

Title: Parity-Violation in Gravity: New Constructions and Inflationary Signals

Speakers: Cyril Creque-Sarbinowski

Series: Quantum Gravity

Date: April 13, 2023 - 2:30 PM

URL: <https://pirsa.org/23040090>

Abstract: I will talk about new constructions to the parity-violating sector of gravity that can yield large observables, in particular those from inflation. Within this framework, I will demonstrate the potential for seeking gravitational parity-violation in uncharted regions of parameter space and give hints to novel methods of probing baryogenesis. I will also comment on a few interesting theoretical aspects regarding inflationary correlators.

Zoom link: <https://pitp.zoom.us/j/95924890817?pwd=MkRyWmV2a3g1dndPNFhoVlQvb3Ixdz09>

# Parity Violation in Gravity: New Constructions and Inflationary Signals

Cyril Creque-Sarbinowski

Quantum Gravity Seminar - Perimeter 2023

Based on 2207.05094 w/ Stephon Alexander  
and 2303.04815 with Stephon, Marc Kamionkowski, & Oliver Philcox



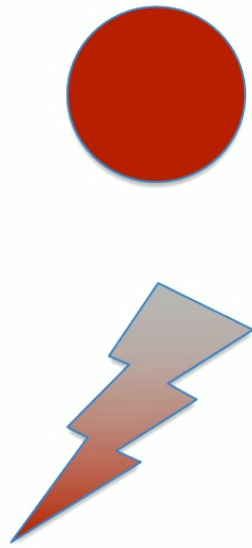
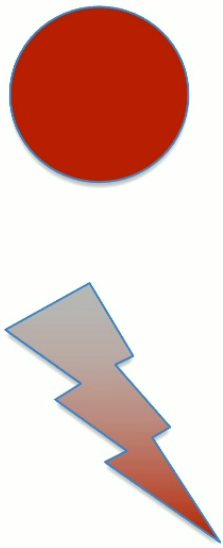
# Roadmap

I will...

- I) Demonstrate two new constructions (i.e types of UV completions) of dynamical Chern-Simons (dCS) gravity that encompass a wide range of EFT scales
- II) Show that dCS gravity yields a parity-violating primordial scalar trispectrum that can be made large.
- III) Demonstrate that ratios of odd and even trispectra contain direct information about spin.
- IV) Quantify how well upcoming and future experiments can constrain scalar trispectra.
- V) Hint to how parity-violating trispectra can constrain baryogenesis.

# What is Parity? How can Gravity violate it?

$$\mathbb{P}\mathbf{x} = -\mathbf{x}, \mathbb{P}\varphi(\mathbf{x}) = \pm\varphi(-\mathbf{x})$$



$$A_L - A_R$$

# What is [dynamical] Chern-Simons (dCS) Gravity?

$$S = \int d^4x \sqrt{-g} \left[ -\frac{1}{2} g^{\mu\nu} (\partial_\mu \phi) (\partial_\nu \phi) + \frac{M_{\text{pl}}^2}{2} R - \frac{\phi}{4f} *RR \right]$$

↑
↑
↑
↑

Spacetime  
Metric,  
Lorentzian  
Signature

dCS pseudo-scalar

Einstein-  
Hilbert Term

dCS Decay  
Constant

Pontryagin  
Density

$*RR = *R^\rho{}_\sigma{}^{\mu\nu} R^\sigma{}_{\rho\mu\nu}$

$*R^\rho{}_\sigma{}^{\mu\nu} = \frac{1}{2} \epsilon^{\mu\nu\alpha\beta} R^\rho{}_{\sigma\alpha\beta}$

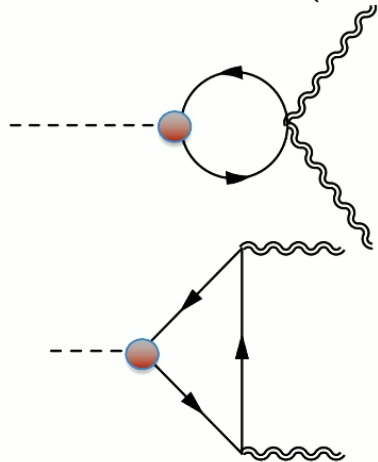
# New Construction I: Spontaneous Symmetry Breaking

$$\Phi = (1/\sqrt{2})(F + \sigma) \exp[ia(x)/F]$$



$$\mathcal{L}_\Phi = \sum_{j=L,R} i\bar{\Psi}_j \gamma^\mu \partial_\mu \Psi_j - \tilde{m}_\Psi (\bar{\Psi}_R \Psi_L + \bar{\Psi}_L \Psi_R) + \partial_\mu \Phi \partial^\mu \Phi^* - V(|\Phi|^2) - (y_\Phi \Phi \bar{\Psi}_L \Psi_R + \text{h.c.})$$

$$\mathcal{L}_\Phi = \bar{\Psi} \left[ i\gamma^\mu \partial_\mu - \left( \tilde{m}_\Psi + \frac{y_\Phi F}{\sqrt{2}} \right) \right] \Psi + \frac{1}{2} \partial_\mu a \partial^\mu a - \frac{y_\Phi}{\sqrt{2}} a \bar{\Psi} \gamma^5 \Psi,$$



$$\mathcal{L}_g^{\text{eff}} = -\frac{g}{384\pi^2} \frac{a}{2m_\Psi} * RR$$

$$f = 192\sqrt{2}\pi^2 \frac{m_\Psi}{y_\Phi}$$



# New Construction 2: Dynamical Symmetry Breaking

$$\mathcal{L}_\Psi = \bar{\Psi} (i\gamma^\mu \partial_\mu - \tilde{m}_\Psi) \Psi - \lambda \bar{\Psi} \Psi \bar{\Psi} \Psi$$

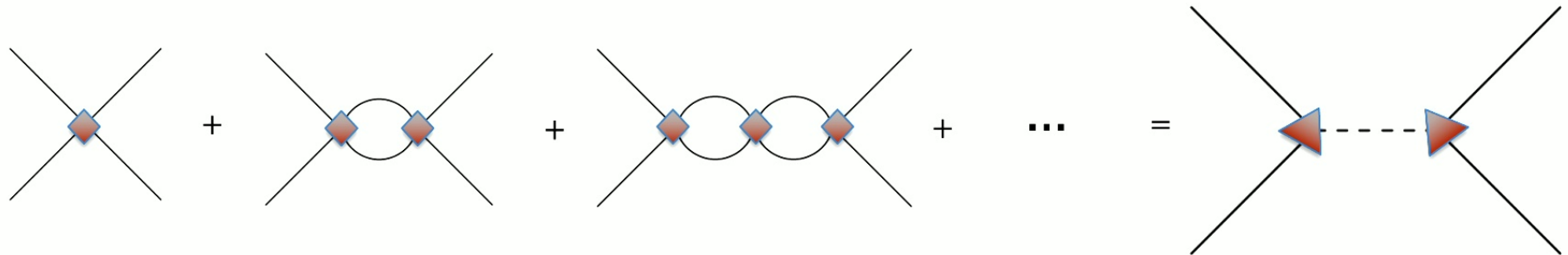
★ Cutoff  $\Lambda$

↑  
Dirac Gamma  
Matrices

↑  
"bare" mass

↑  
Dirac  
Fermion

↑  
(Attractive) Fermion  
Self-Coupling Constant



# New Construction 2: Dynamical Symmetry Breaking

$$Z_\alpha = \int \mathcal{D}\alpha \mathcal{D}\bar{\alpha} \exp\left(-\int d^4x \tilde{m}_\Phi^2 \bar{\alpha}\alpha\right)$$

+

$$\Psi = \Psi_\ell + \Psi_s$$

$$\Phi = \alpha - \tilde{m}_\Phi^{-2} \bar{\Psi}_s \Psi_s$$

+

$\Rightarrow$

$$\begin{aligned} \tilde{\mathcal{L}}_\Psi &= \bar{\Psi}(i\gamma^\mu \partial_\mu - \tilde{m}_\Psi)\Psi - \lambda \bar{\Psi}\Psi\bar{\Psi}\Psi \\ &+ (\partial_\mu \Phi^*) (\partial^\mu \Phi) - \underline{y_\Phi (\Phi \bar{\Psi}\Psi + \text{h.c.})} \\ &+ m_\Phi^2 |\Phi|^2 - \frac{\lambda_\Phi}{4} |\Phi|^4, \end{aligned}$$



Integrate Out Short – Scale Modes

$$y_\Phi \sim \lambda \Lambda^2$$



# What are Observable Signatures of dCS?

- 1) Compact Binary Coalescence (CBC) Waveform-Modification
  - 2) Birefringence of CBC Gravitational Waves
  - 3) Modification of Orbits around Earth
- $f \ll eV$
- } Late-Time
- 
- 4) Polarized Primordial Gravitational Waves
  - 5) Primordial Imprints on Galaxy Clustering and the CMB
- } Early-Time

# How is Galaxy Clustering Described?

## Galaxy N-Point Functions!

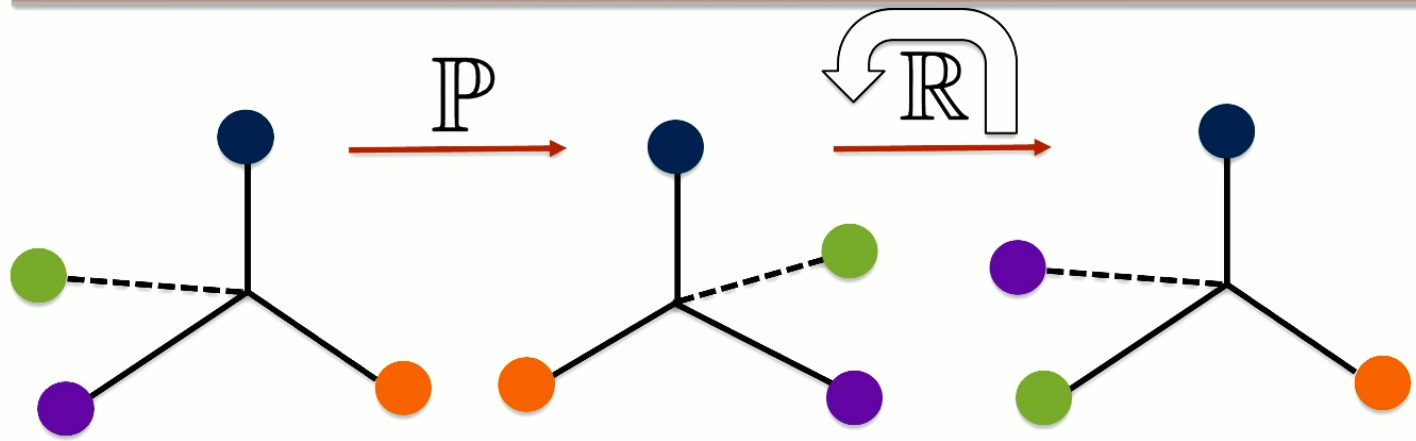
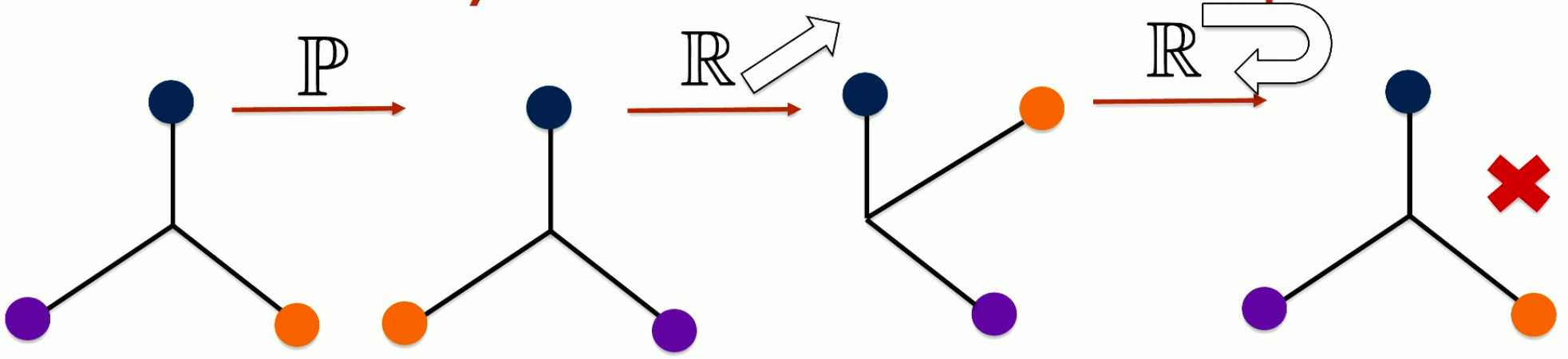
$$\delta_g(\mathbf{x}) = \frac{n_g(\mathbf{x}) - \bar{n}_g}{\bar{n}_g}$$

$$\langle \delta(\mathbf{x}) \rangle_c = 0$$

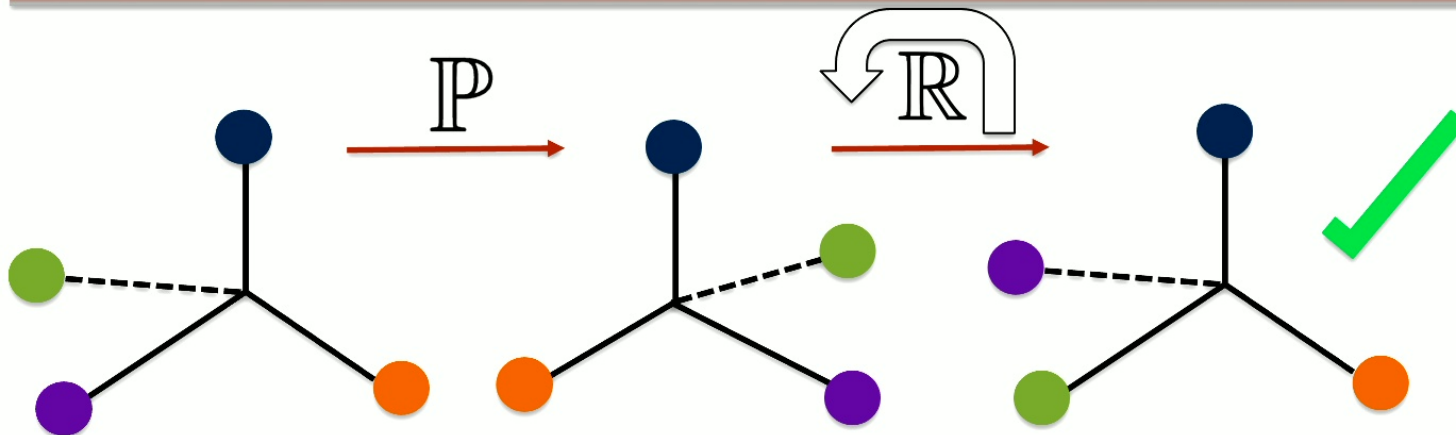
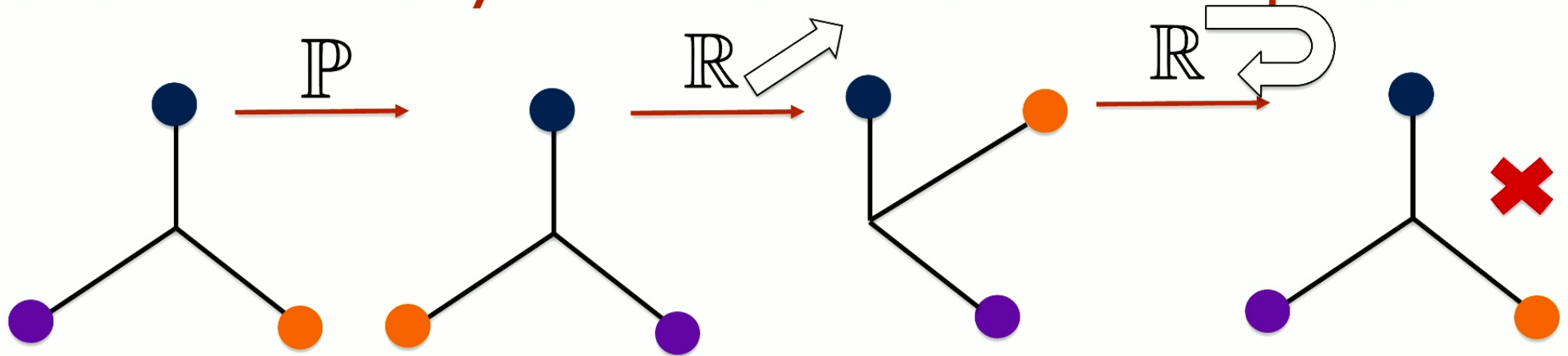
$$\langle \delta(\mathbf{x}_1) \delta(\mathbf{x}_2) \rangle_c = \xi_2(\mathbf{x}_1, \mathbf{x}_2) = \xi_2(|\mathbf{x}_1 - \mathbf{x}_2|)$$

$$\langle \delta(\mathbf{x}_1) \dots \delta(\mathbf{x}_N) \rangle_c = \xi_N(\mathbf{x}_1, \dots, \mathbf{x}_N)$$

# Enter: The Galaxy Four Point Function / Trispectrum



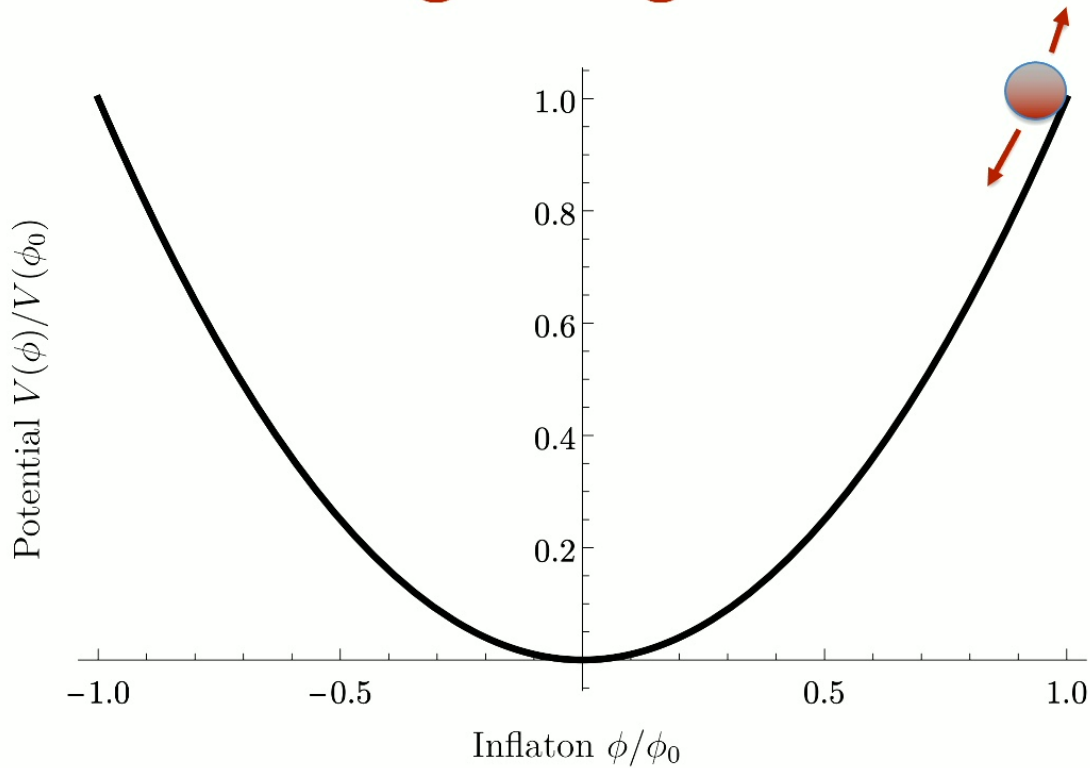
# Enter: The Galaxy Four Point Function / Trispectrum



P-odd Trispectrum  
 =  
 Im (Full Trispectrum)



# Inflation: A Lightning Review



Slow-Roll Parameter

$$\epsilon = \frac{1}{2} \frac{\dot{\phi}^2}{M_{\text{pl}}^2 H^2}$$

Solves:

- 1) The Horizon Problem
- 2) The Flatness Problem
- 3) Magnetic Monopole Problem

Quantum Inflaton Fluctuations Seed  
Density Perturbations



# dCS Inflation: Full Action

$$S = \int d^4x \sqrt{-g} \left[ -\frac{1}{2} g^{\mu\nu} (\partial_\mu \phi) (\partial_\nu \phi) + \frac{M_{\text{pl}}^2}{2} R - \frac{\phi}{4f} *RR \right]$$

↑
↑
↑
↑
↑
↑
↑
↑

Spacetime  
Metric,  
Lorentzian  
Signature

dCS pseudoscalar/Inflaton

Einstein-  
Hilbert Term

Pontryagin  
Density

$$g_{\mu\nu} = a^2(\tau) \begin{pmatrix} -1 & & 0 \\ & \delta_{ij} & \\ 0 & & + h_{ij} \end{pmatrix}$$

dCS Decay  
Constant

$*RR = *R^\rho_{\sigma\mu\nu} R^\sigma_{\rho\alpha\beta}$   
 $*R^\rho_{\sigma\mu\nu} = \frac{1}{2} \epsilon^{\mu\nu\alpha\beta} R^\rho_{\sigma\alpha\beta}$

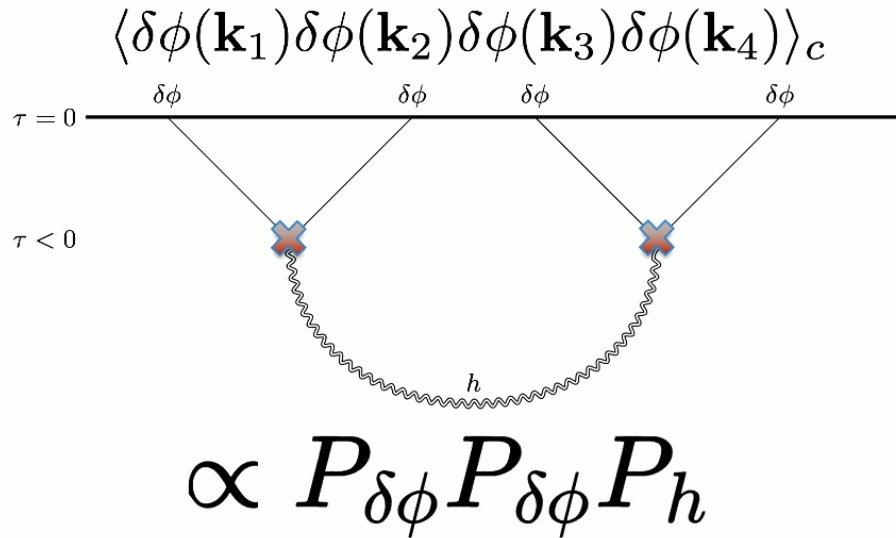
Inflaton Potential:  $V \sim \Lambda^4$

$$\Lambda \sim \sqrt{HM_{\text{pl}}} = 10^{16} \text{ GeV} [H/(10^{14} \text{ GeV})]^{1/2}$$



arXiv:2303.04815, Collaborators: Oliver Philcox, Stephon Alexander, Marc Kamionkowski

# Primordial Scalar Trispectrum from dCS



## Perturbed Inflaton Kinetic Term

$$\frac{1}{2}\sqrt{-g}g^{\mu\nu}\partial_\mu\delta\phi\partial_\nu\delta\phi = \frac{1}{2}a^2\eta^{\mu\nu}\partial_\mu\delta\phi\partial_\nu\delta\phi + \boxed{\frac{1}{2}a^2h^{ij}\partial_i\delta\phi\partial_j\delta\phi}$$

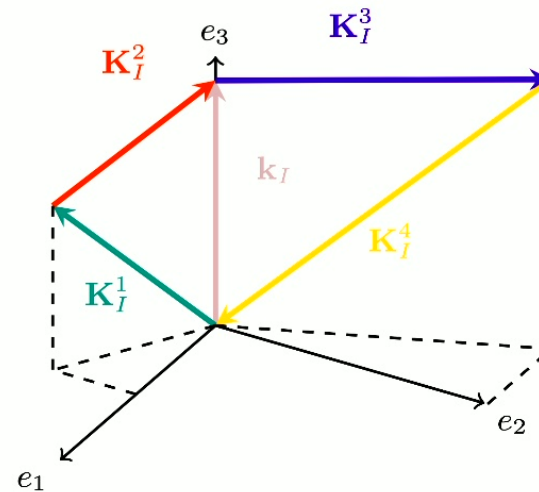
$$\mathbf{K}_s = \{\mathbf{k}_1, \mathbf{k}_2, \mathbf{k}_3, \mathbf{k}_4\}$$

$$\mathbf{K}_t = \{\mathbf{k}_1, \mathbf{k}_3, \mathbf{k}_2, \mathbf{k}_4\}$$

$$\mathbf{K}_u = \{\mathbf{k}_1, \mathbf{k}_4, \mathbf{k}_3, \mathbf{k}_2\}$$

$$\mathbf{K}_I^1 + \mathbf{K}_I^2 = -\mathbf{K}_I^3 - \mathbf{K}_I^4 = \mathbf{k}_I$$

$$I \in \{s, t, u\}$$



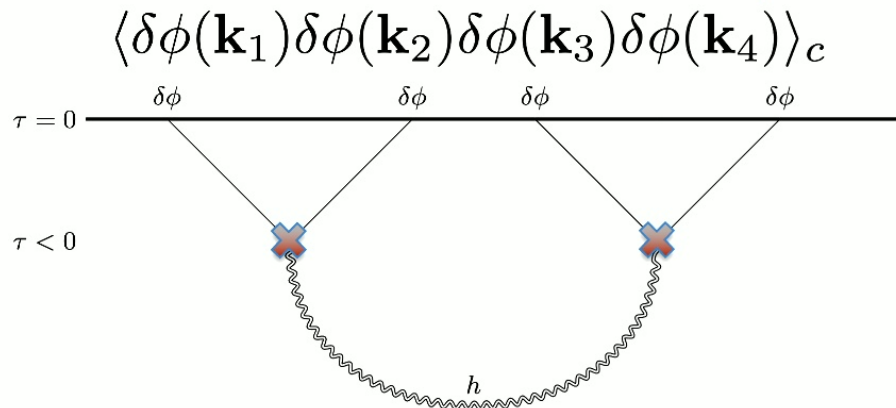
$$\begin{aligned} \hat{\mathbf{e}}_1 &= R_I \hat{\mathbf{x}} \\ \hat{\mathbf{e}}_2 &= R_I \hat{\mathbf{y}} \\ \hat{\mathbf{e}}_3 &= R_I \hat{\mathbf{z}} \end{aligned}$$

$$\hat{\mathbf{e}}_3 = \hat{\mathbf{k}}_I$$

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# Primordial Scalar Trispectrum from dCS



$$\propto P_{\delta\phi} P_{\delta\phi} P_h$$

$$\zeta(\mathbf{k}) = -H[\delta\phi(\mathbf{k})/\dot{\phi}]$$

## Perturbed Inflaton Kinetic Term

$$\frac{1}{2}\sqrt{-g}g^{\mu\nu}\partial_\mu\delta\phi\partial_\nu\delta\phi = \frac{1}{2}a^2\eta^{\mu\nu}\partial_\mu\delta\phi\partial_\nu\delta\phi + \frac{1}{2}a^2h^{ij}\partial_i\delta\phi\partial_j\delta\phi$$

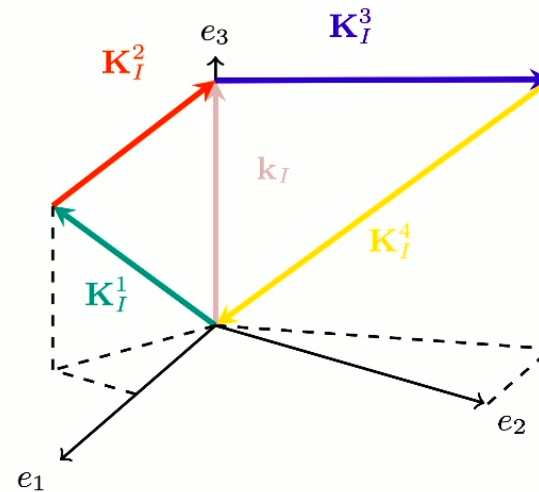
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$$I \in \{s, t, u\}$$



$$\begin{aligned} \hat{\mathbf{e}}_1 &= R_I \hat{\mathbf{x}} \\ \hat{\mathbf{e}}_2 &= R_I \hat{\mathbf{y}} \\ \hat{\mathbf{e}}_3 &= R_I \hat{\mathbf{z}} \end{aligned}$$

$$\hat{\mathbf{e}}_3 = \hat{\mathbf{k}}_I$$

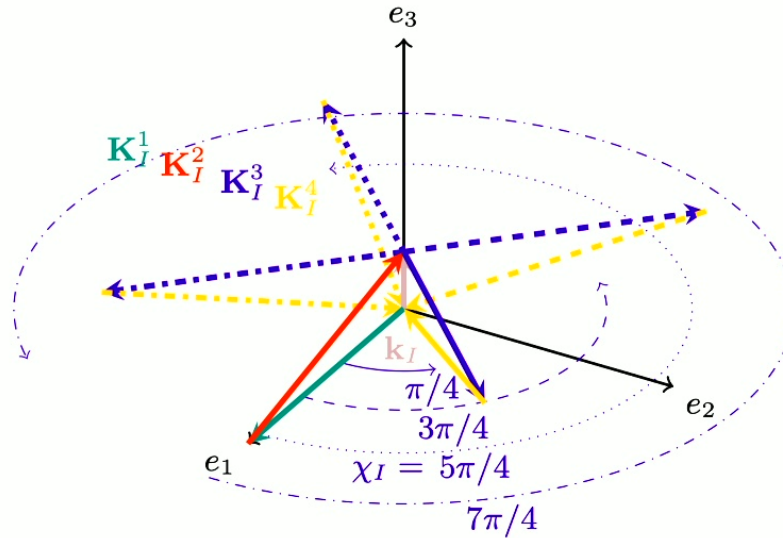
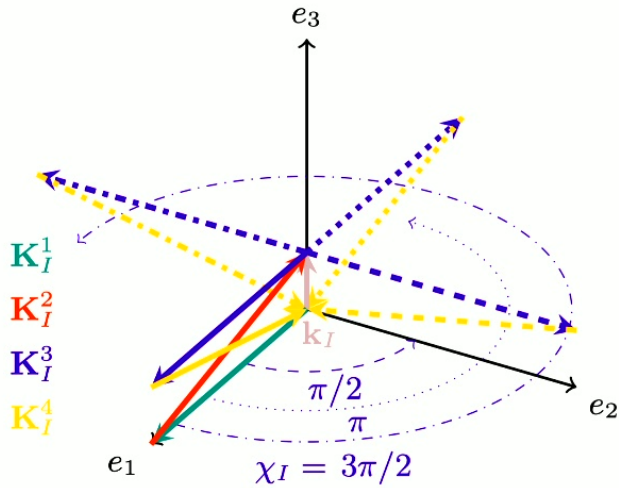
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# The Collapsed Limit Scalar Trispectra

$$\text{Re} \langle \zeta(\mathbf{k}_1) \zeta(\mathbf{k}_1) \zeta(\mathbf{k}_3) \zeta(\mathbf{k}_4) \rangle_c^{\text{GE}} \Big|_{I\text{-Coll.}} = \frac{9}{16} r \cos(2\chi_I) \sin^2(\theta_I^1) \sin^2(\theta_I^3) P_\zeta(k_I) P_\zeta(K_I^1) P_\zeta(K_I^3)$$

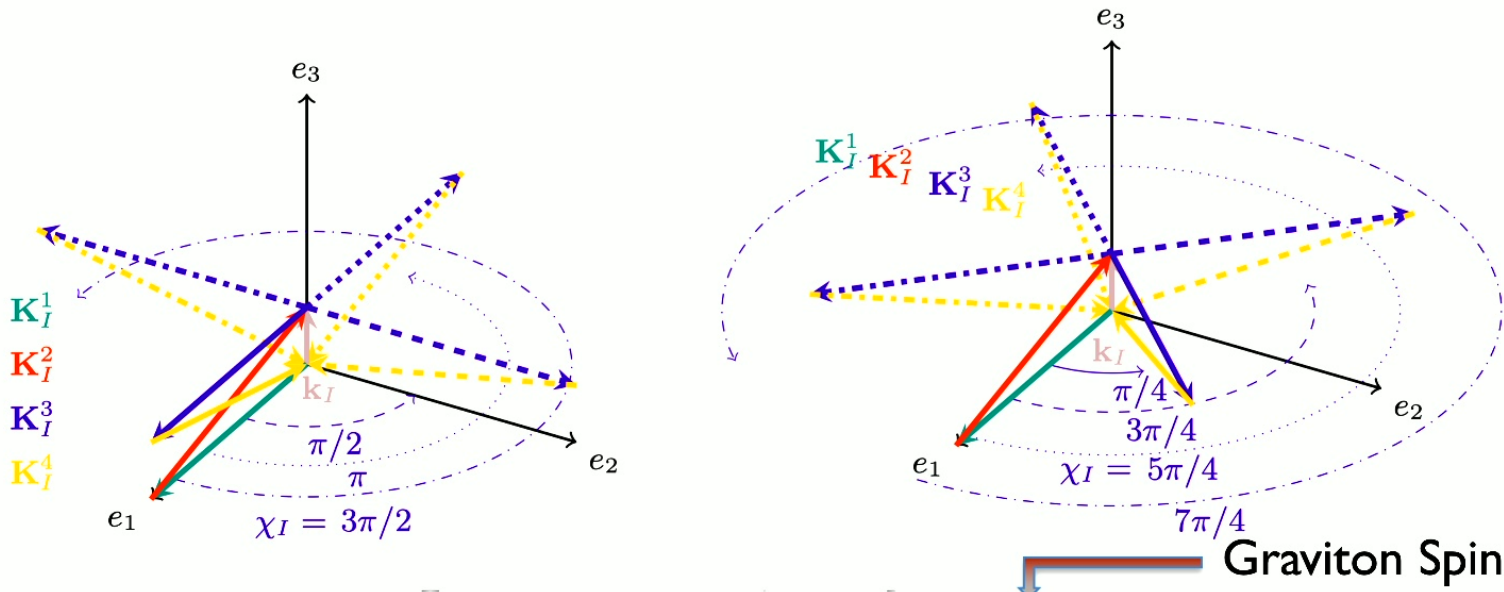
$$\text{Im} \langle \zeta(\mathbf{k}_1) \zeta(\mathbf{k}_1) \zeta(\mathbf{k}_3) \zeta(\mathbf{k}_4) \rangle_c^{\text{GE}} \Big|_{I\text{-Coll.}} = \frac{9}{16} \Pi_{\text{circ}} r \sin(2\chi_I) \sin^2(\theta_I^1) \sin^2(\theta_I^3) P_\zeta(k_I) P_\zeta(K_I^1) P_\zeta(K_I^3)$$



# The Collapsed Limit Scalar Trispectra

$$\text{Re} \langle \zeta(\mathbf{k}_1) \zeta(\mathbf{k}_1) \zeta(\mathbf{k}_3) \zeta(\mathbf{k}_4) \rangle_c^{\prime \text{GE}} \Big|_{I\text{-Coll.}} = \frac{9}{16} r \cos(2\chi_I) \sin^2(\theta_I^1) \sin^2(\theta_I^3) P_\zeta(k_I) P_\zeta(K_I^1) P_\zeta(K_I^3)$$

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$$\text{Odd/Even} = \Pi_{\text{circ}} \cot(2\chi_I)$$

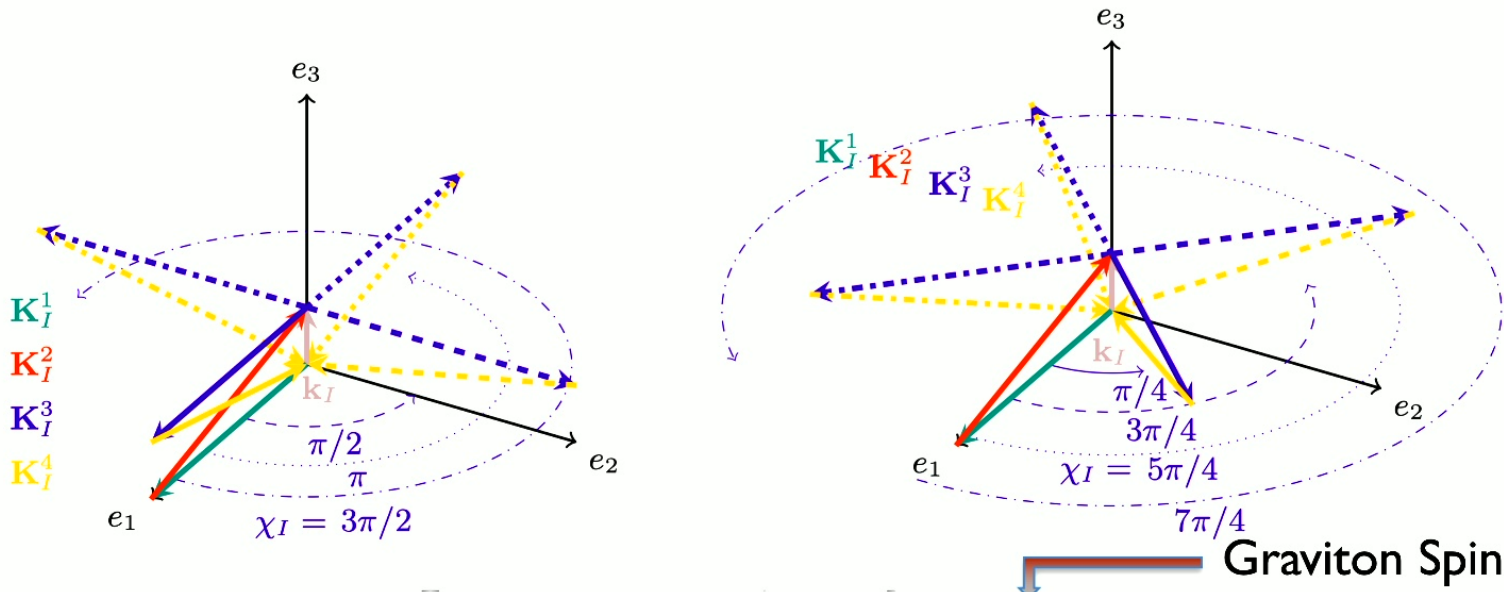


arXiv:2303.04815, Collaborators: Oliver Philcox, Stephon Alexander, Marc Kamionkowski

# The Collapsed Limit Scalar Trispectra

$$\text{Re} \langle \zeta(\mathbf{k}_1) \zeta(\mathbf{k}_1) \zeta(\mathbf{k}_3) \zeta(\mathbf{k}_4) \rangle_c'^{\text{GE}} \Big|_{I\text{-Coll.}} = \frac{9}{16} r \cos(2\chi_I) \sin^2(\theta_I^1) \sin^2(\theta_I^3) P_\zeta(k_I) P_\zeta(K_I^1) P_\zeta(K_I^3) \times \underline{\mathbf{F}}$$

$$\text{Im} \langle \zeta(\mathbf{k}_1) \zeta(\mathbf{k}_1) \zeta(\mathbf{k}_3) \zeta(\mathbf{k}_4) \rangle_c'^{\text{GE}} \Big|_{I\text{-Coll.}} = \frac{9}{16} \Pi_{\text{circ}} r \sin(2\chi_I) \sin^2(\theta_I^1) \sin^2(\theta_I^3) P_\zeta(k_I) P_\zeta(K_I^1) P_\zeta(K_I^3) \times \underline{\mathbf{F}}$$



$$\text{Odd/Even} = \Pi_{\text{circ}} \cot(2\chi_I)$$

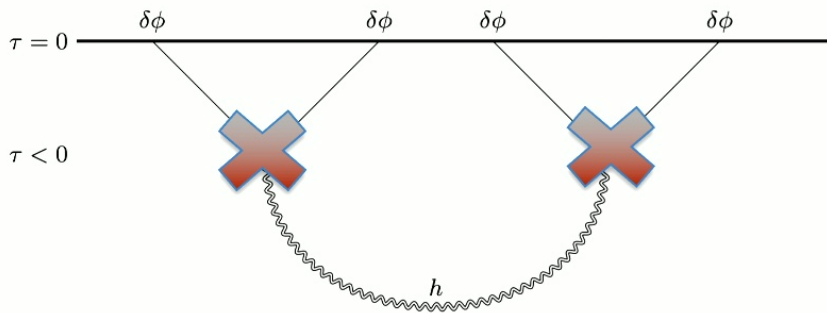


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# Two Examples Beyond Vanilla dCS Inflation

Superluminal Scalar Sound Speed

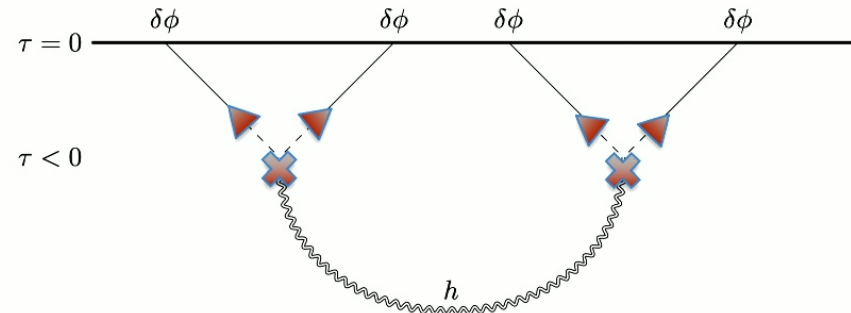
$$c_s^2 g^{ij} (\partial_i \phi) (\partial_j \phi) \quad \times$$



$$F = 2 \times 10^6 (c_s/6)^8$$

Quasi-Single Field/Multi-Field

$$a^3 \lambda \delta \sigma \dot{\delta} \phi \quad \triangle \quad h^{ij} \partial_i (\delta \sigma) \partial_j (\delta \sigma) \quad \times$$



$$F \sim 10^4 [w(\nu)/20]$$

$$\nu = \sqrt{9/4 - (m/H)^2} \lesssim 0.5$$

arXiv: 1504.05993  
(Dimastrogiovanni, Fasiello, Kamionkowski)

$$F \sim 10^7 (N/60)^4$$

$$\nu \sim 3/2$$

arXiv: 0911.3380  
(Xingang Chen, Yi Wang)



The Trispectra can be Large

arXiv: 2303.04815, Collaborators: Oliver Philcox, Stephon Alexander, Marc Kamionkowski



# How Sensitive are we to Collapsed Trispectra?

$$[1 + \delta_{\text{odd}}^P (\Pi_{\text{circ}} - 1)] r \gtrsim 0.04 \left(\frac{n}{3}\right) \left(\frac{8 \times 10^5}{F}\right) \left(\frac{10^6}{N_{\text{modes}}}\right)^{1/2} \left(\frac{k_{\text{min}}}{0.003 \text{ h/Mpc}}\right)^{3/2} \left(\frac{0.3 \text{ h/Mpc}}{k_{\text{max}}}\right)^{3/2}$$

$$f \lesssim 4 \times 10^9 \text{ GeV} \left(\frac{3}{n}\right) \left(\frac{F}{8 \times 10^5}\right) \left(\frac{\varepsilon}{10^{-2}}\right)^{3/2} \left(\frac{H}{10^{14} \text{ GeV}}\right)^2 \left(\frac{N_{\text{modes}}}{10^6}\right)^{1/2} \left(\frac{0.003 \text{ h/Mpc}}{k_{\text{min}}}\right)^{3/2} \left(\frac{k_{\text{max}}}{0.3 \text{ h/Mpc}}\right)^{3/2}$$

$$N_{\text{modes}} = V \int_{k_{\text{min}}}^{k_{\text{max}}} \frac{d^3 \mathbf{k}}{(2\pi)^3} W(\mathbf{k}) \frac{[G^2(\mathbf{k}, \bar{z}) P_L^{gg}(\mathbf{k}, \bar{z})]^2}{[P_{\text{NL}}^{gg}(\mathbf{k}, \bar{z}) + \bar{n}^{-1}]^2}$$

# What about Experiments?

$$[1 + \delta_{\text{odd}}^P (\Pi_{\text{circ}} - 1)] r \gtrsim 0.04 \left(\frac{n}{3}\right) \left(\frac{8 \times 10^5}{F}\right) \left(\frac{10^6}{N_{\text{modes}}}\right)^{1/2} \left(\frac{k_{\text{min}}}{0.003 \text{ h/Mpc}}\right)^{3/2} \left(\frac{0.3 \text{ h/Mpc}}{k_{\text{max}}}\right)^{3/2}$$

Experiments									
Name	$[z_{\text{min}}, z_{\text{max}}]$	$f_{\text{sky}}$	$[k_{\text{min}}, k_{\text{max}}] [h/\text{Mpc}]$	$\bar{n} [(h/\text{Mpc})^3]$	$\bar{b}$	$V [(\text{Gpc}/h)^3]$	$N_{\text{modes}}$	$F$	
BOSS <sup>a</sup> [112, 113]	[0.43, 0.7]	0.23	[0.01, 0.24]	$3 \times 10^{-4}$	2.0	3.6	$10^5$	$2 \times 10^7$	
DESI <sup>a</sup> [115]	[0.6, 1.7]	0.34	[0.003, 0.31]	$3.8 \times 10^{-4}$	1.4	45	$8 \times 10^5$	$8 \times 10^5$	
Euclid <sup>a</sup> [114]	[0.9, 1.8]	0.36	[0.003, 0.34]	$4.3 \times 10^{-4}$	1.7	44	$10^6$	$6 \times 10^5$	
MegaMapper <sup>a</sup> [116, 117]	[2, 5]	0.34	[0.003, 0.64]	$2.5 \times 10^{-4}$	3.8	155	$7 \times 10^6$	$10^5$	
MSE <sup>a</sup> [120]	[1.6, 4]	0.24	[0.003, 0.54]	$2.8 \times 10^{-4}$	3.7	91	$6 \times 10^6$	$10^5$	
SPHEREx <sup>a</sup> [118, 119]	[0.1, 4.3]	0.65	[0.003, 0.46]	$2.0 \times 10^{-3}$	1.1	360	$2 \times 10^7$	$10^5$	
HIRAX <sup>b</sup> [121]	[0.8, 2.5]	0.36	[(0.01, 0.1), 0.38]	$10^{-3}$	1.9	88	$3 \times 10^6$	$(2 \times 10^6, 7 \times 10^7)$	
PUMA-32K <sup>b</sup> [122]	[2, 6]	0.5	[(0.01, 0.1), 0.71]	$7.6 \times 10^{-3}$	6.3	290	$7 \times 10^7$	$(2 \times 10^5, 5 \times 10^6)$	

<sup>a</sup>Spectroscopic Galaxy Survey

<sup>b</sup>21-cm Experiment



arXiv:2303.04815, Collaborators: Oliver Philcox, Stephon Alexander, Marc Kamionkowski

# dCS Inflation: Dynamics and Primordial Power Spectra

$$u''(\tau, \mathbf{k}) - \frac{2}{\tau} u'(\tau, \mathbf{k}) + k^2 u(\tau, \mathbf{k}) = 0.$$

$$u''_{\pm}(\tau, \mathbf{k}) - \frac{2}{\tau} u'_{\pm}(\tau, \mathbf{k}) + \left( k^2 \mp \frac{2k\mu}{\tau} \right) u_{\pm}(\tau, \mathbf{k}) = 0.$$

---

$$u(\tau, \mathbf{k}) = \frac{H}{\sqrt{2k^3}} (1 + ik\tau) e^{-ik\tau}$$

$$u_{\pm}(\tau, \mathbf{k}) = \frac{H}{M_{\text{pl}} \sqrt{k^3}} (1 + ik\tau) e^{-ik\tau} \left[ \frac{U(2 \mp i\mu, 4, 2ik\tau)}{U(2, 4, 2ik\tau)} \right] \exp\left(\pm \frac{\pi}{2} \mu\right)$$

---

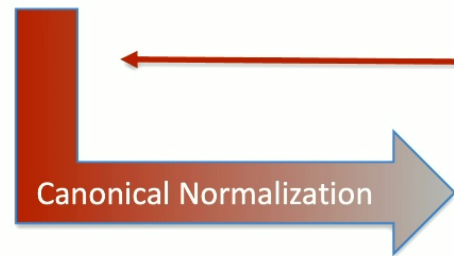
$$\Pi_{\text{circ}} = 0.9 \left( \frac{\varepsilon}{10^{-2}} \right)^{1/2} \left( \frac{H}{10^{14} \text{ GeV}} \right)^2 \left( \frac{10^9 \text{ GeV}}{f} \right)$$



# dCS Inflation: Lepto/Baryogenesis

$$a^4 \frac{\phi(\tau)}{4f} * RR$$

$$\frac{\dot{\phi}}{4f} = \frac{1}{2} \frac{M_{\text{pl}}}{H} \left[ \sqrt{\frac{\epsilon}{2}} \frac{H^2}{M_{\text{pl}} f} \right] = \frac{1}{2} \frac{M_{\text{pl}}^2}{H} \underline{\mu}$$



Canonical Normalization

$$\partial_\mu J_5^\mu = *RR / (384\pi^2) \quad J_5^0 = n_L - n_R$$

$$\sim H\mu(n_L - n_R)$$



# Conclusions/Future Extensions

- 1) Two new constructions of dCS, yielding a large range of decay constants
- 2) First example of massless spin-exchange yielding a parity-violating scalar trispectrum
- 3) Scalar trispectra can probe dCS decay constants at very high values
- 4) Odd to even ratios, in collapsed limit, (i.e. the phase of the trispectrum)  
give direct information about spin
- 5) Potential new probes of baryogenesis

# Conclusions/Future Extensions

- 1) Two new constructions of dCS, yielding a large range of decay constants
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- 4) Odd to even ratios, in collapsed limit, (i.e. the phase of the trispectrum)  
give direct information about spin
- 5) Potential new probes of baryogenesis

## dCS Extensions

- 1) Removal of the dCS Ghost
- 2) Parity-Odd Signal for all configurations
- 3) Where do the trispectra peak?
- 4) CMB signal
- 5) Intermediary dCS constants?  
(Gravi-axion DM/DE)?

## Baryogenesis Extensions

- 1) What type of spin-1 baryogenesis  
can be probed? (What vertex interactions,  
shapes, sensitivities, UV completions etc)
- 2) Other CMB/LSS probes?



# dCS Inflation: Quadratic Action

$${}^*RR = \frac{2\tilde{\epsilon}^{ijk}}{a^4(\tau)} \left[ h''_{il} (\partial_j h_k^l)' + (\partial_m h'_{li}) (\partial_j \partial^l h_k^m + \partial_k \partial^m h_j^l) \right]$$

# UV Completions for Fermion SI

(Scalar) Yukawa

$$\mathcal{L} \supset g_\chi \chi \bar{\Psi} \Psi$$

$$\Lambda \sim m_\chi$$

$$\lambda = (g_\chi / m_\chi)^2$$

$$f = 34 \text{ eV } g_\chi^{-2} \left( \frac{m_\Psi}{10^{-3} \text{ eV}} \right)$$

Gravitational Torsion

$$\mathcal{L} \supset (2/\gamma) R^{\mu\nu} \wedge e_\mu \wedge e_\nu$$

$$\Lambda \sim M_{\text{pl}}$$

$$\lambda = 3\pi / \Lambda_T^2$$

