

Title: Gravitational-wave astronomy and fundamental physics with charged black holes

Speakers: Gabriele Bozzola

Series: Strong Gravity

Date: December 08, 2022 - 1:00 PM

URL: <https://pirsa.org/22120059>

Abstract: Charge (electric, magnetic, or any $U(1)$ charge) is a parameter often neglected in simulations of black holes. As a result, little is known about the dynamics of charged binaries. In this talk, I will highlight the importance of understanding the non-linear interaction of charged black holes for astrophysics and fundamental physics. I will show results from fully self-consistent general-relativistic simulations of merging black holes, touching upon the challenges faced in performing such calculations and the improvements that enabled successful long-term evolution. I will discuss general features of quasi-circular inspirals, and present constraints on the charge of astrophysical black holes and deviation from general relativity obtained from the gravitational-wave event GW150914. Finally, I will highlight the relevance of this line of research in the context of the upcoming gravitational-wave detectors.

Zoom link: <https://pitp.zoom.us/j/93809443805?pwd=bmcvd3NZWjUraERBcGdtL2Y3WTI6QT09>

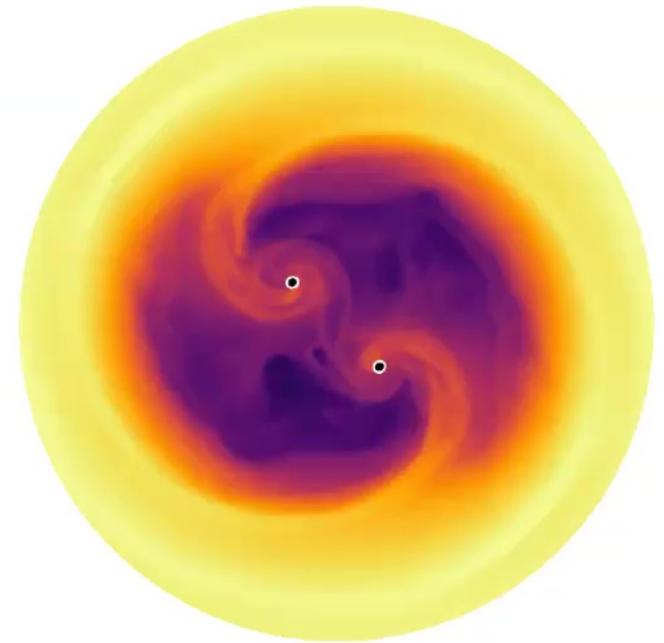
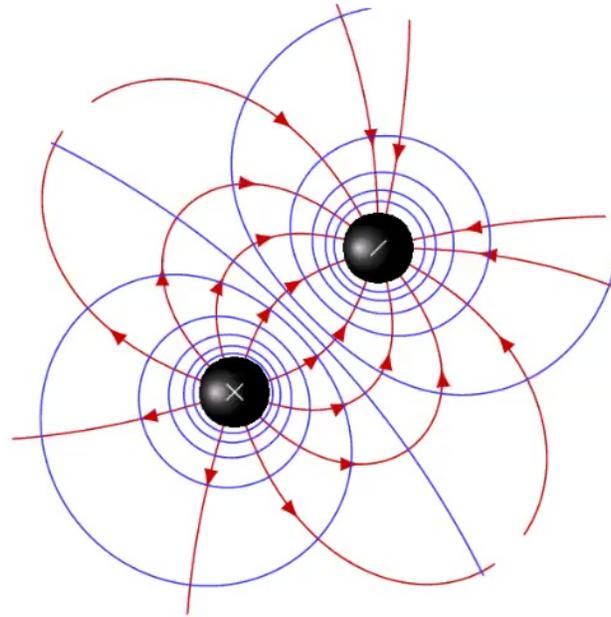


NUMERICAL RELATIVITY, GRAVITATIONAL WAVES, AND FUNDAMENTAL PHYSICS WITH CHARGED BLACK HOLES

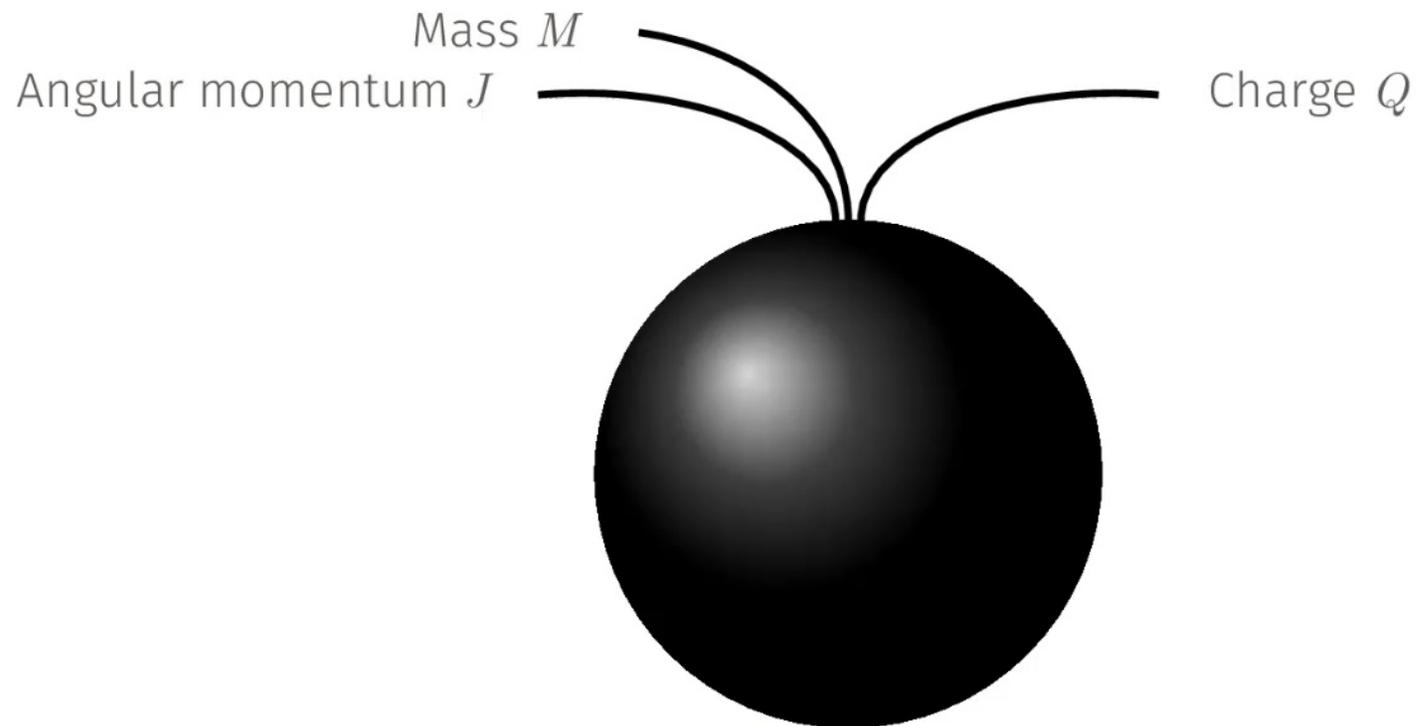
December 8, 2022

Gabriele Bozzola,
Vasileios Paschalidis

Department of Astronomy,
University of Arizona

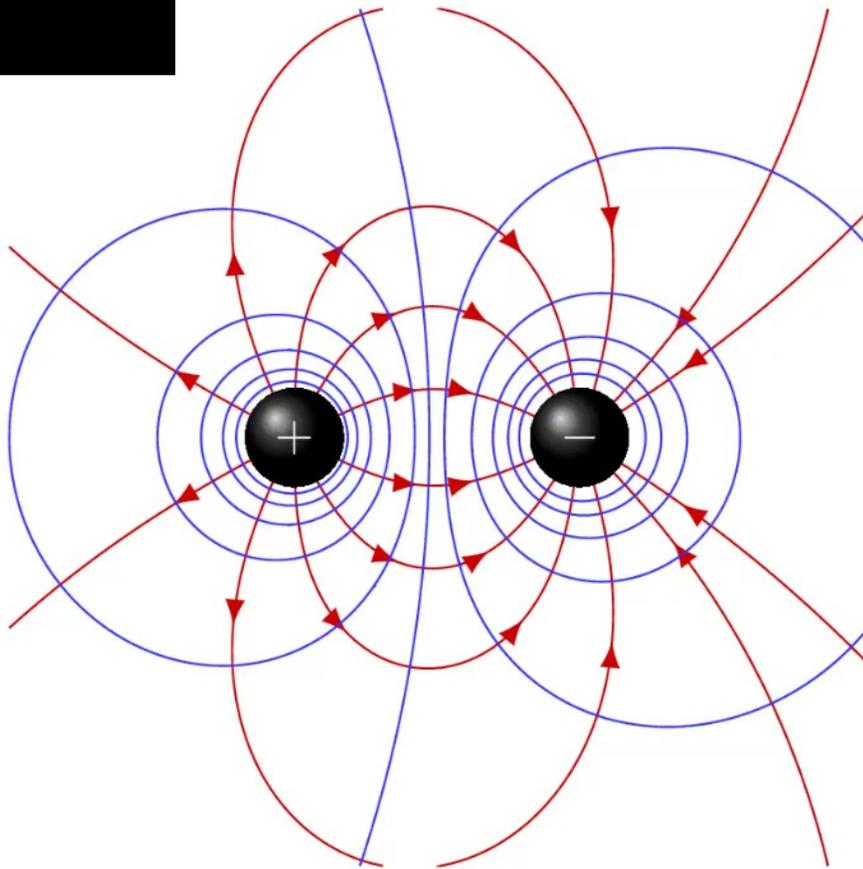


Black holes are simple objects: the *no hair* theorem



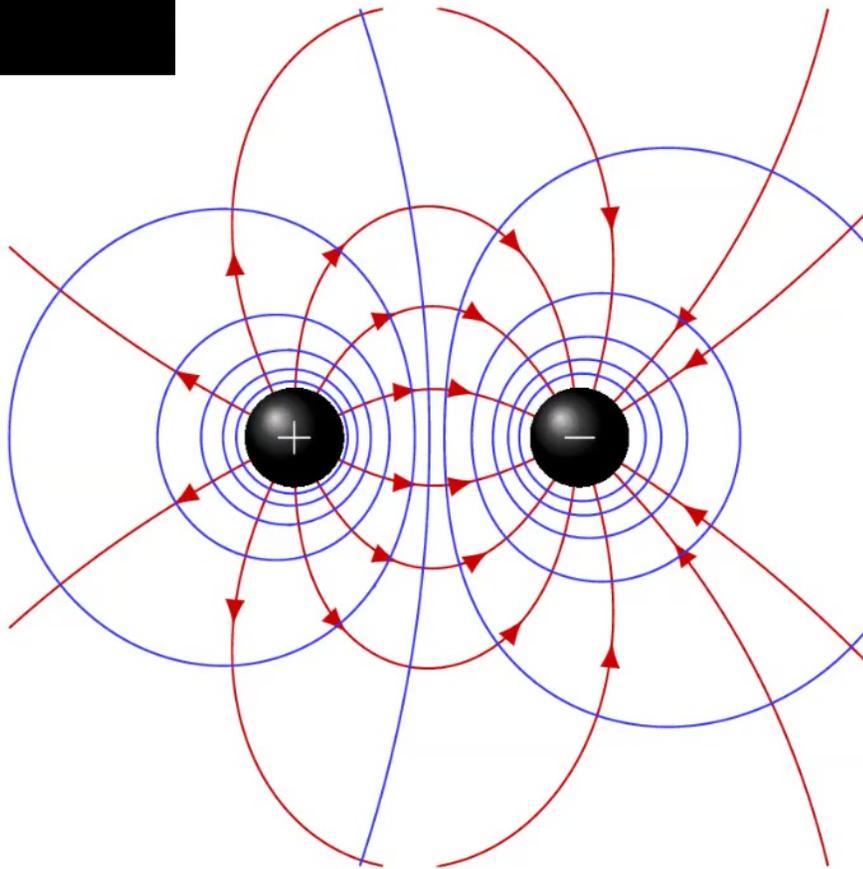
Israel (1967), Israel (1968), Carter (1971), ...

2



The nonlinear interaction between charged binary black holes is **unexplored**

Orbits? Emission? ...?



The nonlinear interaction between charged binary black holes is **unexplored**

Orbits? Emission? ...?

We assume

$$G = c = (4\pi\epsilon_0)^{-1} = 1$$

$$(q_{\text{proton}}/m_{\text{proton}} \approx 10^{18})$$

1. Results apply to stellar-mass, supermassive, and microscopic BHs
2. Charge does not have to be electromagnetic

Non-linear interaction is important for

Astrophysics



Constraints on charge
GW templates for LIGO-Virgo



- BH charge is largely unconstrained
- GW models do include charge

Much more on this later!

Non-linear interaction is important for

Exotic Astrophysics



Constraints on charged dark matter
and magnetic monopoles



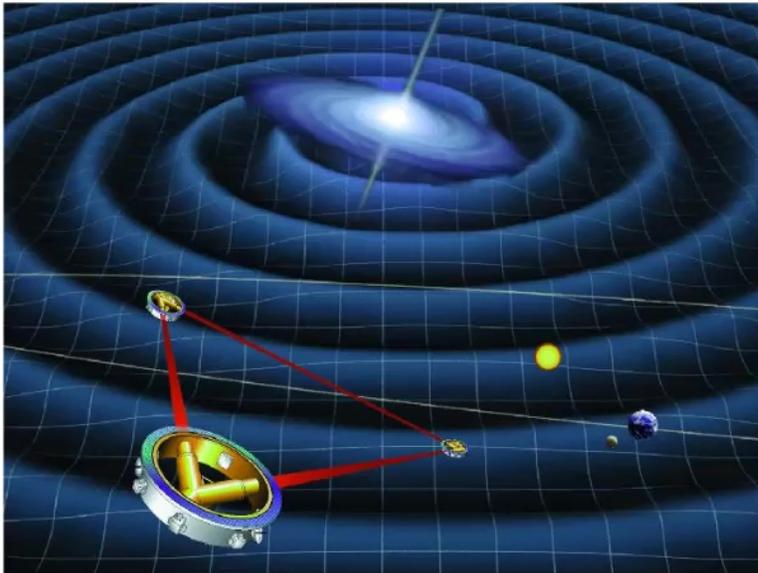
- Primordial BH could have (electric or magnetic) charge
- Dark matter could be (darkly) charged

Non-linear interaction is important for

Modified Gravity



Springboard and proxy for other theories
(e.g. Einstein-Maxwell-Scalar)



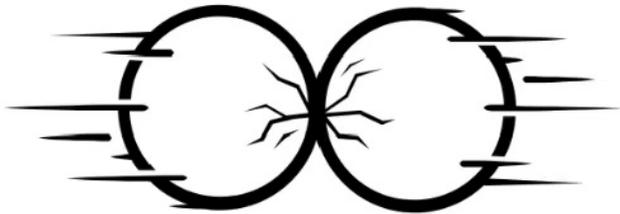
- Well defined way to go beyond GR
- Some theories are mathematically identical in specific limits
- Better understanding for future facilities

Non-linear interaction is important for

First Principles



Ultra-relativistic collisions
Scattering
Cosmic censorship



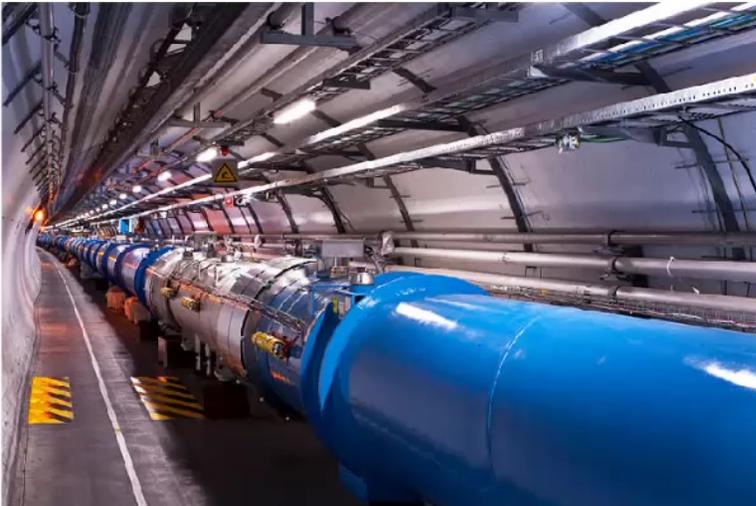
- Exceptional laboratories for controlled numerical experiments
- Interplay between extreme electromagnetic and gravitational fields

Non-linear interaction is important for

Energetic Particles

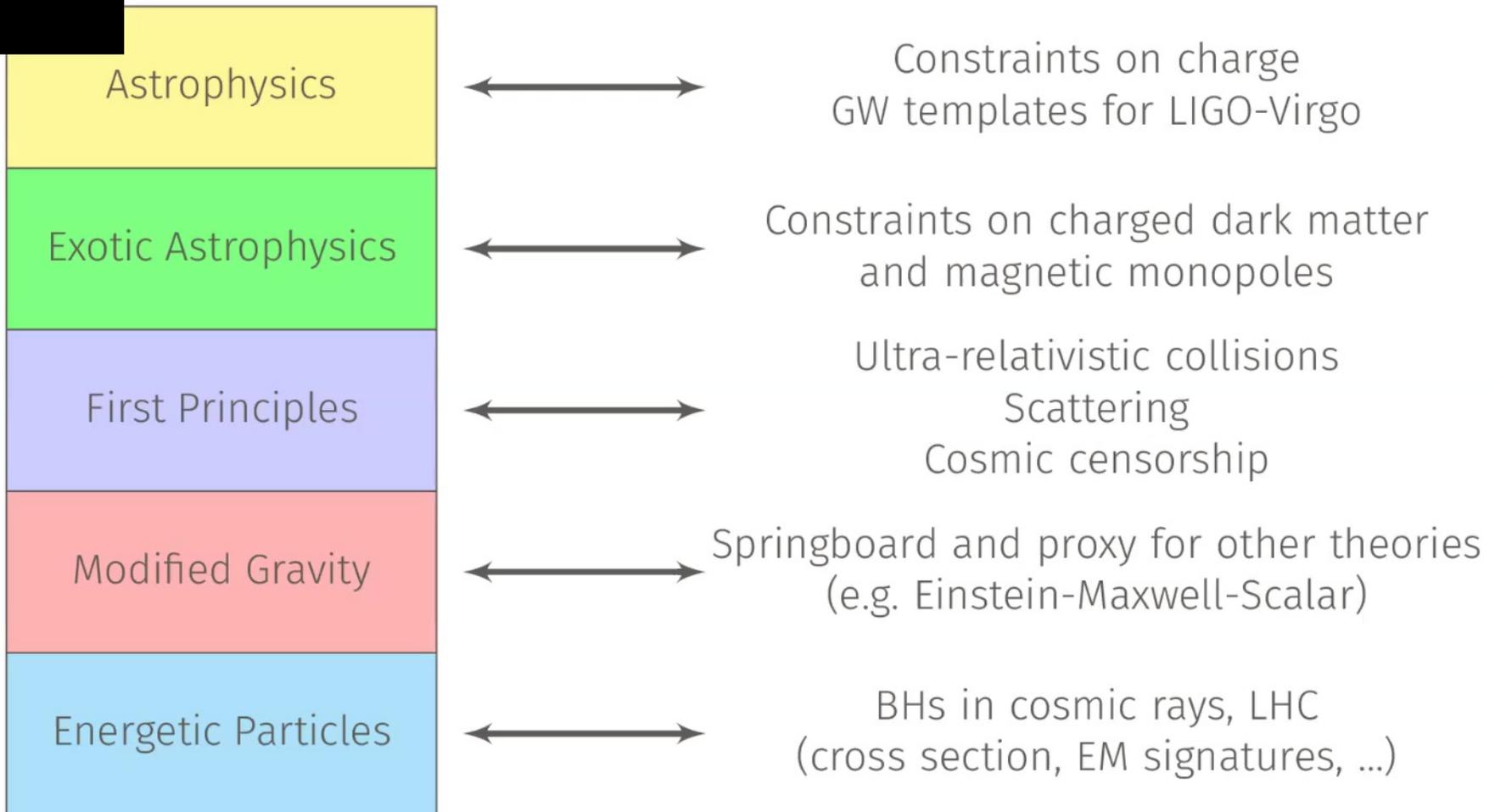


BHs in cosmic rays, LHC
(cross section, EM signatures, ...)



- Microscopic black hole production and detection
- Tests of specific grand unified theories

Non-linear interaction is important for



This list is incomplete; you can help by expanding it.

Einstein-Maxwell

$$R_{\mu\nu} - \frac{1}{2}Rg_{\mu\nu} = 8\pi T_{\mu\nu}^{\text{EM}}$$

$$\nabla_{\mu}F^{\mu\nu} = 4\pi J^{\nu}$$

$$\nabla_{\mu}{}^{\star}F^{\mu\nu} = 0$$

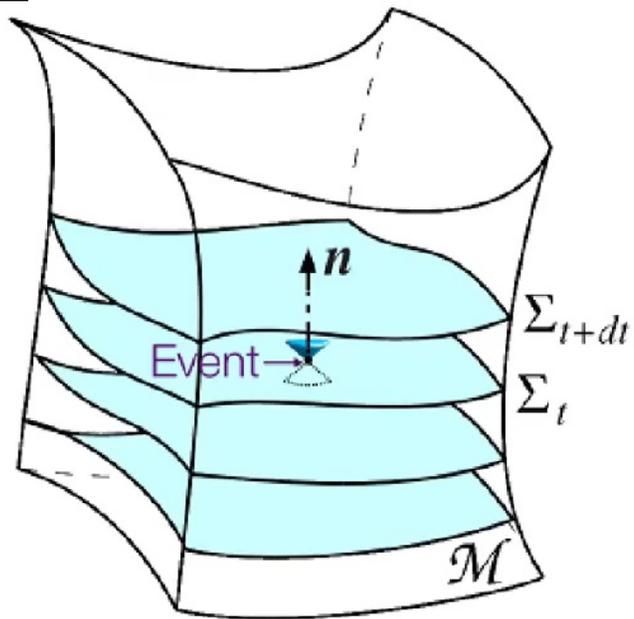
for binary black holes

- No symmetries
- Non perturbative
- Highly non-linear coupled PDEs



NUMERICAL RELATIVITY

How to solve Einstein-Maxwell's equations

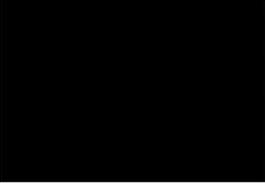


Numerical solution of
Einstein-Maxwell's equations as
initial-value-problem

Issues with stability, gauge, initial
data, ...

HARD PROBLEM!

Numerical relativity = *know-how* to be successful in this feat



Today's talk

Numerical Relativity

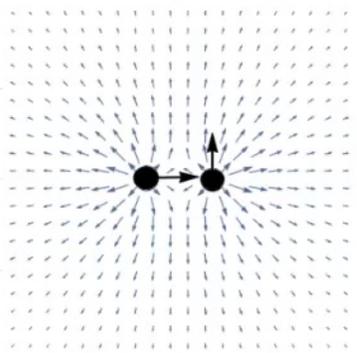
First numerical-relativity simulations of
quasi-circular mergers

Astrophysics

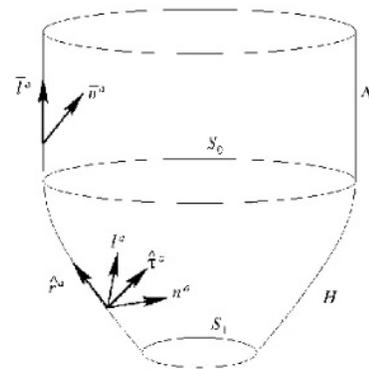
First constraints on BH charge from
gravitational waves

Ask questions or reach out for more!

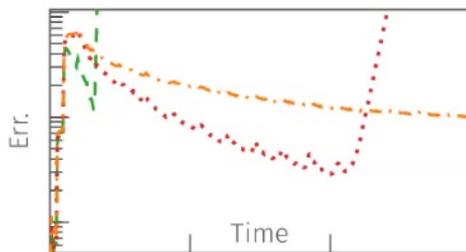
Novel/improved methods



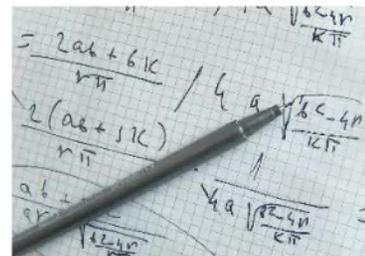
Initial data
TwoChargedPunctures



Interpretation
QuasiLocalMeasuresEM



Stability
Continuous Dissipation



Analysis
Newman-Penrose

Bozzola&Paschalidis (2019, 2021b), Luna&Bozzola+ (2022)

Full non-linear simulations with the Einstein Toolkit

Initial data	TwoChargedPunctures
Adaptive mesh refinement	Carpet
Evolution	ProcaEvolve + Lean
Diagnostics	ProcaConstraints ProcaDiagnostics
Wave extraction	Proca_NPScalars
Horizons properties	AHFinderDirect + QuasiLocalMeasuresEM
Postprocessing	kuibit



[Bozzola&Paschalidis \(2019, 2020, 2021a, 2021b\)](#) [Bozzola \(2021\)](#) [Sperhake \(2006\)](#)
[Thornburg \(2003\)](#) [Dreyer et al \(2002\)](#) [Schnetter \(2003\)](#) [Zilhão et al \(2015\)](#)

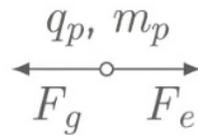


- Mostly multi-resolution Cartesian grids
- (Non-)linear elliptic and hyperbolic PDEs
- Pseudo-spectral solvers, finite differences, method of the lines
- Volume/surface integrals with Lagrange/Hermite interpolations
- C, C++, Fortran, OpenMP, MPI (no GPU)
- Weeks/months on supercomputers



[Bozzola&Paschalidis \(2019, 2020, 2021a, 2021b\)](#) [Bozzola \(2021\)](#) [Sperhake \(2006\)](#)
[Thornburg \(2003\)](#) [Dreyer et al \(2002\)](#) [Schnetter \(2003\)](#) [Zilhão et al \(2015\)](#)

Charge of astrophysical black holes

 $Q_{\text{BH}}, M_{\text{BH}}$ Accretion $\Leftrightarrow F_e + F_g < 0$

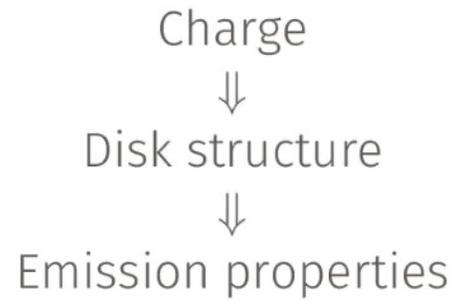
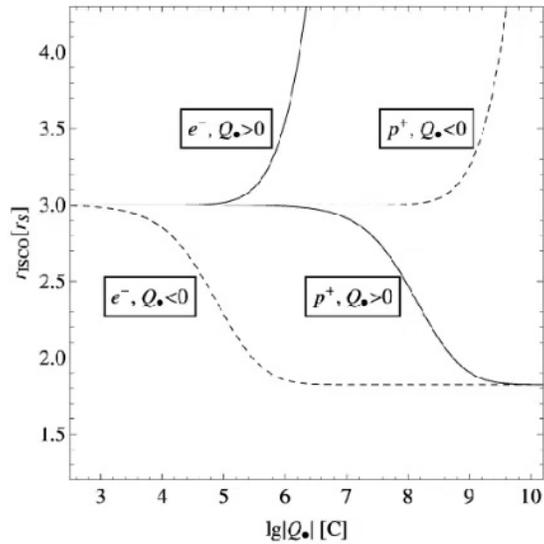
$$\frac{q_p Q}{r^2} - \frac{m_p M}{r^2} < 0$$

$$\frac{Q_{\text{BH}}}{M_{\text{BH}}} < \frac{m_{\text{proton}}}{q_{\text{proton}}} \sim 10^{-18}$$

BH with $Q/M > 10^{-18}$ Accretes opposite charges
Does not accrete same charge

Discharge

Measurements of BH charge are model-dependent



Extensive non-trivial modeling effort

Plot from Zajacek+ (2018)

Discharge limits rely on $\frac{Q_{\text{BH}}}{M_{\text{BH}}} < \frac{m_{\text{proton}}}{q_{\text{proton}}} \sim 10^{-18}$

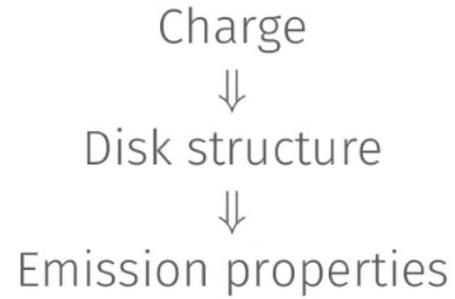
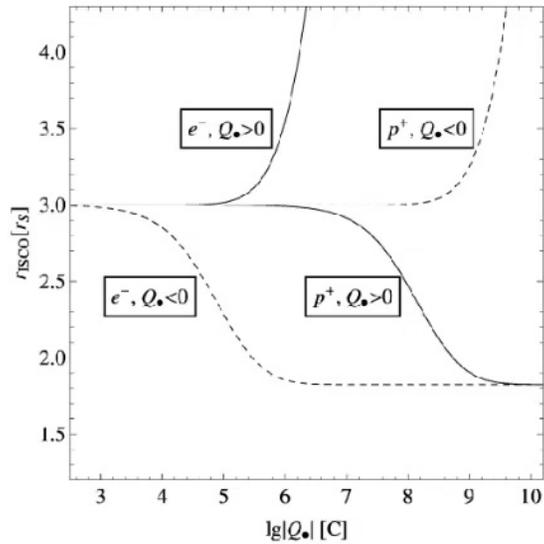
Evaded if:

- Mini-charged dark matter ($m/q \sim 1$)
- Dark electromagnetism (different EM coupling)
- Magnetic monopoles (no discharge)
- Gravitational charge (STVG*, $Q/M = \sqrt{\alpha/(1 + \alpha)}$, with α coupling)

* = Scalar-Tensor-Vector Gravity

Preskill+ (1984), Moffatt (2006), Feng+ (2009), Cardoso+ (2016), ...

Measurements of BH charge are model-dependent

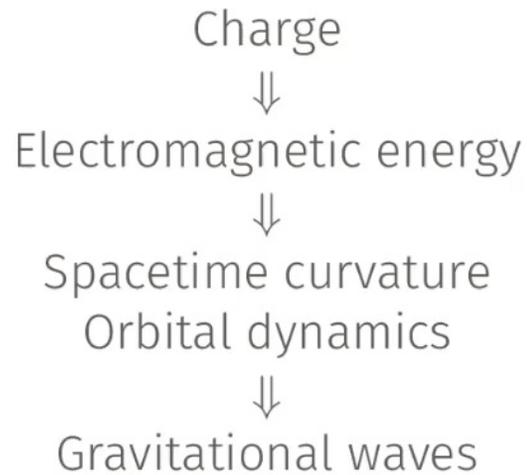


Extensive non-trivial modeling effort

Non-zero charge would challenge current understanding of BH formation and environments

Plot from Zajacek+ (2018)

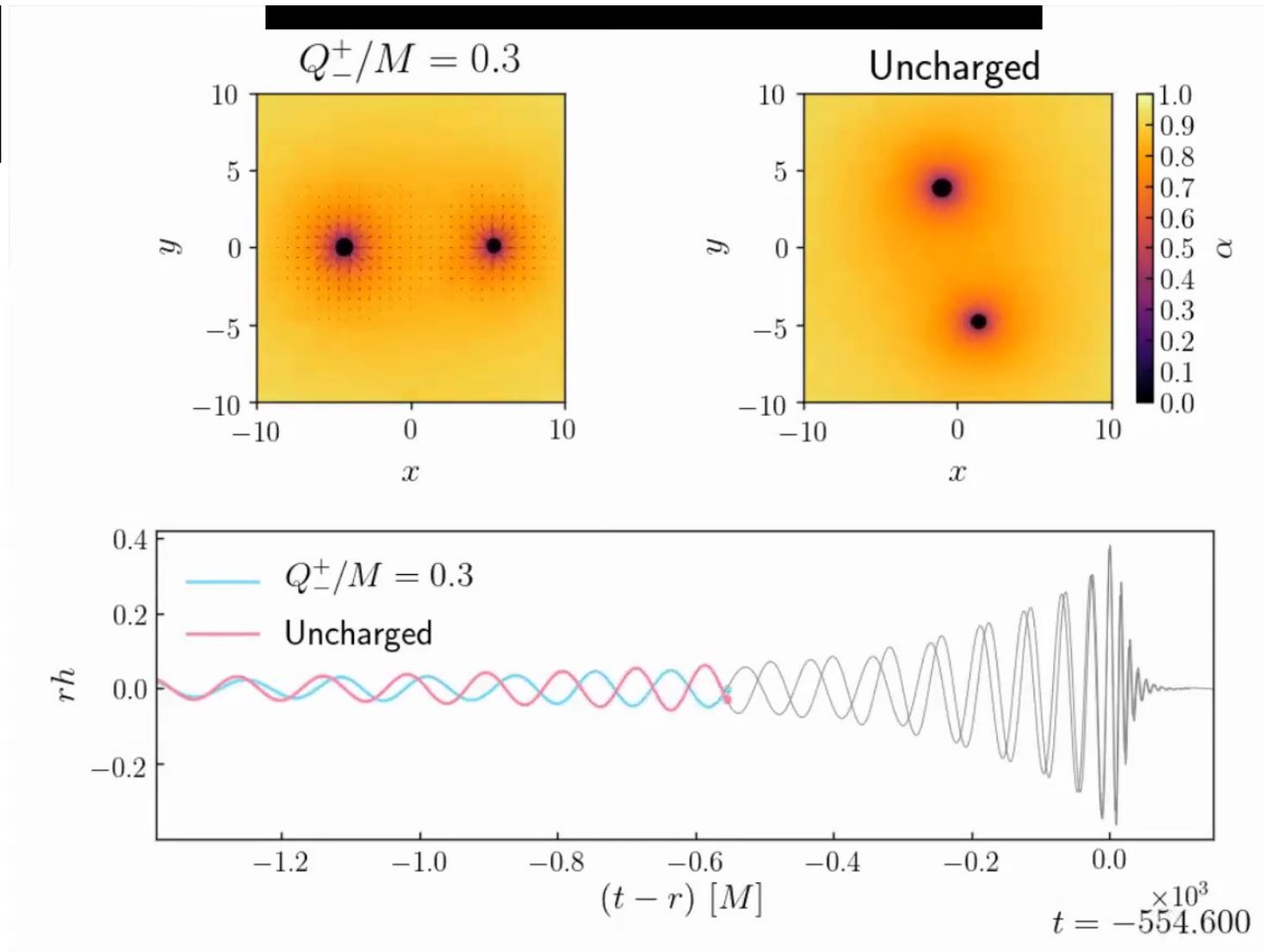
How can we test this?



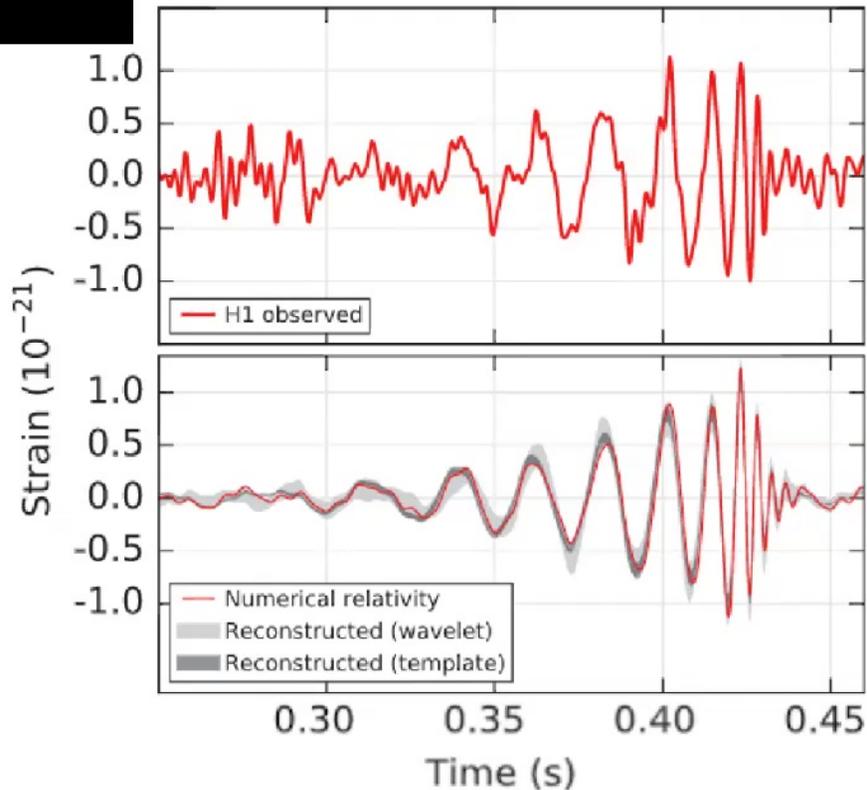
GWs know about charge

Strengths:

- Little modeling required
- Applicable to exotic astrophysics



Hanford, Washington (H1)

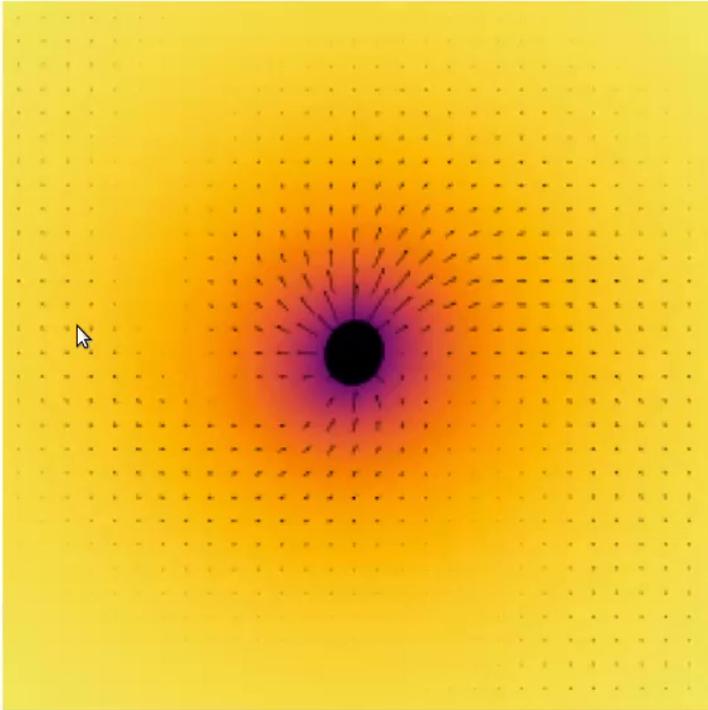


GW150914:

- First confirmed GW event
- Loud and well-characterized (SNR = 25)
- Mass-ratio 29/36
- (Disputed) coincident EM observations

Full Bayesian analysis requires GW templates (= *currently impossible*)

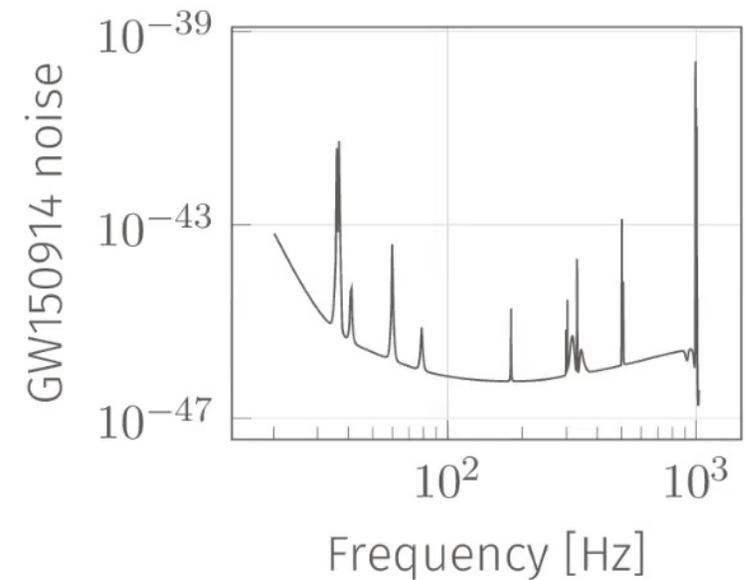
Electromagnetic emission is model dependent



- Propagation/interaction with plasma
- Different for non-EM theories

Constraint from mismatch analysis

1. Perform array of charged simulations
2. Compute SNR needed to distinguish them from GW150914
3. If SNR marginalized over time, phase, and mass $< 25 \rightarrow$ ruled out
4. Find maximum charge still compatible



First constraints on:

Black hole charge from GW150914:

Opposite charge: $Q/M < 0.2$

Single charge: $Q/M < 0.35$

Same charge: $Q/M < 0.4$

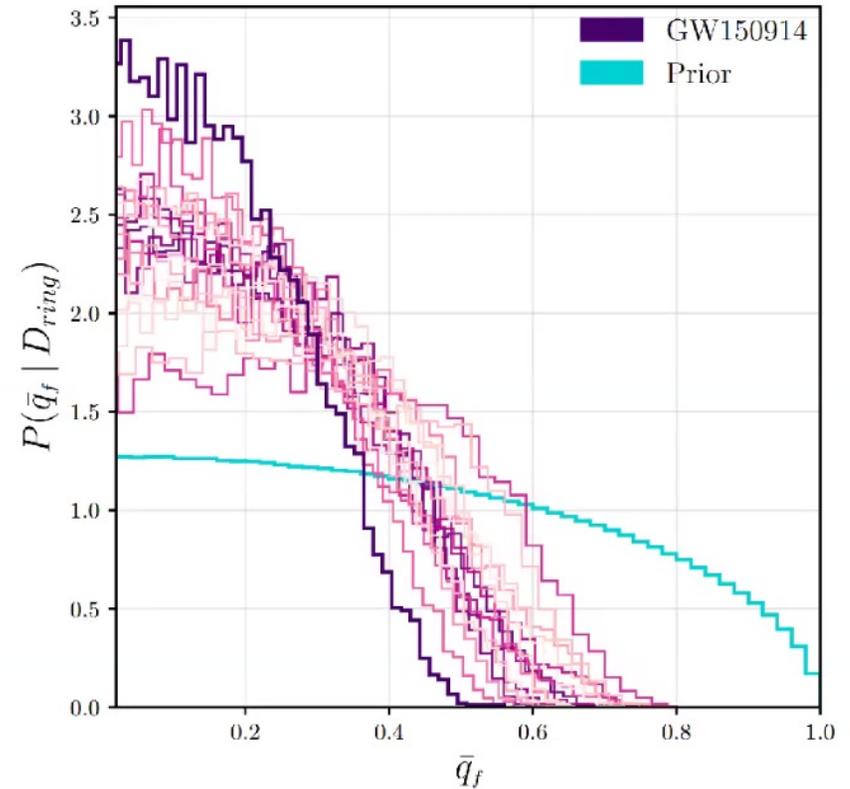
(charge imbalance $< 10^{-17} M_{\odot}$)

Tightest constraint on STVG
(Scalar-Tensor-Vector Gravity)
in strong-field:

$\alpha < 0.19$ (before $\alpha \lesssim 9$)

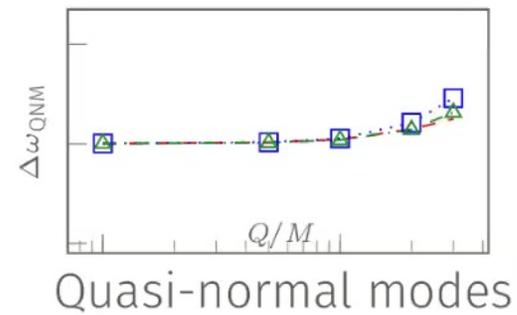
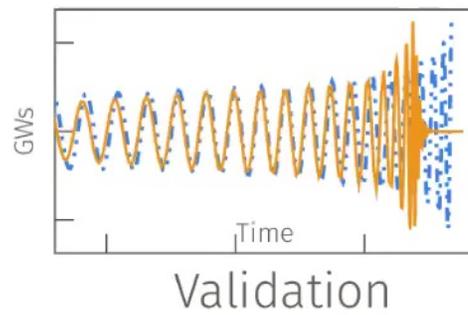
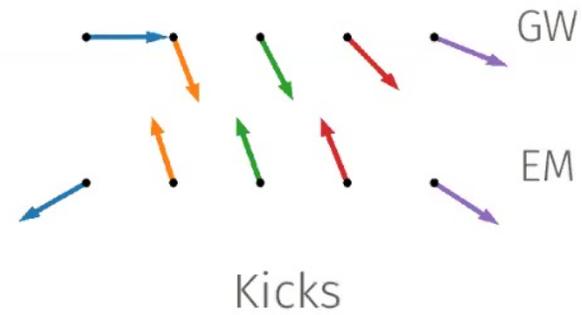
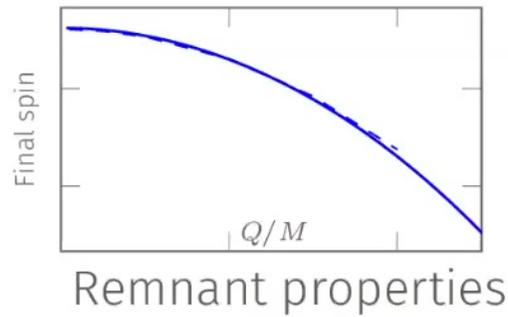
[Bozzola&Paschalidis](#) (PRL 2021)

Later confirmed by Carullo+, 2022:

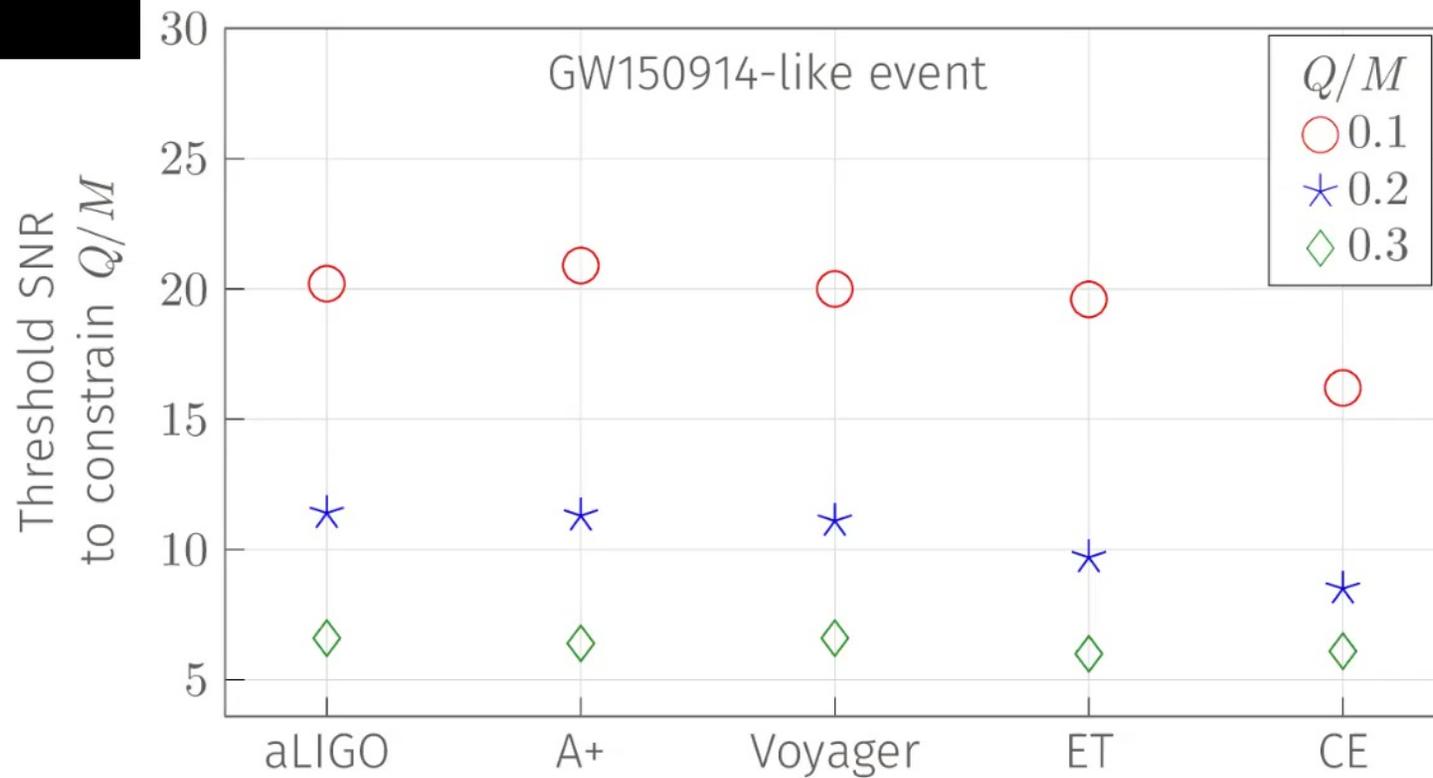


\bar{q}_f = final charge-to-mass-ratio

More advancements:

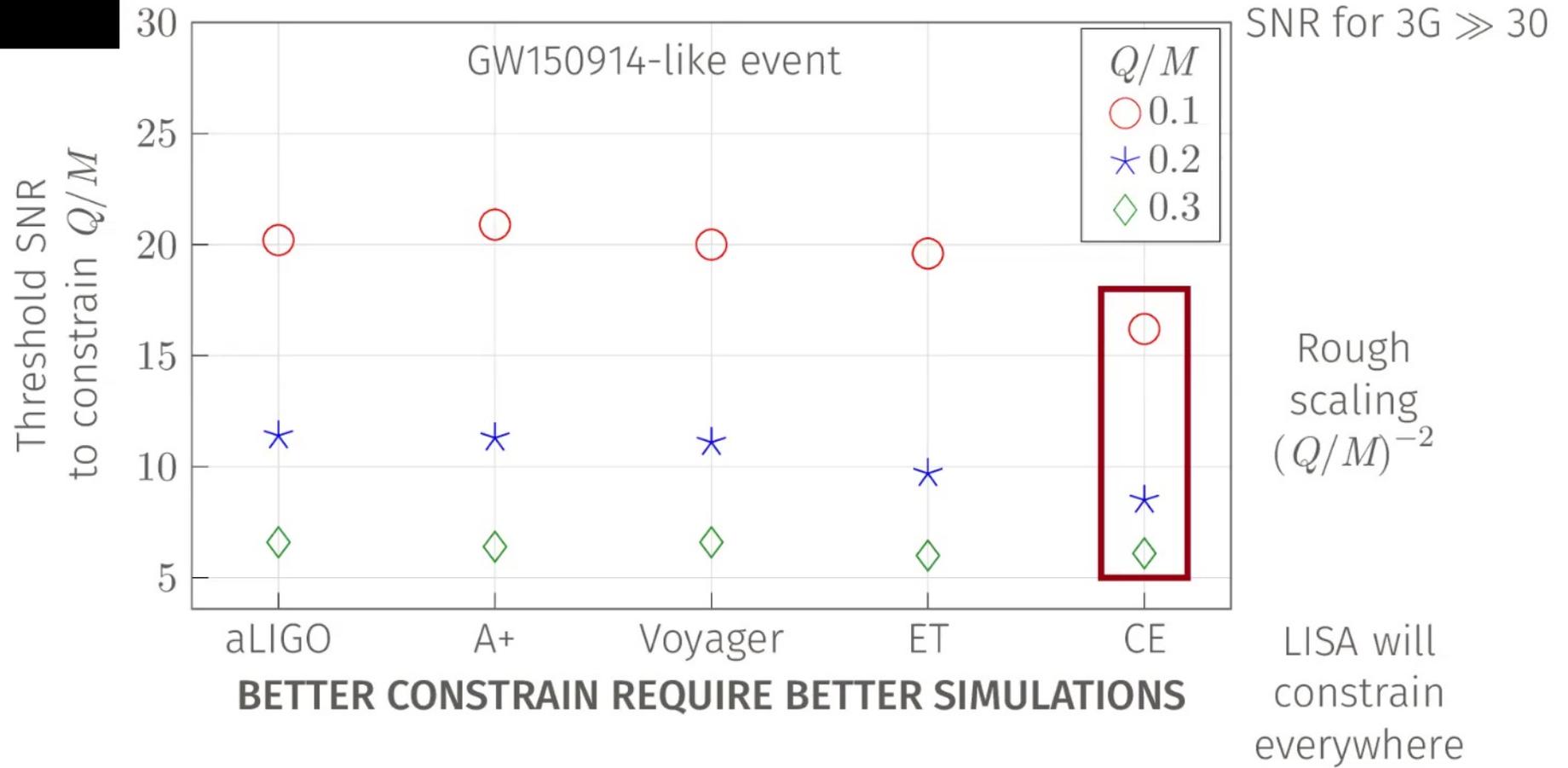


Expected constraints with future detectors



Bozzola&Paschalidis (2021b)

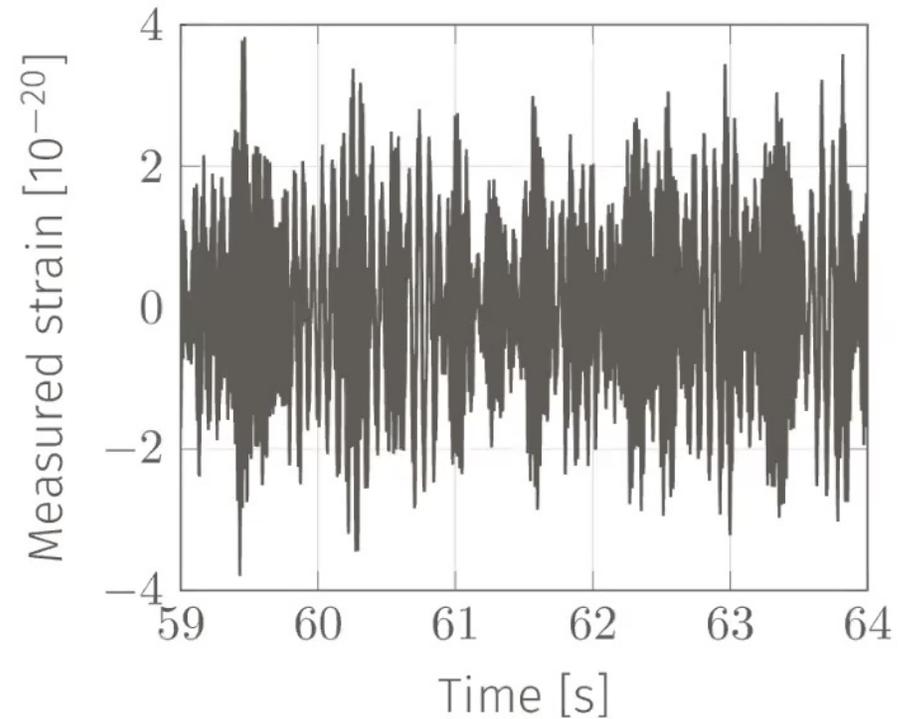
Expected constraints with future detectors



Bozzola&Paschalidis (2021b)

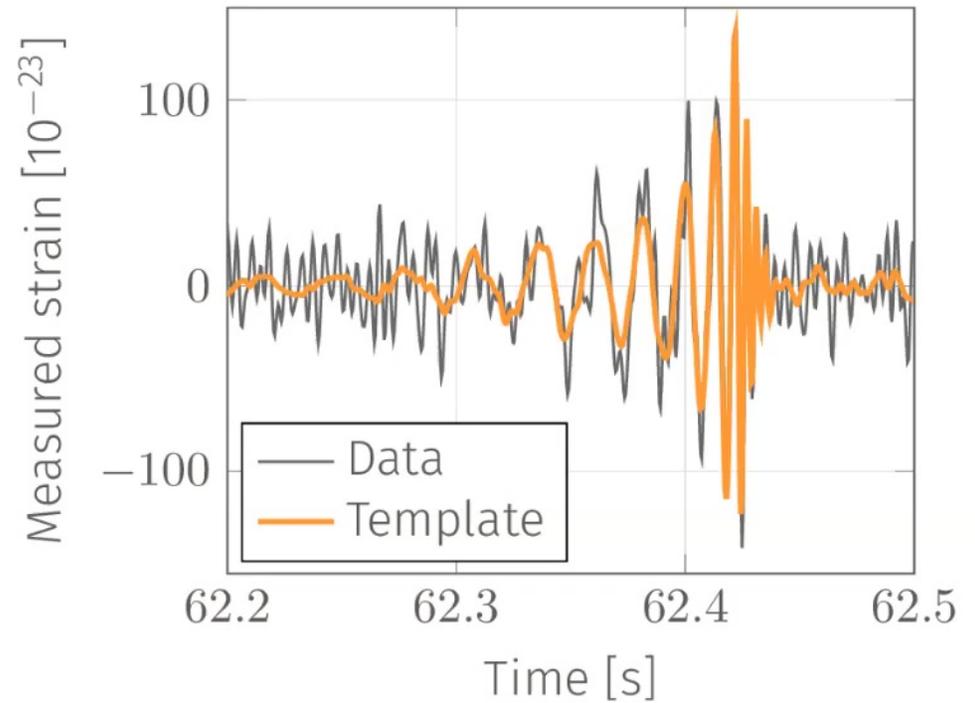
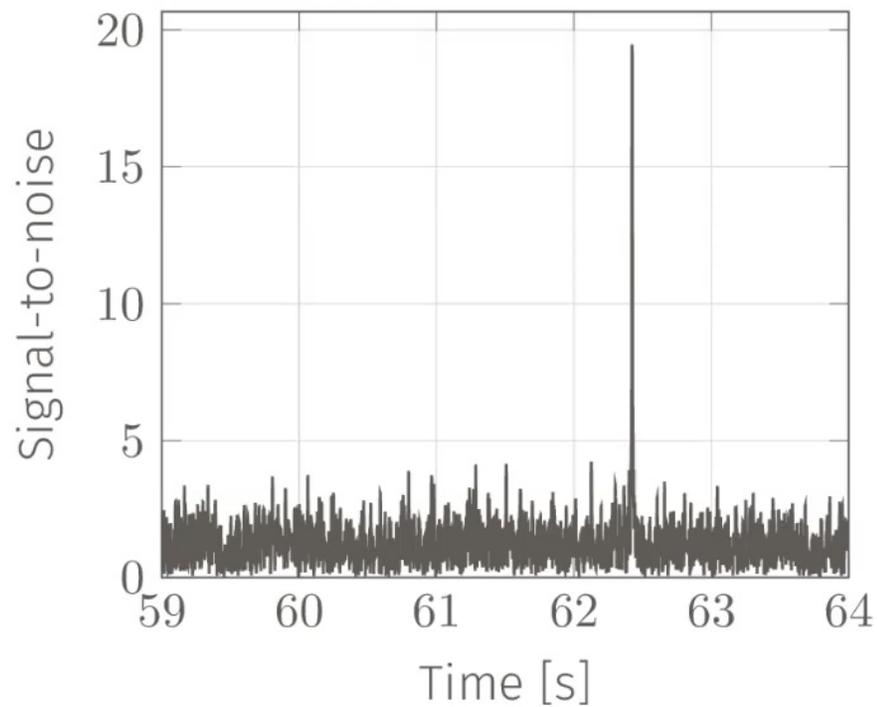
GW detectors look for needles in haystacks

- Noise \gg signal
- GW detection relies on matched filtering

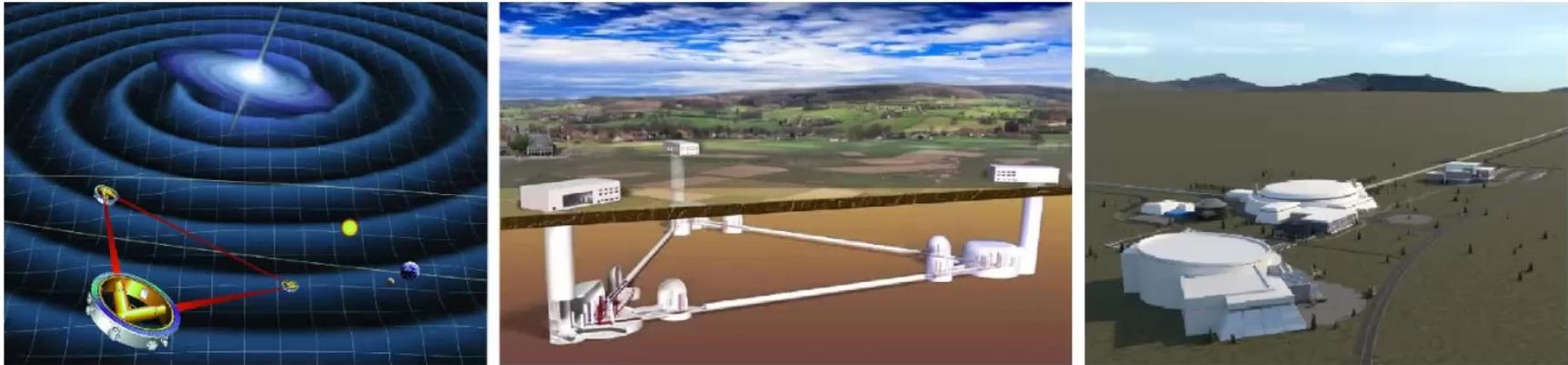


W models are critical for detection and interpretation!

Compare NR-calibrated models with data



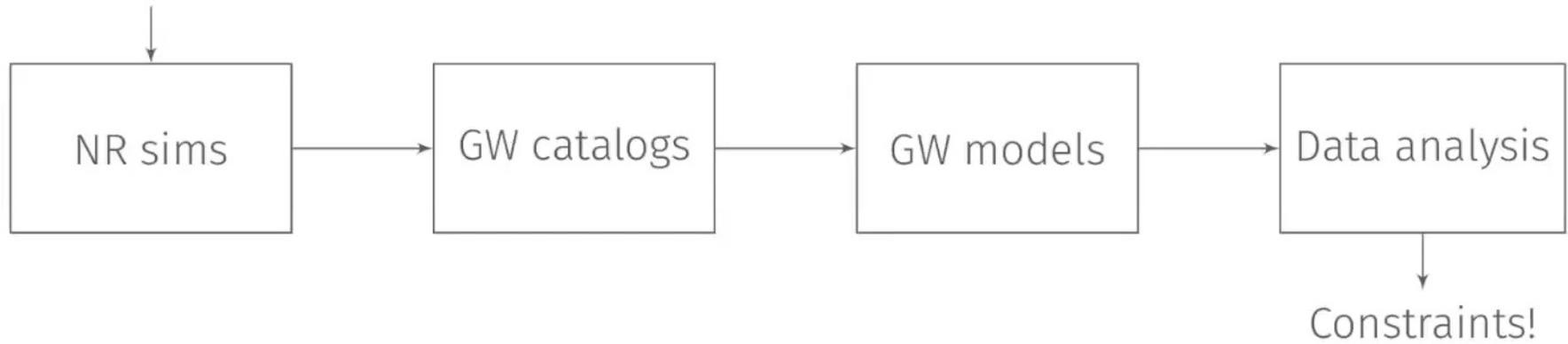
Future detectors require better and more varied simulations



- Vacuum GR waveforms are not accurate enough
- Going off the beaten path to understand biases and degeneracies
- Einstein-Maxwell is a great way to develop tools and techniques for gravitational-wave astronomy beyond vacuum GR

I am building...

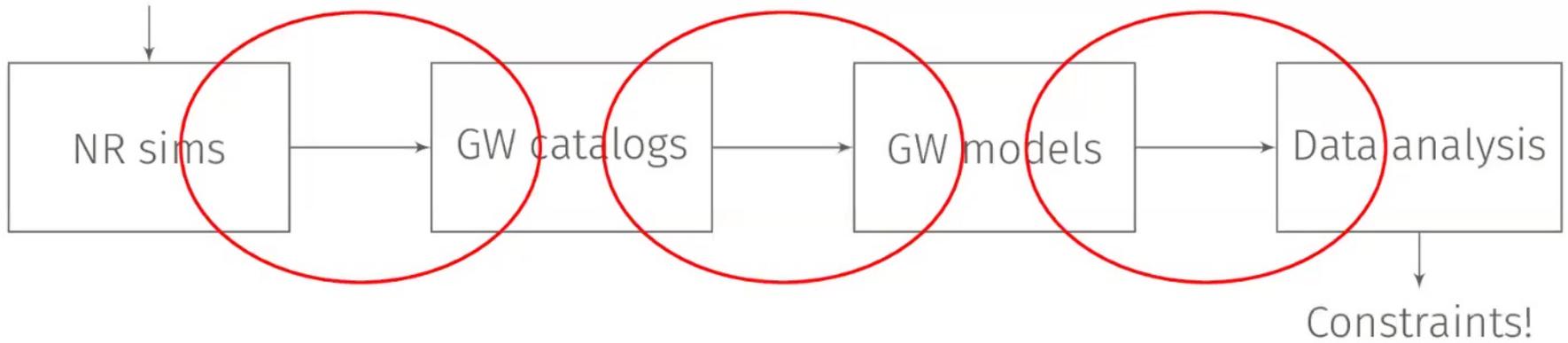
Unmodeled
merger effects



Skip Ad ▶

I am building...

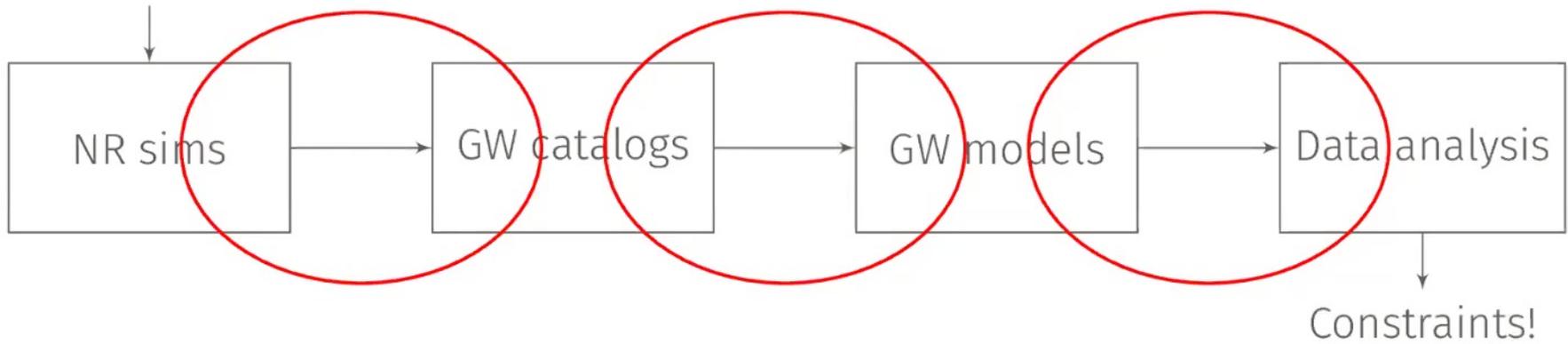
Unmodeled
merger effects



Skip Ad ▶

I am building...

Unmodeled
merger effects

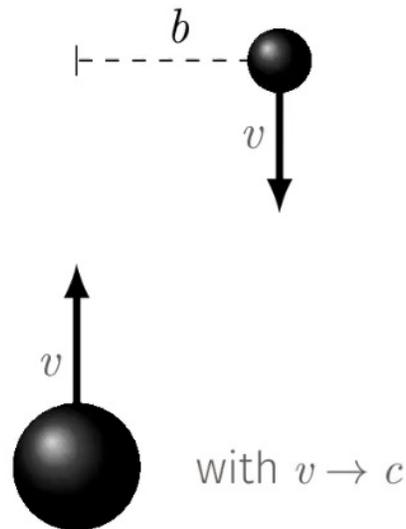


and **theoretical/numerical experiments...**

Skip Ad ▶

29

erged BHs are optimal laboratories for testing conjectures



Energetic events, a few parameters



Extreme conditions in a **controlled** environment

Numerical Relativity

Astrophysics

Modified Gravity

→ Initial data

→ Stable evolutions

→ New formalisms

→ Full NR waveforms

→ Bounds on Q/M

→ Future bounds

→ Bound on STVG α

Current work

→ GW pipeline beyond vacuum GR

→ Cosmic censorship in quasi-circular mergers (+ C. Worley)

→ Quasi-normal-modes and 3G detectors (+ G. Carullo, M. de Amicis, V. Cardoso)

→ Hyperbolic encounters (+ M. Smith)

→ Effects of plasma (+ R. Luna, M. Zilhão, V. Cardoso)



Ask me about:

$$R_{\mu\nu} - \frac{1}{2}Rg_{\mu\nu} = 8\pi T_{\mu\nu}^{EM}$$

$$\nabla_{\mu} F^{\mu\nu} = 4\pi J^{\nu}$$

$$\nabla_{\mu} \star F^{\mu\nu} = 0$$

Einstein-Maxwell

Bozzola&Paschalidis (2019, 2021, 2021b, in prep), Bozzola (2021, 2022), Luna&Bozzola+ (2022)