

Title: Strong Gravity Lecture - 230302

Speakers: William East

Collection: Strong Gravity (2022/2023)

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URL: <https://pirsa.org/23030042>

$$\Delta = 0$$

$$r_{\pm} = M \pm \sqrt{M^2 - a^2}$$

Kerr Schild Cartesian coordinates

$$ds^2 = -dt^2 + dx^2 + dy^2 + dz^2 + \frac{2Mr^3}{r^4 + a^2 z^2} \left[dt + \frac{r(xdx + ydy)}{a^2 + r^2} + \frac{a(ydx - xdy)}{a^2 + r^2} + \frac{z}{r} dz \right]$$

$$x^2 + y^2 + z^2 = r^2 + a^2 \left(1 - \frac{z^2}{r^2} \right)$$

$$t_{KS} = t_{BL} + 2M \int \frac{r}{\Delta} dr$$

$$\phi_{KS} = \phi_{BL} + a \int \frac{r}{\Delta} dr$$

$$x + iy = (r - ia) e^{i\phi_{KS}}$$

$$z = r \cos \theta$$

$$\Delta = 0$$

$$r_{\pm} = M \pm \sqrt{M^2 - a^2}$$

coordinates

$$z^2 + \frac{2Mr^3}{r^4 + a^2 z^2} \left[dt + \frac{r(xdx + ydy)}{a^2 + r^2} + \frac{a(ydx - xdy)}{a^2 + r^2} + \frac{z}{r} dz \right]^2$$

$$x^2 + y^2 + z^2 = r^2 + a^2 \left(1 - \frac{z^2}{r^2} \right)$$

$$+ 2M \int \frac{r}{\Delta} dr$$

$$+ a \int \frac{r}{\Delta} dr$$

$$x + iy = (r - ia) e^{i\phi_K S}$$

$$z = r \cos \theta$$

$$r = 0, \theta = \frac{\pi}{2}$$

$$x^2 + y^2 = a^2$$

Metric in Boyer Lindquist coordinates

$$ds^2 = - \left(\frac{\Delta - a^2 \sin^2 \theta}{\Sigma} \right) dt^2 - \frac{2a \sin^2 \theta (r^2 + a^2 - \Delta)}{\Sigma} dt d\phi + \left(\frac{(r^2 + a^2)^2 - \Delta a^2 \sin^2 \theta}{\Sigma} \right) d\theta^2 + \frac{\Delta}{\Sigma} dr^2$$

$$\Sigma = r^2 + a^2 \cos^2 \theta$$

$$\Delta = r^2 + a^2 - 2Mr$$

$$\Delta = 0$$

$$r_{\pm} = M \pm \sqrt{M^2 - a^2}$$

Kerr Schild Cartesian coordinates

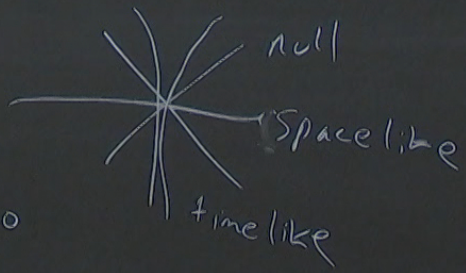
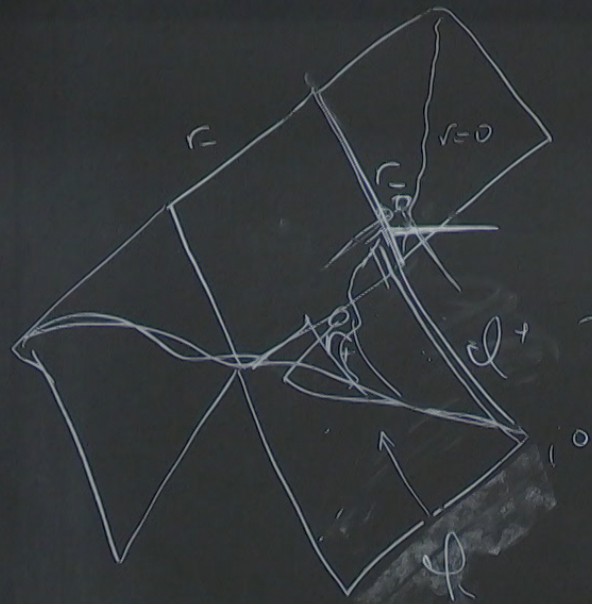
ates

$$\frac{(r^2 + a^2 - \Delta)}{\Sigma} dt d\phi + \left(\frac{(r^2 + a^2)^2 - \Delta a^2 \sin^2 \theta}{\Sigma} \right) \sin^2 \theta d\phi^2 + \frac{\Sigma}{\Delta} dr^2 + \Sigma d\theta^2$$

$$r^2 + a^2 - 2Mr$$

$$r_{\pm} = M \pm \sqrt{M^2 - a^2}$$

$$\left[dt + \frac{r(xdx + ydy)}{a^2 + r^2} + \frac{a(ydx - xdy)}{a^2 + r^2} + \frac{\Sigma}{r} dz \right]^2$$



Ergos

$$r = r_+$$

$$\uparrow_a \quad \uparrow_a$$

$$\uparrow_c \left(\begin{matrix} a \\ + \end{matrix} \right)_a$$

null
 spacelike
 timelike

Ergosphere

$$r = r_+$$

$$\hat{t}^a = \begin{pmatrix} 1 \\ 0 \\ 0 \\ 0 \end{pmatrix}$$

$$\hat{t}^a \hat{t}_a = g_{tt} = \frac{(1 - a^2 \sin^2 \theta)}{\Sigma} \geq 0$$

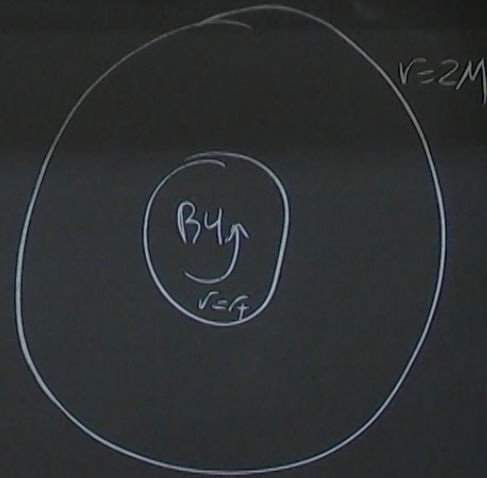
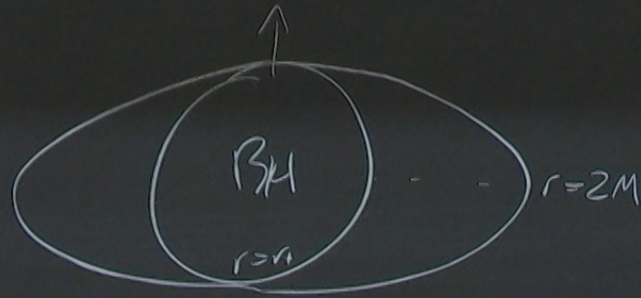
spacelike

Null

$$\hat{t}^a \hat{t}_a = 0 \Rightarrow g_{tt} = 0 \Rightarrow \Delta = a^2 \sin^2 \theta$$

$$\frac{(1 - a^2 \sin^2 \theta)}{\Sigma} \geq 0 \quad \text{spacelike}$$

$$0 \Rightarrow g_{tt} = 0 \Rightarrow \Delta = a^2 \sin^2 \theta \Rightarrow r^2 + a^2 \cos^2 \theta - 2Mr = 0$$



spacelike

$$\Rightarrow \Delta = a^2 \sin^2 \theta \Rightarrow r^2 - 4a^2 \cos^2 \theta - 2Mr = 0$$

$$\begin{aligned}
 -1 = \underline{g_{ab} U^a U^b} &= \underbrace{g_{tt}(U^t)^2 + g_{rr}(U^r)^2 + g_{\theta\theta}(U^\theta)^2 + g_{\phi\phi}(U^\phi)^2}_{\geq 0} \\
 &+ 2g_{t\phi} U^t U^\phi \\
 &\leq -1 \\
 g_{t\phi} < 0 \quad U^t = \frac{dt}{d\tau} > 0, \quad U^\phi = \frac{d\phi}{d\tau} > 0
 \end{aligned}$$

$$g_{tt}(v^t)^2 + g_{rr}(v^r)^2 + g_{\theta\theta}(v^\theta)^2 + g_{\phi\phi}(v^\phi)^2 \geq 0$$

$$+ 2g_{t\phi}v^t v^\phi$$

$$g_{t\phi} < 0 \quad v^t = \frac{dt}{d\tau} > 0 \quad v^\phi = \frac{d\phi}{d\tau} > 0$$

$$\alpha \left(\hat{t}^a + \Omega \hat{\phi}^a \right)$$

$$\hat{\phi}^a \hat{\chi}_a = 0$$

$$g_{\phi t} + g_{\phi\phi}\Omega = 0$$

$$\Omega = \frac{d\phi}{dt} = -\frac{g_{\phi t}}{g_{\phi\phi}}$$

$$r \rightarrow r_+ \quad \Omega = \frac{a}{2r_+} = \Omega_H$$

$$\hat{\chi}^a = \hat{t}^a + \Omega_H \hat{\phi}^a$$

$$\text{For } \bar{a} \ll 1 \quad r_+ = \left(2 - \frac{1}{2}\bar{a}^2\right)M$$

$$r_{\text{H}} = \bar{a}/4M$$

$$\bar{a} = 1 - \varepsilon, \quad \varepsilon \ll 1, \quad r_+ = (1 + \sqrt{2\varepsilon})M$$

$$r_{\text{H}} = \frac{1 - \sqrt{2\varepsilon}}{2M}$$

Geodesics of Kerr

Geodesic Eqn.

$$U^a \nabla_a U^b = 0$$

$$U^a = \frac{dx^a}{d\tau} \quad \tau = \text{proper time}$$

$$U^a U_a = -1$$

$$\nabla_c K_{ab} = 0$$

$$\frac{d^2 x^a}{d\tau^2} + \Gamma^a_{bc} U^b U^c = 0$$

Constants:

$$\begin{aligned} \tilde{E} &= -\dot{\phi}^a U_a = -U_t = -(g_{tt} U^t + g_{t\phi} U^\phi) \\ \tilde{J} &= \dot{\phi}^c U_c = U_\phi = g_{\phi\phi} U^\phi + g_{\phi t} U^t \end{aligned}$$

$$\nabla_a K_{ab} = 0$$

$$U^a \nabla_a (K^b{}_c U^c) = 0$$

$= -1$

$$C = K_{ab} u^a u^b$$

(t, ϕ, θ)

$g_{\mu\nu} u^\mu$

$$= \frac{(r^2 + a^2)^2 - \Delta a^2 \sin^2 \theta}{\Sigma} \sin^2 \theta \frac{d\phi}{dt} - \frac{2M a r \sin^2 \theta}{\Sigma} \frac{dt}{dt}$$

$$U^a \nabla_a (K^b U_b) = 0$$

② Restrict Equatorial orbits $U^\theta = 0, \theta = \frac{\pi}{2}$

$$-1 = g^{ab} U_a U_b = g^{tt} \tilde{E}^2 + g^{rr} (U_r)^2 + g^{\phi\phi} \tilde{J}^2 - 2g^{t\phi} \tilde{E} \tilde{J}$$

$$U_r = g_{rr} U^r = \frac{r^2}{\Delta} \frac{dr}{d\tau}$$

$$g^{tt} = \frac{-((r^2 + a^2)^2 - a^2 \Delta)}{\Delta r^2}, \quad g^{rr} = \frac{\Delta}{r^2}, \quad g^{\phi\phi} = \frac{\Delta}{r^2 \sin^2 \theta}$$

$$= 0, \quad \theta = \frac{\pi}{2}$$

$$r^2 + g^{\phi\phi} \tilde{J}^2 - 2g^{t\phi} \tilde{E} \tilde{J}$$

$$\Delta = r^2 + a^2 - 2Mr$$

$$\frac{(r^2 + a^2)^2 - a^2 \Delta}{\Delta r^2}, \quad g^{rr} = \frac{\Delta}{r^2}, \quad g^{\phi\phi} = \frac{\Delta - a^2}{\Delta r^2}, \quad g^{t\phi} = \frac{-2Ma}{\Delta r}$$

$$(*) \Rightarrow -1 = \frac{\Delta}{r^2} u_r^2 + \frac{a^2 \Delta - (r^2 + a^2)^2}{\Delta r^2} \tilde{E}^2 + \left(\frac{\Delta - a^2}{\Delta r^2} \right) \tilde{J}^2 + \frac{4Ma}{r} \tilde{E} \tilde{J}$$

$$-\Delta^2 u_r^2 = (1 - \tilde{E}^2) r^4 - 2Mr^3 + [a^2(1 - \tilde{E}^2) + \tilde{J}^2] r^2 - 2M(a$$

$$:= V(\tilde{E}, \tilde{J}, r)$$

$$\leq 0$$

Bound $\tilde{E} < 1$

$$+ \frac{a^2 \Delta - (r^2 + a^2)^2}{\Delta r^2} \tilde{E}^2 + \left(\frac{\Delta - a^2}{\Delta r^2} \right) \tilde{J}^2 + \frac{4Ma}{r} \tilde{E} \tilde{J}$$

$$- \tilde{E}^2 r^4 - 2Mr^3 + \left[a^2 (1 - \tilde{E}^2) + \tilde{J}^2 \right] r^2 - 2M(a\tilde{E} - \tilde{J})^2 r$$

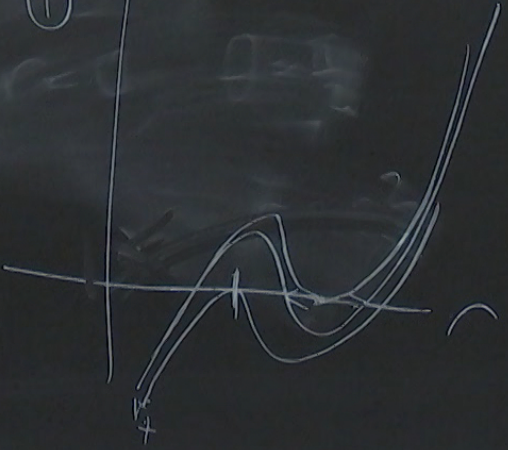
$(\tilde{E}, \tilde{J}, r)$

Bound $\tilde{E} < 1$

$\frac{J}{2M(a\tilde{E} - J)^2 r}$

$V(r_+) = - (2M r_+ \tilde{E})^2 - (aJ)^2 + 4M a r_+ \tilde{E} J$

$= - (2M r_+ \tilde{E} - aJ)^2 \leq 0$

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