

Title: Quantum Fields and Strings Lecture - 230308

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Collection: Quantum Fields and Strings (2022/2023)

Date: March 08, 2023 - 10:15 AM

URL: <https://pirsa.org/23030018>

Why CFTs?

- Asymptotic IR limit of QFT } TOFT
- slogan: a QFT is definition of a CFT
- Directly relevant for nature

- Thermal phase transitions

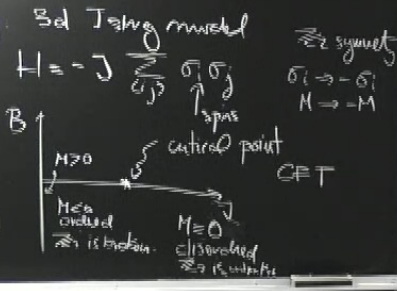
- Quantum phase transitions

- Quantum gravity is asymptotically AdS

= string theory

2d CFT

λ such that $\Delta(\lambda) = 0$
 $\Rightarrow$ CFT



- Thermal phase transitions
 - Quantum phase transitions
 - Quantum gravity in asymptotically AdS
 - string theory

$H = -J \sum_{\langle i,j \rangle} \sigma_i \sigma_j$
 $\sigma_i \rightarrow -\sigma_i$
 $H \rightarrow -H$
 $\mathbb{Z}_2$ symmetry

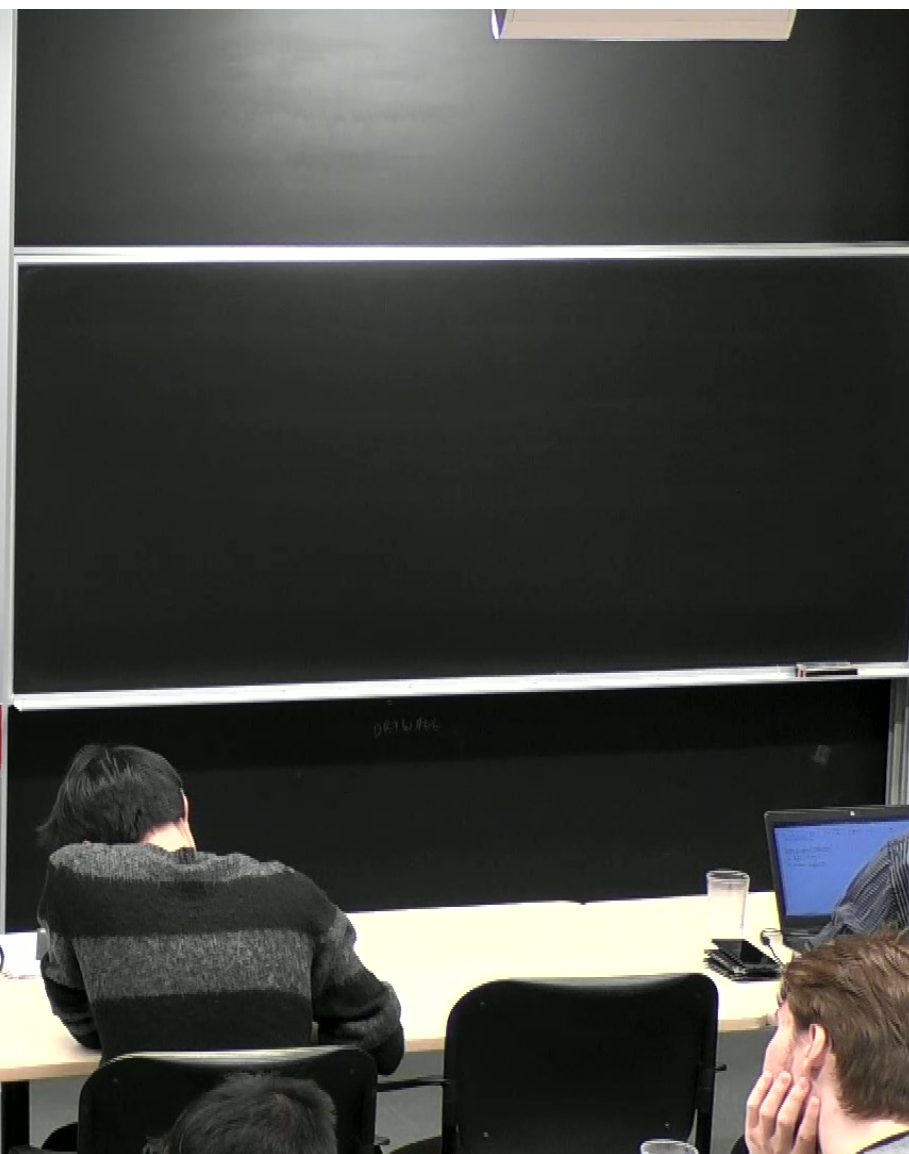
B
 $M > 0$
 $M = 0$
 $M < 0$
 critical point
 CFT

λ^* such that $\Delta(\lambda^*) = 0$
 $\Delta(\lambda) \rightarrow \Delta_0$
 $E_1(\lambda)$
 E_0
 2d CFT
 CFT

Kinematics
 Dynamics

1. Geometry
 2. Realize symmetry on quantum fields

- Geometry of conformal transformations
 - Define a Lie group G
 - Continuous transformations (vs discrete)



- Thermal phase transitions
 - Quantum phase transitions
 - Quantum gravity in asymptotically AdS
 - string theory

$H = -J \sum_{\langle i,j \rangle} \sigma_i \sigma_j$ ↑ spins
 $\mathbb{Z}_2$ symmetry
 $\sigma_i \rightarrow -\sigma_i$
 $H \rightarrow -H$

$\Delta(X^*) = 0$
 $\Rightarrow$ CFT

$\Delta(\lambda)$
 $E_1(\lambda)$
 E_0

$M > 0$
 $M = 0$
 $M < 0$

critical point
 CFT

Dynamics (2. Realize symmetry on quantum fields)

- Geometry of conformal transformations
 $g(x)g(p) = g(f(x,p))$
 $\alpha = 1 + \alpha \gamma^a \gamma_a$
 $\beta =$

- Define a Lie group G .
 $[T^a, T^b] = i f^{ab}_c T^c$

- Continuous transformations (vs discrete)
 - All transformations are completely determined by a Lie algebra

Conformal Transform

DETAIL

- Thermal phase
 - Quantum phase transitions
 - Quantum gravity in asymptotically AdS
 - string theory

$H = - \int \frac{1}{4} \sum_{i,j} \sigma_i \sigma_j$ Ising model
 $\sigma_i \rightarrow -\sigma_i$ $\mathbb{Z}_2$ symmetry
 $H \rightarrow -H$

B
 $M > 0$ critical point
 $M = 0$ CFT
 $\mathbb{Z}_2$ is broken CFT
 $\mathbb{Z}_2$ is unbroken CFT

λ^* such that $\Delta(\lambda^*) = 0$
 $\Delta(\lambda) \rightarrow \Delta_0$
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 $\Delta(\lambda) \rightarrow \Delta_0$

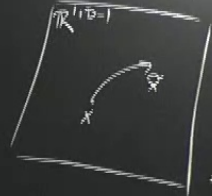
Dynamics (2. Realize symmetry on quantum fields)

- Geometry of conformal transformations
 - Define a Lie group G .
 - Continuous transformations (vs discrete)
 - All transformations are completely determined by a

$g(x)g(\beta) = (f(\alpha/\beta))$
 $\alpha = 1 +$
 $\beta =$
 $[f] = i f \frac{\partial h}{\partial c} T^c$

Conformal Transformation / Minkowski (1, D-1)

1. $x^\mu \rightarrow \hat{x}^\mu(x)$ spacetime transformation that preserve angles & lightcones



$ds^2 \rightarrow e^{2\sigma(x)} ds^2$ $\eta_{\mu\nu}$: Minkowski metric
 $(- + \dots +)$
 $\eta_{\rho\sigma} d\hat{x}^\rho d\hat{x}^\sigma = e^{2\sigma(x)} \eta_{\mu\nu} dx^\mu dx^\nu$
 $\Rightarrow \eta_{\rho\sigma} \frac{\partial \hat{x}^\rho}{\partial x^\mu} \frac{\partial \hat{x}^\sigma}{\partial x^\nu} = e^{2\sigma(x)} \eta_{\mu\nu}$
 $R^\rho_\mu = \frac{\partial \hat{x}^\rho}{\partial x^\mu} e^{-\sigma(x)}$
 $R^\rho_\mu R^\mu_\nu = \delta^\rho_\nu$

- Thermal phase
 - Quantum phase transitions
 - Quantum gravity in asymptotically AdS
 - string theory

2d CFT $\Rightarrow$ CFT

λ^* such that $\Delta(\lambda^*) = 0$
 $\Rightarrow$ CFT

$H = -J \sum_{\langle ij \rangle} \sigma_i^z \sigma_j^z$
 $\sigma_i^z \rightarrow -\sigma_i^z$
 $H \rightarrow -H$

B
 $M > 0$
 $M = 0$
 critical point
 CFT

Dynamics (2. Realize symmetry on quantum fields)

- Geometry of conformal transformations
 $g(x)g(p) = g(f(\alpha, \beta))$
 $\alpha = 1 + \alpha^a T_a$
 $\beta =$

- Define a Lie group G .
 $[T^a, T^b] = i f_{ab}^c T^c$

- Continuous transformations (vs discrete)
 - All transformations are completely determined by a Lie algebra

$\int d\tilde{x}^\mu d\tilde{x}^\nu = e^{2\sigma(x)} \eta_{\mu\nu} dx^\mu dx^\nu$
 $R^\mu_\nu = \frac{\partial \tilde{x}^\mu}{\partial x^\nu} e^{-\sigma(x)} \Rightarrow \eta_{\mu\nu} \frac{\partial \tilde{x}^\mu}{\partial x^\alpha} \frac{\partial \tilde{x}^\nu}{\partial x^\beta} = e^{2\sigma(x)} \eta_{\alpha\beta}$
 $R^T \eta R = \eta$

$x^\mu \rightarrow \tilde{x}^\mu = x^\mu + \xi^\mu(x)$
 ξ^μ small
 spacetime vector field. $\xi = \xi^\mu \partial_\mu$

$\eta_{\mu\nu} (\delta^\mu_\alpha + \partial_\alpha \xi^\mu) (\delta^\nu_\beta + \partial_\beta \xi^\nu) = \left| \frac{\partial \tilde{x}}{\partial x} \right|^{2/D} \eta_{\mu\nu}$
 $\parallel \frac{2}{D} \partial \cdot \xi$

$$R^{\rho}_{\mu} = \frac{\partial \tilde{x}^{\rho}}{\partial x^{\mu}} e^{-\sigma(x)} \eta_{\rho\sigma} \frac{\partial x^{\sigma}}{\partial x^{\mu}} = e^{-\sigma(x)} \eta_{\mu\nu}$$

$$\boxed{R^T \eta R = \eta}$$

$$x^M \rightarrow \tilde{x}^M = x^M + \xi^M(x)$$

↑ small
spacetime vector field.

$$\xi = \xi^M \partial_M$$

$$\eta_{\rho\sigma} (\delta^{\rho}_{\mu} + \partial_{\mu} \xi^{\rho}) (\delta^{\sigma}_{\nu} + \partial_{\nu} \xi^{\sigma}) = \left| \frac{\partial \tilde{x}}{\partial x} \right| \eta_{\mu\nu}$$

$$\parallel \frac{2}{1 + \frac{2}{D} \partial \cdot \xi}$$

$$0(\xi) \quad \boxed{\partial_{\mu} \xi_{\nu} + \partial_{\nu} \xi_{\mu} = \frac{2}{D} \eta_{\mu\nu} \partial \cdot \xi}$$

conformal Killing vector equation

$= 0$ isometry

$$\text{cond} \quad \nabla_{\mu} \xi_{\nu} + \nabla_{\nu} \xi_{\mu} = \frac{2}{D} g_{\mu\nu} (\nabla \cdot \xi)$$

$$\partial_{\rho} \partial_{\rho} \xi_{\nu} = 0$$

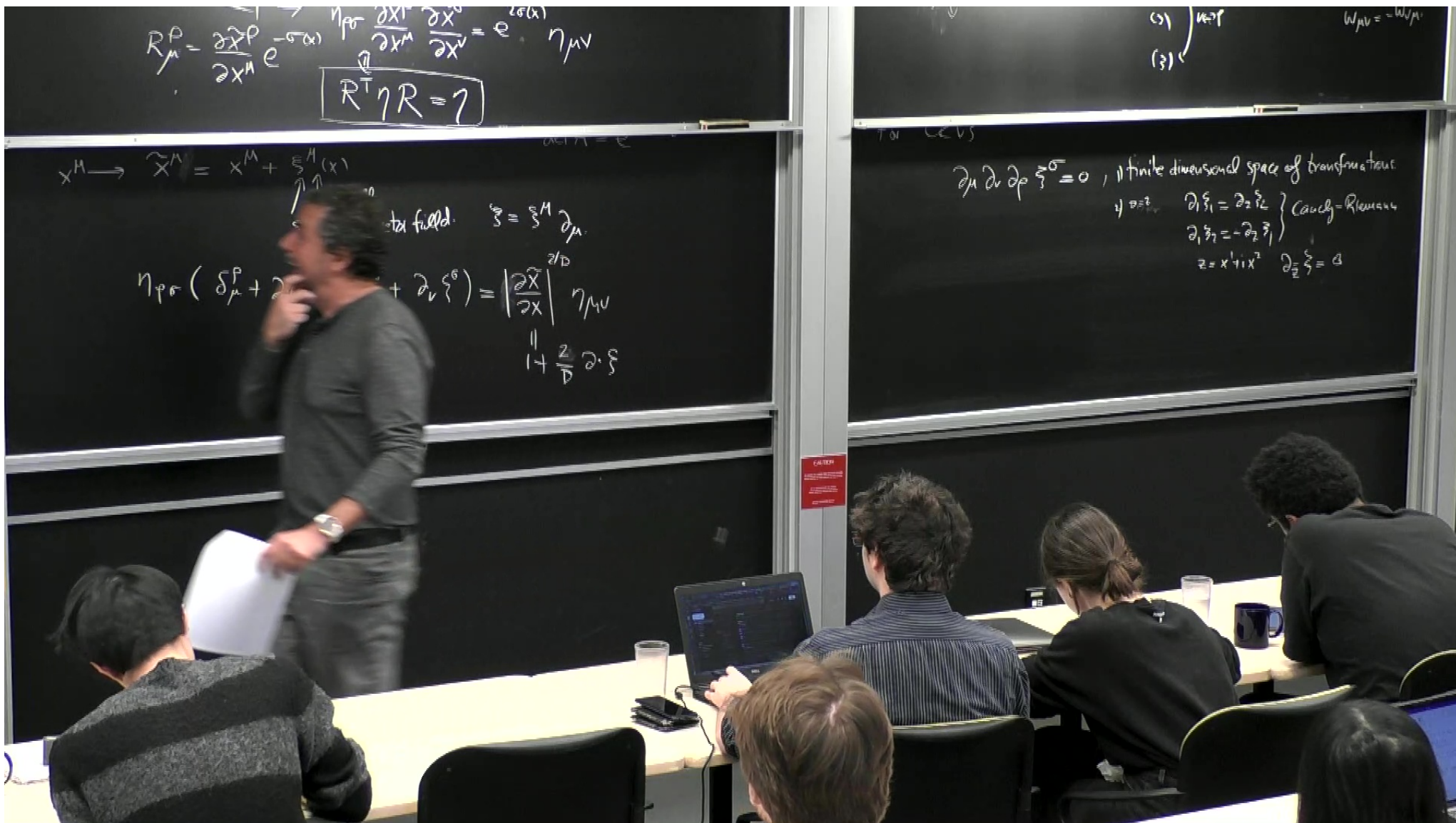
$$1 \text{ Poincaré } \quad \partial_{\mu} \xi_{\nu} + \partial_{\nu} \xi_{\mu} = 0$$

$$(1) + (2) - (3)$$

$$\partial_{\mu} \xi_{\nu} + \partial_{\nu} \xi_{\mu} = 0$$

$$\xi^M = a^M + \omega^M_{\nu} x^{\nu}$$

$$\omega_{\mu\nu} = -\omega_{\nu\mu}$$



$$\xi^M = a^M + \omega^\mu_\nu x^\nu + \lambda x^M + b^M x^2 - 2x^M b \cdot x$$

↓
 P_μ
 D

↓
 $M_{\mu\nu}$
 $\frac{D(D-1)}{2}$

↓
 $\hat{D}$ (dilatation)
1

↓
 K_μ (special conformal)
 D

$\frac{(D+2)(D+1)}{2}$ is
dimension of conformal
algebra.

$$\xi(1), \xi(2)$$

$$[\xi(1), \xi(2)] = \xi(3)$$

$$\xi = \xi^\mu \partial_\mu$$

$$[M^{\mu\nu}, M^{\rho\sigma}] = i (\eta^{\mu\rho} M^{\nu\sigma} - \eta^{\nu\rho} M^{\mu\sigma} - \eta^{\mu\sigma} M^{\nu\rho} + \eta^{\nu\sigma} M^{\mu\rho})$$

$$[M^{\mu\nu}, P^\rho] = \eta^{\mu\rho} P^\nu - \eta^{\nu\rho} P^\mu$$

$$[M^{\mu\nu}, K^\rho] = \eta^{\mu\rho} K^\nu - \eta^{\nu\rho} K^\mu$$

$$z=0$$

$$[D, P^\mu] = -i P^\mu$$

$x^\nu \partial_\nu$ " ∂_μ

$$[D, K^\mu] = i K^\mu$$

$$[P^\mu, K^\nu] = -2i(\eta^{\mu\nu} D + M^{\mu\nu})$$

$$X^M = (X^{\mu}, X^D, X^{D+1}) \quad M=0,1,2,\dots,D-1,D,D+1$$

$$[L_{MN}, L_{PQ}] = i \hat{\eta}_{MP} L_{NQ} + \dots$$

Minkowski:

$$L_{MN} = \begin{cases} L_{\mu\nu} = M_{\mu\nu} \\ L_{DD+1} = D \\ L_{\mu D} = \frac{1}{2} (P_{\mu} + K_{\mu}) \\ L_{\mu D+1} = \frac{1}{2} (P_{\mu} - K_{\mu}) \end{cases}$$

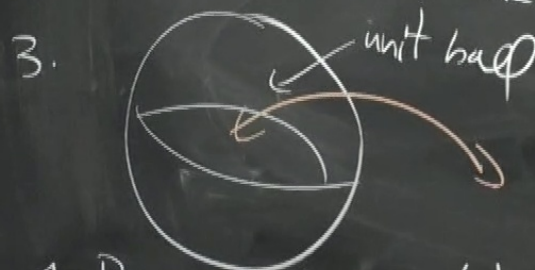
$$\begin{aligned}
 & X^M = (X^{\mu}, X^D, X^{D+1}) \quad \mu = 0, 1, 2, \dots, (D-1), D, D+1 \\
 & \quad \uparrow \\
 & \quad \text{Minkowski} \\
 & L_{MN} = \begin{cases} L_{\mu\nu} = M_{\mu\nu} \\ L_{DD+1} = D \\ L_{\mu D} = \frac{1}{2} (P_{\mu} + K_{\mu}) \\ L_{\mu D+1} = \frac{1}{2} (P_{\mu} - K_{\mu}) \end{cases} \\
 & [L_{MN}, L_{PQ}] = i \hat{H}_{MP} L_{NQ} + \dots \\
 & \hat{H} = \begin{pmatrix} - & + & & \\ & + & + & \\ & & + & + \\ & & & + & - \end{pmatrix} \\
 & \text{conformal algebra in } D \text{ dim is } SO(2, D) \\
 & \quad \uparrow \text{CFT}_D \text{ and } AdS_{D+1}
 \end{aligned}$$

I
↑
inversion

$$x^M \rightarrow x^M / x^2$$

1. check it is conformal $e^{2\sigma(x)} = \frac{1}{(x^2)^2}$

2. $\mathbb{Z}_2$ transformation $I^2 = 1$



4. Reverses the orientation of spacetime.

4. Reverses the orientation of spacetime.

$$\Lambda = \Lambda_{\mu}$$

$$\Lambda = \begin{bmatrix} \end{bmatrix}$$

Finite transformation for k_{μ} :

$$x^{\mu} \rightarrow \frac{x^{\mu} + a^{\mu} x^2}{1 + a^2 x^2 + 2a \cdot x}$$

$$\approx x^{\mu} + a^{\mu} x^2 - 2a \cdot x$$

Discrete conformal

I
↑
inversion

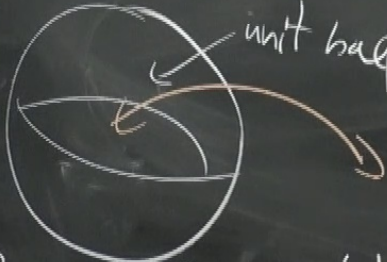
$$x^M \rightarrow x^M / x^2$$

Group of conformal transformations
 $SO(2,D)$

$P \equiv I \wedge$
↑
connected component

1. check it is conformal $e^{2\sigma(x)} = \frac{1}{(x^2)^2}$

2. $\mathbb{Z}_2$ transformation $I^2 = 1$

3.  $I(dx^0 \wedge dx^1 \wedge \dots \wedge dx^{D-1})$
 $= - (dx^0 \wedge \dots \wedge dx^{D-1})$

4. Reverses the orientation of spacetime.