

Title: Machine Learning Lecture - 230309

Speakers: Lauren Hayward

Collection: Machine Learning for Many-Body Physics (2022/2023)

Date: March 09, 2023 - 9:00 AM

URL: <https://pirsa.org/23030031>

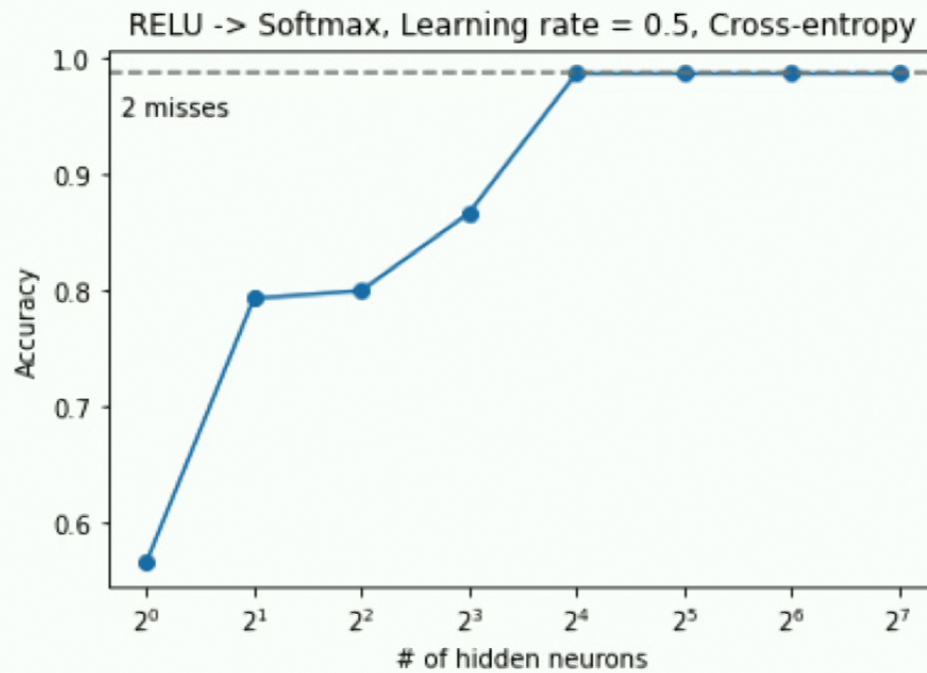
Machine Learning for Many-Body Physics

Lecture 6



Tutorial 2, Problem 2

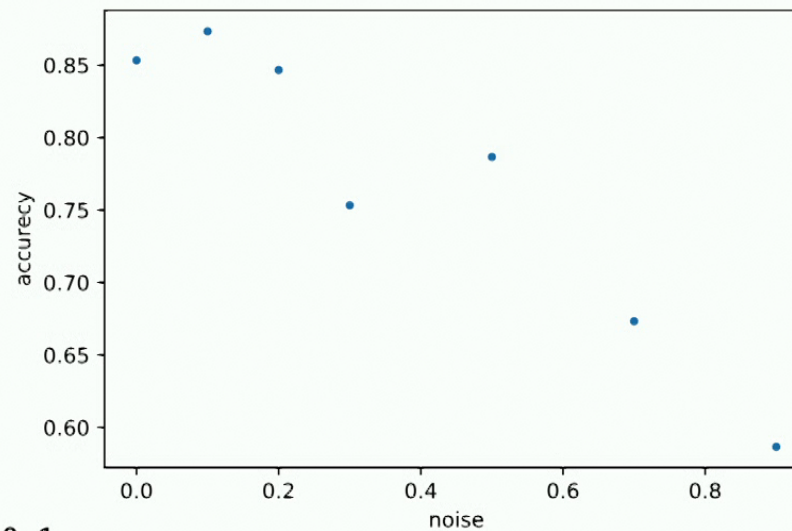
Accuracy vs. the number of neurons in the hidden layer



Tutorial 2, Problem 2

Results from
2022 students

Accuracy vs. the magnitude of noise in the data

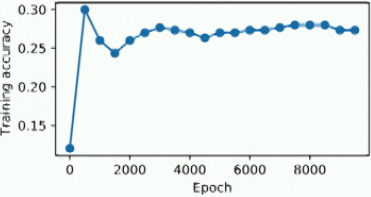
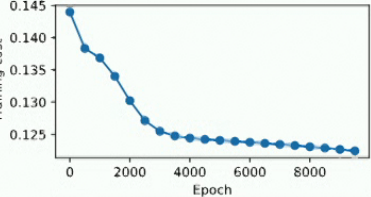
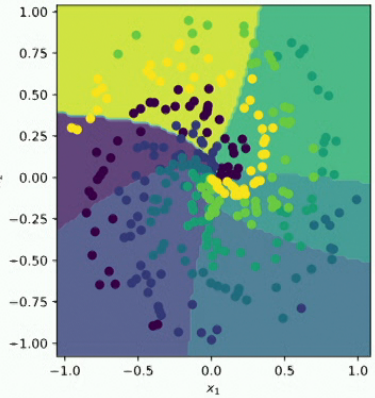
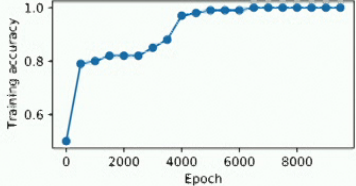
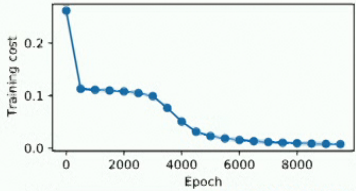
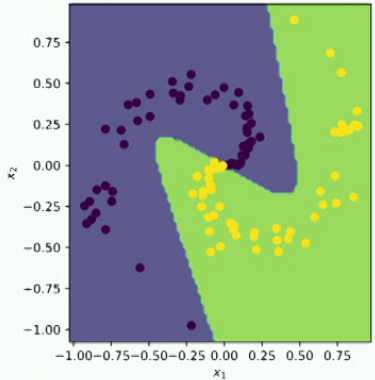


- learning_rate = 0.1
- cost_func = torch.nn.MSELoss()
- activation_function = sigmoid
- number_of_layer and neurons: 1, 4
- N_epochs = 100000

Tutorial 2, Problem 2

Results from 2022 students

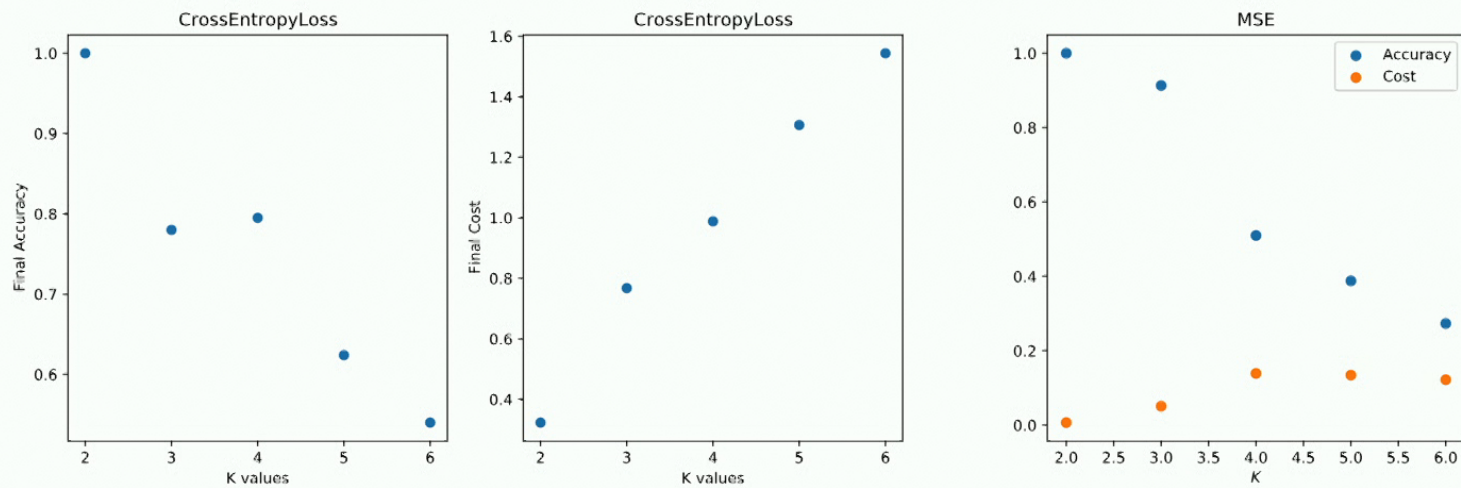
Varying the number of different labels (branches)



Tutorial 2, Problem 2

Results from
2022 students

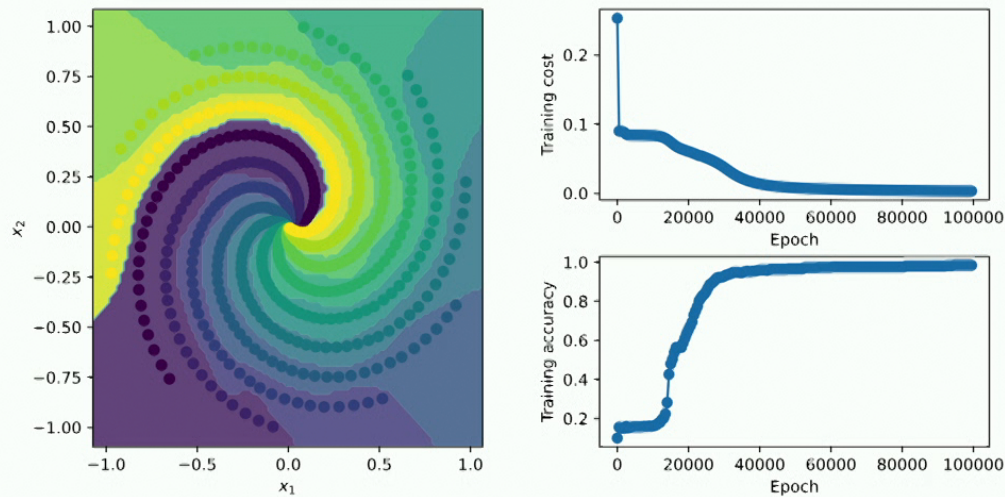
Accuracy vs. the number of different labels (branches)



- learning rate = 1
- 4 neurons in the hidden layer
- first layer uses sigmoid
- outer layer uses softmax

Tutorial 2, Problem 2

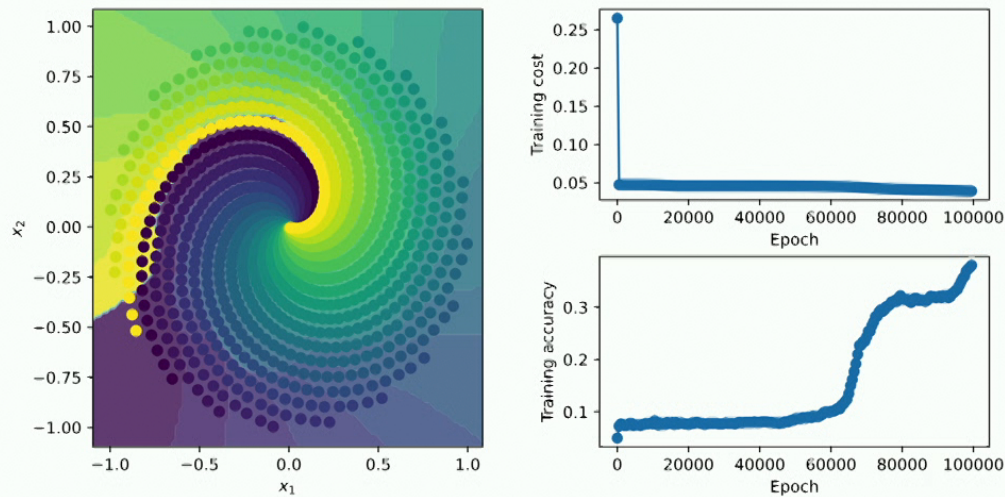
Group 3



- 4 layers with 2, 20, 40, and k neurons
- Adam optimization wasn't working!
- Sigmoid activation function
- MSELoss

Tutorial 2, Problem 2

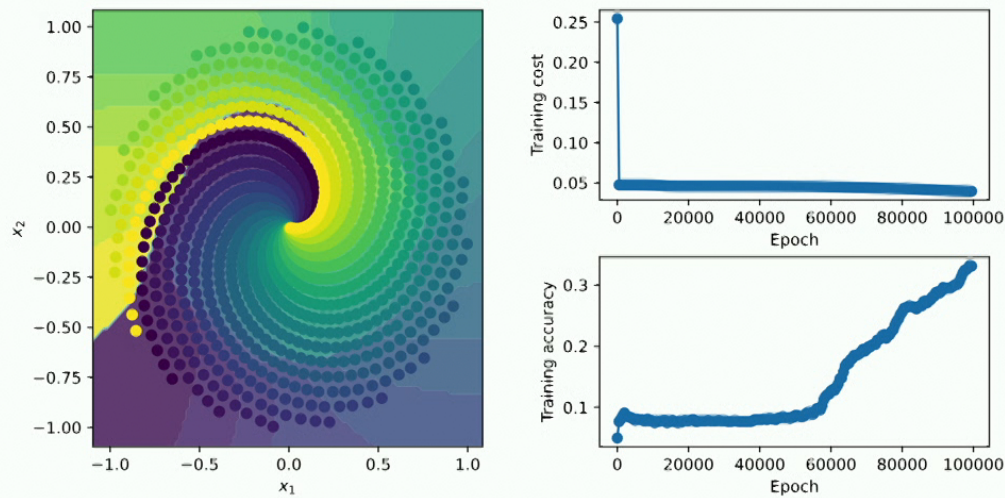
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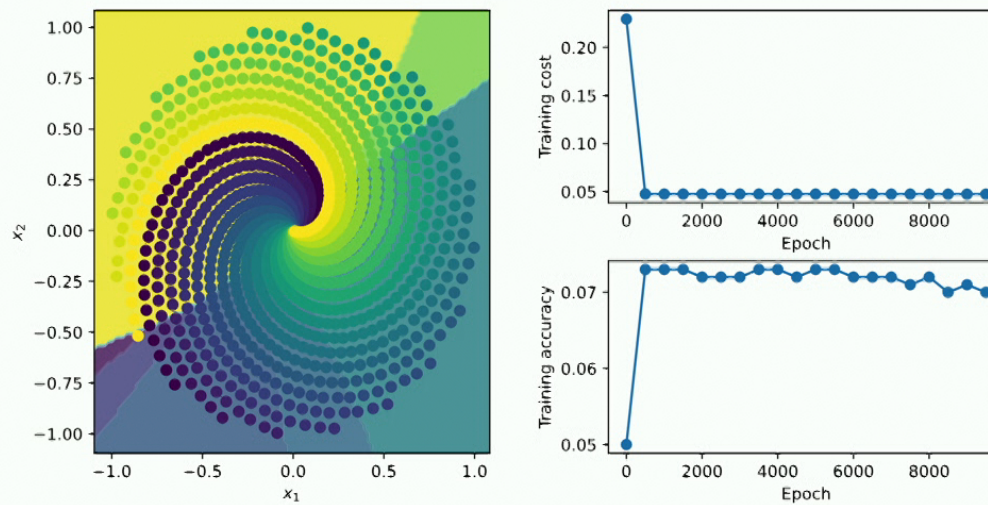
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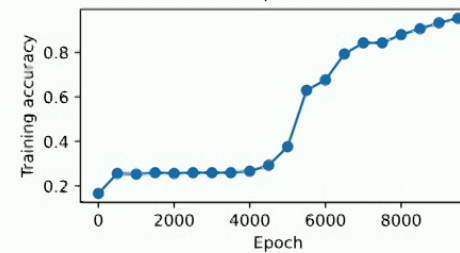
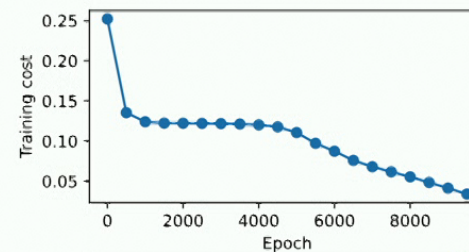
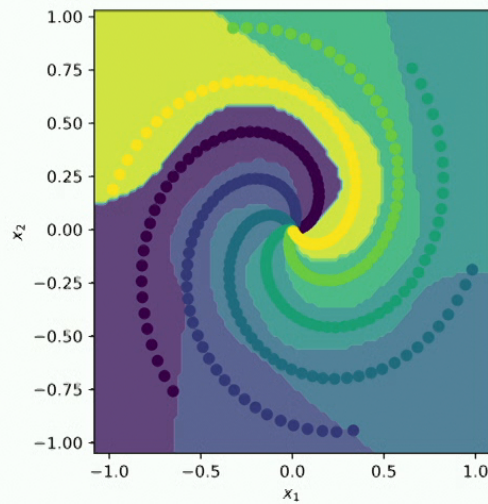
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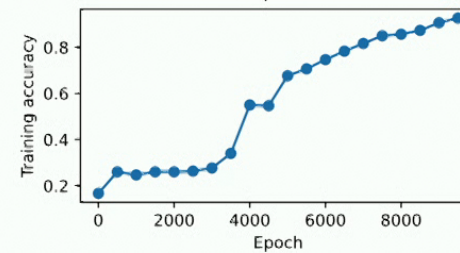
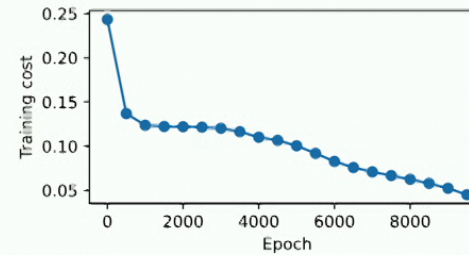
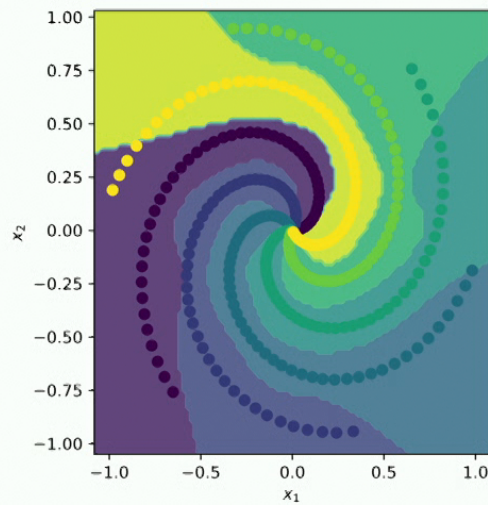
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Tutorial 2, Problem 2

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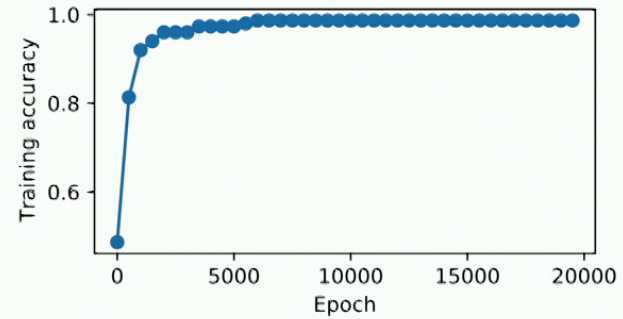
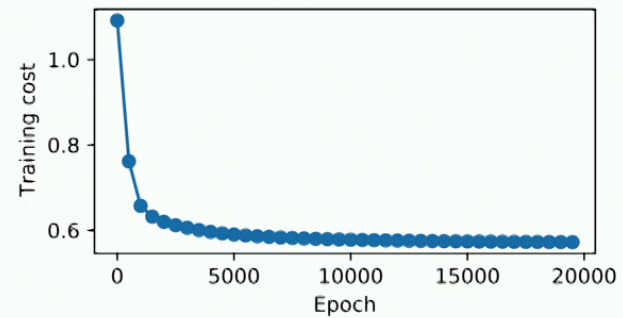
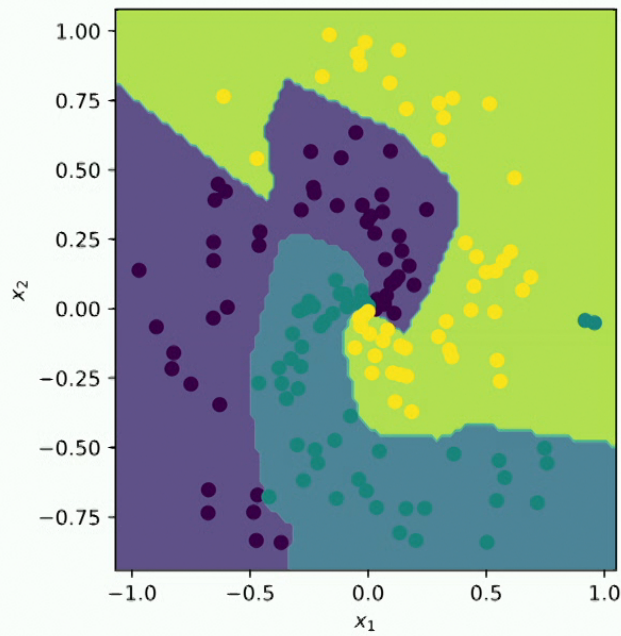


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Overfitting

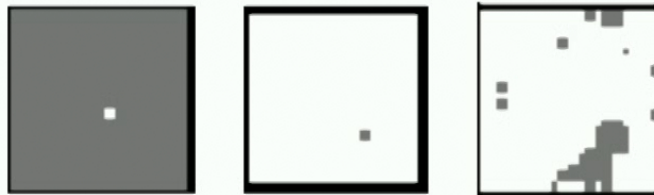
Results from
2022 students

mag_noise = 0.5
hidden_size = 100



Ising model classification

Ferromagnetic phase:



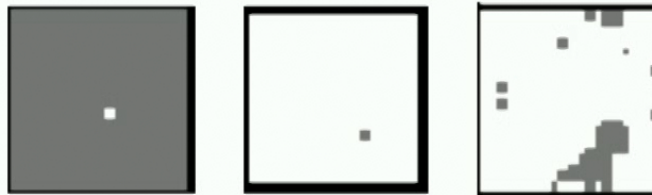
Paramagnetic phase:



Carrasquilla and Melko, arXiv:1605.01735

Ising model classification

Ferromagnetic phase:



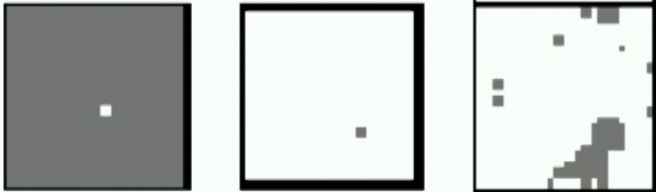
Paramagnetic phase:



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Ising model classification

Ferromagnetic phase:

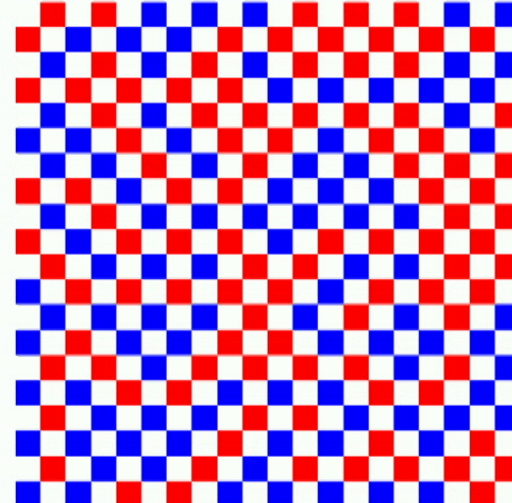
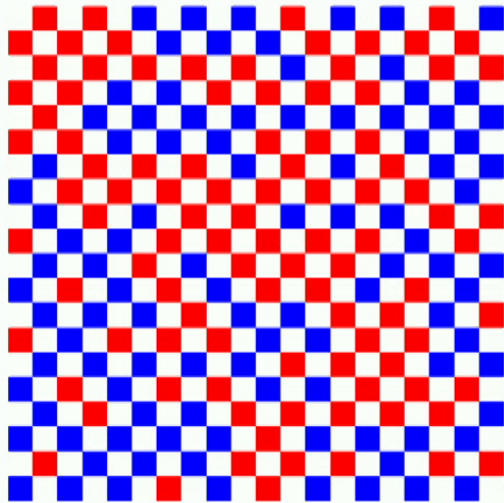


Paramagnetic phase:

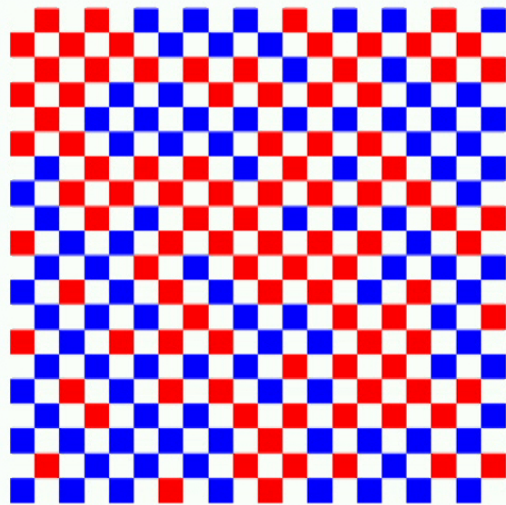


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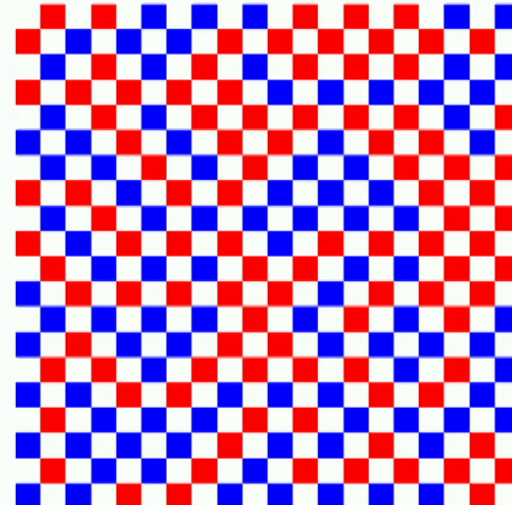
Ising gauge theory classification



Ising gauge theory classification



Random state



Ordered state

Outline for today:

- Using supervised neural networks (NNs) to learn about phases of matter
 - ↳ Ising \mathbb{Z}_2 gauge theory
- Recap of hyperparams. in NNs
- Monte Carlo sampling

Distinguishing phases of the classical Ising \mathbb{Z}_2 gauge theory

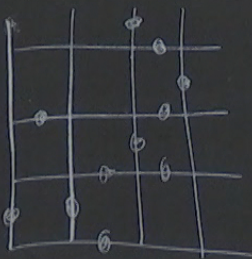
$$H = -J \sum_P \left(\prod_{i \in P} s_i \right) \quad \text{with } s_i = +1 \text{ or } -1$$

The d.o.f. are on the bonds of the lattice

On a torus:



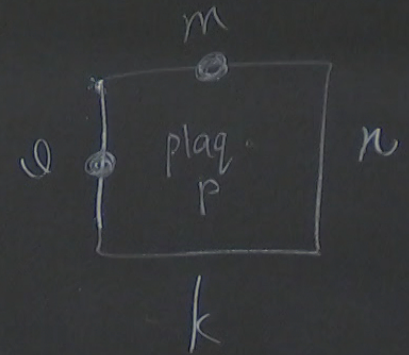
with
periodic
BC



$$s_i = +1 \text{ or } -1$$

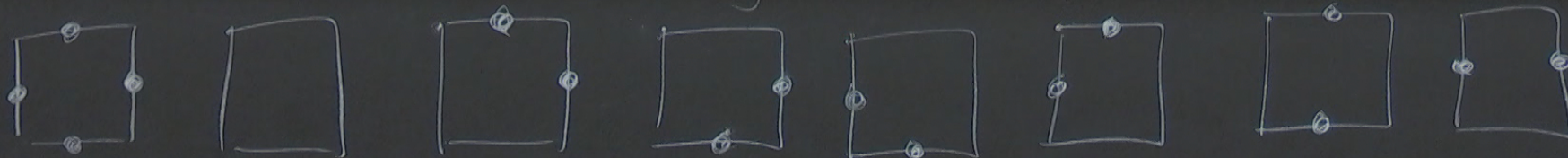
$$\left(\begin{array}{c} | \\ \bullet \\ | \end{array} \text{ or } \begin{array}{c} | \\ | \\ | \end{array} \right)$$

Consider on plaquette p

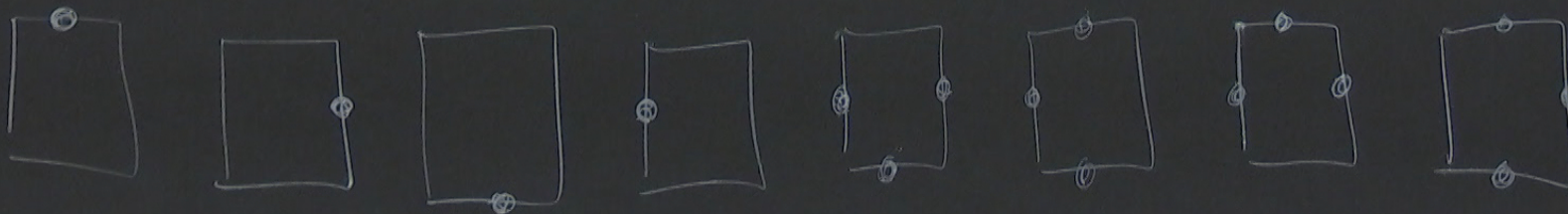


$$\prod_{i \in p} S_i = S_k S_l S_m S_n$$

At low T , there are 8 energetically-favoured plaquettes:



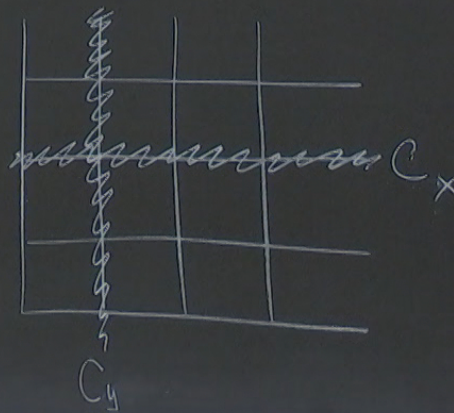
And 8 "excited" plaquettes:



There is topological order exactly at $T=0$ in the thermodynamic limit and no order for $T>0$

We can use the topological Wilson loop W_C as a measure of this order

$$W_C = \prod_{l \in C} S_l = \pm 1$$

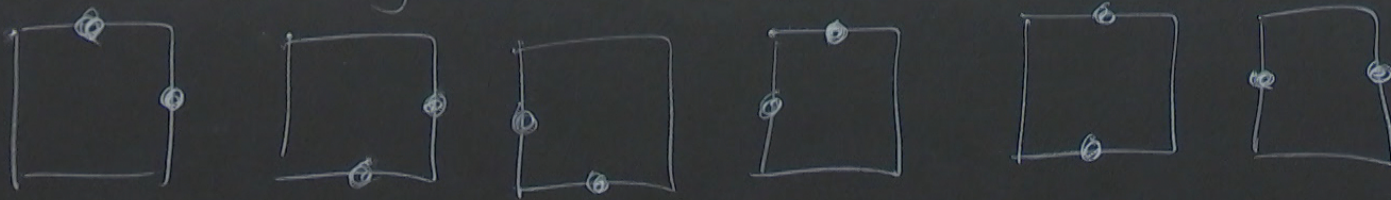


In the topologically-ordered groundstate, W_{C_x} and W_{C_y} are constant for all choices of C_x and C_y

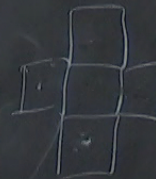
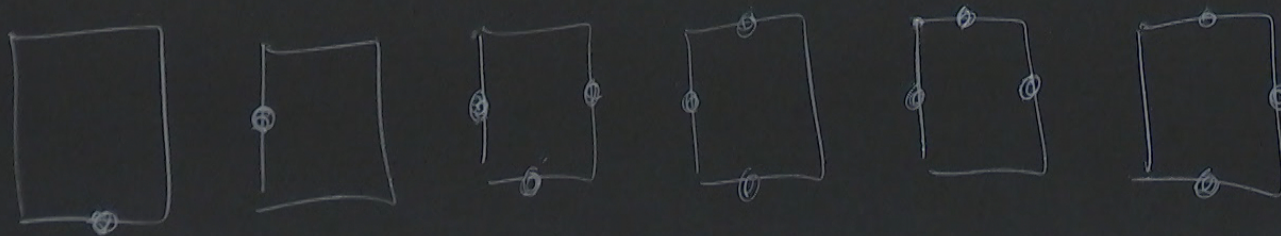
Four possibilities:

W_{C_x}	W_{C_y}
+1	+1
+1	-1
-1	+1
-1	-1

are 8 energetically-favoured plaquettes:



plaquettes:



Hyperparams for feedforward NNs

- # of layers (how "deep" the NN is)
- # of neurons in each hidden layer
- learning rate value, whether it decays with time
- activation function
- cost func.
- training algorithm (GD, SGD, Adam, ...)
- hyperparams. within training alg.: momentum γ , mini-batch size $|B|$, learning rate

- how much regularization is added to the cost (param. λ) and what kind (L1, L2, ...)
- masking prob. if using dropout
- how we partition into training, validation, testing
- whether or not we add noise to our data, what kind and how much
- how the weights and biases are initialized
-
-

Monte Carlo (MC) Sampling (Reference: Newman & Barkema)

(Used to generate data samples for HW1)

Consider a classical system where we want to calculate expectation values of quantities Q

$$\langle Q \rangle = \sum_{\mu} Q_{\mu} P_{\mu}$$

For Boltzmann dist.:

$$P_{\mu} = \frac{1}{Z} e^{-\beta E_{\mu}}, \quad Z = \sum_{\mu} e^{-\beta E_{\mu}}$$

• E_μ is the energy of config μ

• $\beta = \frac{1}{k_B T}$

The sums \sum_μ are over huge # of terms

(2^N terms for classical Ising). Analytically intractable

Idea of MC: estimate $\langle Q \rangle$ from M

States $\{\mu_1, \mu_2, \dots, \mu_M\}$, $M \ll 2^N$

In particular, if each state μ_i is selected from the set of all states according to the desired dist. p_{μ_i} , then

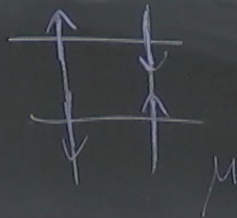
$$\langle Q \rangle \approx Q_M = \frac{1}{M} \sum_{i=1}^M Q_{\mu_i} \quad \left(\begin{array}{l} \text{makes use of} \\ \text{"importance sampling"} \end{array} \right)$$

To implement this sampling, we can use
 Markov chain MC (MCMC)

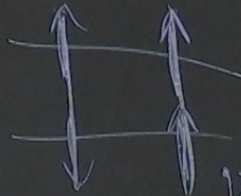
↳ use the current state μ to propose a move
 new related state ν , and then decide to
 accept or reject proposal

eg) single spin flip

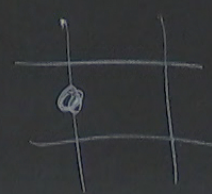
Ising
 model



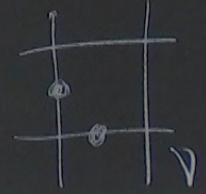
propose
 ~~~~~>



Ising gauge theory:



propose  
 ~~~~~>



In principle, any proposal is allowed Z_2 gauge theory

The prob. of moving from μ to ν is the transition prob. $T(\mu \rightarrow \nu)$, with $\sum_{\nu} T(\mu \rightarrow \nu) = 1$

We can choose any set of transition probs. $\{T(\mu \rightarrow \nu)\}$
as long as we satisfy two conditions:

- ① Ergodicity: there is a non-zero prob. of transitioning (eventually) from μ to ν

$$\mu \rightarrow \dots \rightarrow \nu$$

② Detailed balance (DB)

For the Boltzmann dist.:

$$\frac{T(\mu \rightarrow \nu)}{T(\nu \rightarrow \mu)} = e^{-\beta(E_\nu - E_\mu)}$$