Title: Learning to predict arbitrary quantum processes

Speakers: Hsin-Yuan Huang

Series: Perimeter Institute Quantum Discussions

Date: February 02, 2023 - 11:00 AM

URL: https://pirsa.org/23020031

Abstract: We present an efficient machine learning (ML) algorithm for predicting any unknown quantum process over n qubits. For a wide range of distributions D on arbitrary n-qubit states, we show that this ML algorithm can learn to predict any local property of the output from the unknown process, with a small average error over input states drawn from D. The ML algorithm is computationally efficient even when the unknown process is a quantum circuit with exponentially many gates. Our algorithm combines efficient procedures for learning properties of an unknown state and for learning a low-degree approximation to an unknown observable. The analysis hinges on proving new norm inequalities, including a quantum analogue of the classical Bohnenblust-Hille inequality, which we derive by giving an improved algorithm for optimizing local Hamiltonians. Overall, our results highlight the potential for ML models to predict the output of complex quantum dynamics much faster than the time needed to run the process itself.

Zoom link: https://pitp.zoom.us/j/93857777354?pwd=c044blZuQVhLS200ME4vN25uaGJudz09

Pirsa: 23020031 Page 1/71

Learning to predict arbitrary quantum processes

Credit: DALL-E



Presenter: Hsin-Yuan Huang (Robert) Joint work with Sitan Chen and John Preskill



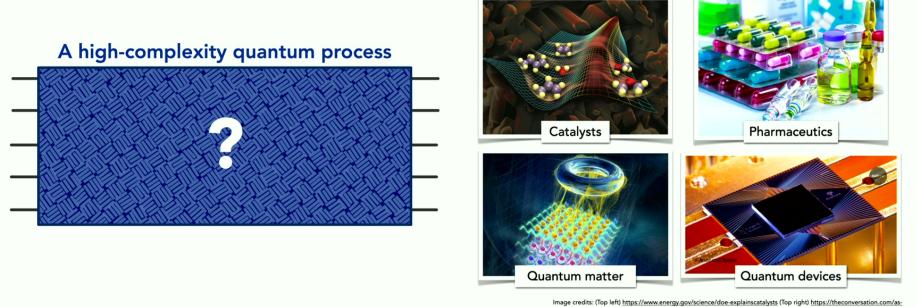




Pirsa: 23020031 Page 2/71

Motivation

- We have seen substantial recent progress on efficiently learning to predict quantum states.
- Are there efficient algorithms for learning to predict quantum circuits / processes?



pharmaceutical-use-continues-to-rise-side-effects-are-becoming-a-costly-health-issue-105494 (Bottom left) https://news.mit.edu/2019/ ultra-quantum-matter-uqm-research-given-8m-boost-0529 (Bottom right) https://www.nature.com/articles/d41586-019-03213-z

Pirsa: 23020031 Page 3/71

The Setting

• In this work, we focus on training an ML model to learn and predict

$$\rho, O \mapsto f_{\mathcal{E}}(\rho, O) = \text{Tr}(O\mathcal{E}(\rho)),$$

where ρ is an input quantum state, $\mathscr E$ is an (unknown) CPTP map, and O is an observable.

• This includes any function computable by a quantum computer (in exponential time).

Pirsa: 23020031 Page 4/71

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Example 1

Predicting outcomes of physical experiments

 ρ : initial state given by classical input x

 ${\mathscr E}$: the physical process in the experiment

O: what the scientist measure



Example 2

Training quantum neural networks

 ρ : input state encoding classical input \boldsymbol{x}

 ${\mathscr E}$: the quantum neural network to learn

 ${\it O}$: a single fixed observable



Example 3

Speeding up complex quantum dynamics

ho : initial state of the physical system

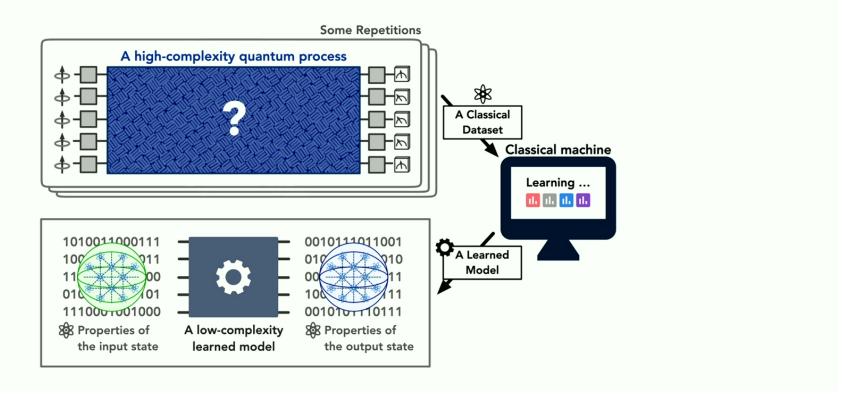
 ${\mathscr E}$: the quantum dynamics to speed up

O: the property we want to predict

Pirsa: 23020031

The goal of this work

Given an n-qubit CPTP map $\mathscr E$ that represents a high-complexity quantum process



Pirsa: 23020031 Page 6/71

A Classical Problem

- Given an unknown classical Boolean circuit C mapping n bits to n bits.
- The input is now an *n*-bit string $x \in \{-1,1\}^n$.
- The 1st output bit of C for input x is equal to $f_C(x) = \text{Tr}(Z_1C(|x\rangle\!\!/x|))$.



Pirsa: 23020031 Page 7/71

Worst-case hardness

- The function f_C is equiv. to an exponentially long vector $\{-1,1\}^{2^n}$ with **no structure**.
- To learn a model h(x), such that $\left|h(x) f_C(x)\right|^2 < 0.5, \forall x \in \{-1,1\}^n$, we must query $f_C(x)$ for all input x. Query complexity: $\Theta(2^n)$.



Pirsa: 23020031 Page 8/71

Average-case hardness

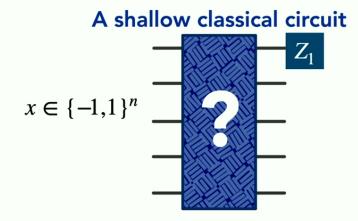
- The function f_C is equiv. to an exponentially long vector $\{-1,1\}^{2^n}$ with **no structure**.
- To learn a model h(x), such that $\mathbb{E}_{x \sim \{-1,1\}^n} \left| h(x) f_C(x) \right|^2 < 0.5$, we must query $f_C(x)$ for at least half of all x. Query complexity: $\Theta(2^n)$.



Pirsa: 23020031 Page 9/71

Average-case hardness for shallow classical circuits

• [AGS19] showed that learning h(x), such that $\mathbb{E}_{x \sim \{-1,1\}^n} \left| h(x) - f_C(x) \right|^2 < 0.5$, is computationally hard (for both classical & quantum computers), even when the classical Boolean circuit is **constant-depth** (with majority gates, i.e., TC_0).



Pirsa: 23020031 Page 10/71

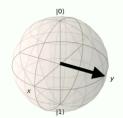
Overview

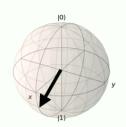
- A classical version of the quantum problem
- A restricted version of the quantum problem
- Generalization to the original quantum problem

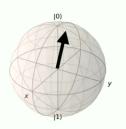
Pirsa: 23020031 Page 11/71

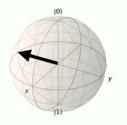
A Quantum Problem

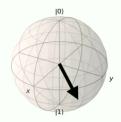
Input:

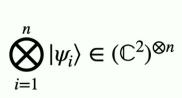


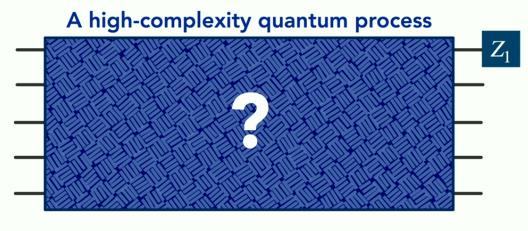










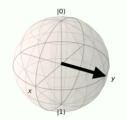


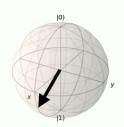
Is this harder?

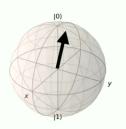
Pirsa: 23020031 Page 12/71

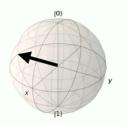
A Quantum Problem

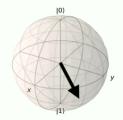
Input:

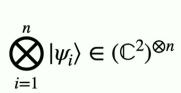


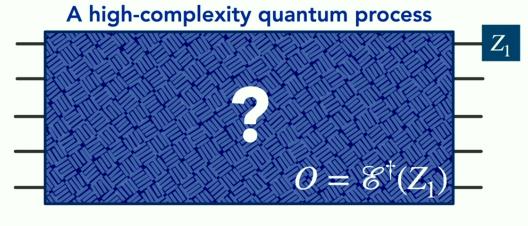












Is this harder?

Pirsa: 23020031 Page 13/71

A Classical Dataset



Classical Dataset about ${\cal O}$

$$\begin{split} |\psi_{\ell}\rangle &= \bigotimes_{i=1}^n |\psi_{\ell,i}\rangle \mapsto y_{\ell}, \quad \mathbb{E}[y_{\ell}] = \langle \psi_{\ell}|O|\psi_{\ell}\rangle \\ & \text{for } \ell=1,...,N. \end{split}$$

Each repetition prepares a random product state, and measures the 1st qubit in the Z basis

Pirsa: 23020031 Page 14/71

The Prediction Task

Classical Dataset about O

$$|\psi_{\ell}\rangle = \bigotimes_{i=1}^{n} |\psi_{\ell,i}\rangle \mapsto y_{\ell}, \quad \mathbb{E}[y_{\ell}] = \langle \psi_{\ell} | O | \psi_{\ell} \rangle$$
 for $\ell = 1, ..., N$.

Given a new state
$$|\psi\rangle = \bigotimes_{i=1}^n |\psi_i\rangle \in (\mathbb{C}^2)^{\otimes n}$$
, how to predict $\langle \psi | O | \psi \rangle$ accurately?

Pirsa: 23020031

Worst-case hardness

- To learn a model $h(|\psi\rangle)$, such that $\left|h(|\psi\rangle) \langle \psi|O|\psi\rangle\right|^2 < 0.5, \forall |\psi\rangle = \bigotimes_{i=1}^n |\psi_i\rangle$, the problem is at least as hard as the classical problem.
- Hence, the query complexity is $\Omega(2^n)$.



Pirsa: 23020031 Page 16/71

Average-case hardness?

- To learn a model $h(|\psi\rangle)$, such that $\mathbb{E}_{|\psi\rangle=\bigotimes_{i=1}^n|\psi_i\rangle}\left|h(|\psi\rangle)-\langle\psi|O|\psi\rangle\right|^2<0.5$, is the problem still exponentially hard?
- Surprisingly, the answer is **no**. The problem can be done in quasi-polynomial time.



Pirsa: 23020031 Page 17/71

$$O = \sum_{P \in \{I, X, Y, Z\}^{\otimes n}} \alpha_P P$$

$$O^{(\mathrm{low})} = \sum_{|P| \le k} \alpha_P P$$

Lemma (Low-weight approximation): $\mathbb{E}_{|\psi\rangle=\bigotimes_{i=1}^n|\psi_i\rangle}\left|\langle\psi|O|\psi\rangle-\langle\psi|O^{(\mathrm{low})}|\psi\rangle\right|^2<\frac{1}{3^k}$.

Pirsa: 23020031

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Interpretation: For most product state $|\psi\rangle = \bigotimes_{i=1}^n |\psi_i\rangle$, $\langle \psi|O|\psi\rangle \approx \langle \psi|O^{(\mathrm{low})}|\psi\rangle$.

Low-weight approximation does not hold in the classical version of this problem

$$O = \sum_{P \in \{I, X, Y, Z\}^{\otimes n}} \alpha_P P$$

$$O^{(\text{low})} = \sum_{|P| \le k} \alpha_P P$$

Classical inputs are perfectly distinguishable.

But quantum state inputs are not.

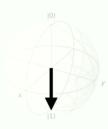
Classical Input:



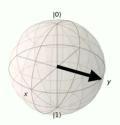


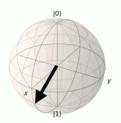


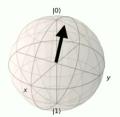


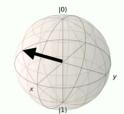


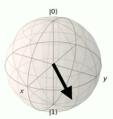
Quantum Input:











$$O = \sum_{P \in \{I, X, Y, Z\}^{\otimes n}} \alpha_P P$$

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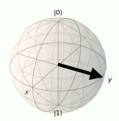


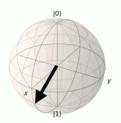


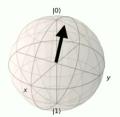


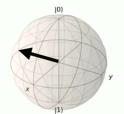


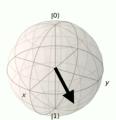
Quantum Input:











Basic Idea for the ML model

Basic idea: Learn the low-weight observable $O^{(low)} = \sum_{|P| \le k} \alpha_P P$ for a small k.

Lemma (Fourier transform):
$$\alpha_P = \mathbb{E}\left[\frac{3^{|P|}}{N}\sum_{\ell=1}^N y_\ell \langle \psi_\ell \,|\, P\,|\, \psi_\ell \rangle\right], \ \forall P \in \{I,X,Y,Z\}^{\otimes n}$$

Classical Dataset

$$\begin{split} |\psi_{\ell}\rangle &= \bigotimes_{i=1}^n |\psi_{\ell,i}\rangle \mapsto y_{\ell}, \quad \mathbb{E}[y_{\ell}] = \langle \psi_{\ell}|O|\psi_{\ell}\rangle \\ \text{for } \ell = 1, \dots, N. \end{split}$$

Pirsa: 23020031 Page 22/71

Basic Idea for the ML model

Basic idea: Learn the low-weight observable $O^{(low)} = \sum \alpha_P P$ for a small k. $|P| \leq k$

Lemma

We only need $N = \mathcal{O}(\log n)$!

Classical Dataset

$$\begin{split} |\psi_{\ell}\rangle &= \bigotimes_{i=1}^n |\psi_{\ell,i}\rangle \mapsto y_{\ell}, \quad \mathbb{E}[y_{\ell}] = \langle \psi_{\ell}|O|\psi_{\ell}\rangle \\ \text{for } \ell = 1, \dots, N. \end{split}$$

Pirsa: 23020031 Page 23/71

An interlude

Optimizing Quantum Hamiltonians

Credit: DALL·E



Presenter: Hsin-Yuan Huang (Robert)

Joint work with Sitan Chen and John Preskill







Pirsa: 23020031 Page 24/71

The Task

Given an n-qubit, k-local Hamiltonian $H = \sum_{|P| \le k} \alpha_P P$.

Find a state $|\psi\rangle$ that maximizes or minimizes $\langle\psi|H|\psi\rangle$.

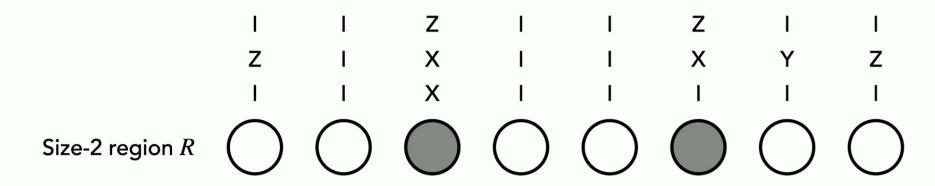
We want a guarantee on $\langle \psi | H | \psi \rangle$ based on the description of $H = \sum_{|P| \leq k} \alpha_P P$

Pirsa: 23020031 Page 25/71

Expansion property

Given an
$$n$$
-qubit, k -local Hamiltonian $H = \sum_{|P| \le k} \alpha_P P$.

H has an expansion coefficient c_e and dimension d_e if for every size- d_e region R, the number of P with $\alpha_P \neq 0$, $\operatorname{dom}(P) \subseteq R$, $R \subseteq \operatorname{dom}(P)$ is at most c_e .



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Example 1

Geometrically-local Hamiltonian

$$c_e = \mathcal{O}(1), d_e = 1$$

Example 2

General k-local Hamiltonian

$$c_e = 4^k, d_e = k$$

Example 3

Degree-d 2-body Hamiltonian

$$c_e = 16d, d_e = 1$$

Theorem

Given an
$$n$$
-qubit, k -local Hamiltonian $H = \sum_{|P| \le k} \alpha_P P$.

If H has an expansion coefficient c_e and dimension d_e , then for $r=2d_e/(d_e+1)\in [1,2)$, we have an algorithm that either finds a maximizing product state $|\psi\rangle$,

$$\langle \psi | H | \psi \rangle \ge \mathbb{E}_{|\phi\rangle: \text{Haar}} \langle \phi | H | \phi \rangle + \frac{1}{c_e^{1/2d_e} 2^{\Theta(k \log k)}} \left(\sum_{P \ne I} |\alpha_P|^r \right)^{1/r},$$

or finds a minimizing product state $|\psi\rangle$ with a similar guarantee (+ \rightarrow -, \geq \rightarrow \leq).

Pirsa: 23020031 Page 28/71

Expansion property

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Pirsa: 23020031 Page 30/71

Theorem

Improved over existing results

Given an
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$$\langle \psi | H | \psi \rangle \ge \mathbb{E}_{|\phi\rangle: \text{Haar}} \langle \phi | H | \phi \rangle + \frac{1}{c_e^{1/2d_e} 2^{\Theta(k \log k)}} \left(\sum_{P \ne I} |\alpha_P|^r \right)^{1/r},$$

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Pirsa: 23020031 Page 31/71

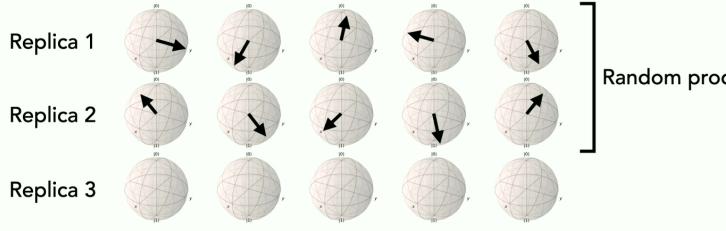
Select a slice with the largest value of α_P

Given an n-qubit, k-local Hamiltonian $H = \sum_{|P|=k} \alpha_P P$.

Find a product state $|\psi\rangle$ that approximately optimizes $\langle\psi|H|\psi\rangle$.

Pirsa: 23020031 Page 32/71

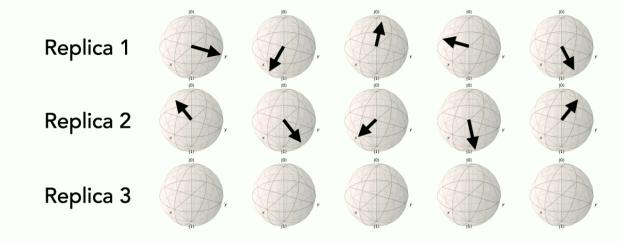
Given an n-qubit, k-local Hamiltonian $H = \sum \alpha_P P$. |P|=k



Random product states

Lift n-qubit H to nk qubits

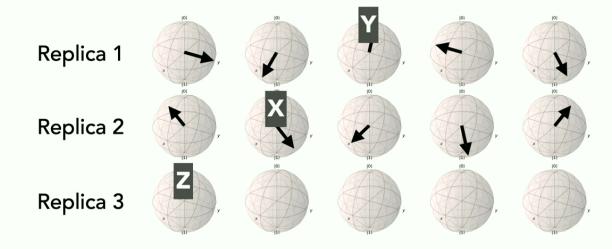
Given an
$$n$$
-qubit, k -local Hamiltonian $H = \sum_{|P|=k} \alpha_P P$.



$$pol(H) = \sum_{|P|=k} \alpha_P pol(P) \in \mathbb{C}^{2^{nk} \times 2^{nk}}$$

Lift n-qubit H to nk qubits

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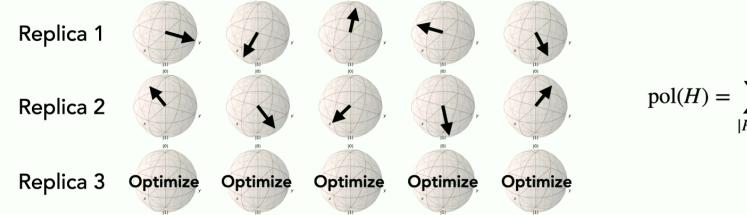


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P = Z X Y I I

Page 35/71

Given an n-qubit, k-local Hamiltonian $H = \sum_{|P|=k} \alpha_P P$.



 $pol(H) = \sum_{|P|=k} \alpha_P pol(P) \in \mathbb{C}^{2^{nk} \times 2^{nk}}$

The Algorithm

Given an
$$n$$
-qubit, k -local Hamiltonian $H = \sum_{|P|=k} \alpha_P P$.











Combine the Bloch vectors using a weighted sum

Given an
$$n$$
-qubit, k -local Hamiltonian $H = \sum_{|P| \le k} \alpha_P P$.

If H has an expansion coefficient c_e and dimension d_e , then for $r=2d_e/(d_e+1)\in [1,2)$, we have an algorithm that either finds a maximizing product state $|\psi\rangle$,

$$\langle \psi | H | \psi \rangle \ge \mathbb{E}_{|\phi\rangle: \text{Haar}} \langle \phi | H | \phi \rangle + \frac{1}{c_e^{1/2d_e} 2^{\Theta(k \log k)}} \left(\sum_{P \ne I} |\alpha_P|^r \right)^{1/r},$$

or finds a minimizing product state $|\psi\rangle$ with a similar guarantee (+ \rightarrow -, \geq \rightarrow \leq).

Pirsa: 23020031 Page 38/71

Another interlude

Generalized Quantum **Bohnenblust-Hille Inequality**

Credit: DALL-E



Presenter: Hsin-Yuan Huang (Robert)

Joint work with Sitan Chen and John Preskill







Pirsa: 23020031 Page 39/71

Given an observable $O = \sum_{|P| \le k} \alpha_P P$ with an expansion coefficient c_e and dimension d_e .

$$||O||_{\infty} \ge \frac{1}{c_e^{1/2d_e} 2^{\Theta(k \log k)}} \left(\sum_{P} |\alpha_P|^r\right)^{1/r} \text{ for } r = \frac{2d_e}{d_e + 1} \in [1, 2).$$

Proof ideas:

- (1) Use the guarantee from the algorithm for optimizing quantum Hamiltonians.
- (2) Adapt by noting that $||O||_{\infty} \ge |\langle \psi | O | \psi \rangle|$, where $|\psi \rangle$ is the state found by the algo.

Pirsa: 23020031 Page 40/71

Given an observable $O = \sum_{e} \alpha_{p}P$ with an expansion coefficient c_{e} and dimension d_{e} .

$$||O||_{\infty} \ge \frac{1}{c_e^{1/2d_e} 2^{\Theta(k \log k)}} \left(\sum_{P} |\alpha_P|^r\right)^{1/r} \text{ for } r = \frac{2d_e}{d_e + 1} \in [1, 2).$$

Example 1

A sum of geometrically-local terms $c_e = \mathcal{O}(1)\text{, } d_e = 1$

$$c_e = \mathcal{O}(1), d_e = 1$$

$$\sum_{P} |\alpha_{P}| \leq \mathcal{O}\left(\|O\|_{\infty}\right)$$

Given an observable $O = \sum_{|P| \le k} \alpha_P P$ with an expansion coefficient c_e and dimension d_e .

$$||O||_{\infty} \ge \frac{1}{c_e^{1/2d_e} 2^{\Theta(k \log k)}} \left(\sum_{P} |\alpha_P|^r\right)^{1/r} \text{ for } r = \frac{2d_e}{d_e + 1} \in [1, 2).$$

Example 2

A sum of k-local terms

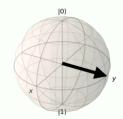
$$c_e = 4^k, d_e = k$$

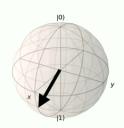
$$\|\overrightarrow{\alpha}\|_{\frac{2k}{k+1}} \le 2^{\mathcal{O}(k\log k)} \|O\|_{\infty}$$

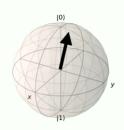
A quantum analogue of the Bohnenblust-Hille inequality

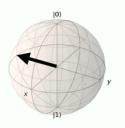
A Quantum Problem

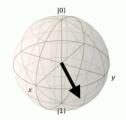
Input:

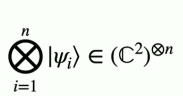


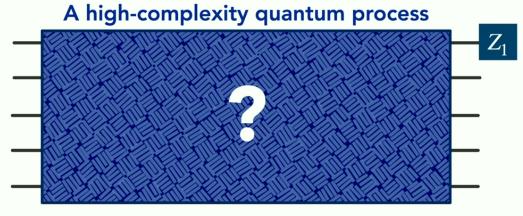












Pirsa: 23020031 Page 43/71

Basic Idea for the ML model

Basic idea: Learn the low-weight observable $O^{(low)} = \sum_{|P| \le k} \alpha_P P$ for a small k.

Lemma (Fourier transform):
$$\alpha_P = \mathbb{E}\left[\frac{3^{|P|}}{N}\sum_{\ell=1}^N y_\ell \langle \psi_\ell \,|\, P\,|\, \psi_\ell \rangle\right], \ \forall P \in \{I,X,Y,Z\}^{\otimes n}$$

Classical Dataset

$$\begin{split} |\psi_{\ell}\rangle &= \bigotimes_{i=1}^n |\psi_{\ell,i}\rangle \mapsto y_{\ell}, \quad \mathbb{E}[y_{\ell}] = \langle \psi_{\ell}|O|\psi_{\ell}\rangle \\ \text{for } \ell = 1, \dots, N. \end{split}$$

Pirsa: 23020031 Page 44/71

Insight from Quantum BH inequality

Insight 1: Learn the low-weight observable $O^{(low)} = \sum \alpha_P P$ for a small k.

Insight 2: The Pauli coef. in $O^{(\text{low})}$ is approximately sparse as $\|\overrightarrow{\alpha}\|_{\frac{2k}{k+1}} \leq 2^{\mathcal{O}(k\log k)} \|O^{(\text{low})}\|_{\infty}$.

This idea is also used in classical learning theory [Al22]



Classical Dataset

$$|\psi_{\ell}\rangle = \bigotimes_{i=1}^{n} |\psi_{\ell,i}\rangle \mapsto y_{\ell}, \quad \mathbb{E}[y_{\ell}] = \langle \psi_{\ell} | O | \psi_{\ell} \rangle$$
 for $\ell = 1, ..., N$.

Pirsa: 23020031 Page 45/71

Basic Idea for the ML model

Basic idea: Learn the low-weight observable $O^{(low)} = \sum \alpha_P P$ for a small k.

Lemma

How large should the data size N be?

Classical Dataset

$$\begin{split} |\psi_{\ell}\rangle &= \bigotimes_{i=1}^n |\psi_{\ell,i}\rangle \mapsto y_{\ell}, \quad \mathbb{E}[y_{\ell}] = \langle \psi_{\ell}|O|\psi_{\ell}\rangle \\ \text{for } \ell = 1, \dots, N. \end{split}$$

Pirsa: 23020031 Page 46/71

Insight 1: Learn the low-weight observable $O^{(low)} = \sum \alpha_P P$ for a small k.

Insight 2: The Pauli coef. in $O^{(\text{low})}$ is approximately sparse as $\|\overrightarrow{\alpha}\|_{\frac{2k}{k+1}} \leq 2^{\mathcal{O}(k\log k)} \|O^{(\text{low})}\|_{\infty}$.



Classical Dataset

$$|\psi_{\ell}\rangle = \bigotimes_{i=1}^{n} |\psi_{\ell,i}\rangle \mapsto y_{\ell}, \quad \mathbb{E}[y_{\ell}] = \langle \psi_{\ell} | O | \psi_{\ell} \rangle$$
for $\ell = 1, ..., N$.

For all $|P| \leq k$,

$$\operatorname{set} \hat{\alpha}_P \leftarrow \frac{3^{|P|}}{N} \sum_{\ell=1}^N y_\ell \langle \psi_\ell \, | \, P \, | \, \psi_\ell \rangle.$$

If $\hat{\alpha}_P$ is small, set $\hat{\alpha}_P \leftarrow 0$.

The learned observable is $\hat{O}^{(\text{low})} = \sum \hat{\alpha}_P P$.

Guarantee for learning O

For any small constant ϵ, ϵ' , given a training set size $N = \mathcal{O}(\log n)$, the prediction error is

$$\mathbb{E}_{|\psi\rangle = \bigotimes_{i=1}^n |\psi_i\rangle} \left| \langle \psi | \hat{O}^{(\text{low})} | \psi \rangle - \langle \psi | O | \psi \rangle \right|^2 < \epsilon + \epsilon' \| O^{(\text{low})} \|_{\infty}^2.$$



Classical Dataset

$$|\psi_{\ell}\rangle = \bigotimes_{i=1}^{n} |\psi_{\ell,i}\rangle \mapsto y_{\ell}, \quad \mathbb{E}[y_{\ell}] = \langle \psi_{\ell} | O | \psi_{\ell} \rangle$$
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The learned observable is $\hat{O}^{(\text{low})} = \sum \hat{\alpha}_P P$.

Guarantee for learning O

For any ϵ, ϵ' , given a training set size $N = \log(n) \, 2^{\tilde{\mathcal{O}}(\log(1/\epsilon)\log(1/\epsilon'))}$, the prediction error is

$$\mathbb{E}_{|\psi\rangle = \bigotimes_{i=1}^n |\psi_i\rangle} \left| \langle \psi | \hat{O}^{(\text{low})} | \psi \rangle - \langle \psi | O | \psi \rangle \right|^2 < \epsilon + \epsilon' \| O^{(\text{low})} \|_{\infty}^2.$$



Classical Dataset

$$|\psi_{\ell}\rangle = \bigotimes_{i=1}^{n} |\psi_{\ell,i}\rangle \mapsto y_{\ell}, \quad \mathbb{E}[y_{\ell}] = \langle \psi_{\ell} | O | \psi_{\ell} \rangle$$
 for $\ell = 1, ..., N$.

For all $|P| \leq k$,

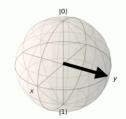
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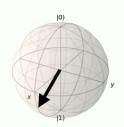
If $\hat{\alpha}_P$ is small, set $\hat{\alpha}_P \leftarrow 0$.

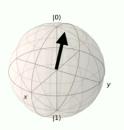
The learned observable is $\hat{O}^{(\text{low})} = \sum \hat{\alpha}_P P$.

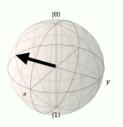
A Quantum Problem

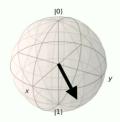


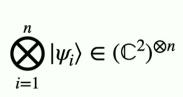


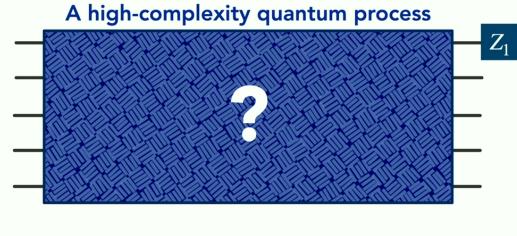












Quasi-polynomially easy!

Pirsa: 23020031 Page 50/71

Guarantee for learning O

For any ϵ, ϵ' , given a training set size $N = \log(n) \, 2^{\tilde{\mathcal{O}}(\log(1/\epsilon)\log(1/\epsilon'))}$, the prediction error is

$$\mathbb{E}_{|\psi\rangle = \bigotimes_{i=1}^n |\psi_i\rangle} \left| \langle \psi | \hat{O}^{(\text{low})} | \psi \rangle - \langle \psi | O | \psi \rangle \right|^2 < \epsilon + \epsilon' \| O^{(\text{low})} \|_{\infty}^2.$$



Classical Dataset

$$|\psi_{\ell}\rangle = \bigotimes_{i=1}^{n} |\psi_{\ell,i}\rangle \mapsto y_{\ell}, \quad \mathbb{E}[y_{\ell}] = \langle \psi_{\ell} | O | \psi_{\ell} \rangle$$
 for $\ell = 1, ..., N$.

For all $|P| \leq k$,

$$\operatorname{set} \hat{\alpha}_P \leftarrow \frac{3^{|P|}}{N} \sum_{\ell=1}^N y_\ell \langle \psi_\ell \, | \, P \, | \, \psi_\ell \rangle.$$

If $\hat{\alpha}_P$ is small, set $\hat{\alpha}_P \leftarrow 0$.

The learned observable is
$$\hat{O}^{(\mathrm{low})} = \sum_{|P| \le k} \hat{\alpha}_P P$$
.

Overview

- A classical version of the quantum problem
- A restricted version of the quantum problem
- Generalization to the original quantum problem

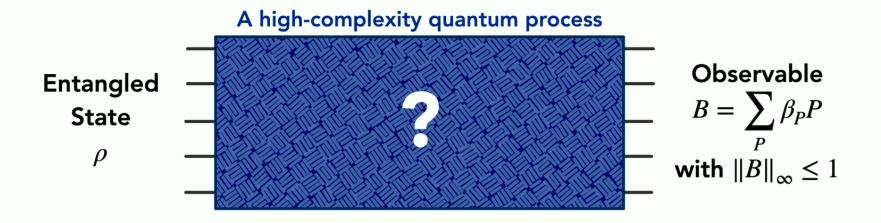
Pirsa: 23020031 Page 52/71

The Restricted Problem



Pirsa: 23020031 Page 53/71

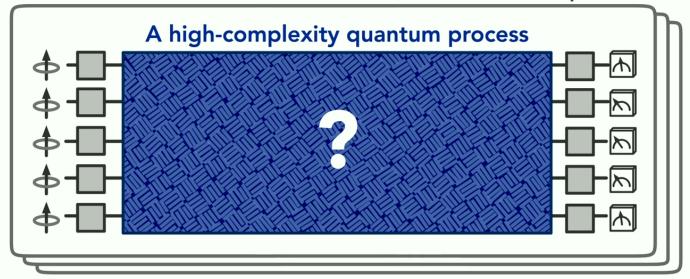
The Original Problem



Pirsa: 23020031 Page 54/71

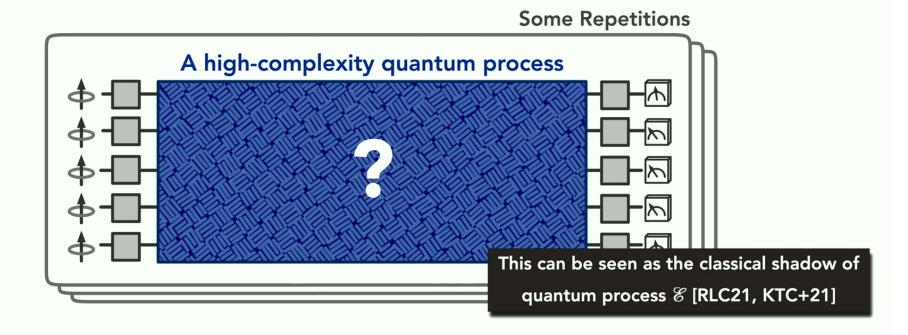
A Classical Dataset for Learning $\mathscr E$

Some Repetitions



Pirsa: 23020031 Page 55/71

A Classical Dataset for Learning $\mathscr E$



Pirsa: 23020031 Page 56/71

A Classical Dataset for Learning $\mathscr E$



Classical Dataset

$$\begin{split} |\psi_{\ell}\rangle &= \bigotimes_{i=1}^n |\psi_{\ell,i}\rangle \mapsto |\phi_{\ell}\rangle = \bigotimes_{i=1}^n |\phi_{\ell,i}\rangle \\ \text{for } \ell = 1, \dots, N. \end{split}$$

This can be seen as the classical shadow of quantum process \mathscr{E} [RLC21, KTC+21]

Pirsa: 23020031 Page 57/71

How to make prediction?



State ρ



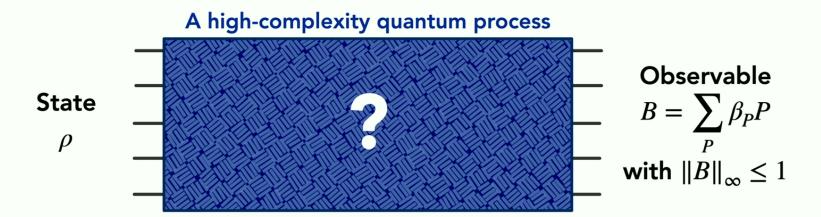
Observable
$$B = \sum_{P} \beta_{P} P$$
 with $\|B\|_{\infty} \leq 1$



Classical Dataset

$$\begin{split} |\psi_{\ell}\rangle &= \bigotimes_{i=1}^n |\psi_{\ell,i}\rangle \mapsto |\phi_{\ell}\rangle = \bigotimes_{i=1}^n |\phi_{\ell,i}\rangle \\ \text{for } \ell = 1, \dots, N. \end{split}$$

Pirsa: 23020031 Page 58/71



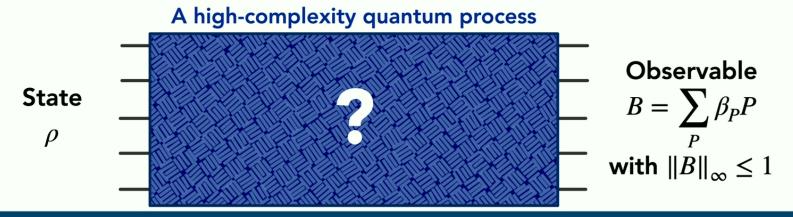
A New Classical Dataset

$$|\psi_{\ell}\rangle = \bigotimes_{i=1}^{n} |\psi_{\ell,i}\rangle \mapsto y_{\ell} = \operatorname{Tr}\left(B\bigotimes_{i=1}^{n} \left(3|\phi_{\ell,i}\rangle\langle\phi_{\ell,i}| - I\right)\right)$$
for $\ell = 1, ..., N$.

Properties [HKP20]:

$$\mathbb{E}[y_{\ell}] = \text{Tr}(B\mathscr{E}(|\psi_{\ell}\rangle\langle\psi_{\ell}|))$$
$$\text{Var}[y_{\ell}] \le ||B||_{\text{shadow}}^{2}$$

Pirsa: 23020031 Page 59/71



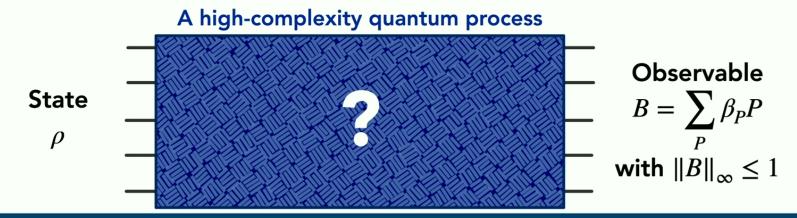
For any sum of local observables B, $\|B\|_{\mathrm{shadow}} \leq \mathcal{O}(\|\vec{\beta}\|_1) \leq \mathcal{O}(\|B\|_{\infty})$ using the generalized quantum BH inequality.

$$|\psi_{\ell}\rangle = \bigotimes_{i=1} |\psi_{\ell,i}\rangle \mapsto y_{\ell} = \operatorname{Tr}\left(B\bigotimes_{i=1} (3|\phi_{\ell,i}\rangle\langle\phi_{\ell,i}| - I)\right)$$
for $\ell = 1, ..., N$.

$$\mathbb{E}[y_{\ell}] = \text{Tr}(B\mathscr{E}(|\psi_{\ell}| | \psi_{\ell}|))$$

$$\text{Var}[y_{\ell}] \leq \|B\|_{\text{shadow}}^{2}$$

Pirsa: 23020031 Page 60/71



For any sum of local observables B, $\|B\|_{\mathrm{shadow}} \leq \mathcal{O}(\|\vec{\beta}\|_1) \leq \mathcal{O}(\|B\|_{\infty})$ using the generalized quantum BH inequality.

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for $\ell = 1, ..., N$.

$$\mathbb{E}[y_{\ell}] = \text{Tr}(B\mathscr{E}(|\psi_{\ell}\rangle\langle\psi_{\ell}|))$$

$$\text{Var}[y_{\ell}] = \mathscr{O}(1)$$

Pirsa: 23020031 Page 61/71

Low-weight approximation

$$O = \sum_{P \in \{I, X, Y, Z\}^{\otimes n}} \alpha_P P$$

$$O^{(\text{low})} = \sum_{|P| \le k} \alpha_P P$$

Lemma (Low-weight approximation): $\mathbb{E}_{\rho \sim \mathcal{D}} \left| \operatorname{Tr}(O\rho) - \operatorname{Tr}(O^{(\text{low})}\rho) \right|^2 < \frac{1}{1.5^k}$.

The lemma holds for any distribution ${\mathcal D}$ over any quantum state ρ as long as ${\mathcal D}$ is flat under single-qubit rotations.

Example: ρ is the ground/thermal state of a generic geometrically-local Hamiltonian.

Pirsa: 23020031 Page 62/71





Classical Dataset

$$|\psi_{\ell}\rangle = \bigotimes_{i=1}^{n} |\psi_{\ell,i}\rangle \mapsto |\phi_{\ell}\rangle = \bigotimes_{i=1}^{n} |\phi_{\ell,i}\rangle$$
 for $\ell = 1, \dots, N$.

Pirsa: 23020031 Page 63/71





Classical Dataset for
$$O = \mathcal{E}^{\dagger}(B)$$

$$|\psi_{\ell}\rangle = \bigotimes_{i=1}^{n} |\psi_{\ell,i}\rangle \mapsto y_{\ell}, \quad \mathbb{E}[y_{\ell}] = \langle \psi_{\ell} | O | \psi_{\ell} \rangle$$
 for $\ell = 1, ..., N$.

For all
$$|P| \le k$$
,
$$\operatorname{set} \hat{\alpha}_P \leftarrow \frac{3^{|P|}}{N} \sum_{\ell=1}^N y_\ell \langle \psi_\ell | P | \psi_\ell \rangle.$$

If $\hat{\alpha}_P$ is small, set $\hat{\alpha}_P \leftarrow 0$.

The learned observable is $\hat{O}^{(\mathrm{low})} = \sum_{|P| \leq k} \hat{\alpha}_P P.$





Classical Dataset for
$$O = \mathcal{E}^{\dagger}(B)$$

$$\begin{split} |\psi_{\ell}\rangle &= \bigotimes_{i=1}^n |\psi_{\ell,i}\rangle \mapsto y_{\ell}, \quad \mathbb{E}[y_{\ell}] = \langle \psi_{\ell}|O|\psi_{\ell}\rangle \\ \text{for } \ell = 1, \dots, N. \end{split}$$

Predict
$$\operatorname{Tr}\left(\hat{O}^{(\mathrm{low})}\rho\right) \approx \operatorname{Tr}\left(B\mathscr{E}(\rho)\right)$$

Pirsa: 23020031

Surprising aspects of the ML algorithm

- We can learn to predict n-qubit exponential-size quantum circuits up to a const. relative error from only $\mathcal{O}(\log n)$ samples.
- The algorithm is computationally efficient (polynomial time for a const. relative error; quasi-polynomial time for a small error).

Pirsa: 23020031 Page 66/71

Surprising aspects of the ML algorithm

- We can learn to predict n-qubit exponential-size quantum circuits up to a const. relative error from only $\mathcal{O}(\log n)$ samples.
- The algorithm is computationally efficient (polynomial time for a const. relative error; quasi-polynomial time for a small error).
- After learning from product state inputs, the algorithm can predict entangled states.
- The entire algorithm can be run on a classical computer.

Pirsa: 23020031 Page 67/71





Classical Dataset for
$$O = \mathcal{E}^{\dagger}(B)$$

$$\begin{split} |\psi_{\ell}\rangle &= \bigotimes_{i=1}^n |\psi_{\ell,i}\rangle \mapsto y_{\ell}, \quad \mathbb{E}[y_{\ell}] = \langle \psi_{\ell}|O|\psi_{\ell}\rangle \\ \text{for } \ell = 1, \dots, N. \end{split}$$

Predict
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Pirsa: 23020031

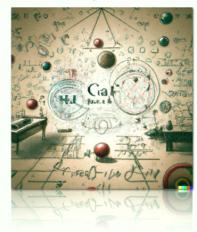
Conclusion

- We give a computationally-efficient ML algorithm that can learn to predict the output of a quantum process with arbitrary complexity.
- Our results highlight the potential that ML models can predict outcomes of a complex quantum dynamics much faster than the process itself.

DALL-E impression of "Predicting quantum processes", "Optimizing quantum Hamiltonians", "Quantum Bohnenblust-Hille"







Pirsa: 23020031 Page 69/71

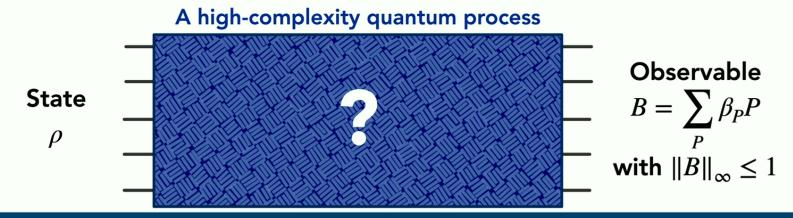
The Task

Given an n-qubit, k-local Hamiltonian $H = \sum_{|P| \le k} \alpha_P P$.

Find a state $|\psi\rangle$ that maximizes or minimizes $\langle\psi|H|\psi\rangle$.

We want a guarantee on $\langle \psi | H | \psi \rangle$ based on the description of $H = \sum_{|P| \leq k} \alpha_P P$

Pirsa: 23020031 Page 70/71



For any sum of local observables B, $\|B\|_{\mathrm{shadow}} \leq \mathcal{O}(\|\vec{\beta}\|_1) \leq \mathcal{O}(\|B\|_{\infty})$ using the generalized quantum BH inequality.

$$|\psi_{\ell}\rangle = \bigotimes_{i=1} |\psi_{\ell,i}\rangle \mapsto y_{\ell} = \operatorname{Tr}\left(B\bigotimes_{i=1} (3|\phi_{\ell,i}\rangle\langle\phi_{\ell,i}| - I)\right)$$
for $\ell = 1, ..., N$.

$$\mathbb{E}[y_{\ell}] = \text{Tr}(B\mathscr{E}(|\psi_{\ell}\rangle\langle\psi_{\ell}|))$$

$$\text{Var}[y_{\ell}] = \mathscr{O}(1)$$

Pirsa: 23020031 Page 71/71