

Title: An operator-algebraic formulation of self-testing

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Abstract: We give a new definition of self-testing for correlations in terms of states on  $C^*$ -algebras. We show that this definition is equivalent to the standard definition for any class of finite-dimensional quantum models which is closed under submodels and direct sums, provided that the correlation is extremal and has a full-rank model in the class. This last condition automatically holds for the class of POVM quantum models, but does not necessarily hold for the class of projective models by a result of Mancinska and Kaniewski. For extremal binary correlations and for extremal synchronous correlations, we show that any self-test for projective models is a self-test for POVM models. The question of whether there is a self-test for projective models which is not a self-test for POVM models remains open. An advantage of our new definition is that it extends naturally to commuting operator models. We show that an extremal correlation is a self-test for finite-dimensional quantum models if and only if it is a self-test for finite-dimensional commuting operator models, and also observe that many known finite-dimensional self-tests are in fact self-tests for infinite-dimensional commuting operator models.

Zoom link: <https://pitp.zoom.us/j/95783943431?pwd=SDFyQVVZR1d4WlVNSDZ4OENzSmJQUT09>

# An operator-algebraic formulation of self-testing

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# Overview

- 1 Correlations and quantum models
- 2 Standard definition of self-testing
- 3  $C^*$ -algebraic formulation of correlations and models
- 4 Abstract state self-testing
- 5 Results and further questions

## Bell scenarios

- Two spatially-separated parties (Alice and Bob) are given measurement settings  $x$  and  $y$  drawn from some finite sets  $X$  and  $Y$  respectively, and return measurement outcomes  $a$  and  $b$  drawn from finite sets  $A$  and  $B$ .
- Alice and Bob's actions are modelled by measurement operators  $\{M_a^x : a \in A\}, x \in X$  and  $\{N_b^y : b \in B\}, y \in Y$  on local Hilbert spaces  $H_A$  and  $H_B$ .
- If the joint system is in the pure state  $|\psi\rangle \in H_A \otimes H_B$ , then the probability that Alice and Bob measure outcomes  $a, b$  on inputs  $x, y$  is

$$p(a, b|x, y) = \langle \psi | M_a^x \otimes N_b^y | \psi \rangle .$$



# Correlations and quantum models

- A collection  $p = \{p(a, b|x, y)\}_{x,y,a,b}$ , such that  $\sum_{a,b} p(a, b|x, y) = 1$ , for all  $x, y \in X \times Y$  is called a **correlation** (or behaviour).
- The collection  

$$S = (H_A, H_B, \{M_a^x : a \in A, x \in X\}, \{N_b^y : b \in B, y \in Y\}, |\psi\rangle)$$
  - i)  $H_A$  and  $H_B$  are finite dimensional Hilbert spaces,
  - ii)  $\{M_a^x : a \in A\}$  is a POVM on  $H_A$  for all  $x \in X$ ,
  - iii)  $\{N_b^y : b \in B\}$  is a POVM on  $H_B$  for all  $y \in Y$ , and
  - iv)  $|\psi\rangle \in H_A \otimes H_B$  is a vector state,
- Such that  $p(a, b|x, y) = \langle \psi | M_a^x \otimes N_b^y | \psi \rangle$  for all  $(a, b, x, y) \in A \times B \times X \times Y$ , is a (POVM) **quantum model** for  $p$ .

# Quantum correlations

- A correlation  $p$  is a **quantum correlation** if it has a quantum model.
- The **set of quantum correlations**  $C_q(X, Y, A, B)$  is a convex subset of  $\mathbb{R}_{\geq 0}^{A \times B \times X \times Y}$ .
- The sets  $C_q$  are not closed in general [Slofstra'19], and the closure of  $C_q$  in  $\mathbb{R}_{\geq 0}^{A \times B \times X \times Y}$  is denoted by  $C_{qa} = C_{qa}(X, Y, A, B)$ .
- While a quantum correlation  $p$  can be directly observed from a Bell scenario, the model  $S$  cannot.
- In fact, there are typically many different models for a given correlation  $p$ .

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- While a quantum correlation  $p$  can be directly observed from a Bell scenario, the model  $S$  cannot.
- In fact, there are typically many different models for a given correlation  $p$ .
- Can we hope to learn anything about the model from observing only its correlations?

### Definition (Local dilation)

A quantum model

$$\tilde{S} = \left( \tilde{H}_A, \tilde{H}_B, \{\tilde{M}_a^x : a \in A, x \in X\}, \{\tilde{N}_b^y : b \in B, y \in Y\}, |\tilde{\psi}\rangle \right),$$

is a **local dilation** of another model

$$S = (H_A, H_B, \{M_a^x : a \in A, x \in X\}, \{N_b^y : b \in B, y \in Y\}, |\psi\rangle).$$

If there are finite dimensional Hilbert spaces  $H_A^{aux}$  and  $H_B^{aux}$ , a vector state  $|aux\rangle \in H_A^{aux} \otimes H_B^{aux}$ , and isometries

$I_A : H_A \rightarrow \tilde{H}_A \otimes H_A^{aux}$  and  $I_B : H_B \rightarrow \tilde{H}_B \otimes H_B^{aux}$  such that

$$(I_A \otimes I_B) \cdot (M_a^x \otimes N_b^y) |\psi\rangle = \left( \tilde{M}_a^x \otimes \tilde{N}_b^y |\tilde{\psi}\rangle \right) \otimes |aux\rangle$$

for all  $(a, b, x, y) \in A \times B \times X \times Y$ .

- We write  $S \succeq \tilde{S}$  to mean that  $\tilde{S}$  is a **local dilation** of  $S$ .
- The relation  $\succeq$  is a preorder i.e. it is transitive and reflexive, but not anti-symmetric.
- The local isometries are necessary to account for the addition of ancilla systems.

### Definition (Self-test for quantum models)

A correlation  $p$  is a **self-test for the class of quantum models** if there is an **ideal** quantum model  $\tilde{S}$  for  $p$ , such that  $S \succeq \tilde{S}$  for any **employed** quantum model  $S$  for  $p$ .

## Summary of Self-Testing

- This definition of a self-test is somewhat ad-hoc: it's clear that some type of equivalence between models is required for the definition, but why exactly this equivalence with local isometries?
- Several definitions of self-testing have appeared since the inception of self-testing in [Mayers-Yao'03], with a rough consensus seeming to form around the above definition only recently, although variants still exist.
- Christandl, Mancinska, and Houghton-Larsen have additionally pointed out a lack of physical or operational interpretation of this definition of self-testing, which they address in their recent work [CMH-L'22].
- Despite this ad-hoc nature, the definition has been very successful [Supic-Bowles'20]. Among other achievements, self-tests have been used in proofs of device-independent cryptography; and self-testing is a key component to the recent proof of  $\text{MIP}^* = \text{RE}$ .



- A quantum model  $S$  for a correlation  $p$  is **projective** (or a **PVM quantum model**) if the operators  $M_a^x$  and  $N_b^y$  are self-adjoint projections for all  $x, y, a, b$ .
- By Naimark dilation, every element of  $C_q$  has a projective quantum model, and by restricting to the support projection, every element of  $C_q$  also has a full-rank (but not necessarily projective) quantum model.
- A quantum model is **full-rank** if  $\dim H_A = \dim H_B$ , and the Schmidt rank of  $|\psi\rangle$  is  $\dim H_A$ .

- A (unital)  $C^*$ -**algebra**  $\mathcal{A}$  is a complex Banach  $*$ -algebra for which the  $C^*$ -identity  $\|a^*a\| = \|a\|^2$  holds, for all  $a \in \mathcal{A}$ .
- An **abstract state** is a linear functional  $f : \mathcal{A} \rightarrow \mathbb{C}$  such that  $f(b) \geq 0$  for all positive elements  $b = a^*a \in \mathcal{A}$  (positivity) and  $f(1) = 1$  (normalized).
- Given a state  $f : \mathcal{A} \rightarrow \mathbb{C}$  the Gel'fand-Naimark-Segal (GNS) construction is: a representation  $\pi_f$  of  $\mathcal{A}$ , a Hilbert space  $H_f$ , and a unique vector state  $|\xi_f\rangle \in H_f$  (called the **cyclic vector** for  $\pi_f$ ), such that  $f(a) = \langle \xi_f | \pi_f(a) | \xi_f \rangle$  for all  $a \in \mathcal{A}$ .
- Conversely, if  $\pi$  is a representation of  $\mathcal{A}$  on Hilbert space  $H$  with a cyclic vector  $|\xi\rangle \in H$  for  $\mathcal{A}$ , then  $f(a) = \langle \xi | \pi(a) | \xi \rangle$  is a positive linear functional on  $\mathcal{A}$ , and  $(\pi, H, |\xi\rangle)$  is unitarily equivalent to the **GNS representation**  $(\pi_f, H_f, |\xi_f\rangle)$  of  $f$ .



- Define the **POVM algebra**  $\mathcal{A}_{POVM}^{X,A}$  to be the universal  $C^*$ -algebra generated by positive contractions  $e_a^x$ ,  $x \in X$ ,  $a \in A$ , subject to the relations:  $\sum_{a \in A} e_a^x = 1$  for all  $x \in X$ .
- Representations  $\phi : \mathcal{A}_{POVM}^{X,A} \rightarrow \mathcal{B}(H)$  on a Hilbert space  $H$ , correspond uniquely to POVMs on  $H$  with  $\phi(e_a^x) = M_a^x$  for  $a \in A$ ,  $x \in X$ .
- Let  $m_a^x := e_a^x \otimes 1$  and  $n_b^y := 1 \otimes e_b^y$  to be the generators for the algebraic tensor product  $\mathcal{A}_{POVM}^{X,A} \otimes_{alg} \mathcal{A}_{POVM}^{Y,B}$ . Completing this using the *min* construction gives us the bipartite POVM  $C^*$ -algebra  $\mathcal{A}_{POVM}^{X,A} \otimes_{min} \mathcal{A}_{POVM}^{Y,B}$ .

- If  $S$  is a quantum model for  $p \in C_q$ , then there is a unique representation  $\phi_A \otimes \phi_B$  of the  $C^*$ -algebra  $\mathcal{A}_{POVM}^{X,A} \otimes_{\min} \mathcal{A}_{POVM}^{Y,B}$  with  $\phi_A(m_a^x) = M_a^x$  and  $\phi_B(n_b^y) = N_b^y$  for all  $(a, b, x, y) \in A \times B \times X \times Y$ .
- Hence, a quantum model  $S$  can be equivalently expressed as  $S = (\phi_A \otimes \phi_B, H_A \otimes H_B, |\psi\rangle)$ .
- The abstract state  $f_S : \mathcal{A}_{POVM}^{X,A} \otimes_{\min} \mathcal{A}_{POVM}^{Y,B} \rightarrow \mathbb{C}$  defined by  $f_S(x) := \langle \psi | (\phi_A \otimes \phi_B)(x) | \psi \rangle$  satisfies

$$f_S(m_a^x \cdot n_b^y) = \langle \psi | \pi_A(m_a^x) \otimes \pi_B(n_b^y) | \psi \rangle = p(a, b | x, y).$$

We refer to  $f_S$  as the **abstract state defined by  $S$** .

- If  $p \in C_q$  then  $f_S$  is **finite dimensional**, since  $H_A \otimes H_B$  is finite dimensional.

- Conversely, if  $f$  is a finite dimensional state on  $\mathcal{A}_{POVM}^{X,A} \otimes_{min} \mathcal{A}_{POVM}^{X,A}$ , then applying the double commutant theorem to a GNS representation of  $f$  yields a quantum model  $S$  such that  $f = f_S$ .
- In particular, a correlation  $p \in \mathbb{R}_{\geq 0}^{A \times B \times X \times Y}$  belongs to  $C_q$  if and only if there is a finite dimensional state  $f$  on  $\mathcal{A}_{POVM}^{X,A} \otimes_{min} \mathcal{A}_{POVM}^{X,A}$  with  $f(m_a^x \cdot n_b^y) = p(a, b|x, y)$  for all  $(a, b, x, y) \in A \times B \times X \times Y$ .
- Consequently, a correlation  $p$  belongs to  $C_{qa}$  if and only if there is a state  $f$  on  $\mathcal{A}_{POVM}^{X,A} \otimes_{min} \mathcal{A}_{POVM}^{X,A}$  with  $f(m_a^x \cdot n_b^y) = p(a, b|x, y)$  for all  $(a, b, x, y) \in A \times B \times X \times Y$ .

- Let  $\mathcal{A}_{PVM}^{X,A}$  be the quotient of  $\mathcal{A}_{POVM}^{X,A}$  by the relations  $(e_a^x)^* = (e_a^x)^2 = e_a^x$  for all  $x \in X, a \in A$ . If

$$q : \mathcal{A}_{POVM}^{X,A} \otimes_{\min} \mathcal{A}_{POVM}^{Y,B} \rightarrow \mathcal{A}_{PVM}^{X,A} \otimes_{\min} \mathcal{A}_{PVM}^{Y,B}$$

is the quotient homomorphism, and  $f$  is a state on  $\mathcal{A}_{PVM}^{X,A} \otimes_{\min} \mathcal{A}_{PVM}^{Y,B}$ , then  $f \circ q$  is a state  $\mathcal{A}_{POVM}^{X,A} \otimes_{\min} \mathcal{A}_{POVM}^{Y,B}$ .

- The pullback map  $q^* : f \mapsto f \circ q$  is an injection, and hence identifies states on  $\mathcal{A}_{PVM}^{X,A} \otimes_{\min} \mathcal{A}_{PVM}^{Y,B}$  with a subset of states on  $\mathcal{A}_{POVM}^{X,A} \otimes_{\min} \mathcal{A}_{POVM}^{Y,B}$ .
- We say that a state on  $\mathcal{A}_{POVM}^{X,A} \otimes_{\min} \mathcal{A}_{POVM}^{Y,B}$  is **projective** if it belongs to the image of  $q^*$ .

- An abstract state  $f$  on  $\mathcal{A}_{POVM}^{X,A} \otimes_{\min} \mathcal{A}_{POVM}^{Y,B}$  is finite dimensional (resp. projective & finite dimensional) if and only if  $f = f_S$  for some quantum model (resp. projective quantum model)  $S$ .
- And, if  $S$  is a quantum model for  $p \in C_q$ , then  $f_S(m_a^x \cdot n_b^y) = \langle \psi | \pi_A(m_a^x) \otimes \pi_B(n_b^y) | \psi \rangle = p(a, b | x, y)$ , for all  $a, b, x, y \in A \times B \times X \times Y$ .

### Definition (Abstract state self-test)

Let  $\mathcal{S}$  be a subset of states on  $\mathcal{A}_{POVM}^{X,A} \otimes_{\min} \mathcal{A}_{POVM}^{Y,B}$ . A correlation  $p$  is an **abstract state self-test for  $\mathcal{S}$**  if there exists a unique abstract state  $f \in \mathcal{S}$  with correlation  $p$ .



## Theorem (PSYY'22)

Suppose  $p \in C_q(X, Y, A, B)$  is an extreme point. Then:

- ①  $p$  is a self-test for the class of quantum models if and only if  $p$  is an abstract state self-test for finite dimensional states.
- ② If  $p$  has a full-rank projective quantum model, then  $p$  is a self-test for projective quantum models if and only if  $p$  is an abstract state self-test for projective finite dimensional states.

- In (2), if  $p$  is an abstract state self-test for projective finite dimensional states, then  $p$  is a self-test for projective quantum models even if  $p$  does not have a full-rank projective quantum model.
- We do not know whether the hypothesis that  $p$  have a full-rank projective quantum model is required for the “projective self-test implies unique projective abstract state” direction.

### Lemma (PSYY'22)

Suppose  $p \in C_q(X, Y, A, B)$  is an extreme point in  $C_q$  and is a self-test for quantum models. Then:

- (a) the unique state  $f$  on  $\mathcal{A}_{\text{POVM}}^{X,A} \otimes_{\min} \mathcal{A}_{\text{POVM}}^{Y,A}$  achieving  $p$  is projective,
- (b)  $p$  is a self-test for projective quantum models,
- (c) any full-rank quantum model for  $p$  is projective,

### Remark

We do not know whether there is a correlation  $p \in C_q$  which is a self-test for projective quantum models but not a self-test for POVM quantum models. Although, the above shows that any such example cannot have a full-rank projective quantum model.

- In general, we do not know if being a self-test for the class of projective models implies being a self-test for the class of (POVM) quantum models.
- Mančinska and Kaniewski have recently shown that there are correlations which **do not have a full-rank projective models**.
- Can this lead to a counterexample? Or perhaps whenever  $p$  is a self-test amongst projective models, there is always a full-rank projective model. This remains unknown.
- General statements aside, we are able to show that in many known cases it suffices to be a self-test amongst projective models.



- A correlation  $p \in C_q(X, Y, A, B)$  is said to be a **synchronous correlation** if  $A = B$ ,  $X = Y$ , and  $p(a, b|x, x) = 0$  whenever  $a \neq b$ .
- A correlation  $p \in C_q(X, Y, A, B)$  is said to a **binary correlation** if  $|A| = |B| = 2$ .

## Theorem (PSYY'22)

*If  $p$  is a synchronous or binary correlation and is an extreme point in  $C_q$ , then the following statements are equivalent:*

- ①  *$p$  is a self-test for quantum models.*
  - ②  *$p$  is an abstract state self-test for finite dimensional states.*
  - ③  *$p$  is a self-test for projective quantum models.*
  - ④  *$p$  is an abstract state self-test for projective finite dimensional states.*
- This implies that many self-tests for projective models, such as in the CHSH game, Mermin-Peres magic square game, etc. are also self-tests for POVM quantum models.

- Recall that a **tensor product (POVM) model** for a correlation  $p$  is a tuple

$$S = (H_A, H_B, \{M_a^x : a \in A, x \in X\}, \{N_b^y : b \in B, y \in Y\}, |\psi\rangle),$$

where the Hilbert spaces  $H_A$  and  $H_B$  are not restricted to be finite dimensional.

- Additionally, we do not know whether the set of states achieving  $p \in C_{qs}$  contains only abstract states with tensor product models
- To our knowledge, it is an open problem to characterize the subset of abstract states on  $\mathcal{A}_{POVM}^{X,A} \otimes_{min} \mathcal{A}_{POVM}^{Y,B}$  which achieve correlations  $p \in C_{qs}(X, Y, A, B)$ . In particular, are these the states with a tensor product representation?

- A **commuting operator POVM model** for a correlation  $p$  is a tuple  $(H, \{M_a^x : a \in A, x \in X\}, \{N_b^y : b \in B, y \in Y\}, |\psi\rangle)$ , where

- ①  $H$  is a Hilbert space,
- ②  $\{M_a^x : a \in A\}, x \in X$  and  $\{N_b^y : b \in B\}, y \in Y$  are POVMs on  $H$  such that

$$M_a^x N_b^y = N_b^y M_a^x$$

for all  $(a, b, x, y) \in A \times B \times X \times Y$ , and

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- i)  $H$  is a Hilbert space,
- ii)  $\{M_a^x : a \in A\}, x \in X$  and  $\{N_b^y : b \in B\}, y \in Y$  are POVMs on  $H$  such that

$$M_a^x N_b^y = N_b^y M_a^x$$

for all  $(a, b, x, y) \in A \times B \times X \times Y$ , and

- iii)  $|\psi\rangle \in H$  is a vector state

such that  $p(a, b|x, y) = \langle \psi | M_a^x \cdot N_b^y | \psi \rangle$  for all  $(a, b, x, y) \in A \times B \times X \times Y$ .

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- ❶  $H$  is a Hilbert space,
- ❷  $\{M_a^x : a \in A\}, x \in X$  and  $\{N_b^y : b \in B\}, y \in Y$  are POVMs on  $H$  such that

$$M_a^x N_b^y = N_b^y M_a^x$$

for all  $(a, b, x, y) \in A \times B \times X \times Y$ , and

- ❸  $|\psi\rangle \in H$  is a vector state

such that  $p(a, b|x, y) = \langle \psi | M_a^x \cdot N_b^y | \psi \rangle$  for all  $(a, b, x, y) \in A \times B \times X \times Y$ .

- Likewise, a projective or **(PVM) commuting operator model** is one where each  $M_a^x$  and  $N_b^y$  is a projection.

- We let  $C_{qc} = C_{qc}(X, Y, A, B)$  be the set of correlations with a commuting operator model.
- It is well-known that  $C_{qc}$  is closed, convex, and contains  $C_{qa}$ , that every correlation in  $C_{qc}$  has a projective commuting operator model, and that the set of correlations with finite dimensional commuting operator models is equal to  $C_q$ .
- It is not apparent from the standard formulation of self-testing, how to give an analogous description of a self-test for correlations with commuting operator models.
- In particular, its not even apparent what a local dilation is for commuting operator models, as there is no longer a tensor product structure.



- However, correlations in  $C_{qc}$  do have a formulation as abstract states on  $C^*$ -algebras.

### Proposition (Fritz'12, JNPPSW'11)

Let  $X, Y, A, B$  be finite sets. A correlation  $p \in \mathbb{R}_{\geq 0}^{A \times B \times X \times Y}$  belongs to  $C_{qc}(X, Y, A, B)$  if and only if there is a state  $f$  on  $\mathcal{A}_{POVM}^{X,A} \otimes_{\max} \mathcal{A}_{POVM}^{Y,B}$  with

$$p(a, b|x, y) = f(m_a^x \cdot n_b^y)$$

for all  $(a, b, x, y) \in A \times B \times X \times Y$ .

- Where  $\mathcal{A}_{POVM}^{X,A} \otimes_{\max} \mathcal{A}_{POVM}^{Y,B}$  is just another way of making the algebraic tensor product of POVM algebras into a  $C^*$ -algebra.



- The definition of an abstract state self-test on  $\mathcal{A}_{POVM}^{X,A} \otimes_{\min} \mathcal{A}_{POVM}^{Y,B}$  suggests a notion of self-test for commuting operator models.

## Definition

Let  $\mathcal{S}$  be a subset of states on  $\mathcal{A}_{POVM}^{X,A} \otimes_{\max} \mathcal{A}_{POVM}^{Y,B}$ . A correlation  $p$  is an **abstract state self-test for  $\mathcal{S}$**  if there exists a unique abstract state  $f \in \mathcal{S}$  with correlation  $p$ . We say that  $p$  is a **commuting operator self-test** if it is an abstract state self-test for all states on  $\mathcal{A}_{POVM}^{X,A} \otimes_{\max} \mathcal{A}_{POVM}^{Y,B}$ .

- There is a surjective homomorphism  $\mathcal{A}_{POVM}^{X,A} \otimes_{\max} \mathcal{A}_{POVM}^{Y,B} \rightarrow \mathcal{A}_{POVM}^{X,A} \otimes_{\min} \mathcal{A}_{POVM}^{Y,B}$ , and this means that states on  $\mathcal{A}_{POVM}^{X,A} \otimes_{\min} \mathcal{A}_{POVM}^{Y,B}$  can be thought of as a subset of states on  $\mathcal{A}_{POVM}^{X,A} \otimes_{\max} \mathcal{A}_{POVM}^{Y,B}$ .

## Definition

Let  $S$ , and  $\tilde{S}$  be two commuting operator models.

- ① We say  $S$  and  $\tilde{S}$  are **equivalent**, and write  $S \cong \tilde{S}$ , if there exists a unitary  $U : H \rightarrow \tilde{H}$  such that
  - ①  $U|\psi\rangle = |\tilde{\psi}\rangle$ , and
  - ②  $UM_a^x U^* = \tilde{M}_a^x$  and  $UN_b^y U^* = \tilde{N}_b^y$  for all  $(a, b, x, y) \in A \times B \times X \times Y$ .
- ②  $S$  is said to be **degenerate** if there exists a non-trivial projection  $\Pi \in \mathcal{B}(H)$  such that  $\Pi|\psi\rangle = |\psi\rangle$  and  $[\Pi, M_a^x] = [\Pi, N_b^y] = 0$  for all  $(a, b, x, y) \in A \times B \times X \times Y$ . In this case, we say  $\tilde{S}$  is a **submodel** of  $S$  if  $\tilde{H} = \Pi H$ ,  $|\tilde{\psi}\rangle = \Pi|\psi\rangle$ , and  $\tilde{M}_a^x = \Pi M_a^x \Pi$ ,  $\tilde{N}_b^y = \Pi N_b^y \Pi$  for all  $(a, b, x, y) \in A \times B \times X \times Y$ . Any commuting operator model is a submodel of itself. A commuting operator model is said to be **nondegenerate** if it is not degenerate.

- A class of commuting operator models  $\mathcal{C}$  is **closed under submodels** if
  - ❶ for any  $S \in \mathcal{C}$ , if  $\tilde{S}$  is a commuting operator model such that  $\tilde{S} \cong S$  then  $\tilde{S} \in \mathcal{C}$ , and
  - ❷ for any  $S \in \mathcal{C}$ , if  $\tilde{S}$  is a submodel of  $S$  then  $\tilde{S} \in \mathcal{C}$ .

### Remark

*For any two commuting operator models  $S$  and  $\tilde{S}$ , we have  $f_S = f_{\tilde{S}}$  whenever  $\tilde{S}$  is equivalent to  $S$ , or  $\tilde{S}$  is a submodel of  $S$ .*

*Moreover, any commuting operator model  $S$  has a nondegenerate submodel  $\tilde{S}$  which is the GNS representation for the abstract state  $f_S$ .*

We can then show that this definition of commuting operator self-test is equivalent to having a unique nondegenerate commuting-operator model.

### Theorem (PSZZ'22)

Let  $\mathcal{C}$  be a class of commuting operator models that is closed under submodels, and let  $\mathcal{S} := \{f_S : S \in \mathcal{C}\}$  be the set of states on  $\mathcal{A}_{\text{POVM}}^{X,A} \otimes_{\max} \mathcal{A}_{\text{POVM}}^{Y,B}$  induced by  $\mathcal{C}$ . Then  $p \in C_{qc}$  is a self-test for  $\mathcal{S}$  if and only if there is a commuting operator model

$$\tilde{S} = (\tilde{H}, \{\tilde{M}_a^x : a \in A, x \in X\}, \{\tilde{N}_b^y : b \in B, y \in Y\}, |\tilde{\psi}\rangle)$$

for  $p$  in  $\mathcal{C}$ , such that for every other commuting operator model

$$S = (H, \{M_a^x : a \in A, x \in X\}, \{N_b^y : b \in B, y \in Y\}, |\psi\rangle),$$

for  $p$  in  $\mathcal{C}$ , there is a submodel of  $S$  which is equivalent to  $\tilde{S}$ .

- It remains to connect this result to something more analogous to the local dilation definition of self-testing.
- This definition for commuting operator self-tests also gives us a definition of self-testing for non-finite dimensional states on  $\mathcal{A}_{POVM}^{X,A} \otimes_{\min} \mathcal{A}_{POVM}^{Y,B}$  (i.e. for correlations  $p \in C_{qa}$  which do not have a finite dimensional model).
- This is potentially useful in examples of self-tests currently being constructed. For instance, Mancinska and Schmidt recently gave an example of a nonlocal game which is a non-robust self-test, by combining a finite dimensional self-test with a game with a perfect  $C_{qa}$  strategy, but no perfect  $C_q$  strategy.
- In our language, this nonlocal game is a finite dimensional abstract-state self-test, but not a self-test for all states on  $\mathcal{A}_{POVM}^{X,A} \otimes_{\min} \mathcal{A}_{POVM}^{Y,B}$ .



## Proposition (PSZZ'22)

Let  $p \in C_q(X, Y, A, B)$ . Then  $p$  is an abstract state self-test for finite dimensional states (resp. projective finite dimensional states) on  $\mathcal{A}_{\text{POVM}}^{X,A} \otimes_{\min} \mathcal{A}_{\text{POVM}}^{Y,B}$  if and only if  $p$  is an abstract state self-test for finite dimensional states (resp. projective finite dimensional states) on  $\mathcal{A}_{\text{POVM}}^{X,A} \otimes_{\max} \mathcal{A}_{\text{POVM}}^{Y,B}$ .

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- As a result, if  $p \in C_q$  is an extreme point, then  $p$  is a self-test for (POVM) quantum models if and only if  $p$  has a unique nondegenerate commuting operator model.
- If, in addition, there exists a projective full-rank quantum model for  $p$ , then  $p$  is a self-test for projective quantum models if and only if  $p$  has a unique nondegenerate commuting operator model.
- This gives a new criterion for  $p \in C_q$  to be a self-test in the finite dimensional case.

- Tsirelson showed that a wide family of correlations in  $C_q$  are in fact commuting operator self-tests.
- To state this result, let  $Cor(X, Y)$  be the set of matrices  $c \in \mathbb{R}^{X \times Y}$  for which there is a Euclidean space  $V$  and vectors  $\{|u_x\rangle\}_{x \in X}$ ,  $\{|v_y\rangle\}_{y \in Y}$  in  $V$  of norm at most 1, such that  $c_{x,y} = \langle u_x | v_y \rangle$  for all  $x, y \in X \times Y$ .
- If  $p \in C_{qc}(X, Y, \mathbb{Z}_2, \mathbb{Z}_2)$  then the matrix  $c$  defined by

$$c_{x,y} = \sum_{a,b \in \mathbb{Z}_2} (-1)^{a+b} p(a, b | x, y)$$

is in  $Cor(X, Y)$ , since if

$S = (H, \{M_a^x : a \in \mathbb{Z}_2, x \in X\}, \{N_b^y : b \in \mathbb{Z}_2, y \in Y\}, |\psi\rangle)$  is a commuting operator model for  $p$ , then

$c_{x,y} = \langle \psi | (M_0^x - M_1^x)(N_0^y - N_1^y) | \psi \rangle$ , where  $\|M_0^x - M_1^x\| \leq 1$  and  $\|N_0^y - N_1^y\| \leq 1$ .



- Let  $C_q^{unbiased} = \{p \in C_q(X, Y, \mathbb{Z}_2, \mathbb{Z}_2) : \sum_{a \in \mathbb{Z}_2} p(0, a | x, y) = \sum_{b \in \mathbb{Z}_2} p(b | x, y), x \in X, y \in Y\}$ .
- Tsirelson showed that  $C_q(X, Y, \mathbb{Z}_2, \mathbb{Z}_2) \rightarrow C_q^{unbiased}(X, Y, \mathbb{Z}_2, \mathbb{Z}_2) \cong Cor(X, Y)$ , restricts to an isomorphism

### Theorem (Tsirelson'87)

If  $p$  is an extreme point of  $C_q^{unbiased}$  and the associated XOR correlation  $c$  has even rank, then  $p$  is a commuting operator self-test.

### Proof.

Since  $p$  is an extreme point of  $C_q^{unbiased}$ , the associated XOR correlation  $c$  is an extreme point in  $Cor(X, Y)$ . By a theorem of Tsirelson, if  $c$  has even rank then all nondegenerate commuting operator models for  $p$  are unitarily equivalent. Hence, there is a unique state on  $\mathcal{A}_{POVM}^{X, \mathbb{Z}_2} \otimes_{\max} \mathcal{A}_{POVM}^{Y, \mathbb{Z}_2}$  achieving  $p$ . □

### Example

It is well known that the unique optimal correlation for the CHSH game is an extreme point of  $C_q^{unbiased}$  with associated XOR correlation matrix

$$\begin{pmatrix} \frac{1}{\sqrt{2}} & \frac{1}{\sqrt{2}} \\ \frac{1}{\sqrt{2}} & \frac{-1}{\sqrt{2}} \end{pmatrix}.$$

Since the associated XOR correlation has rank 2, the optimal CHSH correlation is a commuting operator self-test.

The end

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