

Title: Bridging gaps in the spectrum of GWs: new opportunities for fundamental physics.

Speakers: Diego Blas

Series: Particle Physics

Date: January 24, 2023 - 1:00 PM

URL: <https://pirsa.org/23010058>

Abstract: In this talk I'll highlight the existence of gaps in the spectrum of GWs poorly explored by current observations and that may contain information about BSM physics or primordial cosmology. I'll focus on the muHz gap, and explain how to use the resonance absorption of GWs by binary systems (as the Earth-Moon system or binary pulsars) to access this band. I'll also highlight the potential of high frequency ($\omega > \text{kHz}$) GWs. The focus of this second part will be the use of cavities to detect these signals, together with a brief discussion of sources and what can be learned from them.

Zoom Link: <https://pitp.zoom.us/j/94501758748?pwd=amFsTjV5Z1c5SHJrVVBCanVCRHRMUT09>

Bridging gaps in the spectrum of GWS: *new opportunities for fundamental physics*

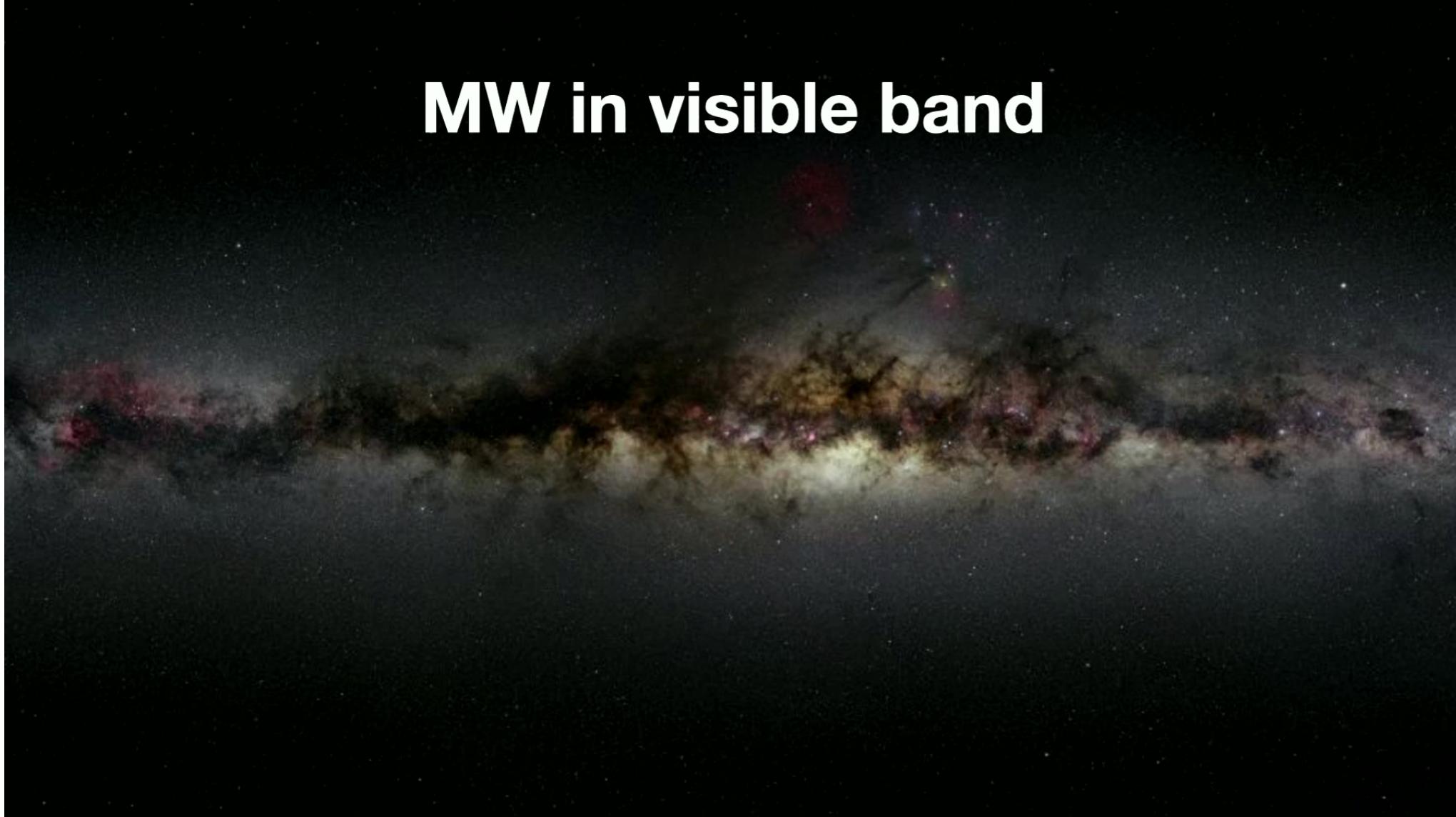
Diego Blas

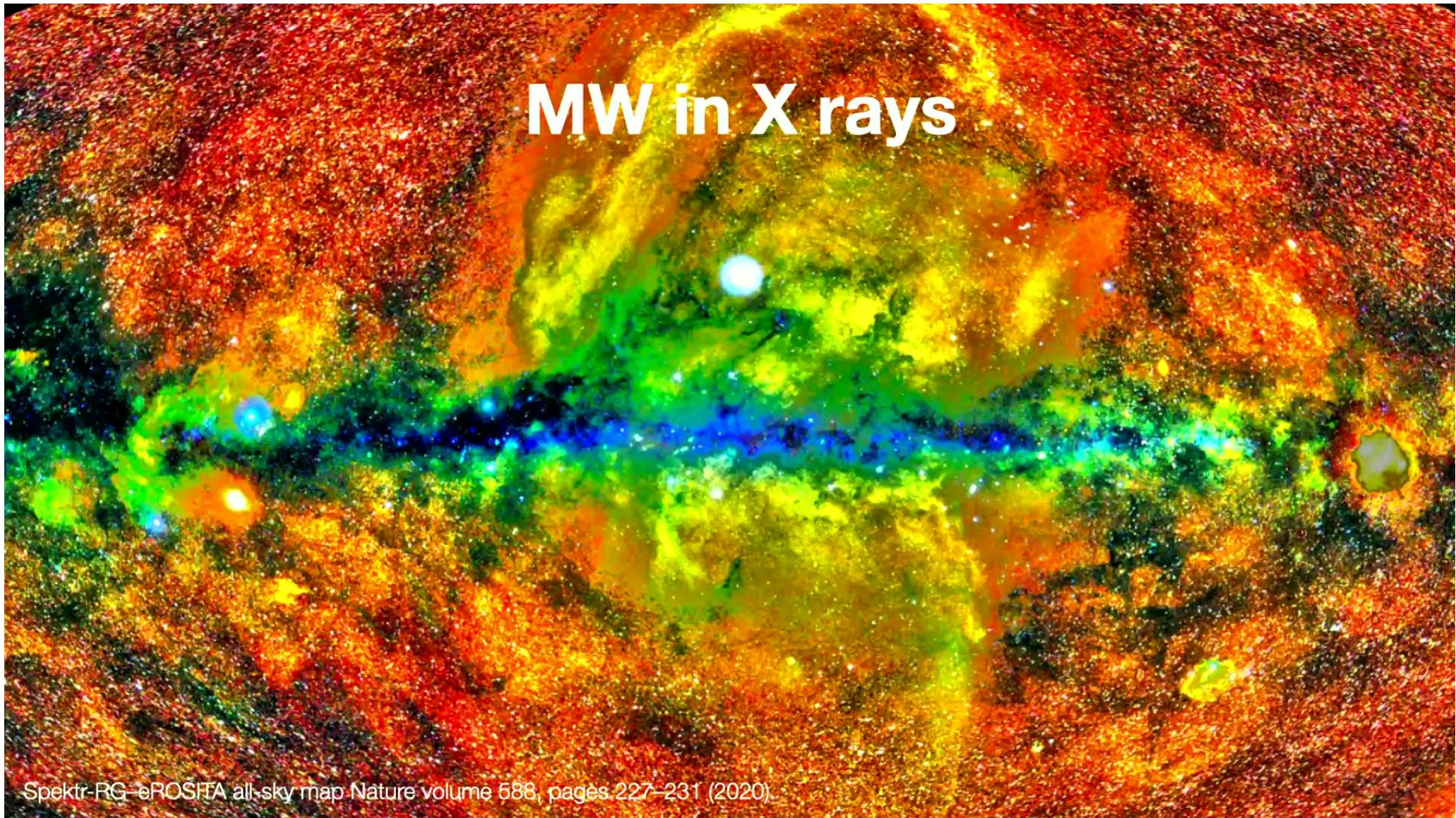
based on 2107.04063/2107.04601 (PRL/PRD22), 2112.11465 (PRD22) + ongoing (2301.xxxx?)

(w. Alex Jenkins // A. Berlin, DB, R. T. D'Agnolo, S. Ellis, R. Harnik, Y. Kahn, J. Schütte-Engel)



MW in visible band



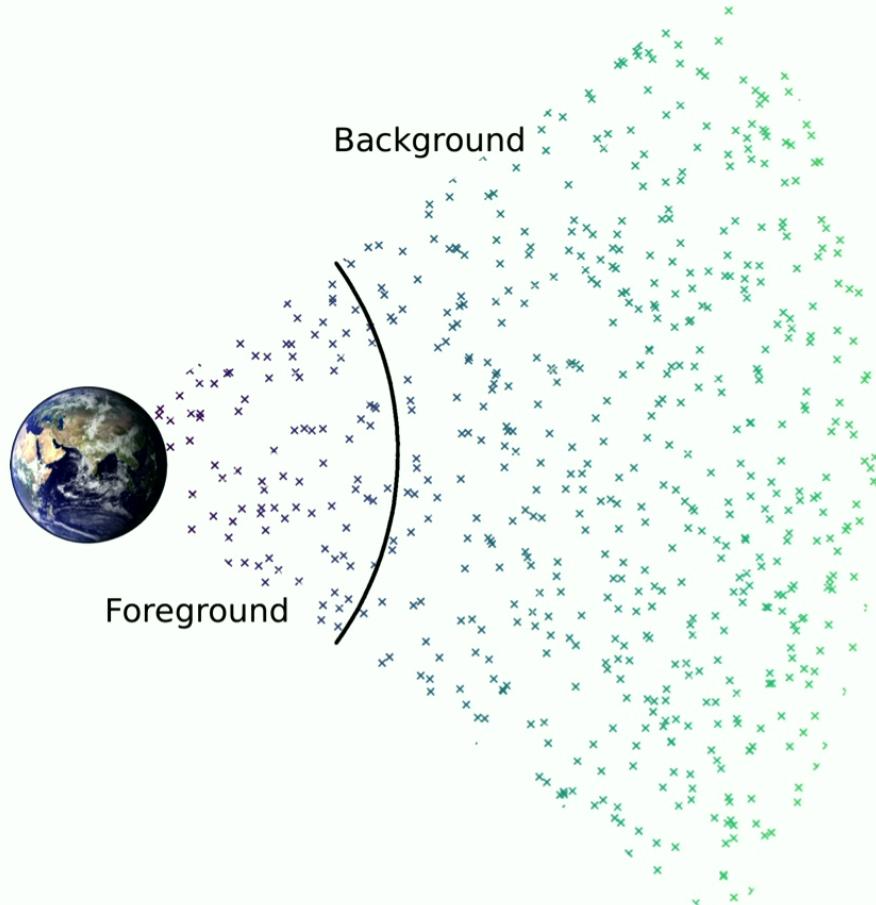


MW in X rays

Spektr-RG-eROSITA all-sky map Nature volume 588, pages 227–231 (2020).

GWs soundscape today

Stochastic gravitational-wave background (SGWB)



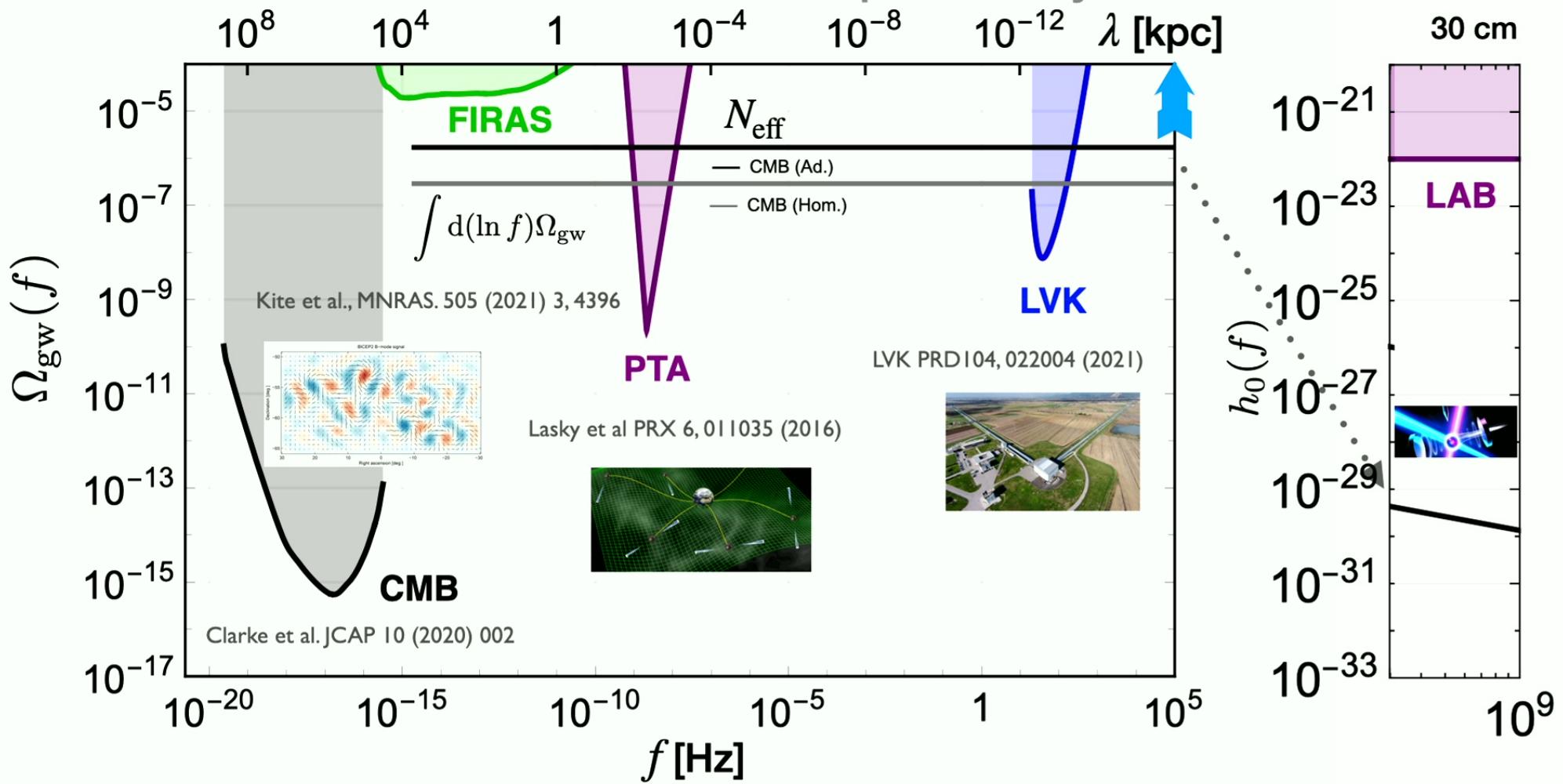
dblasis@ifae.es

- incoherent, persistent GW signal
- faint/numerous sources
- astrophysical and cosmological
- GW density parameter:

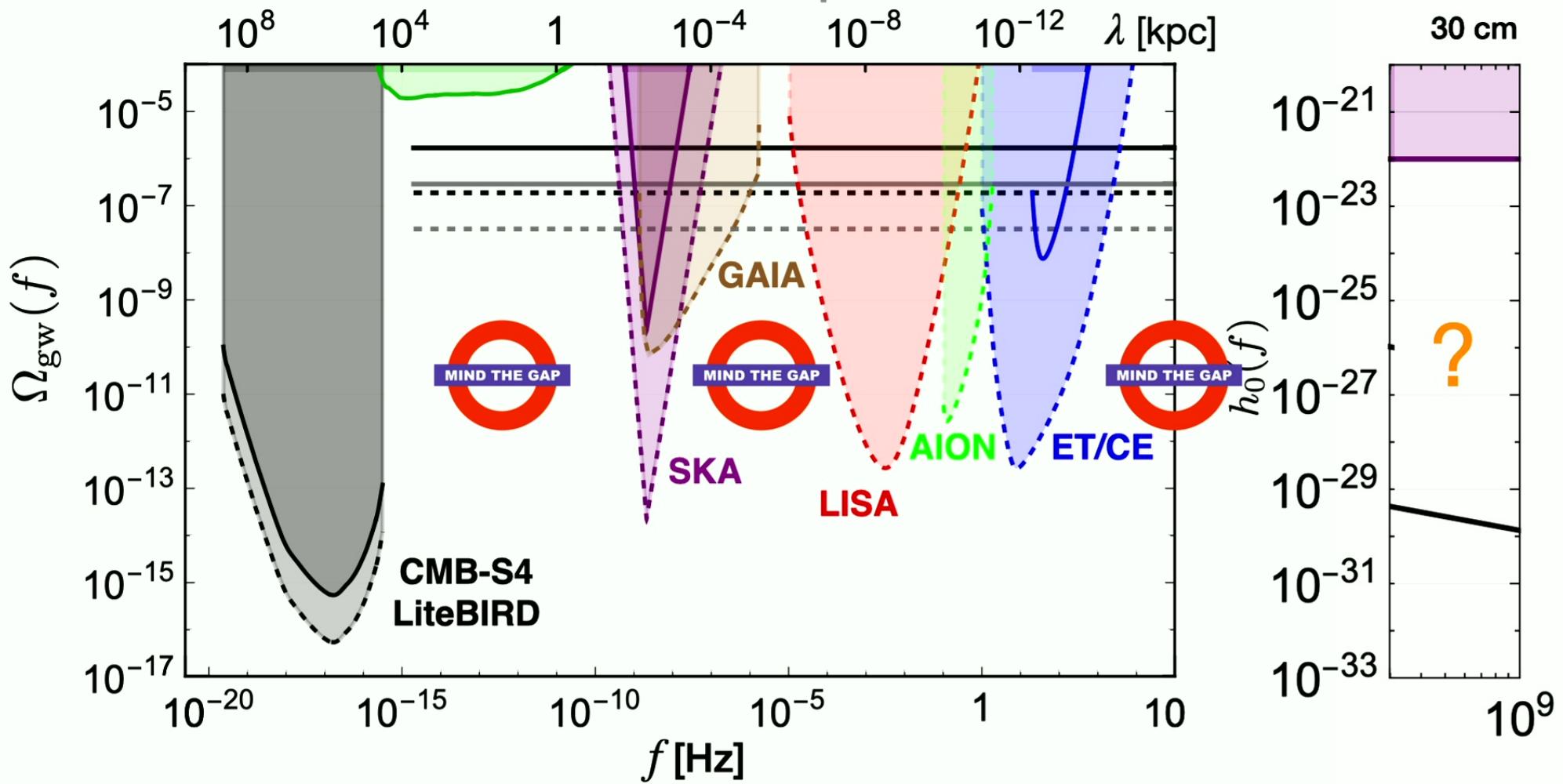
$$\Omega_{\text{GW}}(f) = \frac{1}{\rho_{\text{crit}}} \frac{d\rho_{\text{GW}}}{d(\ln f)}$$
$$\rho_{\text{GW}} \sim M_P^2 \omega^2 h_{\text{GW}}^2$$

$$\rho_c = 1.2 \times 10^{11} M_\odot \text{Mpc}^{-3}$$
$$\sim \text{keV/cm}^3$$

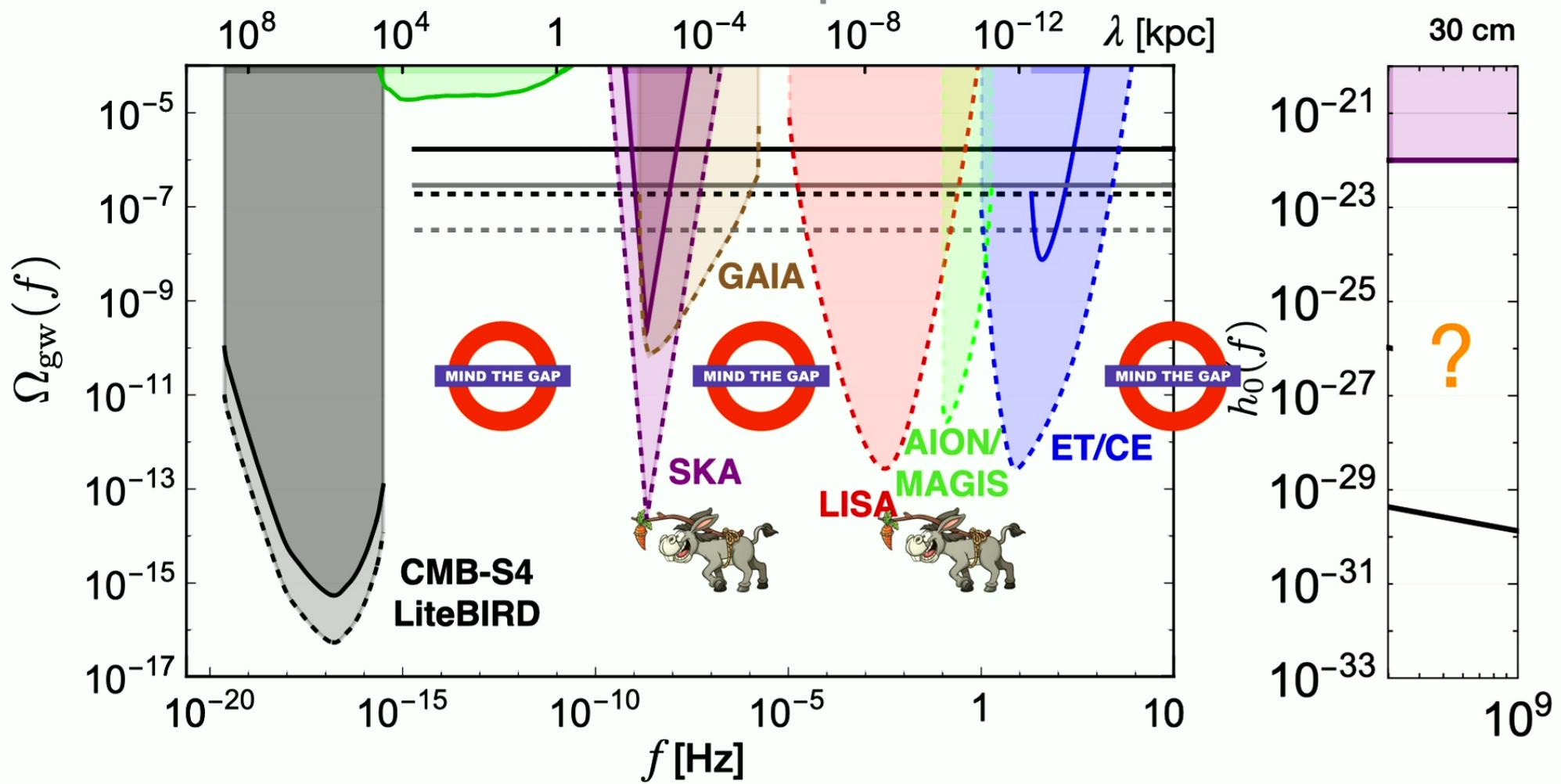
GWs soundscape today



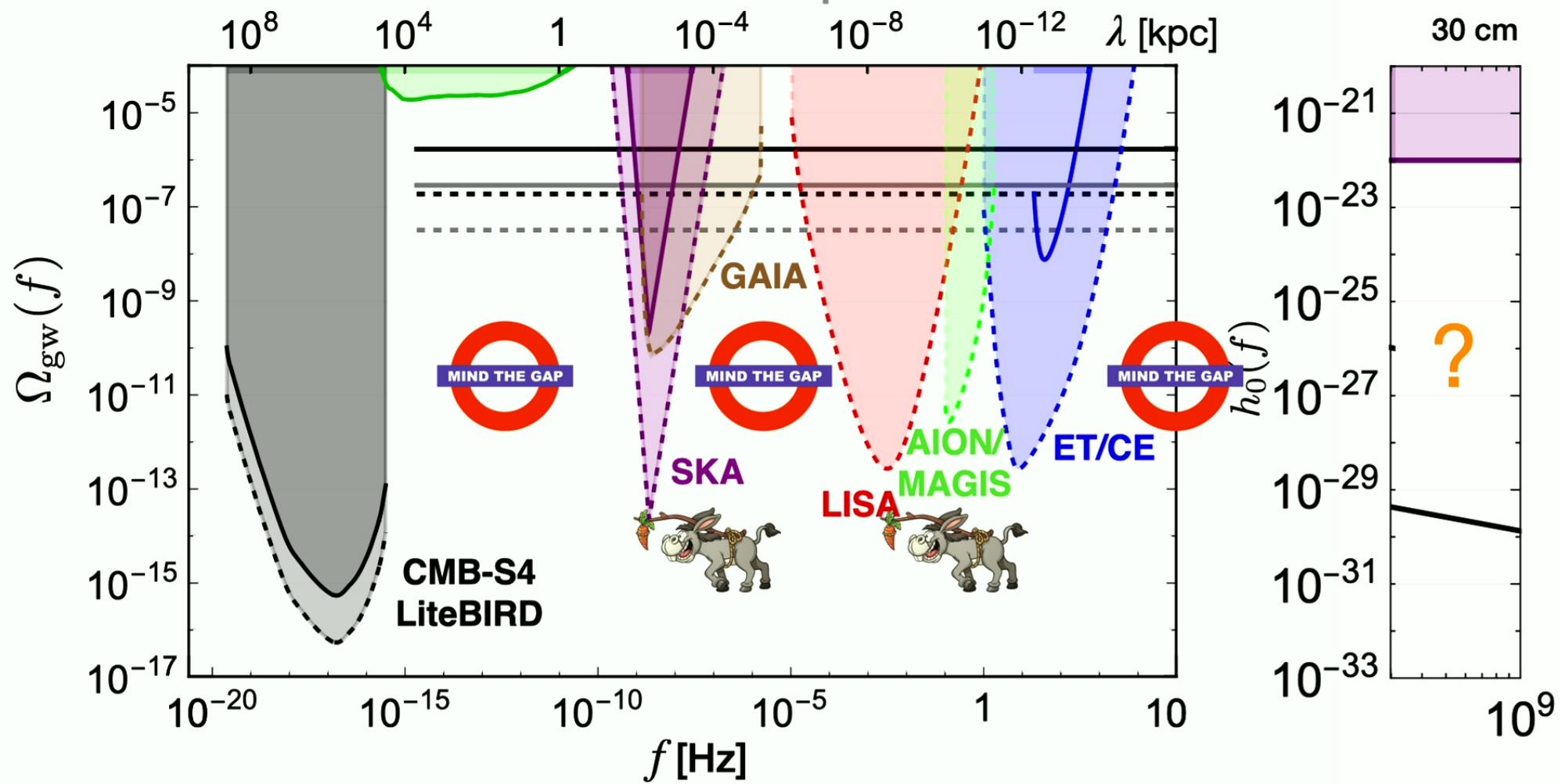
GWs soundscape ca. 2040



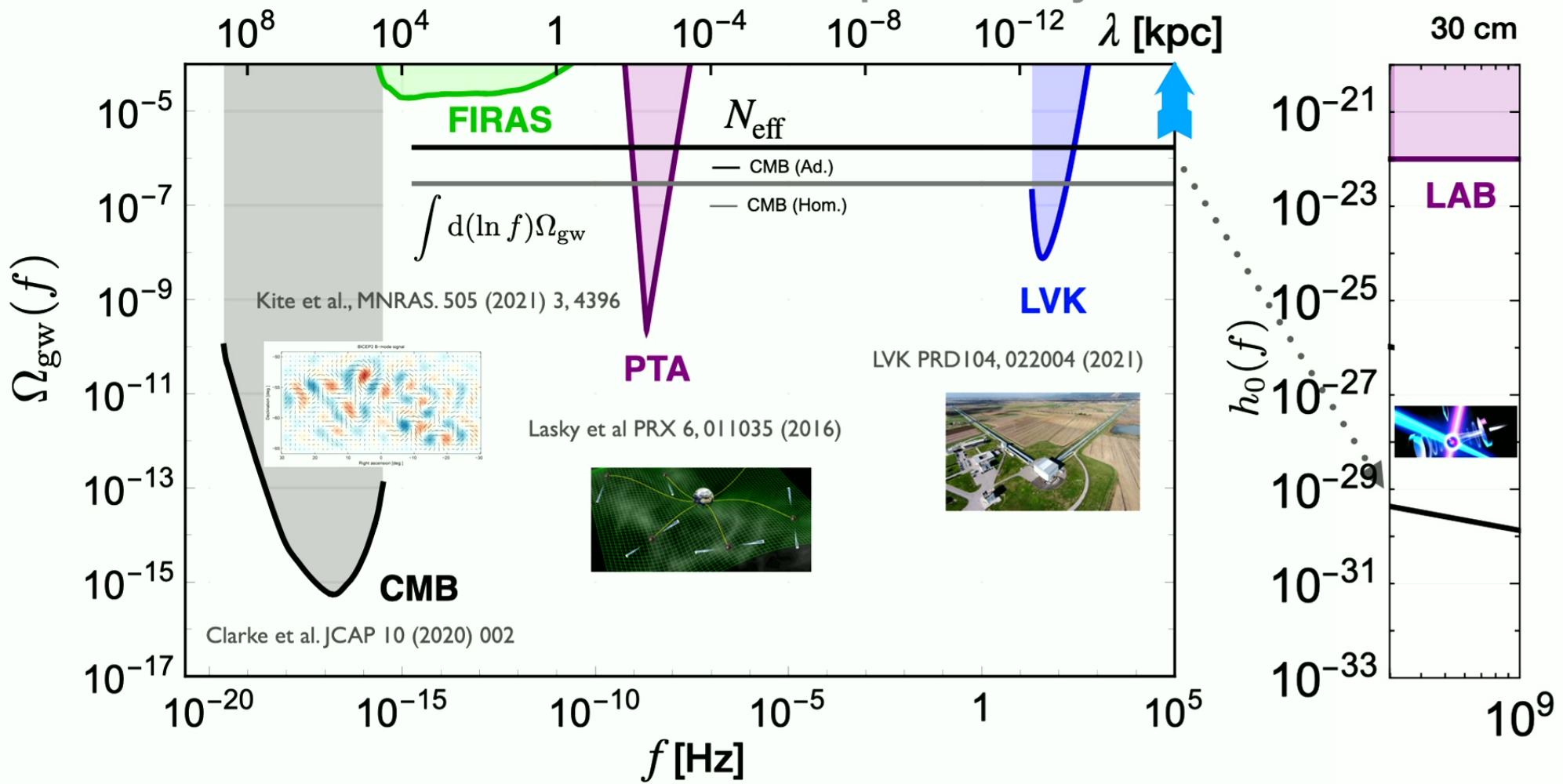
GWs soundscape ca. 2040



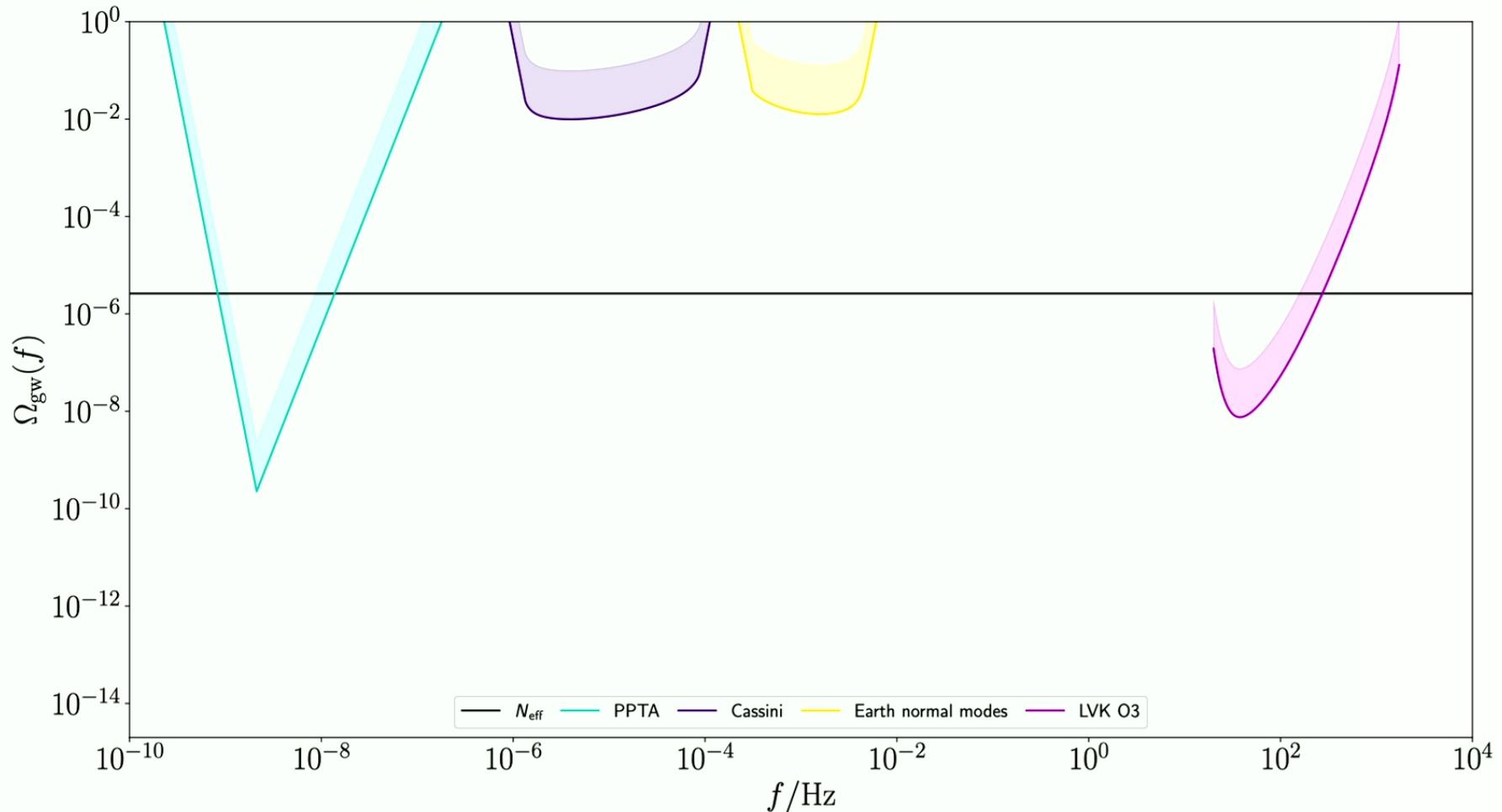
GWs soundscape ca. 2040



GWs soundscape today

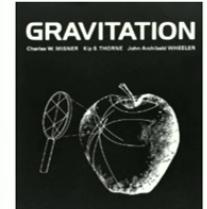


Current SGWB constraints



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Absorption of GWs by binaries



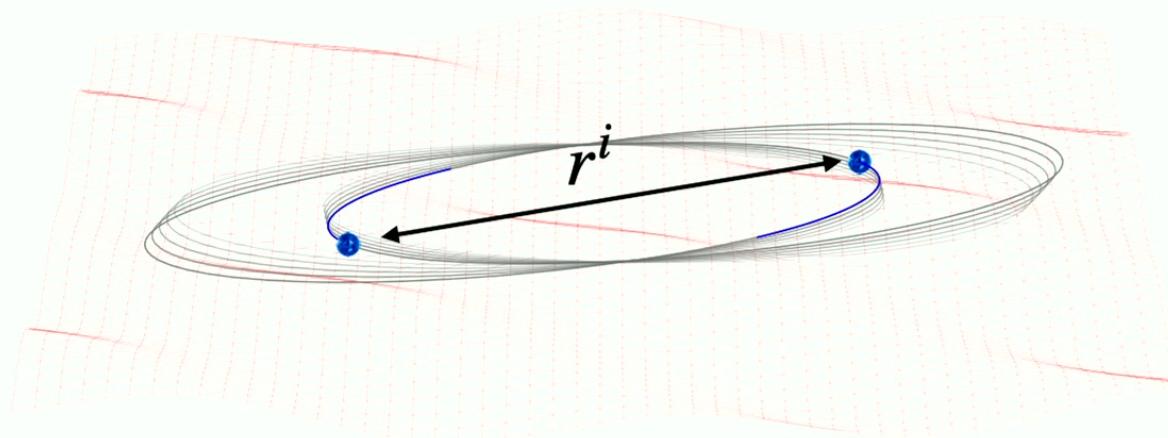
MWT
Hui et al 2013
DB, Lopez-Nacir, Sibiryakov 2017

Intuitive idea (from '60s)

Influence of a GW on a binary system (e.g. non-relativistic)

$$\ddot{r}^i + \frac{GM}{r^3}r^i = \delta^{ik}\frac{1}{2}\ddot{h}_{kj}r^j$$

Newtonian potential ... GW

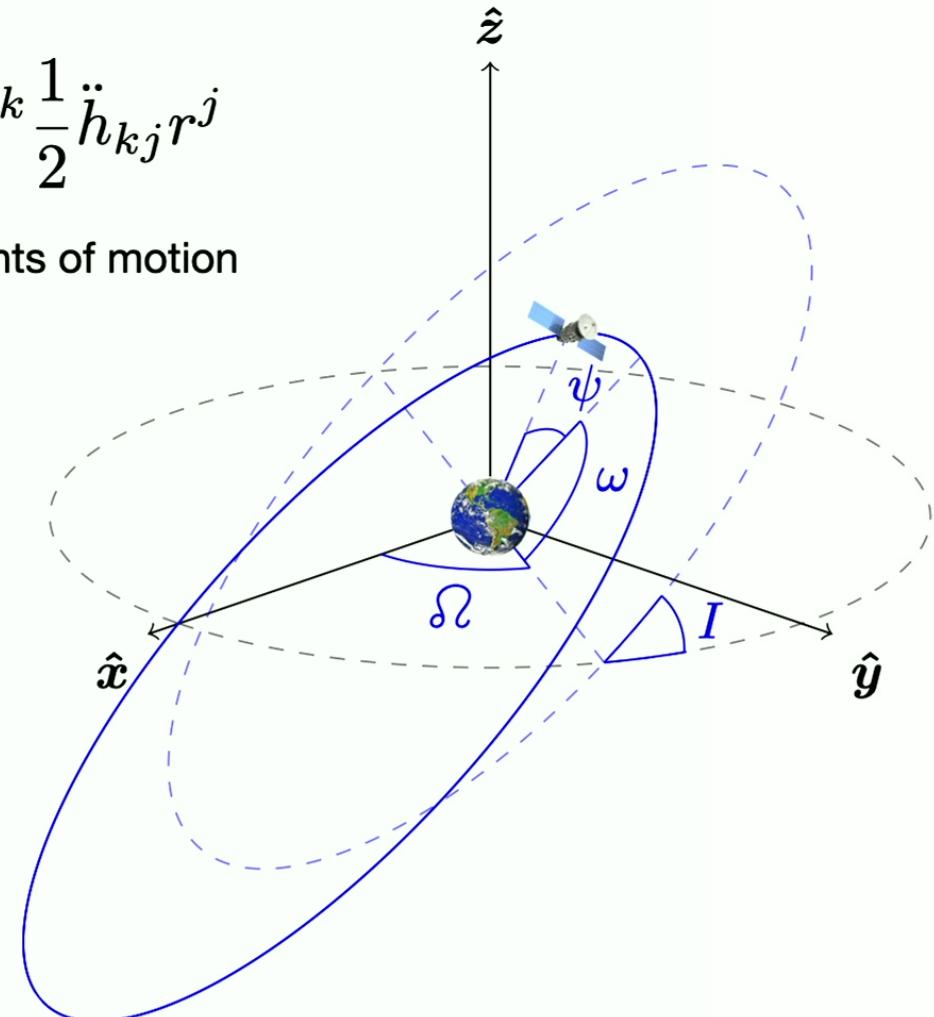


Absorption of GWs by binaries

$$\ddot{r}^i + \frac{GM}{r^3} r^i = \delta^{ik} \frac{1}{2} \ddot{h}_{kj} r^j$$

Better characterised for its 6 Newtonian constants of motion

- **period P , eccentricity e :**
size and shape of orbit
- **inclination I , ascending node Ω :**
orientation in space
- **pericentre ω ,**
mean anomaly at epoch ε :
radial and angular phases



Absorption of GWs by binaries

$$\ddot{\mathbf{r}} + \frac{GM}{r^2} \hat{\mathbf{r}} = \delta \ddot{\mathbf{r}}.$$

■ for generic perturbation:

$$\delta \ddot{\mathbf{r}} = r(\mathcal{F}_r \hat{\mathbf{r}} + \mathcal{F}_\theta \hat{\boldsymbol{\theta}} + \mathcal{F}_\ell \hat{\boldsymbol{\ell}}),$$



$$\begin{aligned}\dot{P} &= \frac{3P^2\gamma}{2\pi} \left[\frac{e \sin \psi \mathcal{F}_r}{1 + e \cos \psi} + \mathcal{F}_\theta \right], \\ \dot{e} &= \frac{\dot{P}\gamma^2}{3Pe} - \frac{P\gamma^5 \mathcal{F}_\theta}{2\pi e(1 + e \cos \psi)^2}, \\ \dot{I} &= \frac{P\gamma^3 \cos \theta \mathcal{F}_\ell}{2\pi(1 + e \cos \psi)^2}, \\ \dot{\Omega} &= \frac{\tan \theta}{\sin I} \dot{I}, \\ \dot{\omega} &= \frac{P\gamma^3}{2\pi e} \left[\frac{(2 + e \cos \psi) \sin \psi \mathcal{F}_\theta}{(1 + e \cos \psi)^2} - \frac{\cos \psi \mathcal{F}_r}{1 + e \cos \psi} \right] - \cos I \dot{\Omega}, \\ \dot{\varepsilon} &= -\frac{P\gamma^4 \mathcal{F}_r}{\pi(1 + e \cos \psi)^2} - \gamma(\cos I \dot{\Omega} + \dot{\omega}),\end{aligned}$$

For the SGWB... Fokker-Planck approach

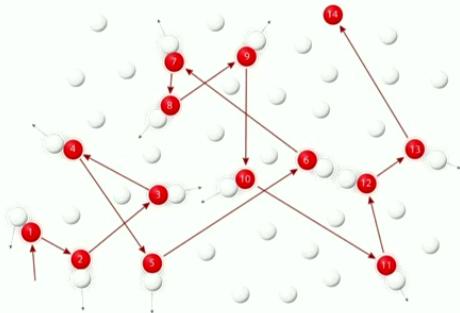
$$\ddot{r}^i + \frac{GM}{r^3}r^i = \delta^{ik}\frac{1}{2}\ddot{h}_{kj}r^j$$

deterministic

$$\dot{X}_i(\mathbf{X}, t) = V_i(\mathbf{X}) + \Gamma_i(\mathbf{X}, t)$$

stochastic

we move from dynamics of the variable to dynamics of the **distribution $W(\mathbf{X})$**



$$\frac{\partial W}{\partial t} = -\partial_i \left(D_i^{(1)} W \right) + \partial_i \partial_j \left(D_{ij}^{(2)} W \right)$$

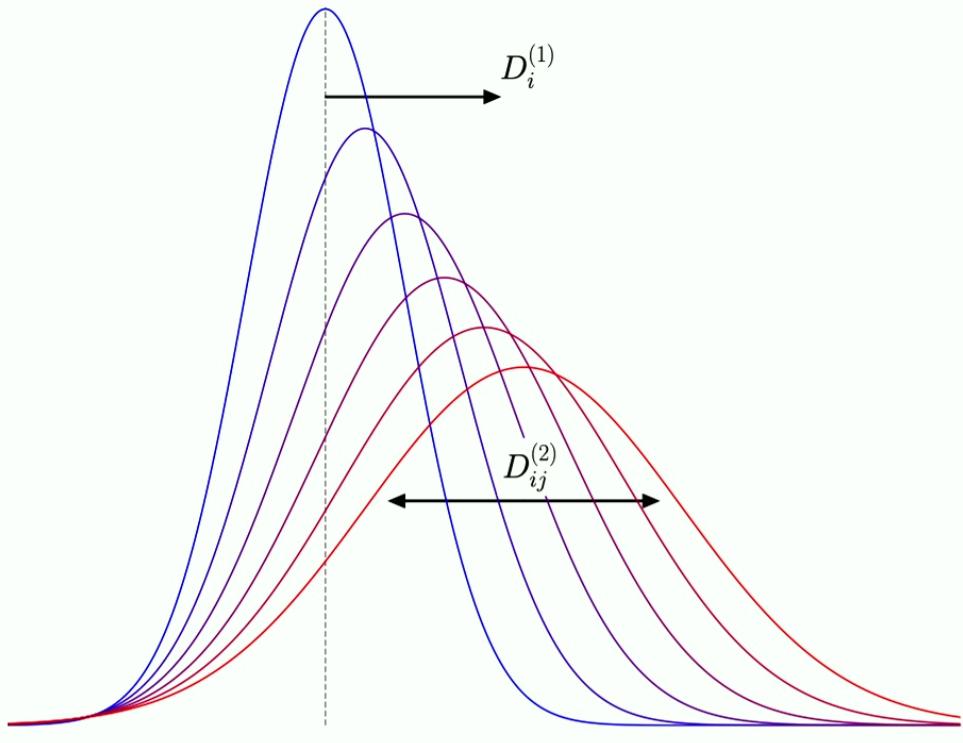
with $\partial_i \equiv \partial/\partial X_i$

$$D_i^{(1)} = V_i + \lim_{\tau \rightarrow 0} \frac{1}{\tau} \int_t^{t+\tau} dt' \int_t^{t'} dt'' \langle \Gamma_j(\mathbf{x}, t'') \partial_j \Gamma_i(\mathbf{x}, t') \rangle.$$

$$D_{ij}^{(2)} = \lim_{\tau \rightarrow 0} \frac{1}{2\tau} \int_t^{t+\tau} dt' \int_t^{t+\tau} dt'' \langle \Gamma_i(\mathbf{x}, t') \Gamma_j(\mathbf{x}, t'') \rangle.$$

Our approach to the problem

Blas&Jenkins Phys.Rev.Lett. 128 (2022) 10, 101103



- track distribution function $W(\mathbf{X}, t)$ of orbital elements $\mathbf{X} = (P, e, I, \delta\Omega, \omega, \varepsilon)$
- evolves through *Fokker-Planck eqn.*

$$\frac{\partial W}{\partial t} = -\frac{\partial}{\partial X_i}(D_i^{(1)} W) + \frac{\partial}{\partial X_i} \frac{\partial}{\partial X_j}(D_{ij}^{(2)} W)$$

- *drift and diffusion coefficients*
(averaged over orbits)

$$D_i^{(1)}(\mathbf{X}) = V_i(\mathbf{X}) + \sum_{n=1}^{\infty} \mathcal{A}_{n,i}(\mathbf{X}) \Omega_{\text{gw}}(n/P)$$

$$D_{ij}^{(2)}(\mathbf{X}) = \sum_{n=1}^{\infty} \mathcal{B}_{n,ij}(\mathbf{X}) \Omega_{\text{gw}}(n/P)$$

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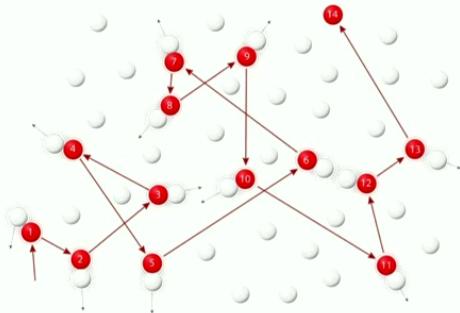
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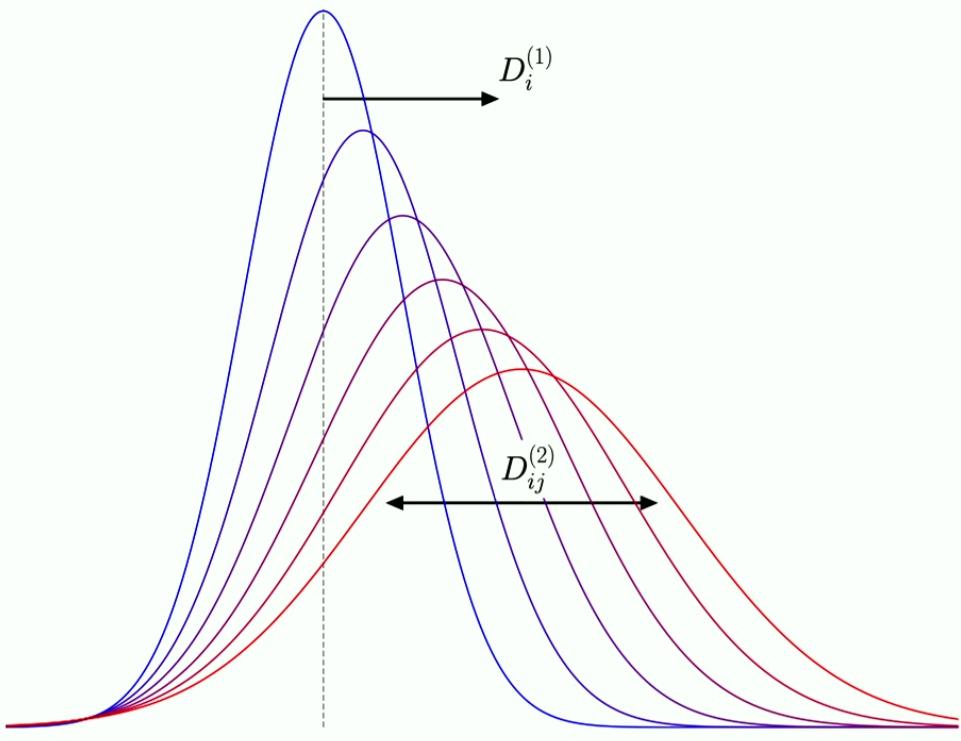
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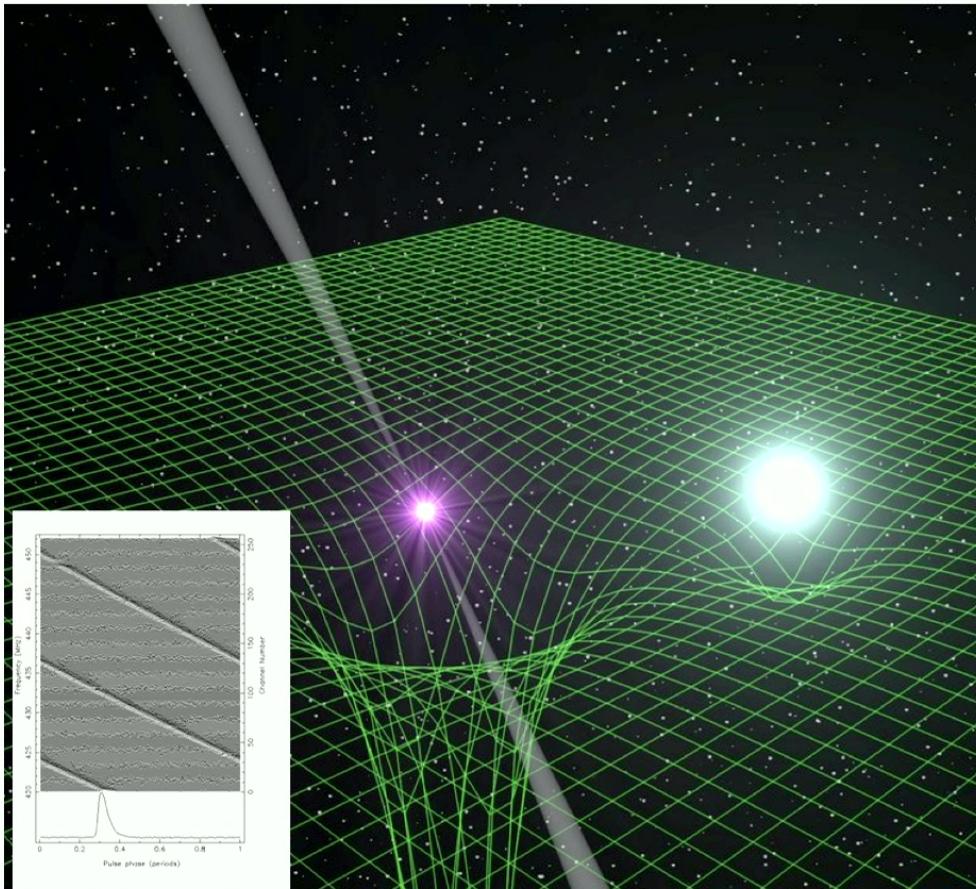
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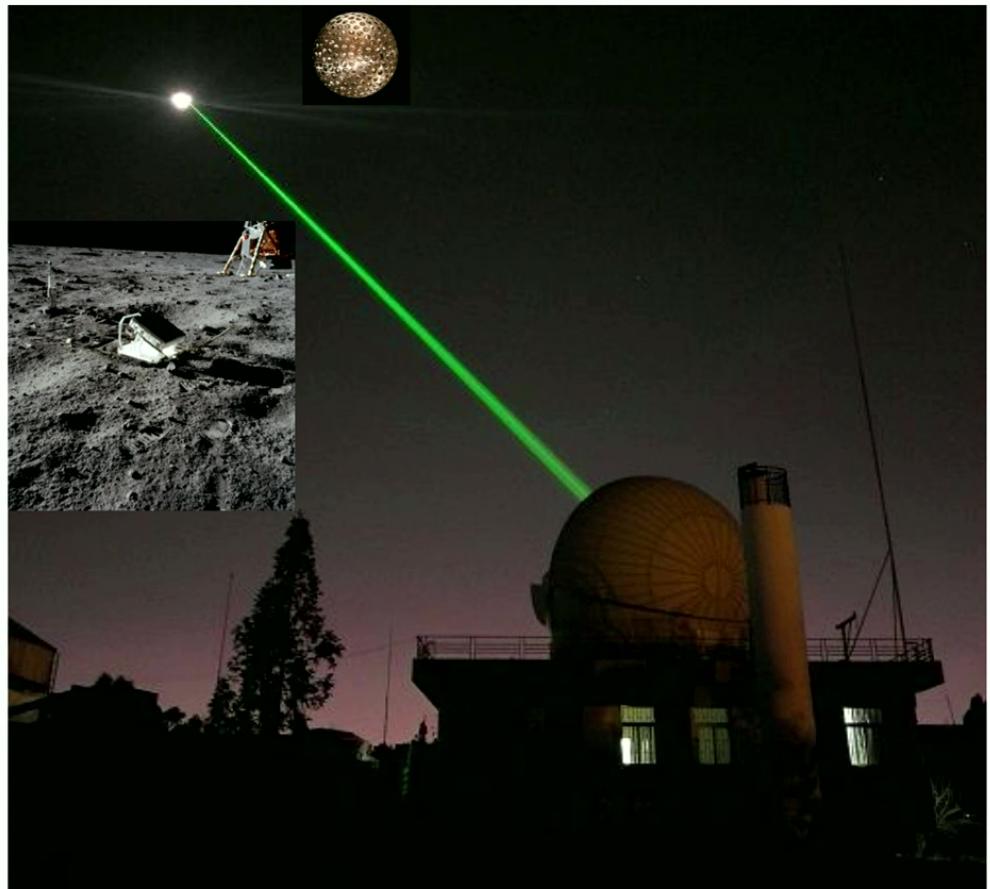
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Two probes

timing of binary pulsars

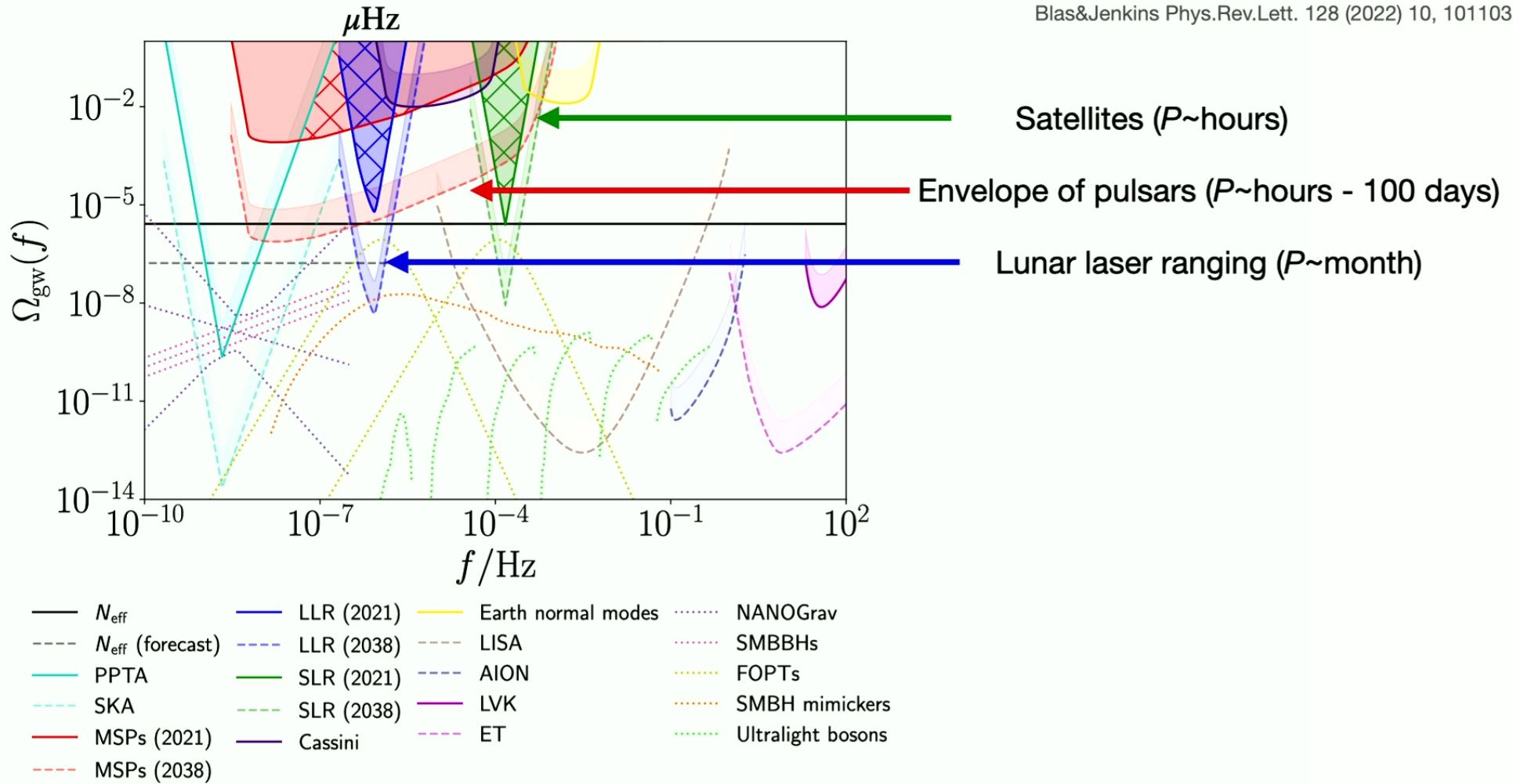


lunar and satellite laser ranging

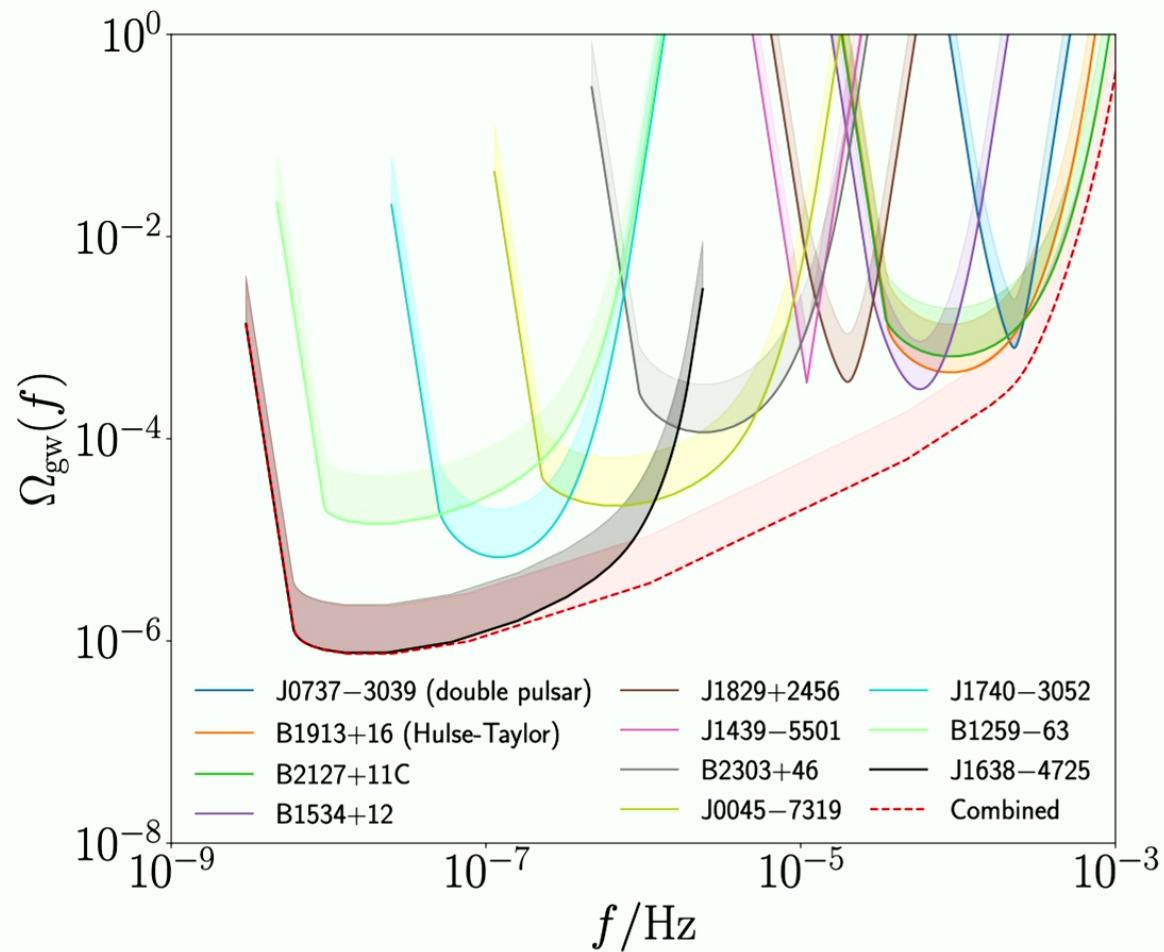


Our estimates

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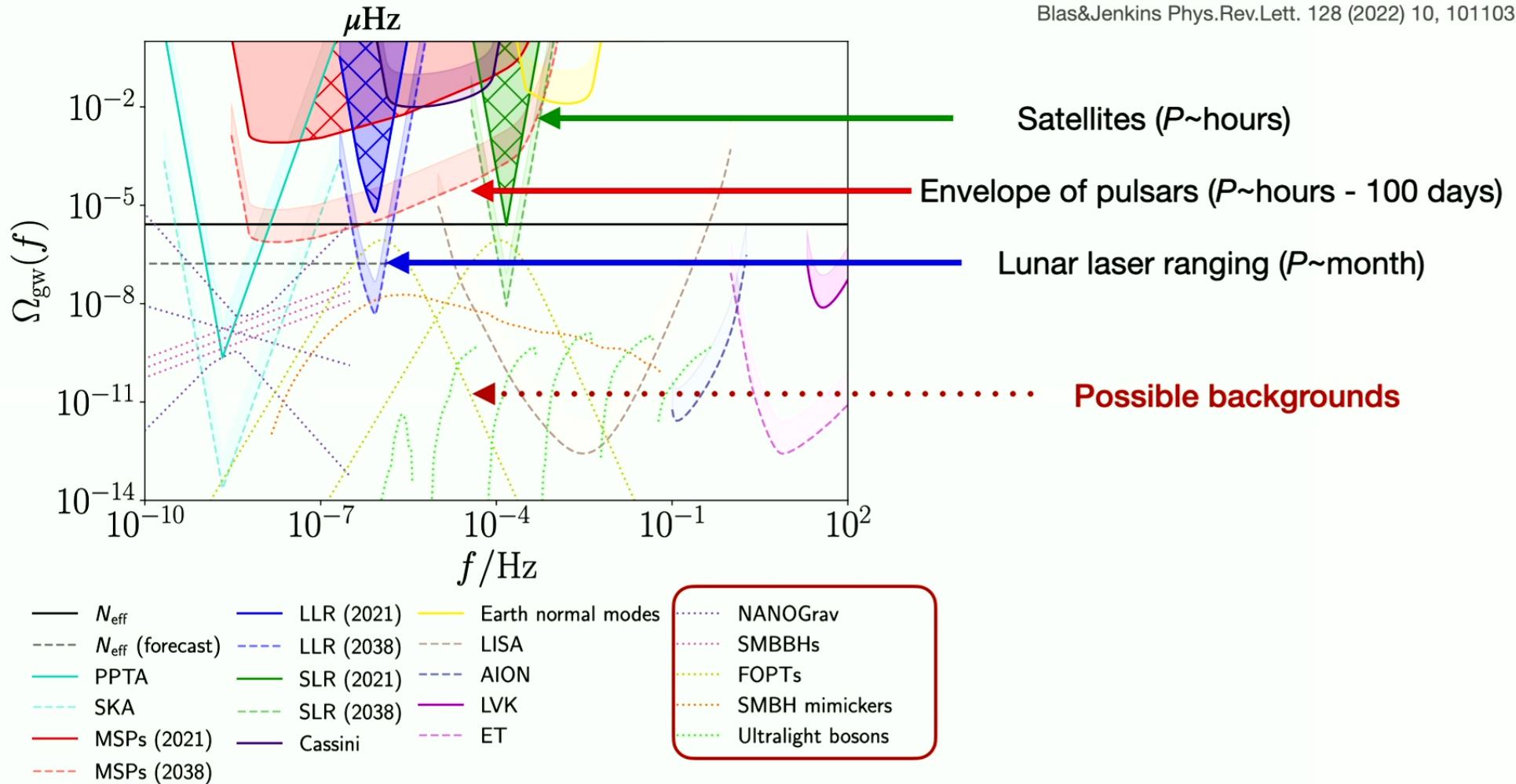


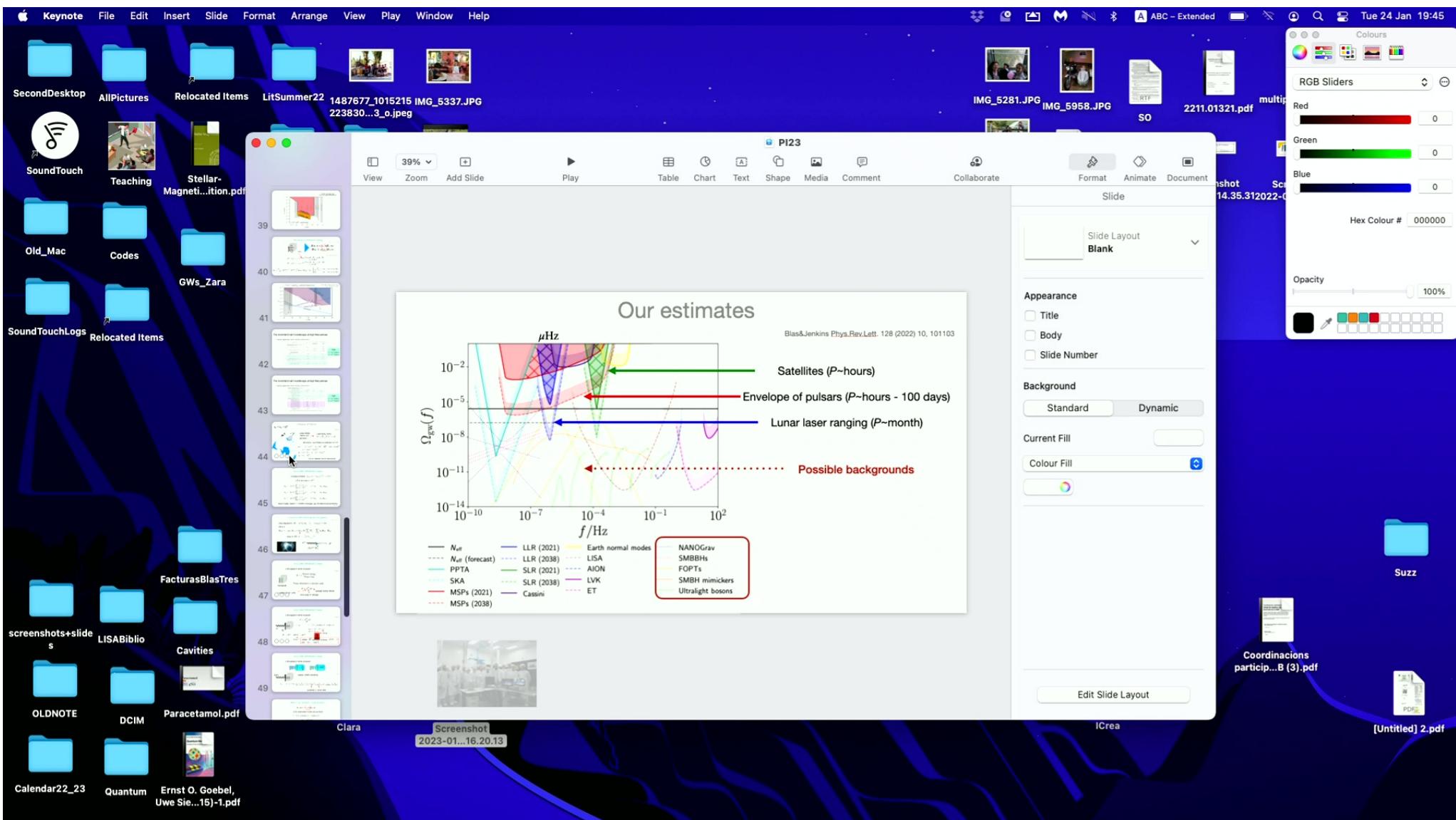
Combining binary pulsar bounds



Our estimates

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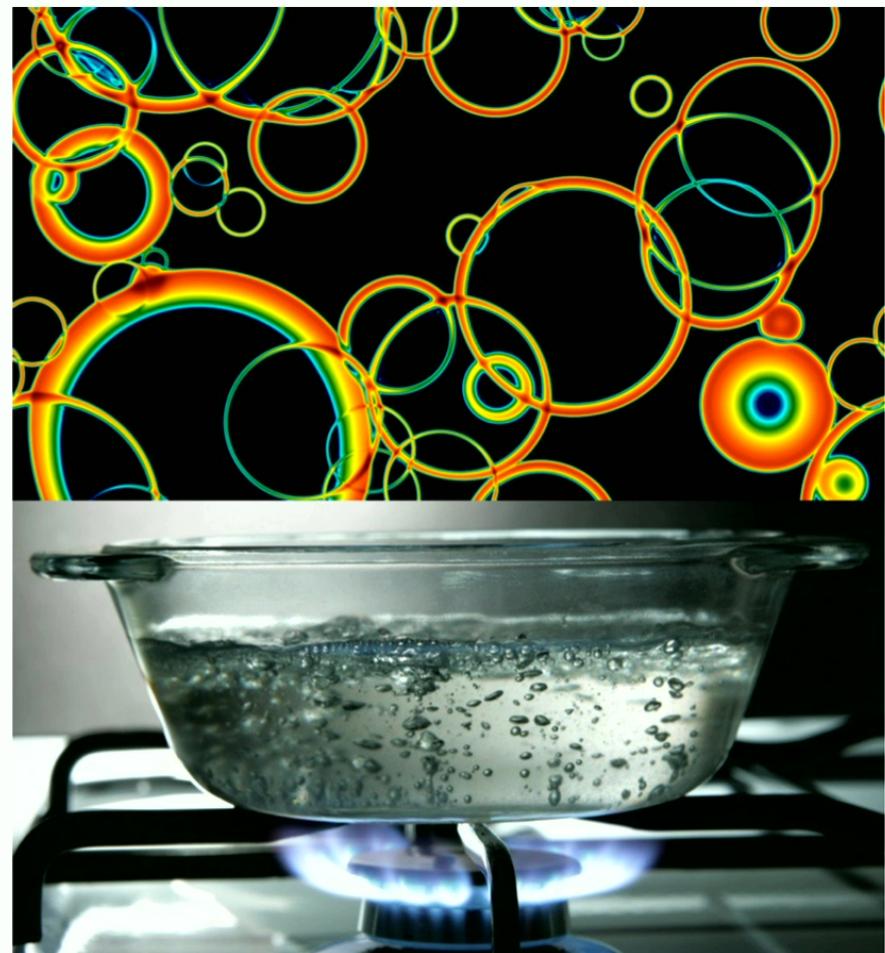




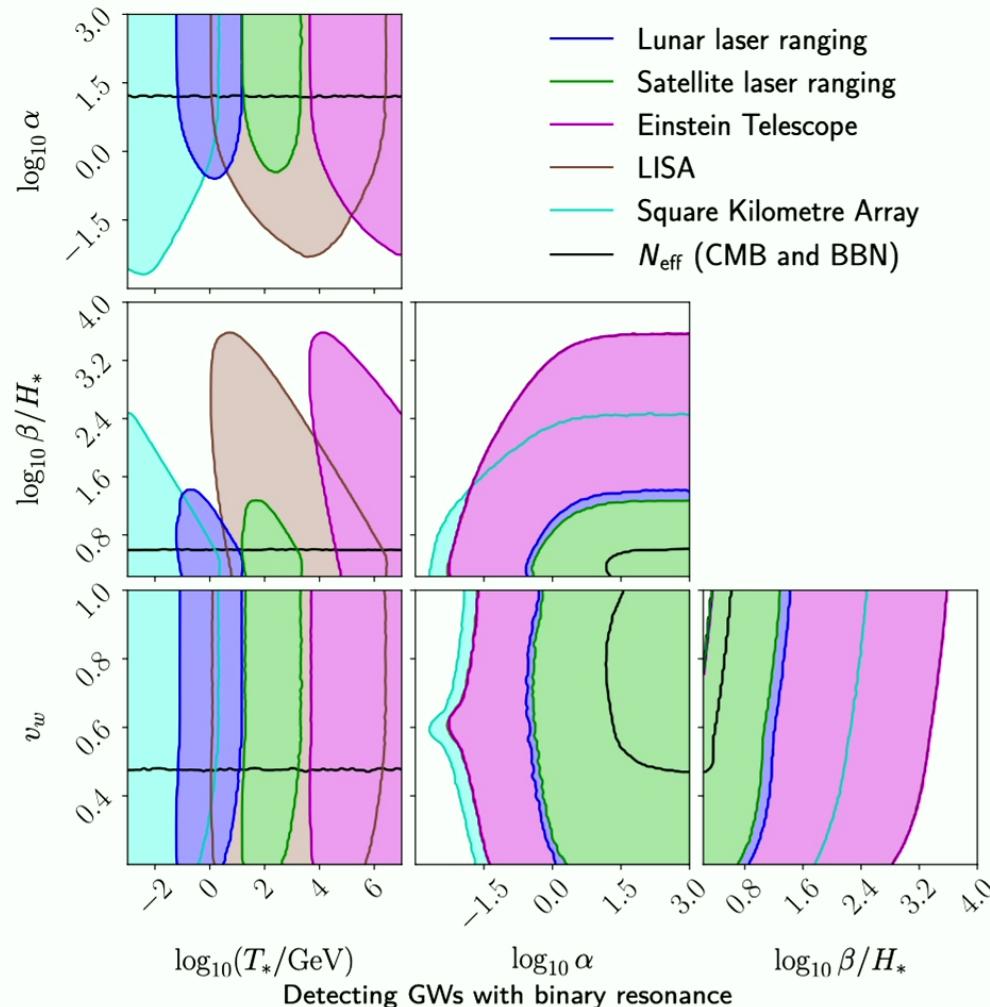
An example signal: cosmological phase transitions

- key prediction of many particle physics models
- four parameters:
 - ▶ temperature T_*
 - ▶ strength α
 - ▶ rate β/H_*
 - ▶ bubble-wall velocity v_w
- peak frequency

$$f_* \approx 19 \mu\text{Hz} \times \frac{T_*}{100 \text{ GeV}} \frac{\beta/H_*}{v_w}$$



Phase transition constraints



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Absorption of GWs by binaries

$$\ddot{r}^i + \frac{GM}{r^3} r^i = \delta^{ik} \frac{1}{2} \dot{h}_{kj} r^j$$

Better characterised for its 6 Newtonian constants of motion

- period P , eccentricity e : size and shape of orbit
- inclination I , ascending node Ω : orientation in space
- pericentre ω , mean anomaly at epoch ε : radial and angular phases

alexander.jenkins@kcl.ac.uk Detecting GWs with binary resonance EPS-HE1, arXiv:1707.08484 [astro-ph.CO] 2/2

Screenshot 2023-01-16 20.13

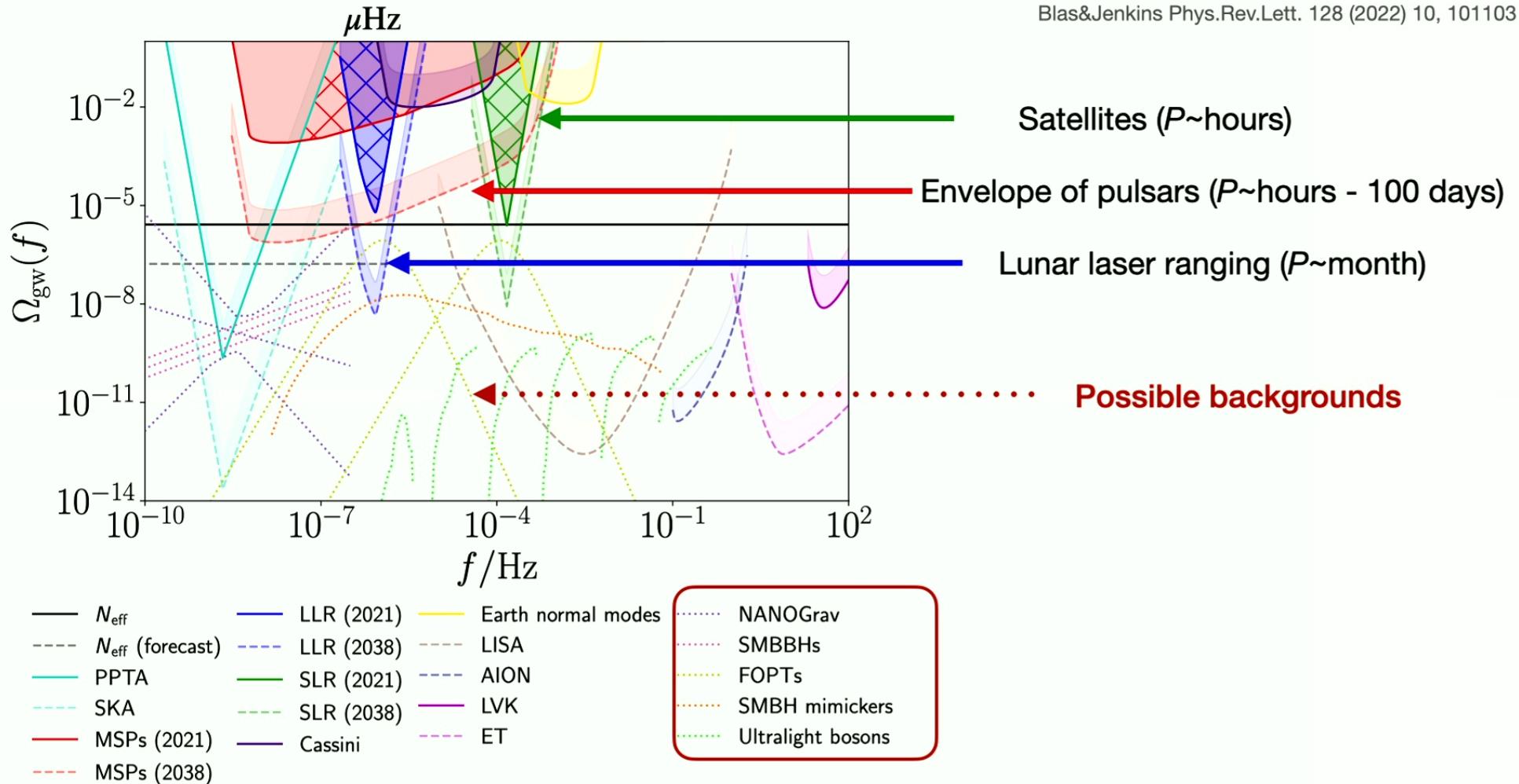
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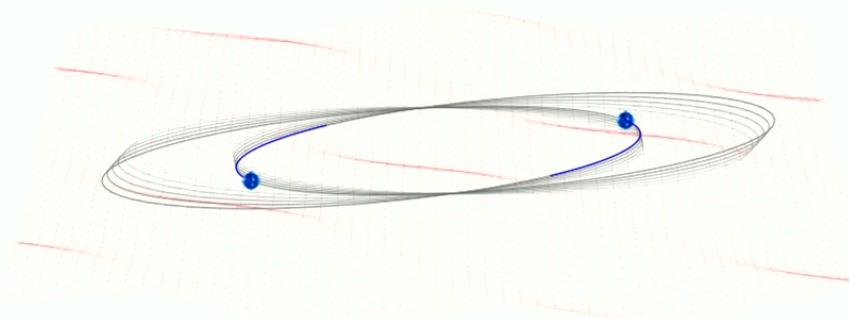
Our estimates

Blas&Jenkins Phys.Rev.Lett. 128 (2022) 10, 101103



Summary and outlook @ μHz

- binary resonance can probe a unique GW frequency band
- we have developed a new formalism (Fokker-Planck treatment)
- unique constraints on phase transitions (and more)
- plenty more work to do! more signals, more systems, plus running on real data



For the SGWB... Fokker-Planck approach

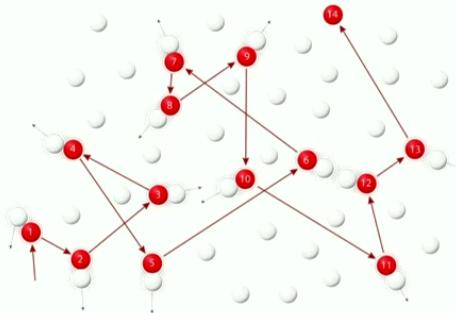
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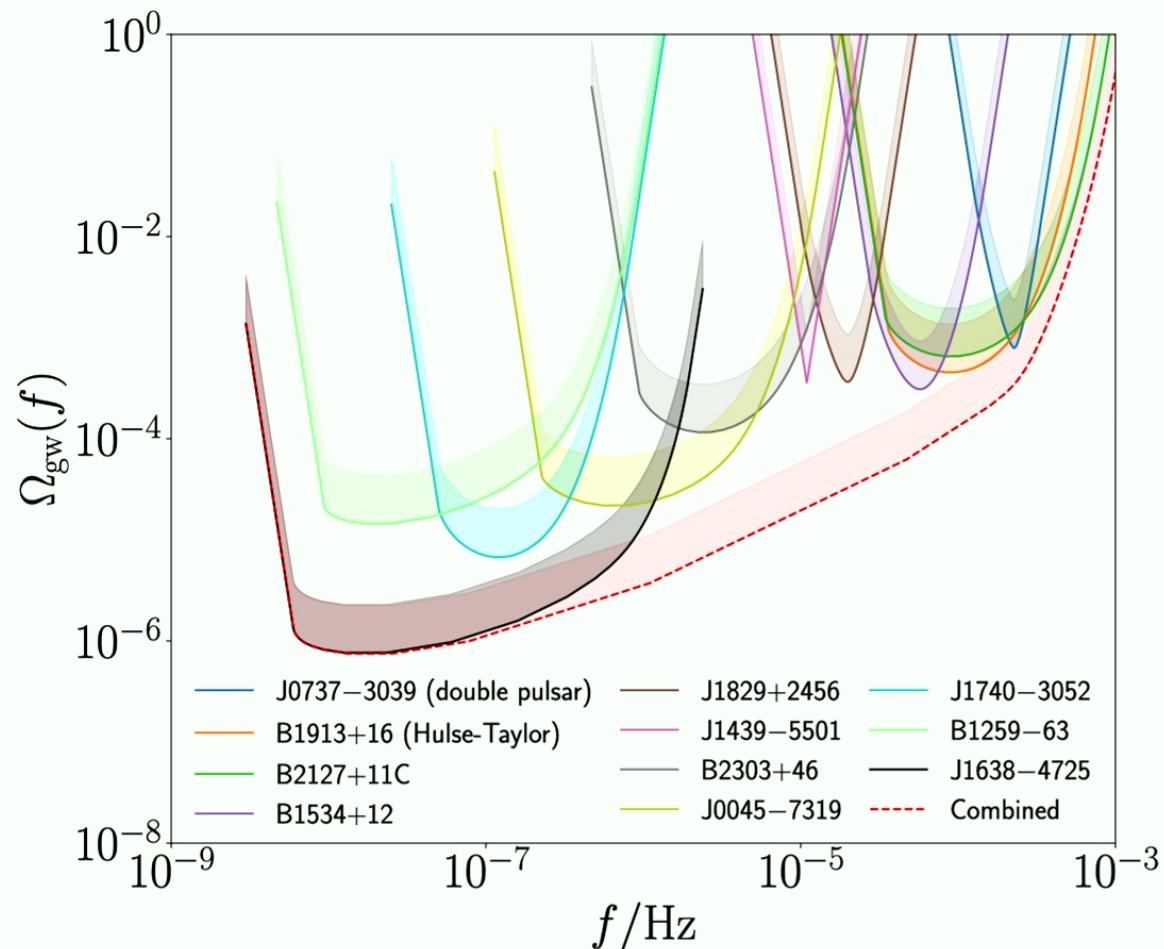
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Combining binary pulsar bounds



Absorption of GWs by binaries

$$\ddot{\mathbf{r}} + \frac{GM}{r^2} \hat{\mathbf{r}} = \delta \ddot{\mathbf{r}}.$$

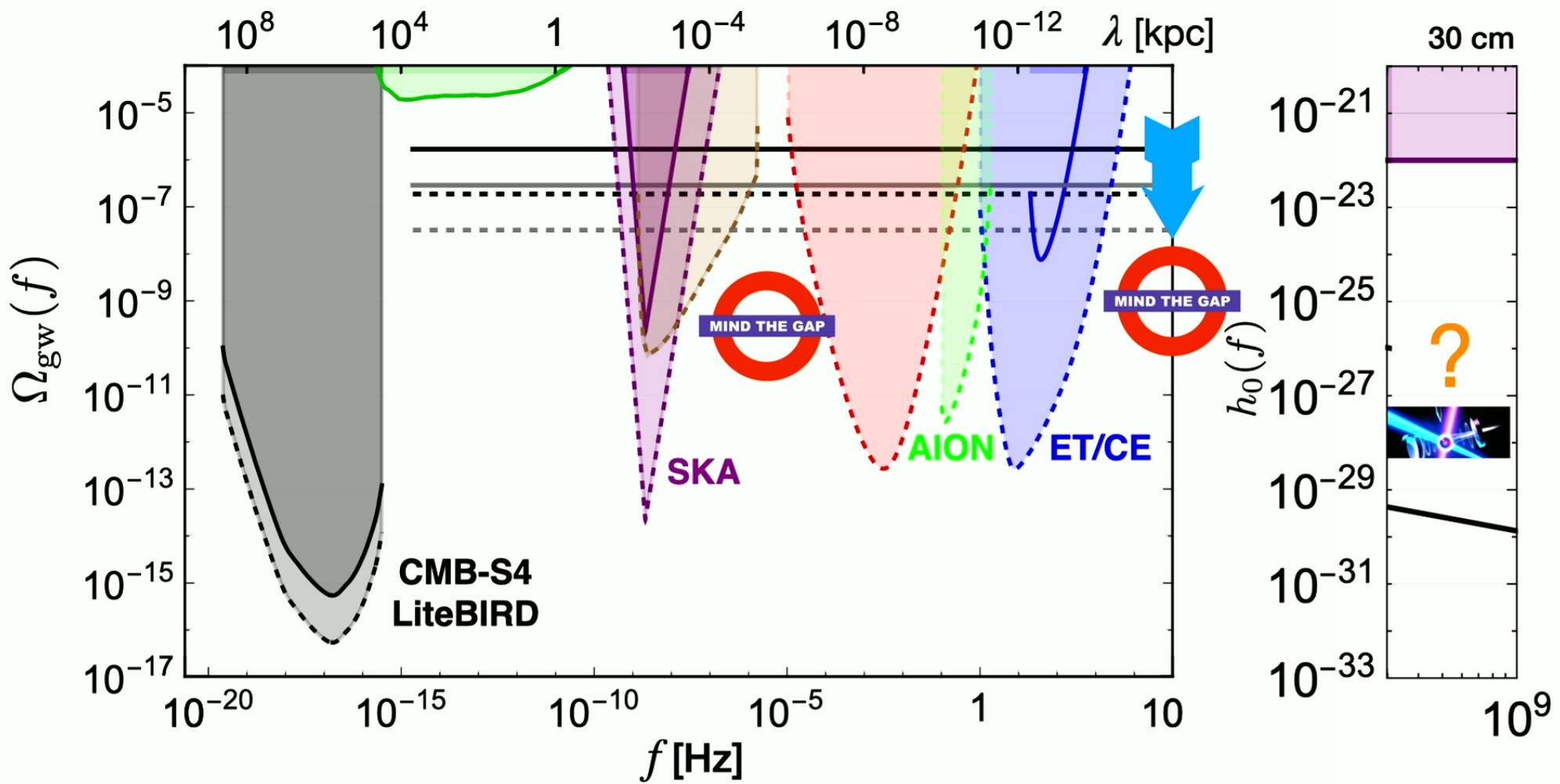
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UHFGWs



Sources of UHFGW

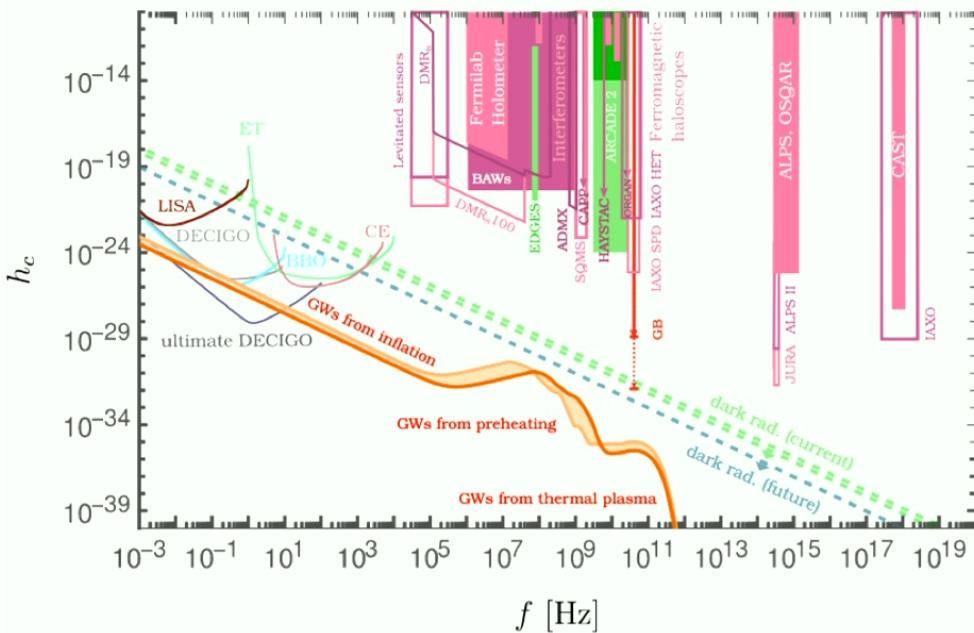
Very different picture comparing to EM spectrum

	Stochastic	Coherent
standard physics	Primordial plasma thermal fluctuations Ghiglieri et al JHEP 07 (2020) 092 Ringwald et al JCAP 03 (2021) 054	?
	Inflation	
new physics	Phase transitions, Cosmic strings Close encounters of BHs	Light PBHB Ultra-light DM Exotic objects

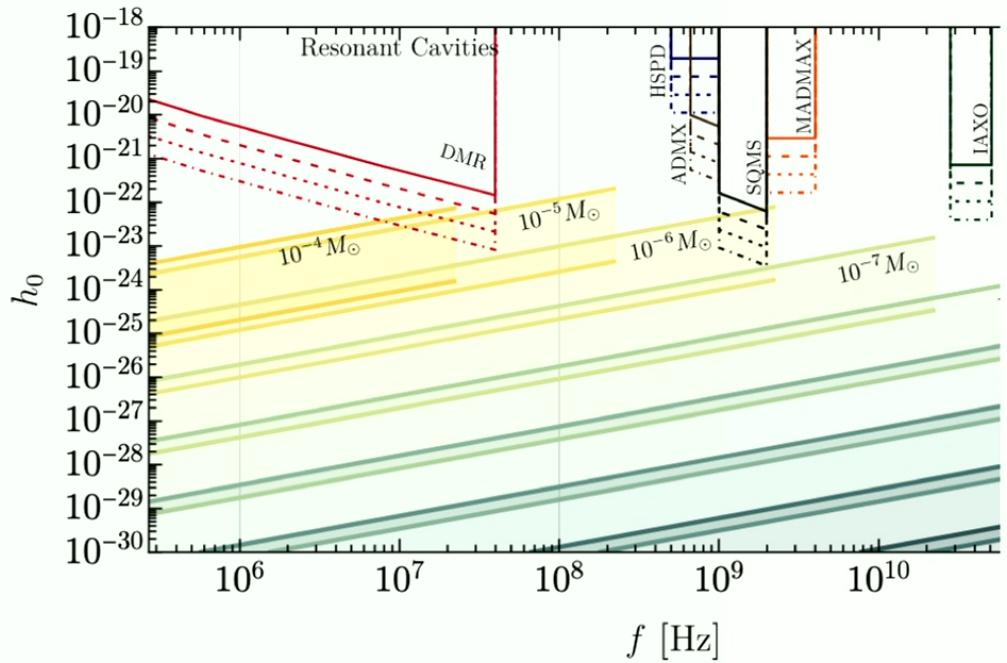
The Gravitational Soundscape at *high frequencies*

$$\Omega_{\text{GW}}(f) = \frac{1}{\rho_{\text{crit}}} \frac{d\rho_{\text{GW}}}{d(\ln f)} \quad \rho_{\text{GW}} \sim M_P^2 \omega^2 h_{\text{GW}}^2$$

SMASH model full spectrum
Ringwald Tamarit 22



GWs from PBHs (rates 1/year)
Franciolini et al 2205.02153

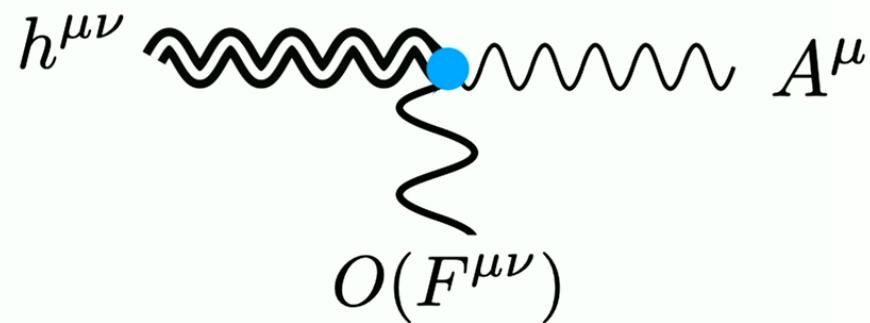


Searching for GWs with light

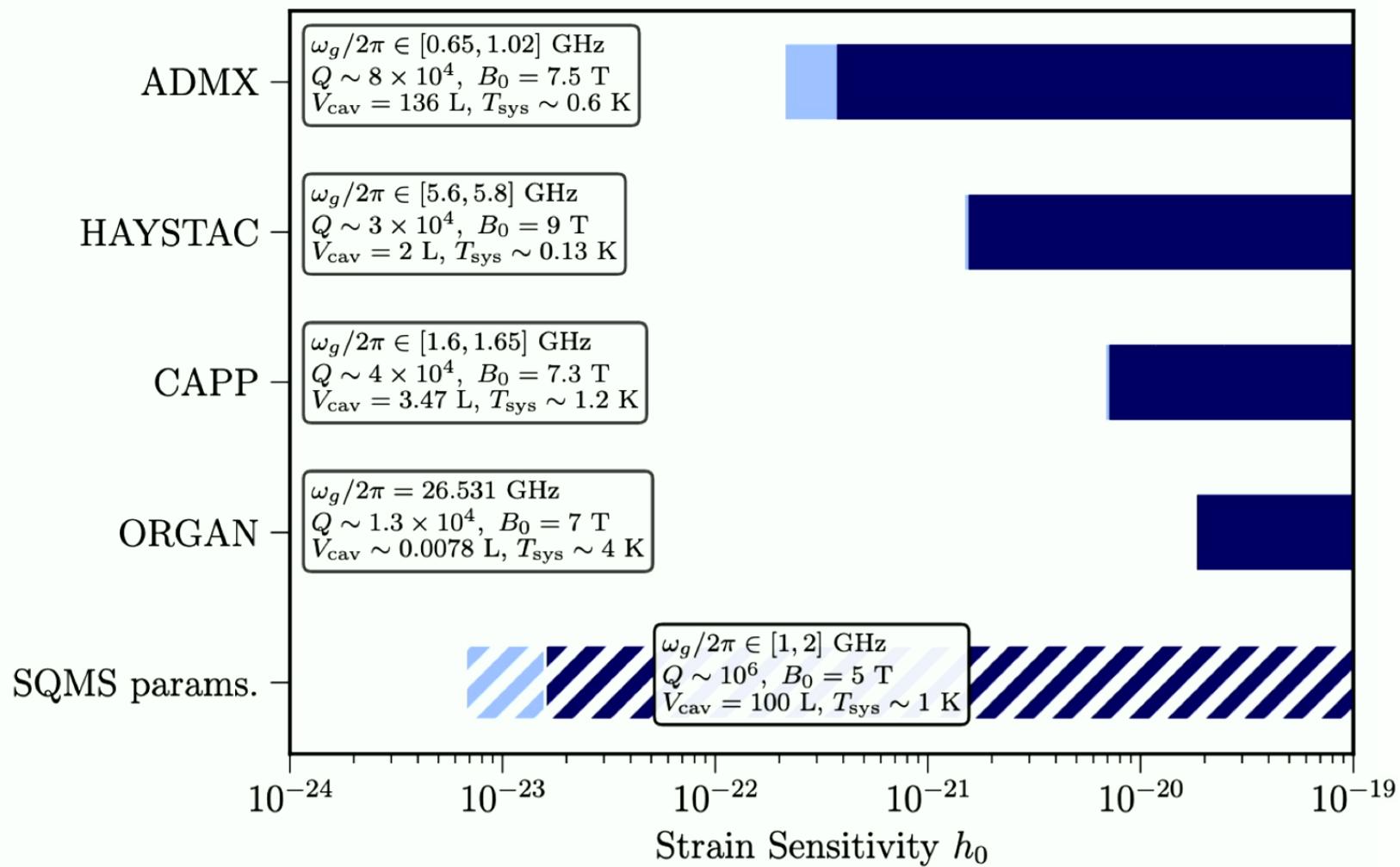
Interaction GWs with light

$$\mathcal{L} = \sqrt{-g} (R + F_{\mu\nu} F^{\mu\nu}) \supset \frac{1}{2} A_\mu j_{\text{eff}}^\mu(h) + \eta^{\mu\alpha} \eta^{\nu\beta} F_{\mu\nu} F_{\alpha\beta} + O(h^2)$$

$$j_{\text{eff}}^\mu = -\partial_\beta \left(\frac{1}{2} h F^{\mu\beta} + h_\alpha^\beta F^{\alpha\mu} - h_\alpha^\mu F^{\alpha\beta} \right)$$



Projected Sensitivities of Axion Experiments



Summary and outlook @ GHz

- SRF cavities are a mature technology to look for GWs at GHz either

- 'ADMX' like
- Heterodyne



- As in any GR calculation: subtleties in working with a consistent gauge

- TT gauge needs to be converted to laboratory frame
- Reach of $h_0 \sim 10^{-23}$ possible, though far from known signals



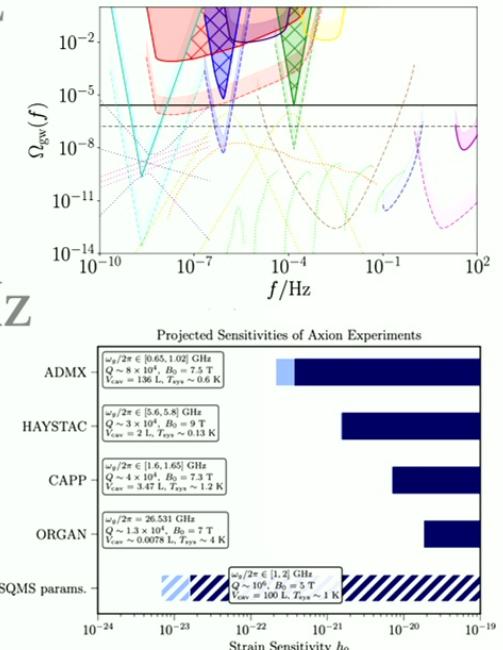
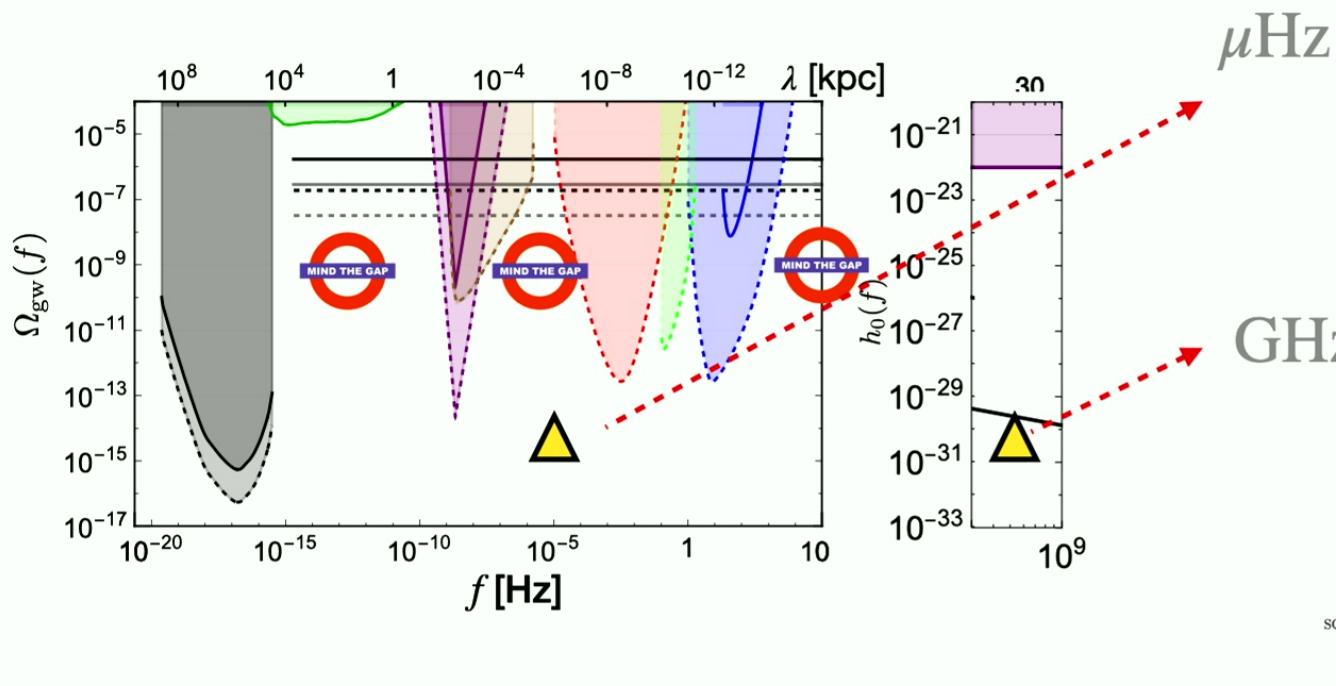
- Stay tuned for the connection to real world... (noise estimates + prospects)

Conclusions

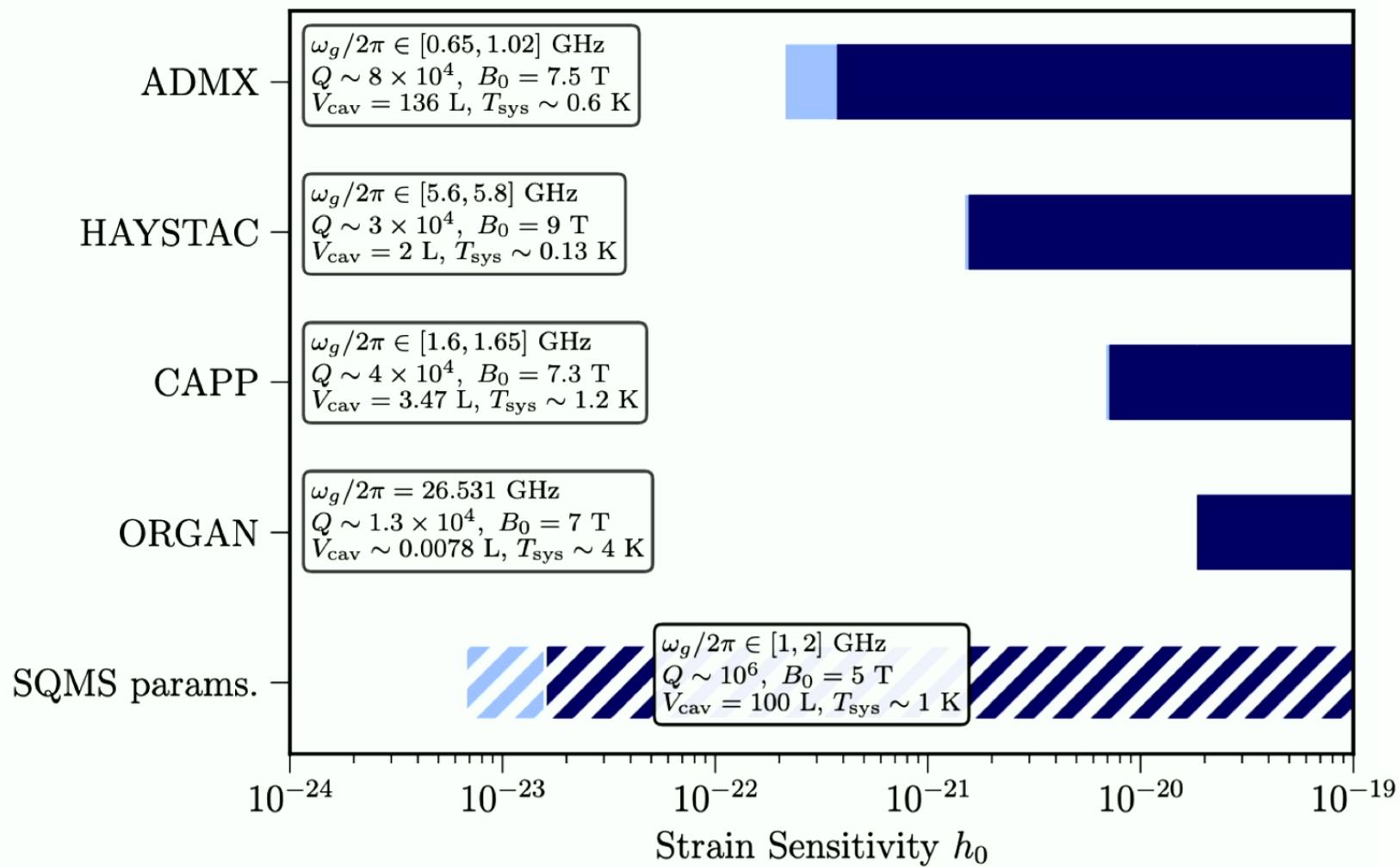
Multiband approach to GWs: great potential to transform both astrophysics and BSM

For this:

- i) there is vigorous effort to in four regions (CMB/PTA/LISA/LVK)
- ii) new ideas are needed in other phenomenologically rich regions



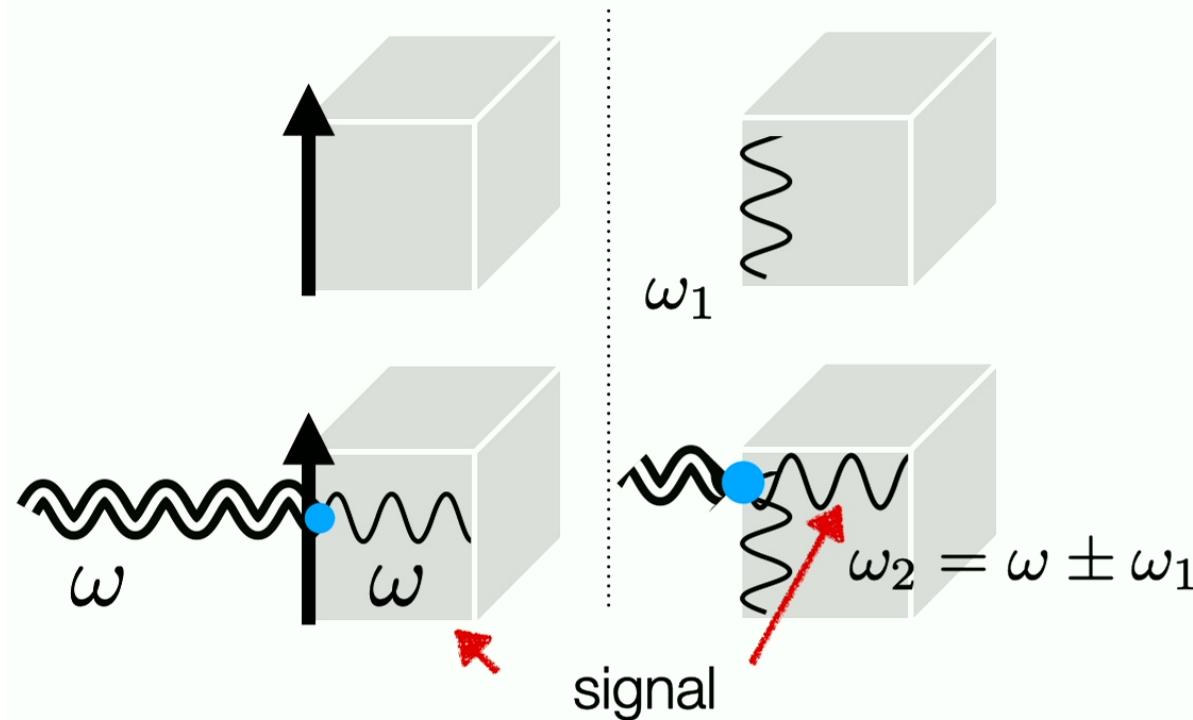
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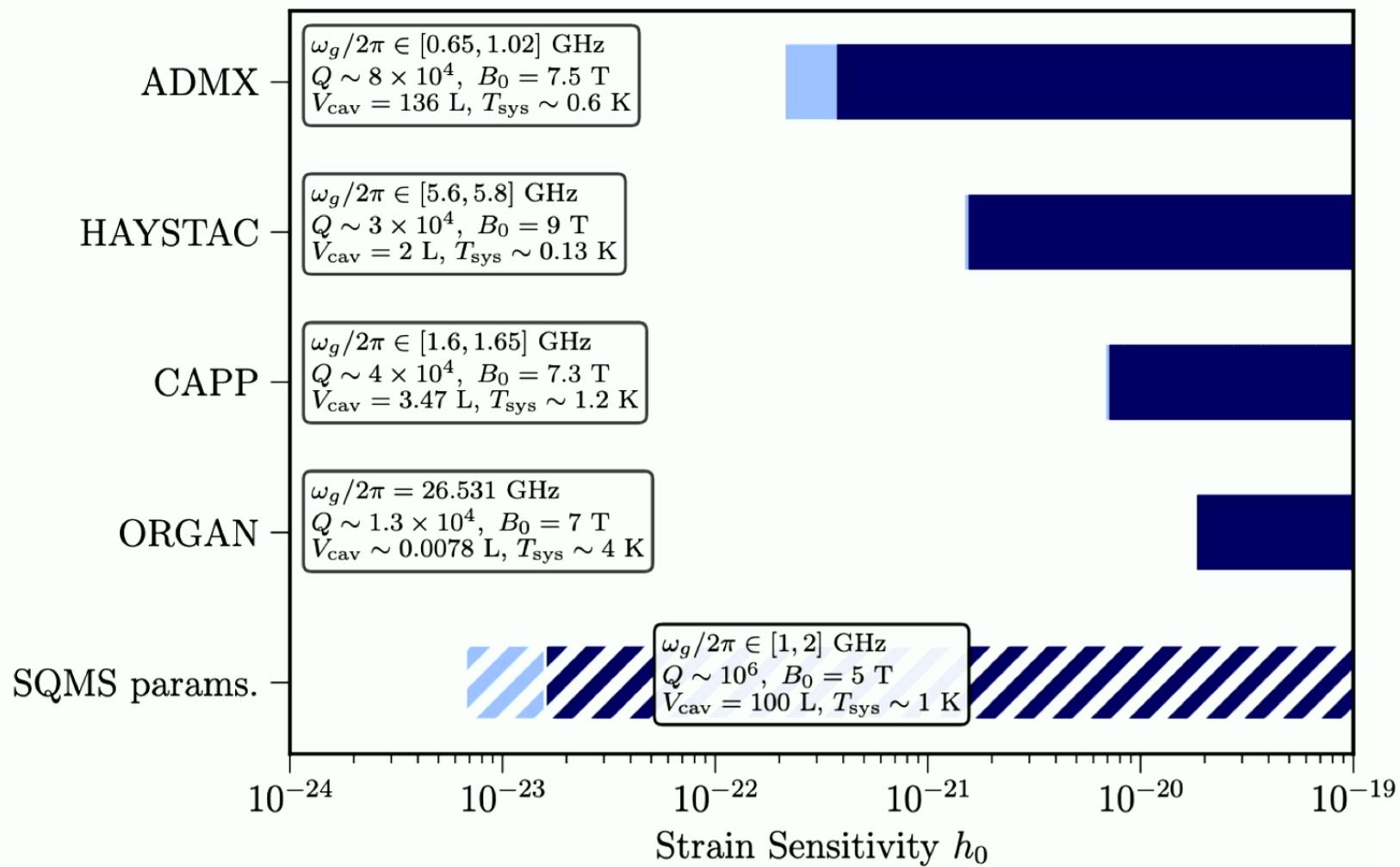
Recycling axion experiments

Cavities (cm -> GHz)

EM-coupling



Projected Sensitivities of Axion Experiments



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- Reach of $h_0 \sim 10^{-23}$ possible, though far from known signals



- Stay tuned for the connection to real world... (noise estimates + prospects)

Recycling axion experiments

Cavities (cm -> GHz)

EM-coupling

