigma \sqrt{|h|} dF^{oi}(A) \wedge dA_i$$

$$F^{\alpha} \sim \dot{A}_i - \nabla^i A_\alpha$$



## Rmks

- 1)  $A_0$  only appears "mixed" with  $\dot{A}_i$
- 2)  $\dot{A}_i$  does not appear in isolation

$$\left[ \text{and } \frac{\partial \mathcal{L}}{\partial (\partial_t A_0)} = 0 \right]$$

## Presymplectic reduction:

we identify all  $A_\mu(t, x)$  according to:

$$[A_\mu] = \left\{ A_\mu(x) : \left. \begin{array}{l} A_i|_\Sigma = a_i \in \Omega^1(\Sigma) \\ \sqrt{h} F^{0i}(A)|_\Sigma = \Sigma^i \end{array} \right\} \right. \left. \begin{array}{l} \text{a vector-valued} \\ \text{density on } \Sigma \end{array} \right\}$$

$\Sigma^i = \sqrt{h} E^i$  densitized electric field.

to split  $A_T = (A_0, \vec{A})$   
 at all "times"

$$\mathcal{P}_{\text{Maxwell}}^{\text{off}} = \left\{ (a_i, \epsilon^i) \right\} = T^* \Omega^1(\Sigma)$$

$$\omega_{\text{Maxwell}}^{\text{off}} = \int_{\Sigma} d\epsilon^i(x) \wedge da_i(x)$$

densitized  
 electric fields

"space like"  
 1-forms on  $\Sigma$   
 (no  $A_0$ !)

vector-valued  
 density on  $\Sigma$

$\sqrt{|h|} E^i$  densitized  
 electric field.

reduction LUCKY.  $(p, q)$  are also in 1-to-1

$(\mathcal{P}^{\text{off}}, \omega^{\text{off}})$  can't be isomorphic to  $(\mathcal{P}, \omega)_{\Sigma}$  on shell!

d.f.c. on shell:  $\sqrt{h} \nabla_i E^i = 0$   
 $\uparrow$   
 $\partial_i \mathcal{E}^i = 0$

(Gauss Constraint)

why?!

presympt  
reduction

$\int_{\Sigma} \rho(\xi) \Omega_{\Sigma} = -d \int_{\Sigma} Q_{\Sigma}(\xi)$

$Q_{\Sigma}(\xi) = \int_{\Sigma} \sqrt{h} (\nabla_i F^{i0}) \xi + \int_{\partial \Sigma} \sqrt{h} F^{i0} \xi$

Noether II  
 $\downarrow$   
 Gauss  $\approx 0$   
 $\nearrow$   $\partial \Sigma = \emptyset$

$\int_{\Sigma} \rho(\xi) \omega^{\text{off}} = -d \int_{\Sigma} (\partial_i \mathcal{E}^i) \xi \approx 0$  by Noether 2

Constraints (e.g. Gauss)

are e.o.m. which have a  
left-over impact on  $(p^{\text{off}}, w^{\text{off}})$

Always present in gauge theories  
because of Noether 2:

$$\underline{\underline{J}}(\xi) = \underbrace{C_a}_{\substack{\approx 0 \\ \boxed{N_2}}} \xi^a + d \dot{\xi}^i$$

$\partial Z = \phi$   
↗  
↖

appears "mixed" with  $\vec{A}_i$   
 do not appear in isolation

$$\left[ \leftarrow \frac{\partial \mathcal{L}}{\partial (\partial_t A_0)} = 0 \rightleftharpoons \text{Lagrange multiplier for Gauss constraint} \right]$$

but for this one needs a foliation to split  $A_\mu = (A_0, \vec{A})$  at all "times"

reduction

$$\underline{\Sigma} = \mathcal{L}_\Sigma^* (*F) \in \Omega^{\text{top}}(\Sigma)$$

pull all  $A_\mu(t, x)$  according to:

$$A_\mu(x) \in \mathcal{M} : \left. \begin{aligned} A_i|_\Sigma &= a_i \in \Omega^1(\Sigma) \\ \sqrt{|g|} F^{0i}(A)|_\Sigma &= \mathcal{E}^i \end{aligned} \right\} \begin{array}{l} \text{vector-valued} \\ \text{density on } \Sigma \end{array}$$

$\mathcal{E}^i = \sqrt{|g|} E^i$  densitized electric field

$$\mathcal{P}_{\text{Maxwell}}^{\text{off}} = \left\{ (a_i, \mathcal{E}^i) \right\} = T^* \Omega^1(\Sigma)$$

$$\omega_{\text{Maxwell}}^{\text{off}} = \int_\Sigma d\mathcal{E}^i(x) \wedge da_i(x)$$

densitized electric fields

"spacelike" 1-forms on  $\Sigma$  (no  $A_0$ !)