

Title: Standard Model Lecture - 230206

Speakers: Latham Boyle

Collection: Standard Model (2022/2023)

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Standard Model

Lecture # 13

Is the Standard Model Exceptional?

Anomalies:

- Classical sym: $\varphi \rightarrow \varphi + \delta\varphi$ leaves $S[\varphi]$ invariant
- Quantum sym: $\varphi \rightarrow \varphi + \delta\varphi$ leaves $\int \mathcal{D}[\varphi] e^{iS[\varphi]}$ invariant
- Anomaly: $\varphi \rightarrow \varphi + \delta\varphi$ leaves $S[\varphi]$ but not $\int \mathcal{D}[\varphi]$ invariant
- Global sym: $\varphi(x) \rightarrow e^{i\alpha} \varphi(x)$ ← anomaly OK!

Gauge sym: $\varphi(x) \rightarrow e^{i\alpha(x)} \varphi(x)$ ← anomaly = inconsistency!

- Example: "Chiral anomaly"

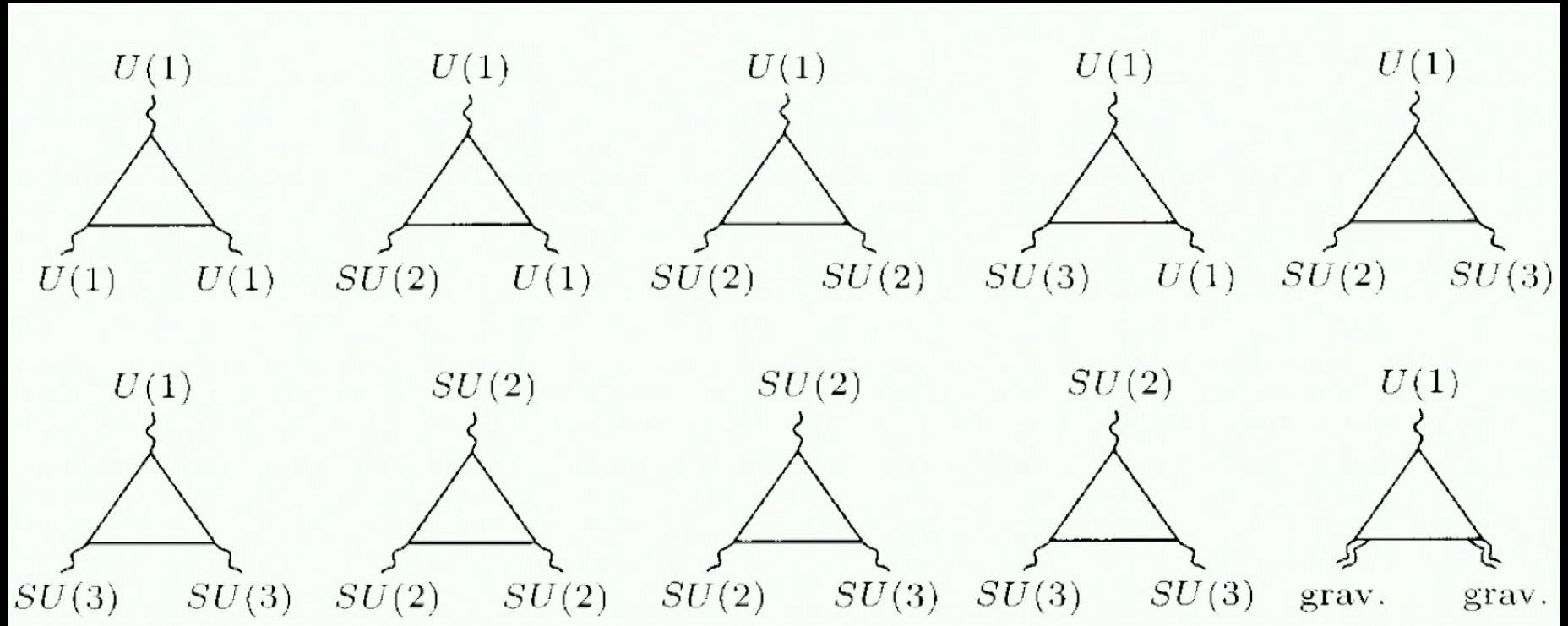
$$Z = \int \mathcal{D}A \mathcal{D}\psi_R \mathcal{D}\bar{\psi}_R e^{iS} \quad S = \int d^4x \left(-\frac{1}{4} \text{Tr} F_{\mu\nu} F^{\mu\nu} + i \bar{\psi}_R \not{D} \psi_R \right)$$

classical sym: $\psi_R \rightarrow e^{i\alpha} \psi_R \Rightarrow S \rightarrow S$

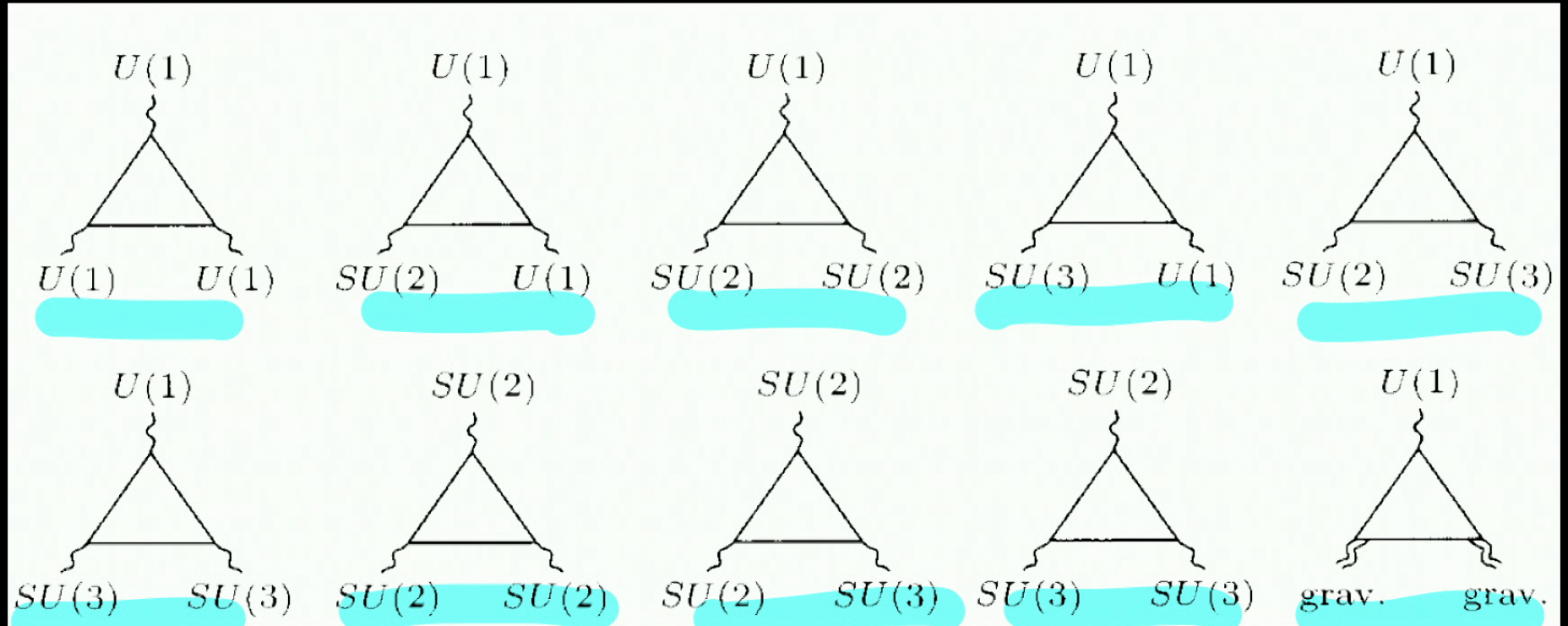
but $\mathcal{D}\psi_R \mathcal{D}\bar{\psi}_R \rightarrow \mathcal{D}\psi_R \mathcal{D}\bar{\psi}_R e^{i \int d^4x \frac{\alpha g^2}{32\pi^2} \text{Tr} F_{\mu\nu} \tilde{F}^{\mu\nu}}$ $\left(\tilde{F}^{\mu\nu} = \epsilon^{\mu\nu\rho\sigma} F_{\rho\sigma} \right)$

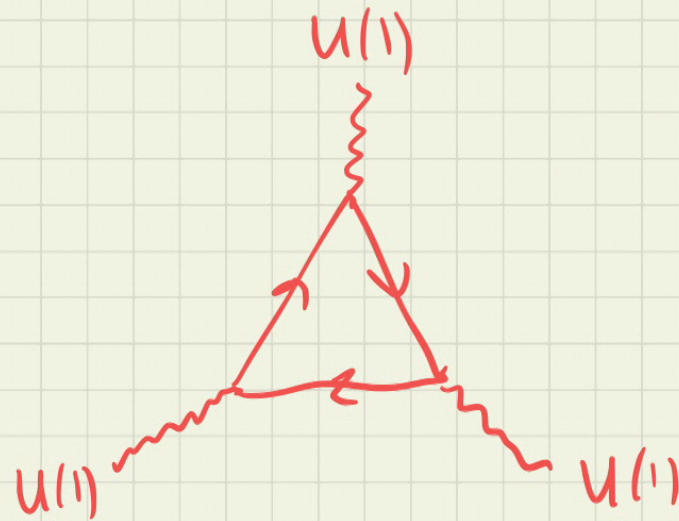
if we replace $\psi_R \rightarrow \psi_L$, then $i \rightarrow -i$

In Standard Model, gauge and gravitational anomalies must cancel:



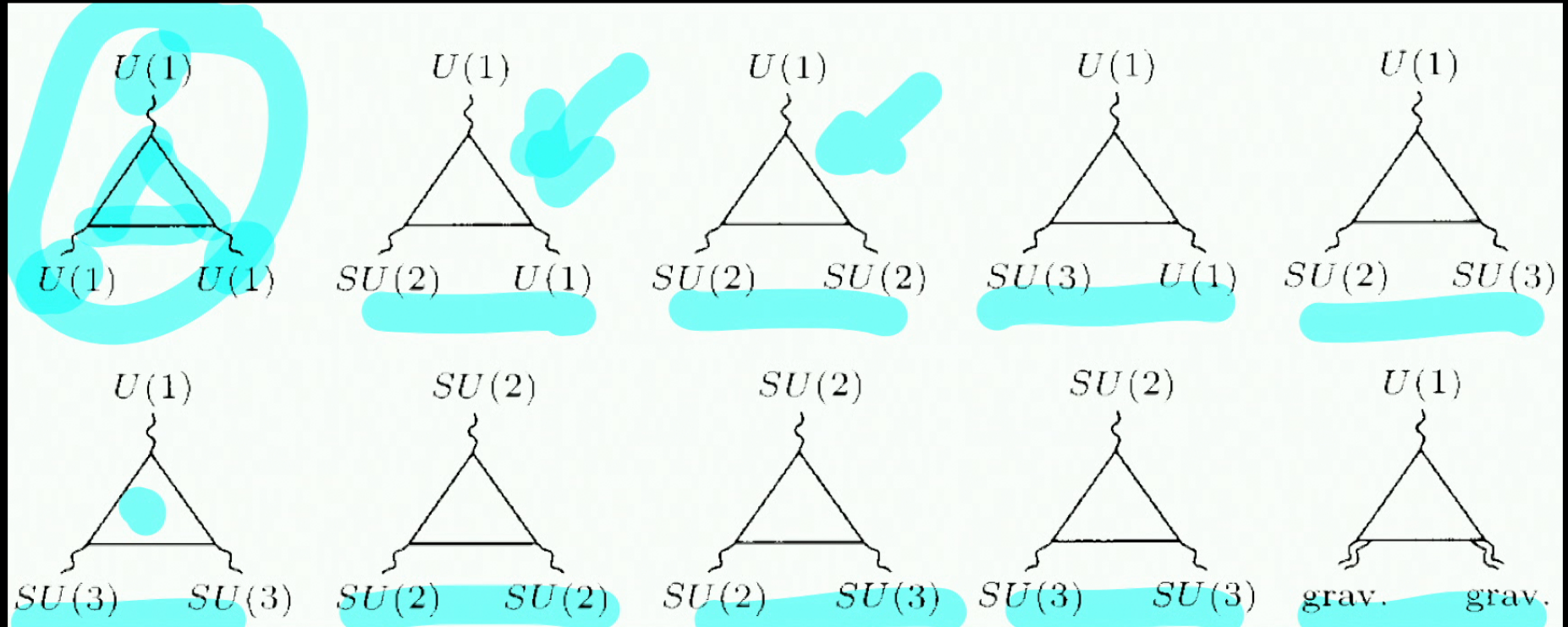
In Standard Model, gauge and gravitational anomalies must cancel:





$$3 \times 2 \times \left(\frac{1}{6}\right)^3 + 3 \times 1 \times \left(-\frac{2}{3}\right)^3 + 3 \times 1 \times \left(\frac{1}{3}\right)^3 \\ + 1 \times 2 \times \left(-\frac{1}{2}\right)^3 + 1 \times 1 \times (1)^3 = 0$$

In Standard Model, gauge and gravitational anomalies must cancel:



Anomaly cancellation: very constraining!

Two fun exercises:

	SU(3)	SU(2)	U(1) _y
q _L	3	2	Y _q
u _R	3	1	Y _u
d _R	3	1	Y _d
l _L	1	2	Y _l
e _R	1	1	Y _e

	SU(3)	SU(2)	U(1) _y	U(1) _x
q _L	3	2	Y _q	X _q
u _R	3	1	Y _u	X _u
d _R	3	1	Y _d	X _d
l _L	1	2	Y _l	X _l
ν _R	1	1	Y _ν	X _ν
e _R	1	1	Y _e	X _e

What about local scale (“Weyl”) invariance?

$$g_{\mu\nu}(x) \rightarrow \Omega^2(x) g_{\mu\nu}(x)$$

- Emphasized by Weyl, Dirac, Dicke, ..., 't Hooft:
 - Natural generalization of diff invariance (gen. covariance)
- Ignoring Higgs, Standard Model is *classically* Weyl invariant.
- But Weyl symmetry is anomalous:

$$\langle T_{\mu}^{\mu} \rangle = c C^2 - a E$$

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Cancelling Weyl anomalies and Vacuum Energy?

$$a = \frac{1}{360(4\pi)^2} [n_0 + \frac{11}{2}n_{1/2} + 62n_1 - 28n'_0]$$
$$c = \frac{1}{120(4\pi)^2} [n_0 + 3n_{1/2} + 12n_1 - 8n'_0]$$

$$E_{\mathbf{k}} = \frac{\hbar\omega}{2} [n_0 - 2n_{1/2} + 2n_1 + 2n'_0]$$

$$S_4[\varphi] = \frac{1}{2} \int d^4x \sqrt{g} \varphi \Delta_4 \varphi$$

$$\Delta_4 = \square^2 + 2R^{\mu\nu} \nabla_\mu \nabla_\nu - \frac{2}{3}R \square + \frac{1}{3}(\nabla^\mu R) \nabla_\mu$$

See Fradkin and Tseytlin, Nucl. Phys. B203 (1982), 157.

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$$E_{\mathbf{k}} = \frac{\hbar\omega}{2} [n_0 - 2n_{1/2} + 2n_1 + 2n'_0]$$

$$n_{1/2} = 4n_1, \quad n'_0 = 3n_1, \quad n_0 = 0.$$

Matches standard model!

$$n_1 = 8 + 3 + 1 = 12$$

$$n_{1/2} = 3 \times 16 = 48$$

Dimension-zero scalars: notable features

$$\langle \varphi(t, \mathbf{x}) \varphi(t, \mathbf{x}') \rangle = \int \frac{d^3 k}{(2\pi)^3} e^{i\mathbf{k} \cdot (\mathbf{x} - \mathbf{x}')} \frac{1}{4k^3}$$

Scale invariant!

No new local degrees of freedom.

*seem to explain spectrum of primordial perturbations in cosmology;
arXiv:2302.00344*

See Bogoliubov, Logunov, Oksak & Todorov (1990) and V.O. Rivelles (2003).

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$$E_{\mathbf{k}} = \frac{\hbar\omega}{2} \left[n_0 - 2n_{1/2} + 2n_1 \right]$$

Standard Model

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Is the Standard Model Exceptional?

Questions:

1) Why $G_{SM} = [SU(3)_c \times SU(2)_L \times U(1)_Y] / \mathbb{Z}_6$?

2) Why $\mathcal{F}_{SM}^{(1)} = (\overset{q_L}{3}, \overset{d_R^c}{2}, \overset{u_R^c}{\frac{1}{6}}) \oplus (\bar{3}, 1, \frac{1}{3}) \oplus (\bar{3}, 1, -\frac{2}{3})$
 $\oplus (\overset{l_L}{1}, \overset{e_R^c}{2}, -\frac{1}{2}) \oplus (1, 1, 1) \oplus (1, 1, 0) \oplus (\overset{\nu_R^c}{1}, 1, 0)$?

3) Why 3 generations:

$$\mathcal{F}_{SM} = \mathcal{F}_{SM}^{(1)} \oplus \mathcal{F}_{SM}^{(1)} \oplus \mathcal{F}_{SM}^{(1)}$$

Old Hints:

1) Grand Unification

i) $SU(5) \rightarrow G_{SM}$

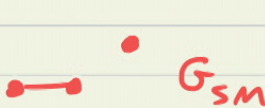
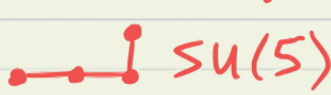
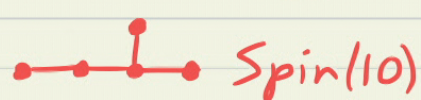
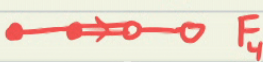
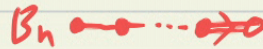
$\Lambda E^5 \rightarrow \rho_{SM}^{(1)}$

ii) $Spin(10) \rightarrow G_{SM}$

$16 \rightarrow \rho_{SM}^{(1)}$

2) Exceptional?

Killing-Cartan:



Magic Square

(Freudenthal, Tits, Gursky, Ramond, Adams)

		\mathbb{K}			
	\mathbb{R}	\mathbb{C}	\mathbb{H}	\mathbb{O}	
\mathbb{R}	$so(3)$	$su(3)$	$sp(3)$	f_4	
\mathbb{C}	$su(3)$	$su(3) + su(3)$	$su(6)$	e_6	
\mathbb{H}	$sp(3)$	$su(6)$	$so(12)$	e_7	
\mathbb{O}	f_4	e_6	e_7	e_8	

$$m(\mathbb{K}, \tilde{\mathbb{K}}) = \underbrace{tr(\mathbb{K}) \oplus tr(\tilde{\mathbb{K}})}_{\text{bosonic}} \oplus \underbrace{(\mathbb{K} \otimes \tilde{\mathbb{K}})_1 \oplus (\mathbb{K} \otimes \tilde{\mathbb{K}})_2 \oplus (\mathbb{K} \otimes \tilde{\mathbb{K}})_3}_{\text{fermionic}} \leftarrow \text{superalgebra}$$

$$m(\mathbb{C}, \mathbb{O}) = \underbrace{tr(\mathbb{C}) \oplus tr(\mathbb{O})}_{e_6} \oplus \underbrace{(\mathbb{C} \otimes \mathbb{O})_1 \oplus (\mathbb{C} \otimes \mathbb{O})_2 \oplus (\mathbb{C} \otimes \mathbb{O})_3}_{spin(10) \oplus u(1)}$$

16: Weyl spinor of spin(10)

$$m(\mathbb{O}, \mathbb{O}') = \underbrace{tr(\mathbb{O}) \oplus tr(\mathbb{O}')}_{e_8} \oplus (\mathbb{O} \otimes \mathbb{O}')_1 \oplus (\mathbb{O} \otimes \mathbb{O}')_2 \oplus (\mathbb{O} \otimes \mathbb{O}')_3$$

see K. Krasnov
arXiv: 2104.01786

right counting for 3 gen (internal + spacetime)

Another Hint:

Dubois-Violette + Todorov (+ Baez) (2018)

$$h_3(\mathbb{D}) \longleftrightarrow \mathcal{G}_{SM}$$

What about \mathcal{G}_{SM} ?