

Title: Mathematical Physics Lecture (230421)

Speakers: Kevin Costello

Collection: Mathematical Physics - Elective (2022/2023)

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Non-linear Graviton
Self-dual (meaning $C = 0$)
Einstein manifolds



\longleftrightarrow { 3 complex dimensional manifold X
with a fibration $X \rightarrow \mathbb{C}P^1$
and a symplectic form ω on fibres,
valued in $\mathfrak{O}(2)$ }

Additional conditions on X .

1) The fibration $X \rightarrow \mathbb{C}P^1$
must admit a hol. section
 $\sigma: \mathbb{C}P^1 \rightarrow X$

with normal bundle $\mathcal{O}(1) \oplus \mathcal{O}(1)$

2) If we want a real Einstein manifold, X must have an anti-holomorphic involution.

M , the Einstein manifold, is

$$M = \left\{ \begin{array}{l} \mathbb{C}P^1 \text{ in } X \\ \text{deforming the section } \sigma \end{array} \right\}$$

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$$M = \left\{ \mathbb{C}P^1 \text{'s in } X \right. \\ \left. \text{deforming the section } \sigma \right\}$$

This is a 4 parameter space

Normal bundle is $\mathcal{O}(1) \oplus \mathcal{O}(1)$

\exists normal vectors n_1, n_2 on a patch of $\mathbb{C}P^1$
($z \neq \infty$)

which vanish as $z \rightarrow \infty$

$\mathcal{O}(1) \oplus \mathcal{O}(1)$ rank 2

change patches, there is a 2×2 matrix

$\mathcal{O}(1) \oplus \mathcal{O}(1)$ means the matrix is

$$\begin{pmatrix} \tilde{z} & 0 \\ 0 & \tilde{z} \end{pmatrix}$$

patch of \mathbb{CP}^1

$$X$$

$$U$$

$$\mathbb{C}P^1 = \text{Im } \sigma$$

On $\mathbb{C}P^1$, have

$$T\mathbb{C}P^1 \subseteq TX$$

$$\xrightarrow{\sigma} \mathbb{C}P^1$$

$$N = \text{quotient}$$

$$TX / T\mathbb{C}P^1$$

On $z \neq \infty$

n_1, n_2 a basis of normal vectors

$z \neq 0$

\tilde{n}_1, \tilde{n}_2 a basis

$$n_j = \frac{1}{z} \tilde{n}_j = \frac{1}{z} \tilde{n}_j$$

Normal vectors which have no poles.

4 of them

n_1, n_2

$z n_1, n_2$

\Rightarrow we can move our $\mathbb{C}P^1$ in a 4 parameter family of holomorphic $\mathbb{C}P^1$'s

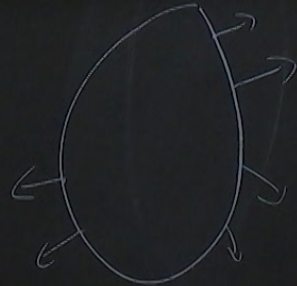
ish as $z \rightarrow \infty$

$k = 2$

there is a 2×2

the matrix is

Picture in 3 real dimensions



$\partial \Omega$

$$n_i = \frac{\partial}{\partial v_i} \text{ on } \mathbb{R}^3$$

Where does the metric come from?

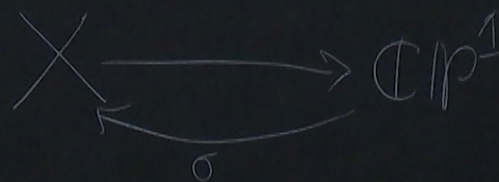
If $\mathbb{C}P^1 \subseteq X$

\Rightarrow it's a point $p \in M$

$T_p M =$ holomorphic sections
of the normal bundle

$n_1, n_2, \bar{z}n_1, \bar{z}n_2$

Want to build \langle, \rangle
inner product



Normal bundle
 $=$ vectors along fibre
of $X \rightarrow \mathbb{C}P^1$

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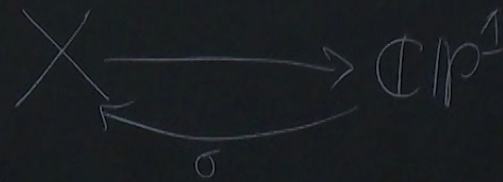
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Want to build \langle, \rangle
inner product.



Normal bundle
= vectors along fibre
of $X \rightarrow \mathbb{C}P^1$

Data on X means there's
a bundle map $\mathbb{R}P^1 \cong \mathcal{O}(2)$
(symplectic form on fibres)

This allows us to define

ε_{ij} an anti-sym pairing
on n_1, n_2

(On flat space had ε_{ij} on $\frac{\partial}{\partial v_a}$)

Inner prod

means there's

$$\Lambda^2 N_0 \cong \mathcal{O}(2)$$

on fibres)

us to define

anti-sym pairing

ad ϵ_{ap} on $\frac{\partial}{\partial v_a}$

Inner product on 4d space with basis

$$n_1, n_2, zn_1, zn_2$$

$$\langle n_i, zn_j \rangle = \omega_{ij}$$

$$\langle zn_i, n_j \rangle = -\omega_{ij}$$

More clever formula

$N_i(z)$ section of normal bundle

$$\langle N_i(z), N_j(z) \rangle = \oint N_i(z) \partial N_j(z) \omega^1 \frac{dz}{z}$$

It is really not obvious that
this solves Einstein equations.

$$H^0(\mathbb{CP}^1, \mathcal{O}(1)) = \{1, z\}$$

\cup
 SL_2

this is fundamental
rep. of SL_2

has natural antisymmetric
inner product.

cut
ms.

N normal bundle

$$= V \otimes \mathcal{O}(1)$$

V symplectic v. space of \dim^2

$$H^0(\mathbb{C}P^1, N)$$

$$= \underbrace{V \otimes H^0(\mathbb{C}P^1, \mathcal{O}(1))}_{\text{Symmetric inner product}}$$

Symmetric inner product.

Role of W_∞

The terminology is
confusing + not consistent.

W_∞ = Lie algebra of
Holomorphic Hamiltonian
symmetries of \mathbb{C}^2

i.e. symmetries like

$$\{h, -\}$$

h is hol. function on \mathbb{C}^2

We can build an X
as above by gluing:

$$X = U \cup V$$

$$U = \mathbb{C}^3 \text{ coords}$$

$$V_{\alpha, z}$$

$$V = \mathbb{C}^3 \text{ coordinates}$$

$$\tilde{V}_{\alpha, z}$$

Given

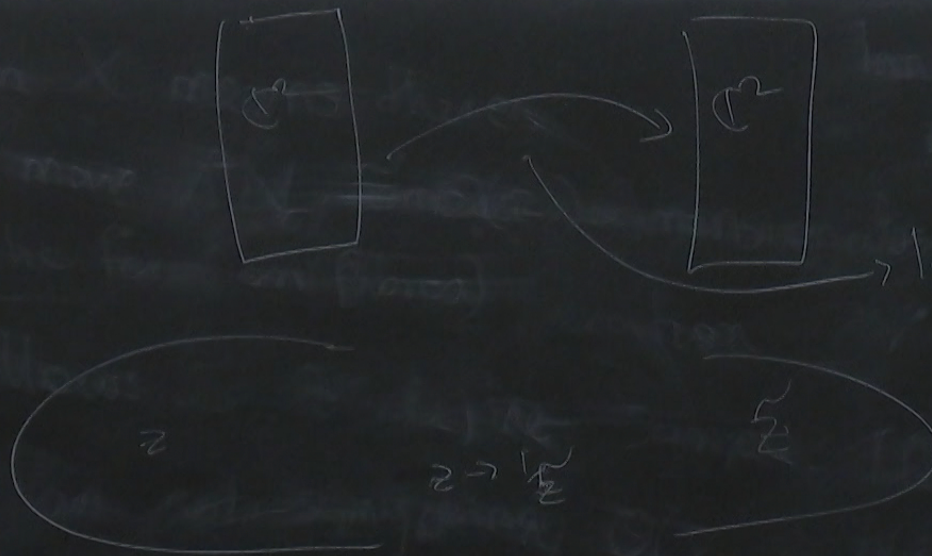
$$h(z) : \mathbb{C}^2 \times \mathbb{C}^* \rightarrow \mathbb{C}$$

holomorphic,

we can identify the coordinates by

$$\tilde{z} = \frac{1}{z}$$

$$\tilde{v}_\alpha = \frac{1}{z} e^{\{h, -\}} v_\alpha$$



Identified by
 a holomorphic transformation
 preserving $\{ \}$ on fibres

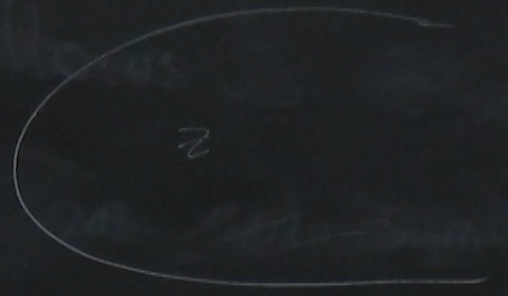
$$Z) : \mathbb{C}^2 \times \mathbb{C}^* \rightarrow \mathbb{C}$$

holomorphic,

identify the coordinates by

$$= \frac{1}{z}$$

$$= \frac{1}{z} e^{\{h, -\}} \quad \forall \alpha \quad \{h, -\} = \epsilon_{\alpha\beta} \frac{\partial h}{\partial v_\alpha} \frac{\partial}{\partial v_\beta}$$



Idea:

SD Gravity has hidden
Symmetry of

Maps $(\mathbb{R}^4, \text{form } (\mathbb{R}^2))$

= Loops into w_∞

$$\left. \begin{array}{l} \frac{\partial}{\partial v_a} \\ \frac{1}{2} \frac{\partial}{\partial v_a} \end{array} \right\} \text{translations}$$

$$v_a \frac{\partial}{\partial v_b} \quad \text{rotations}$$

$$\frac{1}{2} \frac{\partial}{\partial v_a} \quad \text{supertranslations}$$

Also

$$v_a v_b \frac{\partial}{\partial v_c}$$

Also

$$\{ V_\alpha, V_\beta, V_{\partial_t} \}$$

$$V_\nu \frac{\partial}{\partial \nu}$$

not a rotation!

These are part
of the BMS
asymptotic symmetry group

$$V_{(B, \delta)} = \left\{ \frac{\partial}{\partial v} \right\}$$

a rotation!

are part of the BMS asymptotic symmetry group

$$h: \mathbb{P}^2 \times \mathbb{P}^2 \rightarrow \mathbb{C}$$

Gauge trans. on $\mathbb{P}^2 \times \mathbb{P}^2$

a patch on \mathbb{P}^2 which preserve the \mathbb{C} field configuration.

h acts by

$$e^{\{h, -\}} \partial e^{-\{h, -\}} = \bar{\partial}$$

$(\beta, \delta) \sim \}$

$\frac{\partial}{\partial v}$

rotation!

are part

the BMS

asymptotic symmetry group

$$f: \mathbb{S}^2 \times \mathbb{R} \rightarrow \mathbb{C}$$

Gauge trans. on $\mathbb{R}^{1,3}$

a patch on \mathbb{P}^1

which preserve the

field configuration.

f acts by

$$e^{\{f, -\}} \partial e^{-\{f, -\}} = \bar{\partial}$$

$$h \in \mathcal{S}^0(\mathbb{P}^1, \mathbb{C}(z)) \quad \{h, -\}$$