

Title: Standard Model Lecture - 230119

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Collection: Standard Model (2022/2023)

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Standard Model

Lecture # 6

Higgs Mechanism, Electroweak Symmetry Breaking (Part 1)

$$\begin{aligned}
\mathcal{L}_{SM} = & -\frac{1}{2} \text{Tr}[G_{\mu\nu} G^{\mu\nu}] - \frac{1}{2} \text{Tr}[W_{\mu\nu} W^{\mu\nu}] - \frac{1}{4} B_{\mu\nu} B^{\mu\nu} \\
& - (D^\mu h)^\dagger (D_\mu h) + m^2 (h^\dagger h) - \lambda (h^\dagger h)^2 \\
& + i \left[\bar{q}_L^i \not{D} q_L^i + \bar{u}_R^i \not{D} u_R^i + \bar{d}_R^i \not{D} d_R^i + \bar{l}_L^i \not{D} l_L^i + (\bar{\nu}_R^i \not{D} \nu_R^i) + \bar{e}_R^i \not{D} e_R^i \right] \\
& - \left[\bar{q}_L^i \gamma_\mu^{ij} \tilde{h} u_R^j + \bar{q}_L^i \gamma_\mu^{ij} h d_R^j + (\bar{l}_L^i \gamma_\mu^{ij} \tilde{h} \nu_R^j) + \bar{l}_L^i \gamma_\mu^{ij} h e_R^j + \text{h.c.} \right] \\
& - (\bar{\nu}_R^i M^{ij} \nu_R^j) - \theta \text{Tr}[G_{\mu\nu} \tilde{G}^{\mu\nu}]
\end{aligned}$$

where: $G_{\mu\nu} = \partial_\mu G_\nu - \partial_\nu G_\mu + ig_3 [G_\mu, G_\nu]$ etc.
 $\tilde{G}^{\mu\nu} = \epsilon^{\mu\nu\alpha\beta} G_{\alpha\beta}$

$$D_\mu h = (\partial_\mu + ig_2 W_\mu + i(\frac{1}{2})g_1 B_\mu) h$$

$$\not{D} q_L^i = \gamma^\mu (\partial_\mu + ig_3 G_\mu + ig_2 W_\mu + i(\frac{1}{6})g_1 B_\mu) q_L^i$$

$g_1, g_2, g_3, m^2, \lambda, \gamma_u^{ij}, \gamma_d^{ij}, \gamma_\nu^{ij}, \gamma_e^{ij}, M^{ij}$ are constants Yukawa matrices (3x3)

Story so far:

i) Met YM theory

ii) Met SM table (gauge group, matter irreps)

iii) Met L_{SM}

Minimal Standard Model

$$\begin{aligned} L_{SM} = & -\frac{1}{2}\text{Tr}[G_{\mu\nu}G^{\mu\nu}] - \frac{1}{2}\text{Tr}[W_{\mu\nu}W^{\mu\nu}] - \frac{1}{4}B_{\mu\nu}B^{\mu\nu} \\ & - (D_\mu h)^\dagger (D^\mu h) + m^2(h^\dagger h) - \frac{\lambda}{4}(h^\dagger h)^2 \\ & + i[\bar{q}_L \not{D} q_L + \bar{u}_R \not{D} u_R + \bar{d}_R \not{D} d_R + \bar{l}_L \not{D} l_L + \bar{e}_R \not{D} e_R] \\ & - [\bar{q}_L \gamma_\mu \tilde{h} u_R + \bar{q}_L V_{CKM} \gamma_\mu h d_R + \bar{l}_L \gamma_\mu h e_R + \text{h.c.}] \end{aligned}$$

↪ What are the physical particles and interactions?

Focus on Higgs terms: "Higgs potential"

$$L = \dots - (\partial_\mu h)^\dagger (\partial^\mu h) - V(h)$$

$$V(h) = -m^2(h^\dagger h) + \frac{\lambda}{4}(h^\dagger h)^2$$

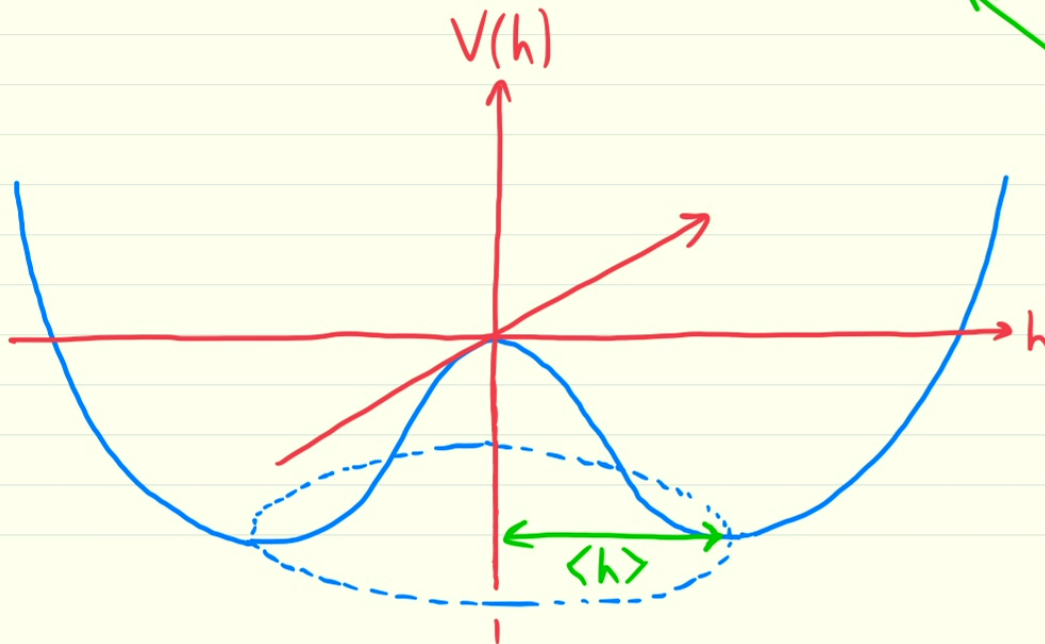
$$\Rightarrow \mathcal{A} = (\partial_t h)^\dagger (\partial_t h) + (\vec{\partial} h)^\dagger (\vec{\partial} h) + V(h)$$

minimized by
 $h(x) = \langle h \rangle = \text{const}$

$\langle h \rangle$ is "vacuum expectation value" ("VEV") of h

$\langle h \rangle \neq 0 \Rightarrow$ "spontaneous symmetry breaking" ("SSB")

In gauge theory:
 SSB $\rightarrow$ "Higgs Mechanism"



Perturbative QFT in brief:

- i) Split $L = L_{\text{free}} + L_{\text{int}}$ ← expanded around vacuum
- L_{free} is quadratic in fields (unperturbed, "free")
 - (read off phys particles, masses + propagators)
 - L_{int} is cubic and higher in fields (small perturbations, "interactions")
 - (read off interaction vertices + Feynman rules for vertices)

ii) $L_{\text{free}}^{\text{scalar}} = -\frac{1}{2}(\partial_\mu \phi)(\partial^\mu \phi) - \frac{1}{2}m^2 \phi^2$

$L_{\text{free}}^{\text{spinor}} = i\bar{\psi} \not{\partial} \psi - m\bar{\psi} \psi$

$\psi = \begin{pmatrix} \psi_L \\ \psi_R \end{pmatrix}$ (Dirac)
 $\psi = \begin{pmatrix} \psi_R^c \\ \psi_R \end{pmatrix}$ (Majorana)
 or $i\bar{\psi}_L \not{\partial} \psi_L$ or $i\bar{\psi}_R \not{\partial} \psi_R$ (Weyl)

$L_{\text{free}}^{\text{vector}} = -\frac{1}{4}F_{\mu\nu}F^{\mu\nu} - \frac{1}{2}m^2 A_\mu A^\mu$

in all cases,
 $e^{ik_\mu x^\mu}$
 with
 $k_\mu k^\mu = -m^2$
 is exact solution

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bosonic terms

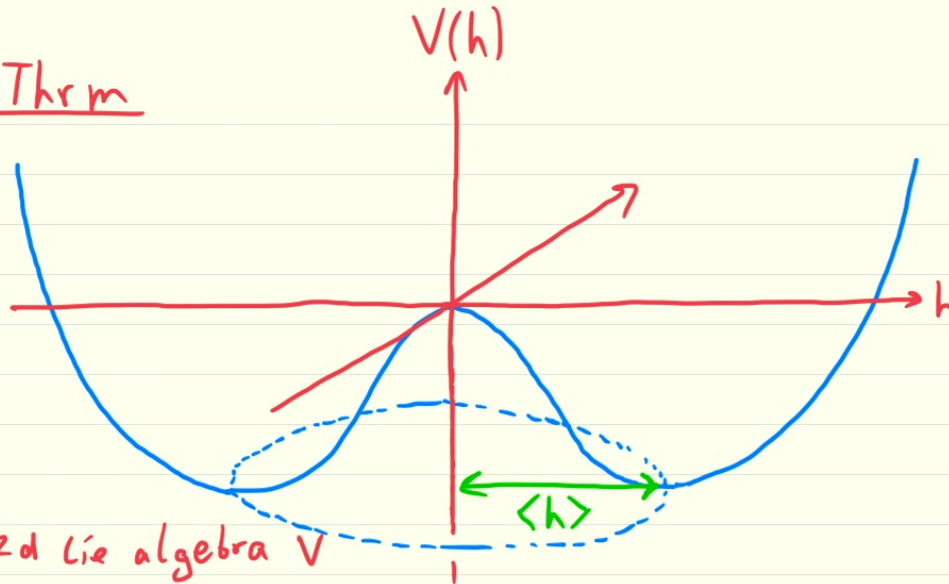
fermionic terms

What are the physical particles and interactions?

- i) Interpret quadratic bosonic terms (SM bosons)
- ii) Interpret quadratic fermionic terms (SM fermions)
- iii) Interpret remaining non-quadratic terms (SM interactions)

SSB (Global): Goldstone's Thrm

- In SM: $h = \begin{pmatrix} \varphi_1 + i\varphi_2 \\ \varphi_3 + i\varphi_4 \end{pmatrix}$
- Vacuum manifold: 3-sphere
 $(h^\dagger h) = (\varphi_1^2 + \varphi_2^2 + \varphi_3^2 + \varphi_4^2) = \text{const}$
- Take VEV $\langle h \rangle = \frac{1}{\sqrt{2}} \begin{pmatrix} 0 \\ v \end{pmatrix}$
- $\text{su}(3)_c \oplus \text{su}(2)_L \oplus \text{u}(1)_Y = 8+3+1 = 12\text{d Lie algebra } V$
- 4 generators act non-triv. on h : $T_1 = \frac{1}{2} \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}$, $T_2 = \frac{1}{2} \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix}$, $T_3 = \frac{1}{2} \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}$, $Y = \frac{1}{2} \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix}$
- But 1d subspace acts trivially on $\langle h \rangle$: $Q = Y + T_3 = \begin{pmatrix} 1 & 0 \\ 0 & 0 \end{pmatrix} \Rightarrow Q\langle h \rangle = 0$
- So $8+1 = 9\text{d}$ sub-algebra $W = \text{su}(3)_c \oplus \text{u}(1)_Q$ annihilates $\langle h \rangle$: $W\langle h \rangle = 0$
 (unbroken generators/gauge sym)
- Remaining 3d quotient space V/W does not: $V/W\langle h \rangle \neq 0$
 (broken generators)
- $\Rightarrow 3$ massless scalar bosons, one for each broken generator in V/W
 (Goldstone bosons) (Goldstone's Thrm)



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- Take VEV $\langle h \rangle = \frac{1}{\sqrt{2}} \begin{pmatrix} 0 \\ v \end{pmatrix}$
- $su(3)_c \oplus su(2)_L \oplus u(1)_Y = 8+3+1=12$ d Lie algebra V
- 4 generators act non-triv. on h : $T_1 = \frac{1}{2} \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}$, $T_2 = \frac{1}{2} \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix}$, $T_3 = \frac{1}{2} \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}$, $Y = \frac{1}{2} \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix}$
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