

Title: Mathematical Physics Lecture - 230202

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Collection: Mathematical Physics (2022/2023)

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States  $\longleftrightarrow$

$(M, \omega)$   
 $f: M \rightarrow B$   
 $\dim B = n = \frac{1}{2} \dim M$   
Further, functions on  $B$

Lagrangian submanifolds

Poisson commute

$H = \{ \mathbb{C}$  valued  
fns on  $B \}$  points  
of  $B \}$

$b \in B$ ,  $b$  is Bohr-Sommerfeld  
 $\delta_b \in$  a state.

Think:

$\delta_b$  is associated to the  
submanifold  $f^{-1}(b) \subseteq M$   
of dimension  $n$

Def'n

A submanifold  $L \subseteq M$   
is Lagrangian if

1)  $\dim L = \frac{1}{2} \dim M$

2)  $\omega|_L = 0$

Equivalently

If, locally  $L$  is cut out by  $f_1, \dots, f_n = 0$

2) is equivalent to

$$\{f_i, f_j\} = 0 \text{ on } L$$

Equivalently

If, locally  $L$  is cut out by  $f_1, \dots, f_n = 0$

2) is equivalent to  
 $\{f_i, f_j\} = 0$  on  $L$

$\delta - f_n$  is a "measure"

$$\int \delta_{x=0} f(x) dx = f(0)$$

Equivalently

If, locally  $L$  is cut out by  $f_1, \dots, f_n = 0$

2) is equivalent to

$$\{f_i, f_j\} = 0 \text{ on } L$$

Eg  $p_1, p_2, q_1, q_2$   $\omega = \sum dq_i \wedge dp_i$

$L = \{p_1 = 0, q_1 = 0\}$  This is not Lagrangian

as,

$\{p_1, q_1\} = 1$  is not zero on  $L$

Or.  $q_2, p_2$  are coords on  $L$

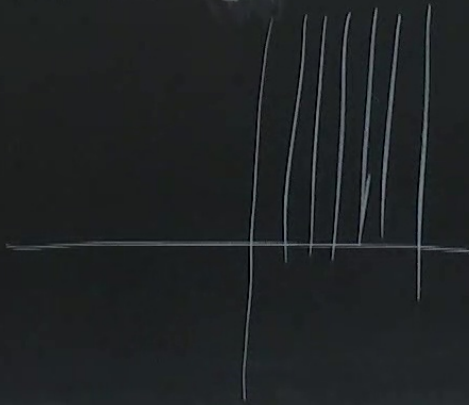
$\omega|_L = dq_2 \wedge dp_2$ , not zero

If  $\pi: M \rightarrow B$

is a polarization

$\pi^{-1}(b) \subseteq M$  is Lagrangian

for generic  $b$



Theorem (Darboux-Weinstein)

If  $L$  is a Lagrangian,

$\exists$  coords locally

$p_i, q_i$  so

$$\omega = \sum dq_i \wedge dp_i$$

and  $L = \{ p_i = 0, i=1..n \}$

Quantum State  $\xrightarrow{\hbar \rightarrow 0}$  Lagrangian

If  $|\psi\rangle$  is a state

$\mathcal{A} =$  def. quantization of  $C^\infty(M)$

Let  $\mathcal{I} \subseteq \mathcal{A}$  of operators  $\Theta$  so that

$$\Theta|\psi\rangle = 0$$

As  
bec  
 $\mathcal{I}_\hbar$



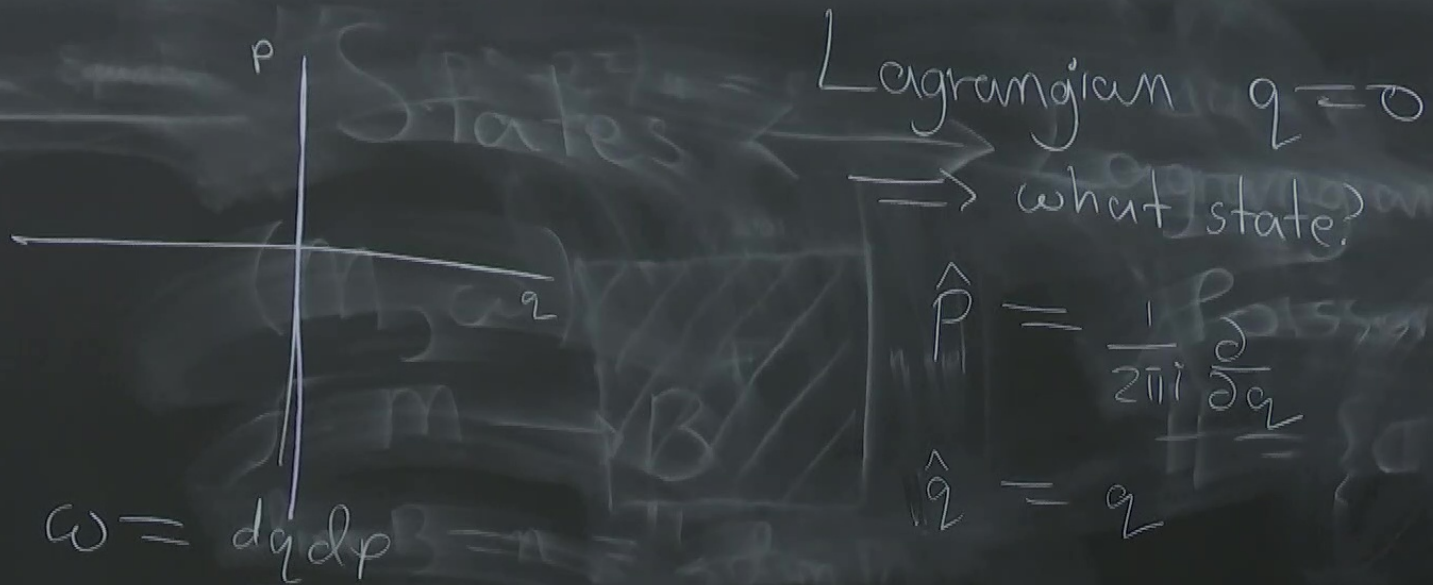
As  $\hbar \rightarrow 0$  this will  
become a subset

$$I_{\hbar=0} \subseteq C^\infty(m)$$

of functions which vanish  
on  $L$ .

$$\text{If } \theta_1, \theta_2 \in I \subseteq A$$

$$[\theta_1, \theta_2]|\psi\rangle = 0 \text{ as well}$$



Want  $\psi(q)$  to be such that Bohr-Sommerfeld

$$q \psi(q) = 0 \quad \psi \in \text{state}$$

Soln  $\psi(q) = \delta_{q=0}$

Definite position,  $S_p$  is associated to the  
undetermined momentum  $f^{-1}(5) \subset \mathbb{R}^n$

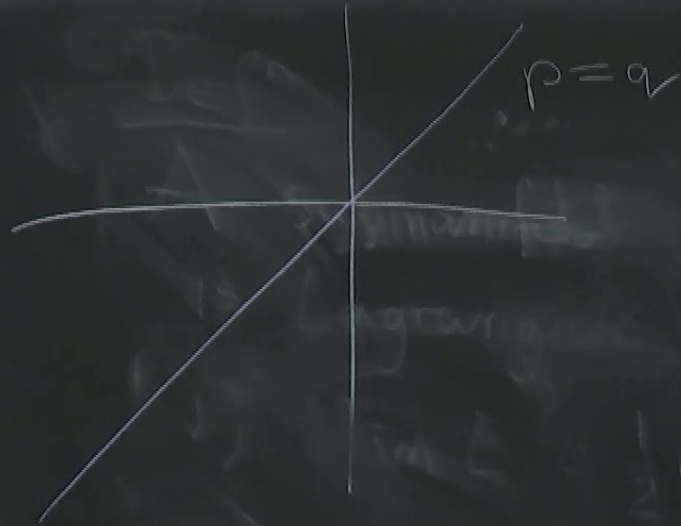
Want  $\psi(q)$  to be such that  $\hat{p} = 0$

$$q \psi(q) = 0 \Rightarrow \psi = 1$$

Sol<sup>n</sup>  $\psi(q) = \delta_{q=0}$

Definite momentum,  
undetermined position

Definite position,  $\delta_{q=0}$  is associated to the  
undetermined momentum  $f^{-1}(5) \subset \mathbb{R}$



$$\hat{p} = \frac{1}{2\pi i} \partial_q$$

$$\left( \frac{1}{2\pi i} \partial_q - a \right) \psi(a) = 0$$

$$\psi = e^{\frac{2\pi i}{2} a^2}$$

$$\int \delta_{q=0} f(q) = f(0)$$

$$\int (\delta_{q=0} g(q)) f(q) = g(0) f(0)$$

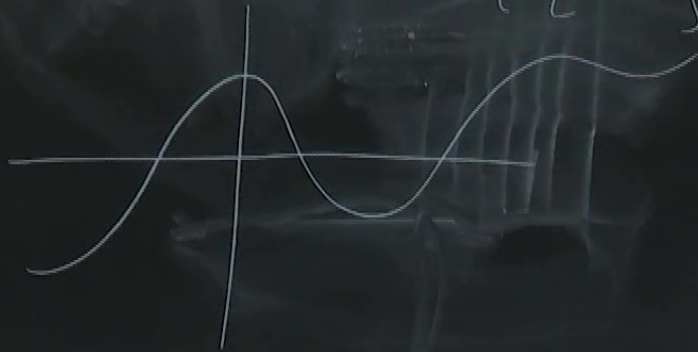
$$g(q) \delta_{q=0} = g(0) \delta_{q=0}$$

is a "measure"  
 $\int \delta_{q=0} = 1$  is not zero on  $\mathbb{R}$   
 $\delta_{q=0}$  are coordinate  
 $dq$  and  $p$  not zero

Given  $L \subseteq \mathbb{R}^2$ , what is  $\psi_L(q)$ ?

Answer: Assume  $\epsilon$  for simplicity

that  $L \cap \{q = c\}$  is always a point.



Then,  $L =$  graph of a function  
 $g$  on  $\mathbb{R}$  " $p = g(q)$ " defines  $L$

State is  $\psi(q)$  satisfies

$$\frac{1}{2\pi i} \partial_q \psi = g(q) \psi$$

which we solve by integrating.



Or:  $\mathcal{L} = 2\pi i q dp$

So  $\omega = \frac{d\alpha}{2\pi i}$  state

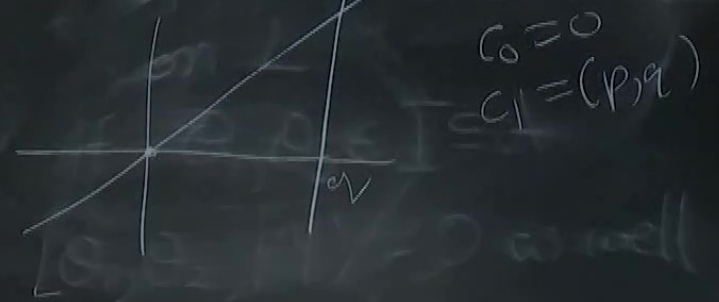
Then, we can compute

$\psi(q_1) \psi(q_0)$

$= \exp\left(\int_{c_0}^{c_1} \alpha\right)$

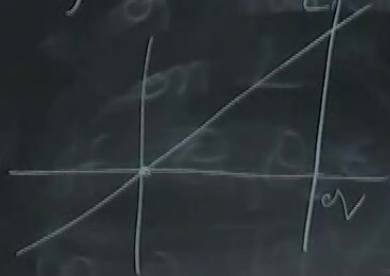
Integral is along the path from  $c_0$  to  $c_1$  in  $L$

If  $L = \{p=q\}$



Integral is along the path  
from  $c_0$  to  $c_1$  in  $L$

$$\text{If } L = \{p=q\}$$



$$c_0=0$$
$$c_1=(p,q)$$

$$\alpha = 2\pi i q dp$$

$$\text{On } L, dp = dq$$

$$\alpha = \frac{2\pi i}{2} dq^2$$

$$\int_0^q \alpha = \frac{2\pi i}{2} q^2$$

$$\psi = \exp\left(\frac{2\pi i}{2} q^2\right)$$

I forgot to mention

Loops in  $L$ .

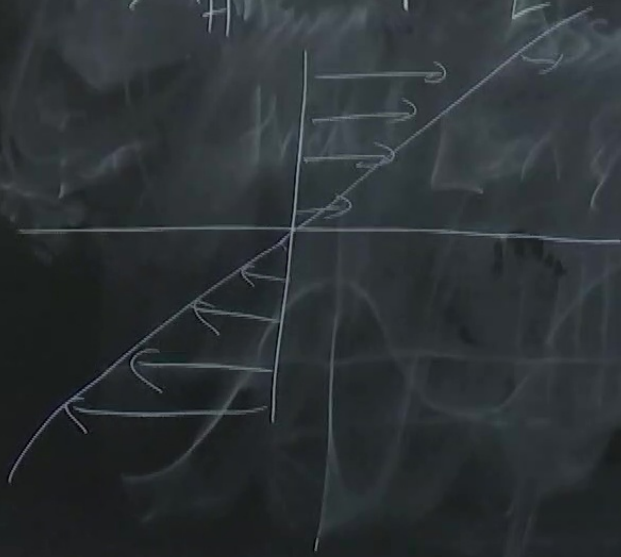
$L \subseteq M$  will correspond  
to a state only if BS  
condition is satisfied.

$$\omega = \frac{1}{2\pi i} \int F(A) \quad A \in U(1)$$

2) gauge field,  $e^{\int A} = 1$  for all

$$H_{lin} = \frac{1}{2} p^2$$

$$X_{Hlin} = p \partial_q$$



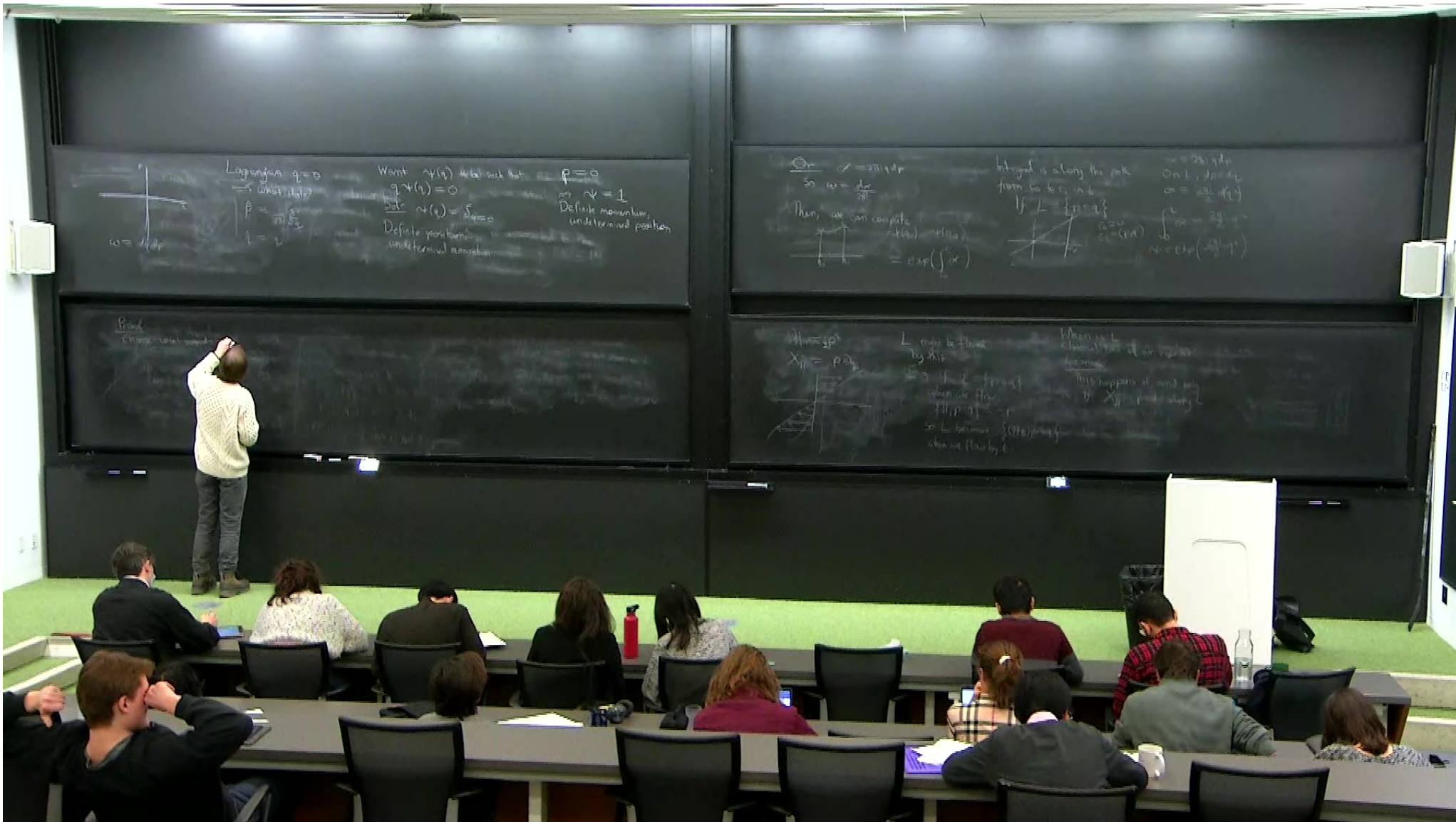
$L$  might be flowed  
by this:

E.g. if  $L = \{p = q\}$

when we flow,

$$\{H, p - q\} = -p$$

so  $L$  becomes  $\{(1+t)p = q\}$   
when we flow by  $t$ .



When is  $L$   
Classical limit of a function?

Lemma

This happens if, and only  
if  $X_{\hbar}(\rho)$  points along  $L$

if we solve by integrating

Proof

Choose local coords  $p_i, q_i$

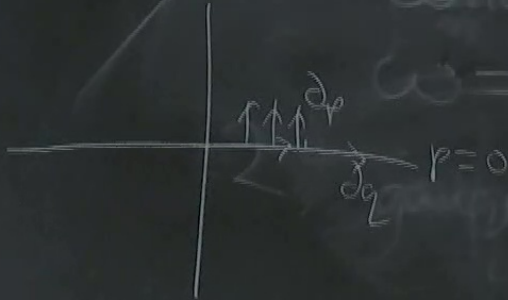
$$\text{so } L = \{p_i = 0\}$$

Then,

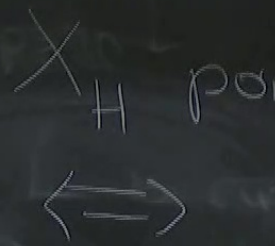
$$X_H = \sum \frac{\partial H}{\partial p_i} \frac{\partial}{\partial q_i}$$

$$= \sum \frac{\partial H}{\partial q_i} \frac{\partial}{\partial p_i}$$

normal to  $L$



along  $L$



$X_H$  points along  $L$

$$\Leftrightarrow \frac{\partial H}{\partial q_i} = 0 \text{ on } L$$

$q_i$  are coords. on  $L$

this means  $H$  is constant on  $L$   
which is semi-classical version of  $L$  being  
an eigenstate.

$$\hat{H}|\psi_L\rangle = c|\psi_L\rangle$$
$$(\hat{H} - c)|\psi_L\rangle = 0$$

In classical limit



$X_H$  points along  $L$

$$\Leftrightarrow \frac{\partial H}{\partial q_i} = 0 \text{ on } L$$

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this means  $H$  is constant on  $L$   
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an eigenstate.

$$\hat{H}|\psi\rangle = c|\psi\rangle$$
$$(H-c)|\psi\rangle = 0$$

In classical limit

$H-c$  vanishes on  $L$

$$H=c \text{ on } L$$

$$(p, q)^2$$

$$(\hat{p}, \hat{q})^2$$

$$\hat{p}^2 \hat{q}^2$$

$$\hat{p} \hat{q} \hat{p} \hat{q}$$

Lagrangian Suppose we have

what is state

Hamiltonian

$\hat{q} = q$

commutator

Suppose we have product  $M_1 \times M_2$  of 2 systems.

e.g.  $M_1 = \mathbb{R}^2$

$M_2 = \mathbb{R}^2$

words  $p_1, q_1$   
"  $p_2, q_2$

A Lagrangian  $L \subset M_1 \times M_2$  is separable

if  $L = L_1 \times L_2$

Otherwise, it's entangled.

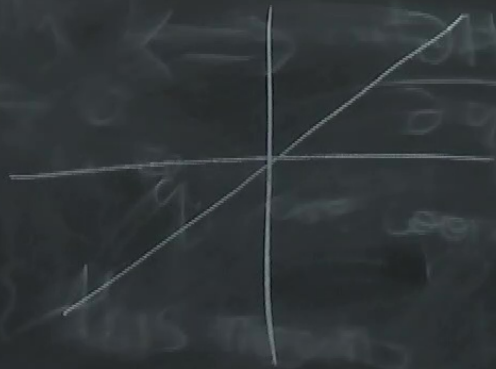
Take

$$L = \begin{cases} q_1 - q_2 \\ p_1 = -p_2 \end{cases}$$

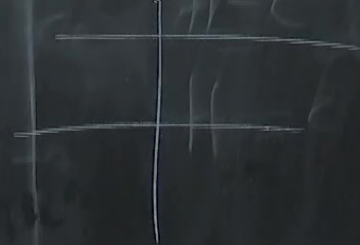
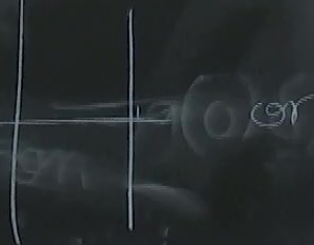
Since  $\{ p_1 + p_2, q_1 - q_2 \} \neq 0$

This is Lagrangian

In  $q_1, q_2$  plane  $L$  is



whereas separable states will be



The quantum state

$\Psi_L(q_1, q_2)$  satisfies

$$q_1 \Psi_L = q_2 \Psi_L$$

$$\left( \frac{\partial}{\partial q_1} + \frac{\partial}{\partial q_2} \right) \Psi_L = 0$$

Solution:

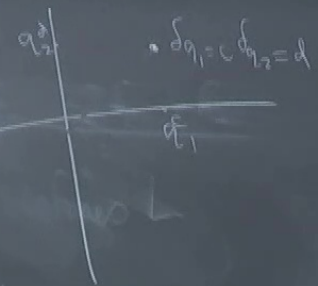
$$\Psi_L(q_1, q_2) = \delta_{q_1=q_2}$$

Solution:

$$\Psi_L(q_1, q_2) = \delta_{q_1=q_2}$$

$\delta_{q_1=c} \otimes \delta_{q_2=c}$  are separable states

$$\delta_{q_1=q_2} = \int_c \delta_{q_1=c} \otimes \delta_{q_2=c} \quad \text{so it is (very) entangled.}$$



$$\alpha = 2\pi i q dp$$

$Q \rightarrow \text{area } 2\pi i q dp$

$P_{3/2} \quad \omega = 2dqdp_{3/2}$  state

Then  $P = e^{2\pi i p}$

$Q = e^{2\pi i q}$  compute

Let

$$PQ = QP e^{2\pi i/n}$$

$$= -QP \text{ on } n=2$$

Hilbert space has 2 states

$$Q|\uparrow\rangle = |\uparrow\rangle$$

$$P|\downarrow\rangle = |\uparrow\rangle$$

$$Q|\downarrow\rangle = -|\downarrow\rangle$$

$$P|\uparrow\rangle = |\downarrow\rangle$$



# Entangled states

As before

2 such systems

$P_1, \Phi_1, P_2, \Phi_2$

$$\Gamma \subseteq (S' \times S') \times (S' \times S')$$

$$\text{is } \left\{ \begin{array}{l} q_1 = q_2 \\ p_1 = p_2 \end{array} \right\}$$

$$Q_1 |\psi\rangle = Q_2 |\psi\rangle$$

$$P_1 |\psi\rangle = P_2 |\psi\rangle$$

$$E_0 \in (S \times S') \times (S' \times S)$$

$$is \left\{ \begin{array}{l} q_1 = q_2 \\ p_1 = -p_2 \end{array} \right\}$$

$$Q_1 |\psi\rangle = Q_2 |\psi\rangle$$

$$P_1 |\psi\rangle = -P_2 |\psi\rangle$$

Write a state as  
 $|\uparrow\rangle \otimes |\downarrow\rangle$  etc

$$\psi = |\uparrow\rangle \otimes |\uparrow\rangle + |\downarrow\rangle \otimes |\downarrow\rangle$$

