

Title: Quantum Foundations Lecture - 230118

Speakers: Lucien Hardy

Collection: Quantum Foundations (2022/2023)

Date: January 18, 2023 - 10:15 AM

URL: <https://pirsa.org/23010050>

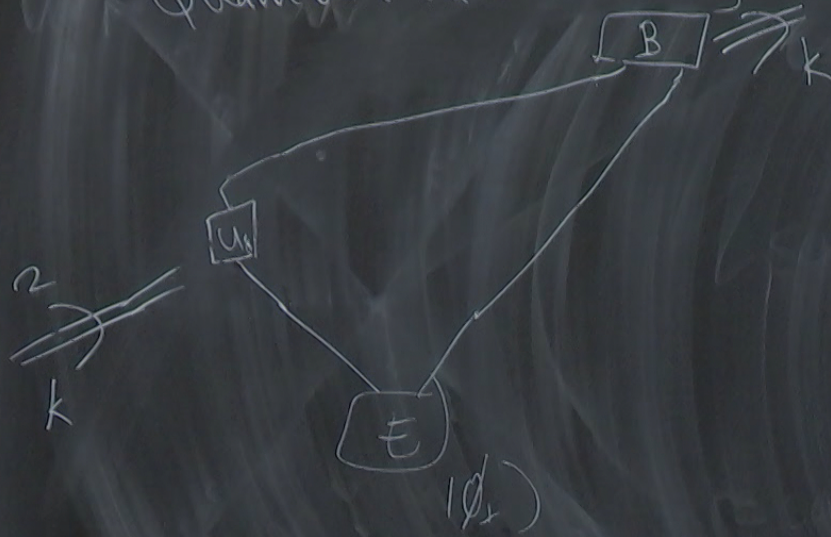
$$\begin{aligned}
 & (\alpha|0\rangle_{a_1} + \beta|1\rangle_{a_1}) |\phi_+\rangle_{a_2 a_3} \\
 & = \frac{1}{2} |\phi_+\rangle_{a_1 a_2} (\alpha|0\rangle_{a_3} + \beta|1\rangle_{a_3}) + \frac{1}{2} |\phi_-\rangle_{a_1 a_2} (\alpha|0\rangle_{a_3} - \beta|1\rangle_{a_3}) \\
 & + \frac{1}{2} |\psi_+\rangle_{a_1 a_2} (\beta|0\rangle_{a_3} + \alpha|1\rangle_{a_3}) + \frac{1}{2} |\psi_-\rangle_{a_1 a_2} (\beta|0\rangle_{a_3} - \alpha|1\rangle_{a_3})
 \end{aligned}$$

Measure in Bell basis. Get on of the four terms.

$$U_0 = \mathbb{1} \text{ (for } \phi_+ \text{ term)} \quad U_1 = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} \text{ for } \phi_- \text{ term} \quad U_2 = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \text{ for } \psi_+ \text{ case} \quad U_3 = \begin{pmatrix} 0 & -1 \\ 1 & 0 \end{pmatrix} \text{ for } \psi_- \text{ case}$$



# Quantum dense coding



Maximally entangled (indistinguishable)



# The Quantum Zeno Effect

Suderhan & Misra (1977)

System in initial state  $|0\rangle_a$   $t=0$

$$|\psi(\delta t)\rangle = e^{i\hat{H}\delta t/\hbar} |0\rangle_a$$

$$\approx \left(1 + \frac{i\hat{H}\delta t}{\hbar}\right) |0\rangle_a = |0\rangle + \frac{i\hat{H}\delta t}{\hbar} |0\rangle$$

Perform a



77)

Perform a projective measurement

$$\hat{P}_0 = |0\rangle\langle 0| \quad \hat{P}_0 = \sum_{n=1}^{N-1} |n\rangle\langle n|$$

$$\hat{P}_0 + \hat{P}_0 = 1$$

$$\text{prob}(\hat{P}_0) = K^2 (\delta t)^2$$

$$= \langle \psi(\delta t) | \sum_{n=1}^{N-1} |n\rangle\langle n| | \psi(\delta t) \rangle$$

$$= \left( \langle 0 | - \frac{i\hat{H}}{\hbar} \delta t | 0 \right) \sum_{n=1}^{N-1} |n\rangle\langle n| \left( |0\rangle + \frac{i\hat{H}}{\hbar} \delta t |0\rangle \right)$$

$$= \sum_{n=1}^{N-1} \frac{(\delta t)^2}{\hbar^2} \langle 0 | \hat{H} | n \rangle \langle n | \hat{H} | 0 \rangle = \sum_{n=1}^{N-1} \frac{|\langle 0 | \hat{H} | n \rangle|^2}{\hbar^2} (\delta t)^2$$



Consider, by contrast, classical exponential decay

$$P_{\text{not decay}} = e^{-dt/\tau} \approx 1 - \frac{dt}{\tau}$$

1 hor.



Q zero effect

Keep measuring system every  $\Delta t$ . do this  $\frac{T}{\Delta t}$  times

$$\text{prob (in } \odot \text{ after } T) = \left(1 - k^2(\Delta t)^2\right)^{\left(\frac{T}{\Delta t}\right)}$$

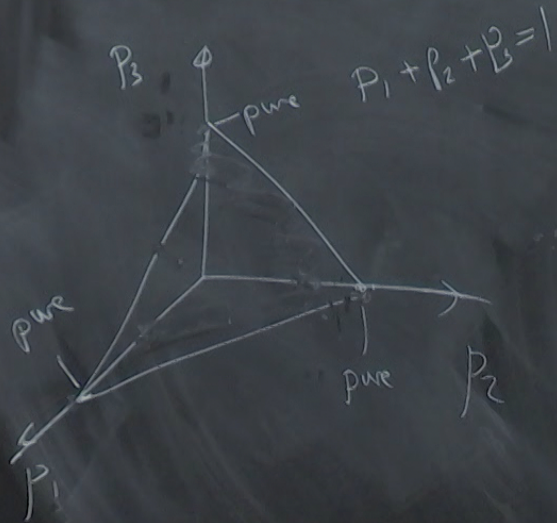
$$= 1 - \frac{T}{\Delta t} k^2(\Delta t)^2 + \dots$$

$$\approx 1 - Tk^2\Delta t$$

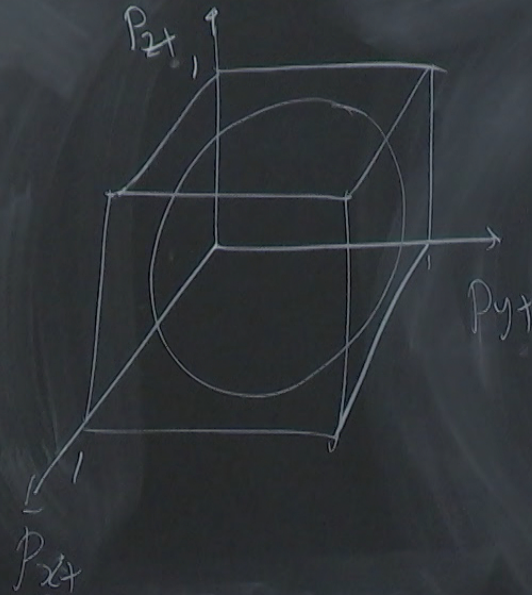
as  $\Delta t \rightarrow 0$  this prob  $\rightarrow 1$



Q zero eff is related to fact that evolution is linear in  $\mathcal{H}$ .



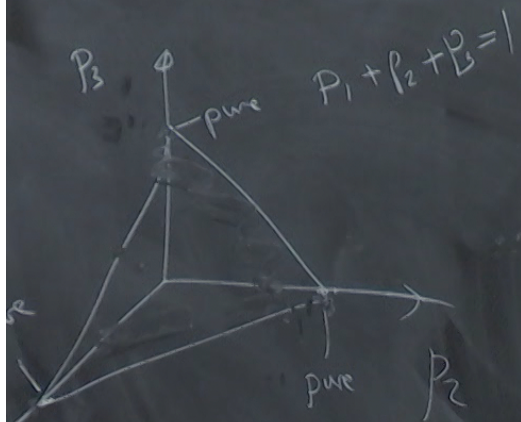
$$\begin{pmatrix} P_1 \\ P_2 \\ P_3 \end{pmatrix}$$



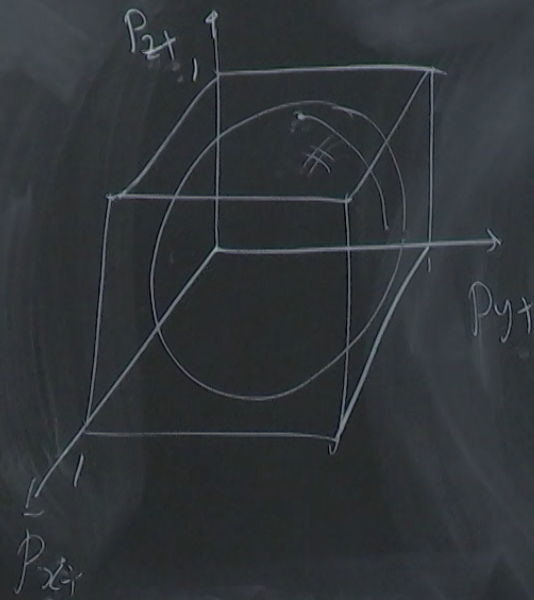
spin  $\frac{1}{2}$



Zero off is related to fact that evolution is linear in  $\mathcal{H}$ .



$$\begin{pmatrix} P_1 \\ P_2 \\ P_3 \end{pmatrix}$$



spin  $\frac{1}{2}$

$$|A\rangle \rightarrow \alpha|A\rangle + \beta|B\rangle$$

