

Title: Quantum Gravity Lecture (230420)

Speakers: Aldo Riello

Collection: Quantum Gravity (2022/2023)

Date: April 20, 2023 - 9:00 AM

URL: <https://pirsa.org/23040023>

Recap

$\underline{L} \in \Omega^{\text{top}, 0}(M \times F)$ Lagrangian density

$$d\underline{L} =: \underline{E}_1 d\varphi^T + d\underline{(\mathbb{1})}$$

\uparrow EOM
 \Downarrow

\uparrow cov. sympl. pot. current

LEL: $\overline{F} = \{\varphi : E_1 \varphi = 0\} \hookrightarrow F$ Euler-Lagrange locus
a.k.a "the shell"

$\underline{\Omega} := d\underline{(\mathbb{1})}$ cov. sympl. pot. current

$$\Sigma \hookrightarrow M$$

$$\underline{\Omega}_\Sigma := \int_\Sigma \underline{\Omega}$$

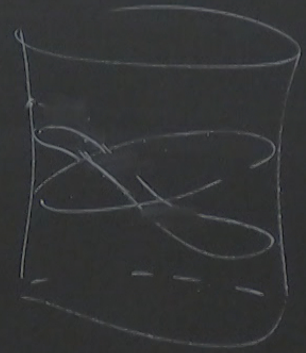
THM

COR



THM $d\Omega \approx 0$
 \uparrow on shell
 (after pullback to $\overline{F} \hookrightarrow \overline{F}$) -

COR: $M = \Sigma \times \mathbb{R}$, $\partial\Sigma = \emptyset \Rightarrow \Sigma$ Cauchy surf
 and $\mathcal{L}_{EL}^* \Omega_\Sigma$ does not
 depend on choice of $\Sigma \hookrightarrow M$ -



$\Rightarrow (\overline{F}, \mathcal{L}_{EL}^* \Omega_\Sigma)$ is a candidate
 (covariant) phase space
 $\begin{cases} \uparrow \\ \text{indep. of choice of } \Sigma \end{cases}$

locus
 shell

M
 \mathbb{R}

SYMMETRIES $\xrightarrow{(N1)}$ Conserved charges

Sym Generators
[?]

$$\rho(\xi) = \int \delta_{\xi} \varphi \frac{\delta \mathcal{L}}{\delta \varphi}$$

off shell

DEF (g, ρ) Lagrangian sym iff $\mathcal{L}_{\rho(\xi)} \mathcal{L} = dR(\xi)$

DEF $\underline{J}(\xi) := \tilde{\rho}(\xi) \underline{\mathcal{L}} - R(\xi)$ Noether current

THM (Noether 1) $d\underline{J}(\xi) = -E_{\xi} \delta_{\xi} \varphi^I \approx 0$

$\partial \Sigma = \phi$, $Q_{\xi}(\Sigma) = \int_{\Sigma} \underline{J}(\xi)$ is indep of $\Sigma \hookrightarrow M$

THM

$\rho(\xi)$ is tangent to $\mathcal{F} \subset \mathcal{F}$

$$D\rho(\xi) E_I \approx 0$$

Generators

"Best case scenario"

PROP: if $R(\xi) = 0$, $D\rho(\xi) \frac{\partial}{\partial \xi} = 0$

then: $\dot{\rho}(\xi) \underline{\Omega} = -d\underline{J}(\xi)$

to $\mathcal{F} \hookrightarrow \mathcal{F}$

ario "

$$p(\vec{z}) = 0, \quad L_{p(\vec{z})} \mathcal{H} = 0$$

$$\Omega = -dJ(\vec{z})$$

Ex particle (1d field th.)

$$M = \mathbb{R} \times \{1pt\}$$

$$\mathcal{F} = C^\infty(M, \mathbb{R}^3) \ni q(t)$$

$$\mathcal{H} = m \dot{\vec{q}} \cdot d\vec{q} = m \dot{q}^i dq^i$$

$$g \text{ rotations } p(\vec{z}) = \int_{\mathbb{R}} \epsilon_{ij}^k \dot{z}^i q^j \frac{\delta}{\delta q^k(t)}$$

$$L_{p(\vec{z})} \mathcal{H} = 0 \quad \checkmark$$

$$J(\vec{z}) = \int_{\mathbb{R}} p(\vec{z}) \mathcal{H} = m \dot{q}^i \delta_{ij} q^j = m \dot{q}^k \underbrace{\epsilon_{ijk}}_{\text{any. man.}} \dot{z}^i q^j$$

THM

$\rho(\xi)$ is tangent to $\mathcal{F} \subset \mathcal{F}$

$$L_{\rho(\xi)} E_I \approx 0$$

then $\underline{J}(\xi)$ cov. sym. generator

Generators

"Best case scenario"

PROP: if $R(\xi) = 0$, $L_{\rho(\xi)} \underline{L} = 0$

then: $\dot{\rho}(\xi) \underline{\Omega} = -d\underline{J}(\xi)$

Ex part

$$M = \mathbb{R} \times \mathbb{S}^1$$

$$\mathcal{F} = C^\infty(M)$$

$$\mathcal{H} = m$$

g rotat

$$L_{\rho(\xi)} \mathcal{H}$$

$$\underline{J}(\xi)$$

Lemma (g, ρ) Logn. sym of L , then

$$\dot{i}_{\rho(\xi)} \underline{\Omega} = -d\underline{J}(\xi) + \underline{r}(\xi)$$

$$\underline{r}(\xi) := \underline{L}_{\rho(\xi)} \textcircled{1} - d\underline{R}(\xi)$$

Thm 1.) OFF-SHELL if either

i) $\underline{r}(\xi) = 0$, or

ii) $\underline{r}(\xi) = d\underline{S}(\xi)$, $\partial \Sigma = \emptyset$

then

$$\dot{i}_{\rho(\xi)} \underline{\Omega}_{\Sigma} = -d\underline{Q}_{\Sigma}(\xi)$$

2) on

then

(test case scenario)

2) ON-SHELL if either

i) $\underline{\Omega}(\Sigma) \approx 0$, or

ii) $\underline{\Omega}(\Sigma) = d\underline{\Omega}(\Sigma)$, $\partial\Sigma = \phi$

← (time-translations
OR
"time" differs in GR)

then

$$\mathbb{P}_\phi(\Sigma) \Omega_\Sigma \approx -\int Q_\Sigma(\Sigma)$$

case scenario

Ex. scalar field + time transl. sym -

$$\underline{L} = \frac{1}{2} \nabla_a \phi \nabla^a \phi \quad \text{on } (M, g_{ab})$$

X is a killing v.f.

$$L_X g_{ab} = \nabla_a X_b + \nabla_b X_a = 0$$

translations
OR
differences in GR

$$\mathbb{D}^a = \nabla^a \phi \, d\phi$$

$$R^a(X) = X^a L = X^a \left(\frac{1}{2} \nabla_c \phi \nabla^c \phi \right)$$

$$J^c(X) = X^b \left(\nabla_b \phi \nabla^c \phi - \frac{1}{2} \delta_b^c \nabla_d \phi \nabla^d \phi \right)$$

$$r^a(X) = X^a \underbrace{\square \phi}_{E} \, d\phi + \nabla_b \left(\underbrace{2 X^{[a} \nabla^{b]} \phi}_{\text{gab}} \, d\phi \right)$$

⇒ case 2ii

off shell obstruction on $\partial \Sigma = \phi$

nonzero only if
 $n \cdot X \neq 0$
i.e. X transverse
to Σ .

$$r_\Sigma(\xi) = \int_\Sigma \sqrt{h} \, (n_a X^a) \, \square \phi \, d\phi$$

of $\nabla \phi$
 $X_b \nabla_b X_a = 0$

Local Sym & Noether 2

• Results so far hold also for local sym's

• In this section we change notation
 \mathfrak{g} = finite dim. (semisimple) Lie alg.

$\mathcal{G} = C^\infty(M, \mathfrak{g}) \leftarrow$ local Lie alg. of sym.

$$\xi, \eta \in \mathcal{G}, \quad [\xi, \eta]_{\mathcal{G}}(x) = [\xi(x), \eta(x)]_{\mathfrak{g}}$$

& Noether 2

for hold also for

ation we change notation
nite dim. (semisimpl) Lie alg.

$(M, g) \leftarrow$ local Lie alg.
of sym.

$$[\xi, \eta]_{\mathfrak{g}}(x) = [\xi(x), \eta(x)]_{\mathfrak{g}}$$

Recall Noether 1 (off shell)

$$dJ(\xi) = -E_I \delta_{\xi} \varphi^I$$

$$\delta_{\xi} \varphi^I = \underbrace{D_{\alpha}^I}_{\text{some } \varphi\text{-dep diff operator}} \xi^{\alpha}$$

some φ -dep diff operator

Ex: YM

$$\delta_{\xi} A_a = \partial_a \xi + [A_a, \xi]$$
$$\delta_{\xi} A_a^{\alpha} = \underbrace{(\delta_{\gamma}^{\alpha} \partial_a + f^{\alpha}_{\beta\gamma} A_a^{\beta})}_{D_{\alpha}^I} \xi^{\gamma}$$

Noether 1 (off shell)

$$\delta S = -E_I \int_{\Sigma} \varphi^I$$

$$p^I = \mathcal{D}_\alpha^I \xi^\alpha$$

some φ -dep diff operator

YM

$$\delta_\xi A_a = \partial_a \xi + [A_a, \xi]$$
$$\delta_\xi A_a^\alpha = \left(\delta_\gamma^\alpha \partial_a + f_{\beta\gamma}^\alpha A_a^\beta \right) \xi^\gamma$$

\mathcal{D}_α^I

$$\partial \Sigma = \phi$$

$$\int_{\Sigma} E_I \mathcal{D}_\alpha^I \xi^\alpha = \text{ibp}$$

adjoint op.

$$\int_{\Sigma} \left(\mathcal{D}_\alpha^{+I} E_I \right) \xi^\alpha$$

HM (Noether 2)

$$\left(\mathcal{D}_\alpha^{+I} E_I \right) = 0$$