

Title: Quantum adiabatic speedup on a class of combinatorial optimization problems

Speakers: Madelyn Cain

Collection: New Frontiers in Machine Learning and Quantum

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Abstract: "One of the central challenges in quantum information science is to design quantum algorithms that outperform their classical counterparts in combinatorial optimization. In this talk, I will describe a modification of the quantum adiabatic algorithm (QAA) [1] that achieves a Grover-type speedup in solving a wide class of combinatorial optimization problem instances. The speedup is obtained over classical Markov chain algorithms including simulated annealing, parallel tempering, and quantum Monte Carlo. I will then introduce a framework to predict the relative performance of the standard QAA and classical Markov chain algorithms, and show problem instances with quantum speedup and slowdown. Finally, I will apply this framework to interpret results from a recent Rydberg atom array experiment [2], which suggest a superlinear speedup in solving the Maximum Independent Set problem on unit-disk graphs.

[1] Farhi et al. (2001) Science 292, 5516

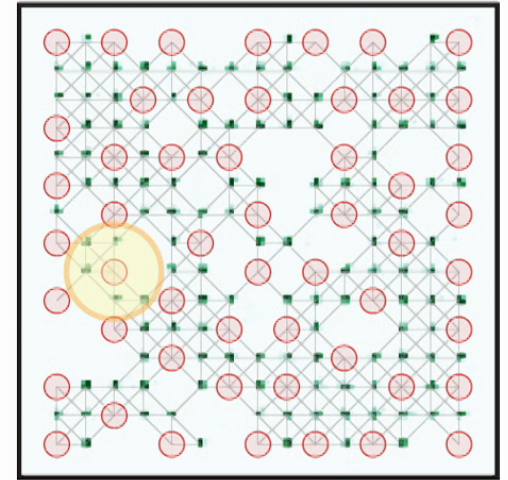
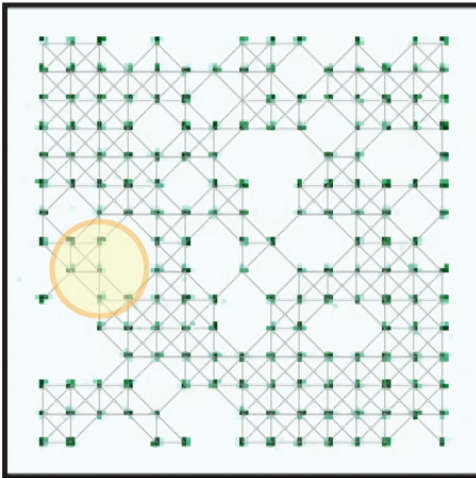
[2] Ebadi et al. (2022) Science 376, 6598"

Quantum adiabatic speedup on a class of combinatorial optimization problems

Madelyn Cain

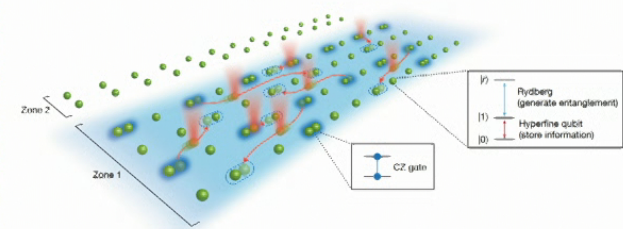
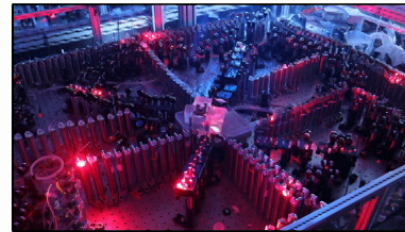
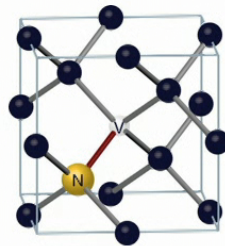
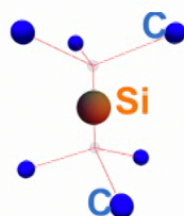
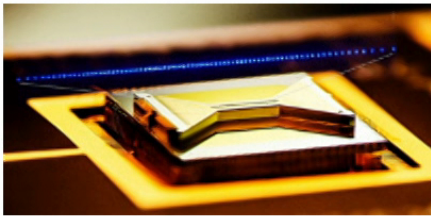
New Frontiers in ML and Quantum

November 22, 2022

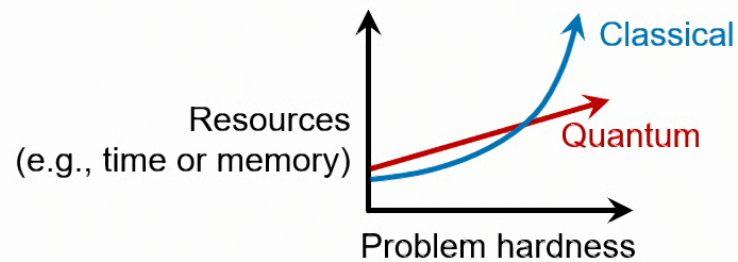


Demonstrating near-term quantum speed-up

Control quantum systems with ≤ 1000 qubits, and no error correction.



Can we demonstrate quantum speed-up on useful problems in the near-term?



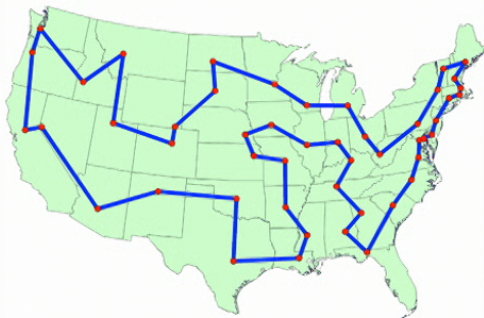
- Shor's algorithm
- Grover's database search
- Solving linear systems
- Quantum dynamics
- Combinatorial optimization etc...

Combinatorial optimization problems

Broad set of problems with the goal: optimize a cost function, subject to constraints.

- NP-hard, intractable to solve in worst-case
- Solving one NP-hard problem is general enough to solve all NP-hard problems

Traveling salesman problem



Minimize: total distance traveled
Constraint: path must include all cities

This talk: Maximum Independent Set



Given a graph (nodes connected by edges),

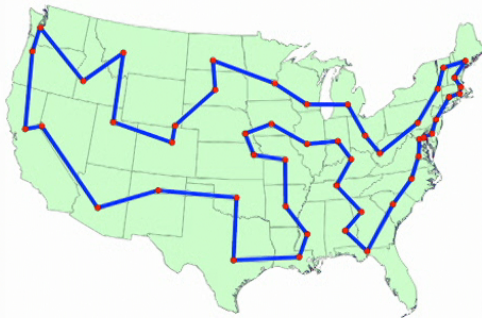
Maximize: size of set of nodes
Independent set constraint: no nodes in set are connected by an edge

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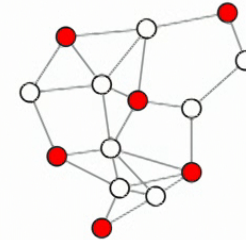
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This talk: Maximum Independent Set

Maximum independent set (MIS) in red



Given a graph (nodes connected by edges),

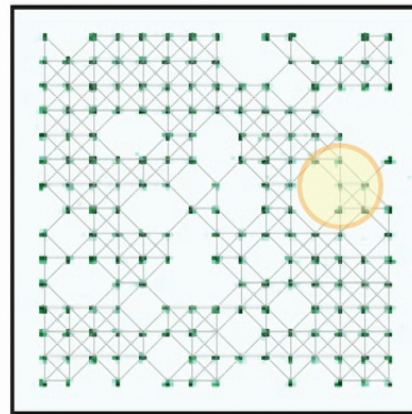
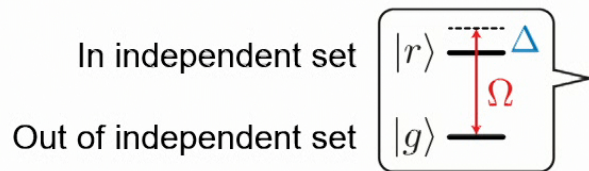
Maximize: size of set of nodes
Independent set constraint: no nodes in set are connected by an edge

Hardware-efficient encoding of Maximum Independent Set

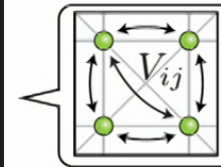
Rydberg atom array platform

Several hundred Rb atoms (qubits) arranged deterministically

Each qubit represents a node in the graph

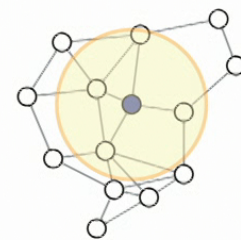


$$V_{ij} = V_0/|\mathbf{x}_i - \mathbf{x}_j|^6$$



Rydberg blockade:
simultaneous Rydberg
excitation forbidden in
blockade radius

**Rydberg interaction naturally encodes the
independent set constraint on unit disk graphs**
(edge between nodes within a unit radius)



**Blockade
radius**

The Quantum Adiabatic Algorithm

System Hamiltonian

Quantum driver

$$H_q = \Omega \sum_{u \in V} |g_u\rangle \langle r_u| + \text{h.c.}$$

Cost Hamiltonian

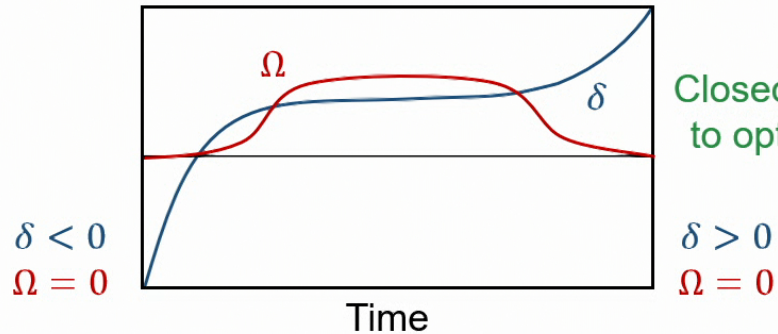
$$H_{\text{cost}} = -\delta \sum_{u \in V} n_u + \sum_{u,v} V_{uv} n_u n_v$$

Favors many nodes in set
Penalizes independent set violations

Notice that ground state of H_{cost} is the MIS!

Adiabatically prepare the ground state of H_{cost}

1. Initialize in easy to prepare many-body ground state
2. Slowly change Hamiltonian
3. Final, hard to prepare, ground state encodes MIS



Closed-loop optimization to optimize δ over time

Initialize atoms in atomic ground state $|g\rangle^N$

Maximize atoms in $|r\rangle$ subject to blockade (corresponds to MIS of graph)

Outline

1. Experimental signatures of a speedup over simulated annealing on certain problem instances
2. Mechanism for observed experimental performance, and a new quantum algorithm with guaranteed speedup

Hard instances for SA

Expected runtime to randomly select MIS $\sim \frac{D_{|MIS|-1}}{D_{|MIS|}}$

D_k = Number (degeneracy) of independent sets of size k

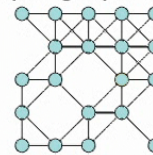
Runtime of any SA algorithm (no. proposed updates)

$$\gtrsim \max_k \frac{D_{k-1}}{nD_k}$$

Among states with the same energy, SA dynamics are a random walk!

Hard MIS unit-disk graphs for SA: exponentially many suboptimal solutions with same energy

Example graph instance:

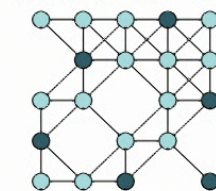


Configuration graph:

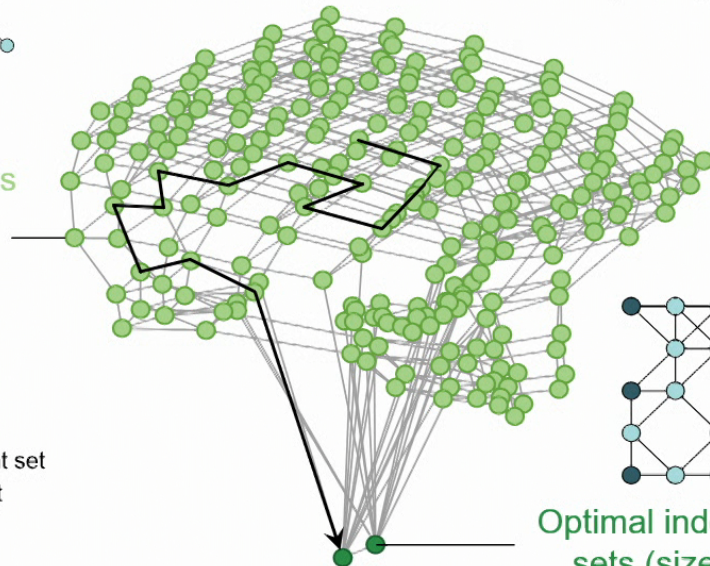
Nodes = independent sets

Edges = sets connected by SA update

Suboptimal independent sets (size $|MIS|-1$)



● Out of independent set
● In independent set



Optimal independent sets (size $|MIS|$)

Hard instances for SA

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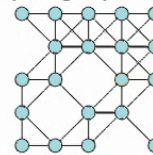
$$\gtrsim \max_k \frac{D_{k-1}}{n D_k}$$

\equiv SA hardness parameter

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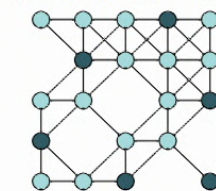


Configuration graph:

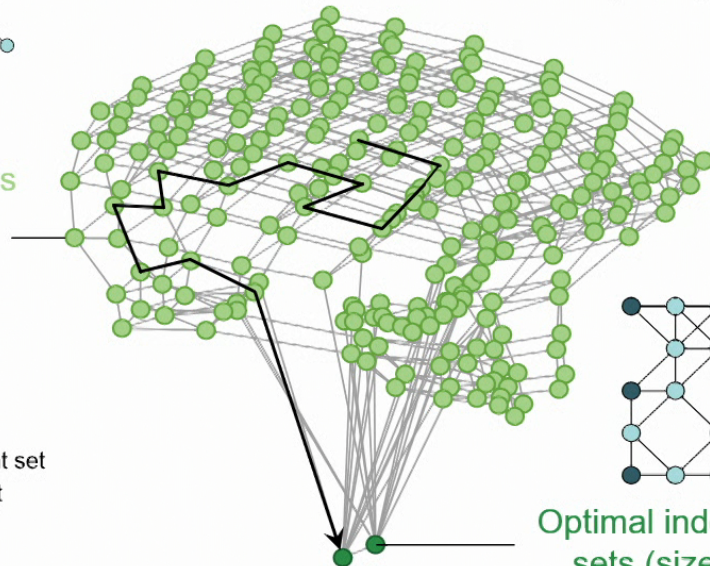
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Hard instances for SA

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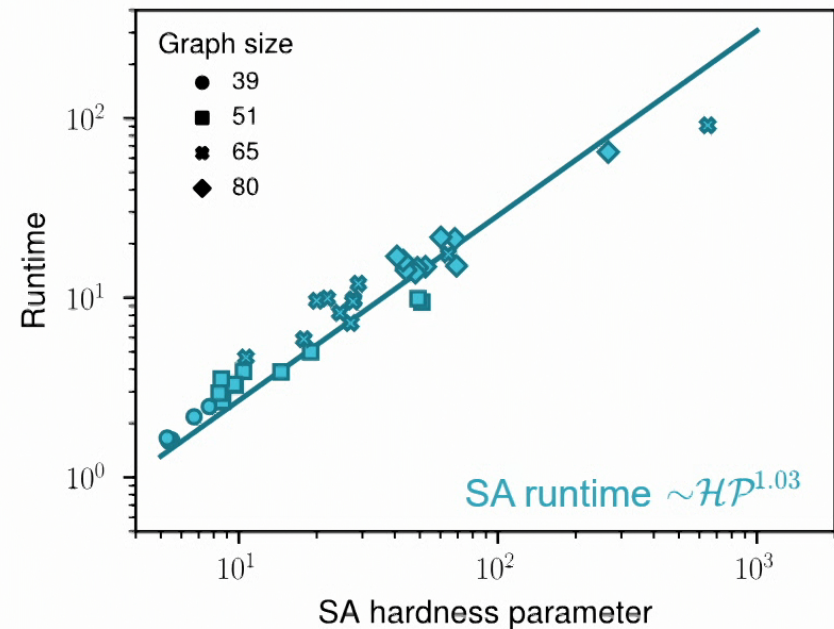
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\equiv **SA hardness parameter**

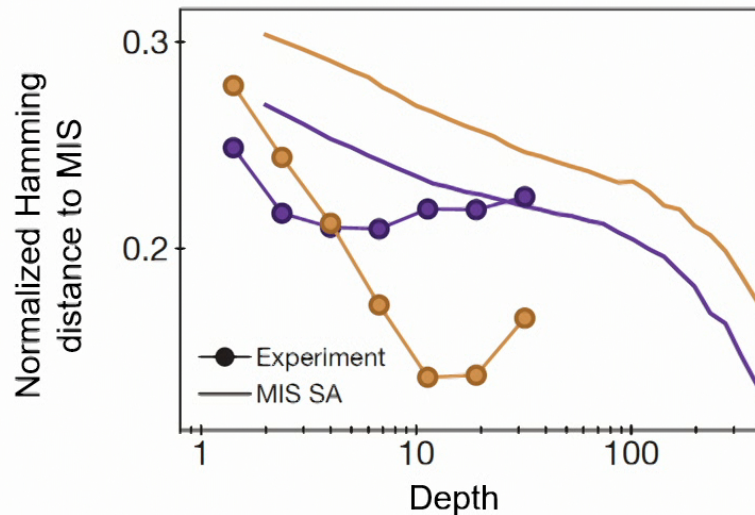
**Generate top 2% hardest instances
maximizing SA hardness parameter**



Experimentally test for quantum speedup

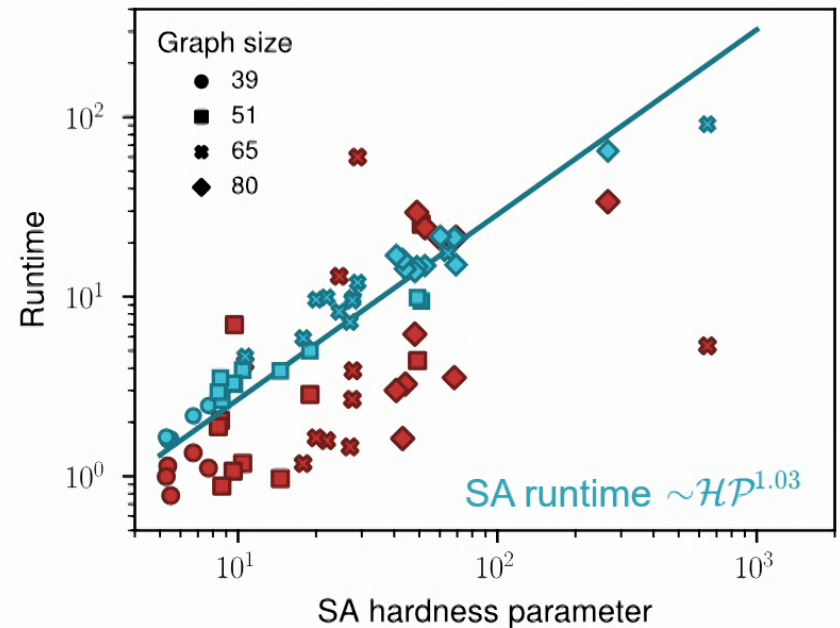
Quantum performance varies on instances with similar SA hardness

- Easy graph for quantum
- Hard graph for quantum



Quantum depth = evolution time / time to flip spin

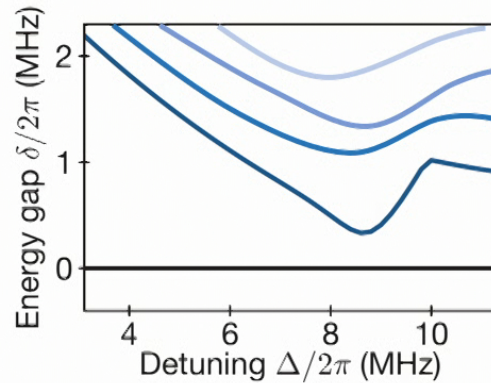
Generate top 2% hardest instances maximizing SA hardness parameter



Quantum hardness not determined by SA hardness parameter!

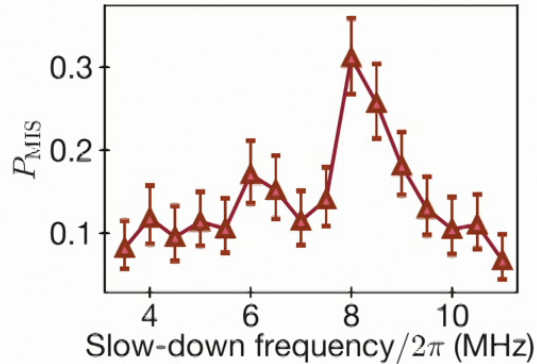
Quantum performance controlled by minimum energy gap

Adiabaticity is limited by the minimum energy gap

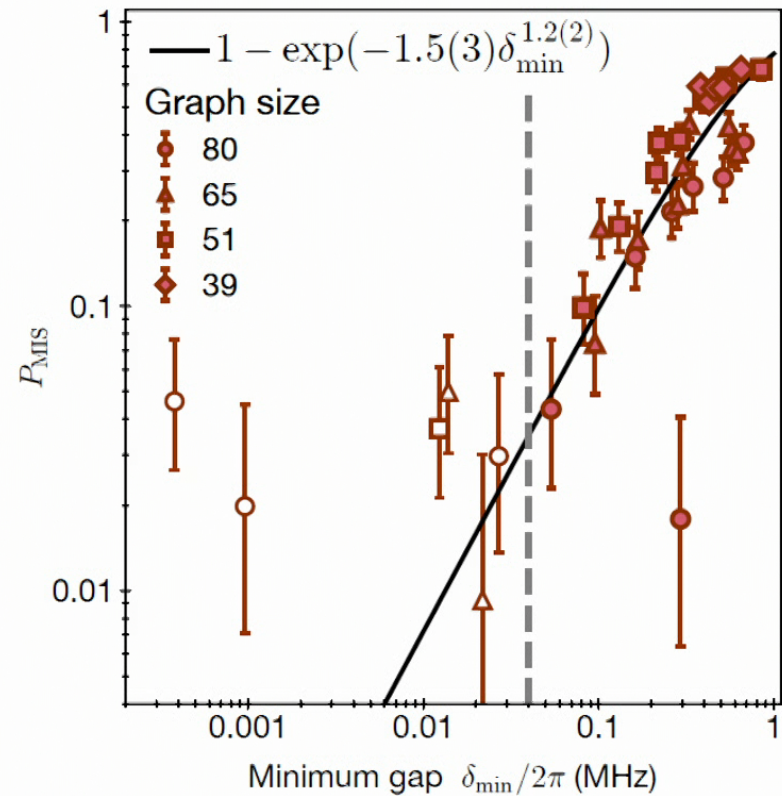


Experiment sensitive to many-body adiabatic gap

MIS probability increases when detuning sweep slowed at gap



Hardness controlled by adiabatic gap
Dependence on gap near-optimal (linear)

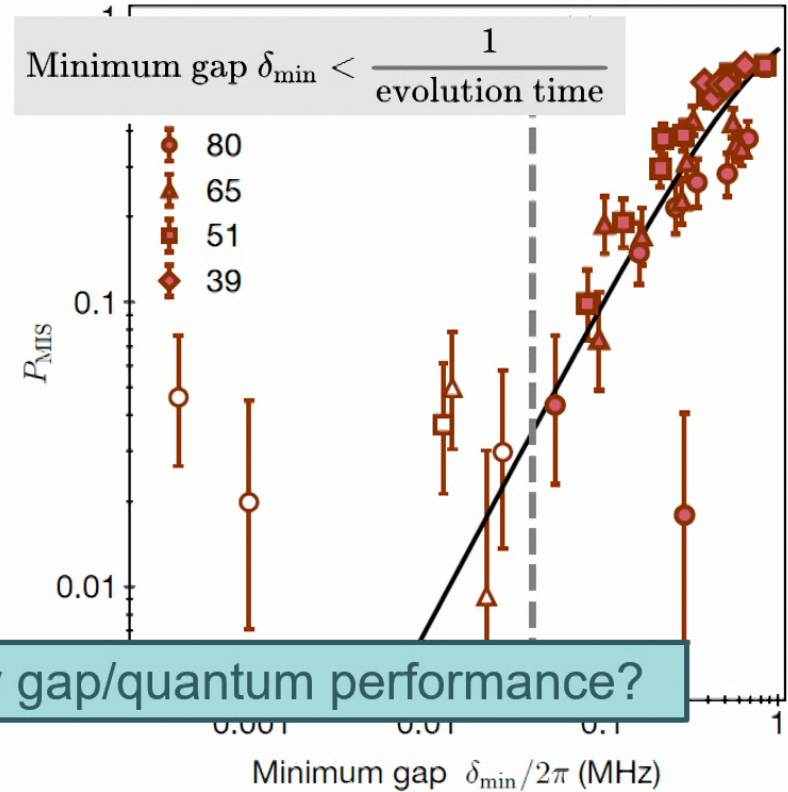
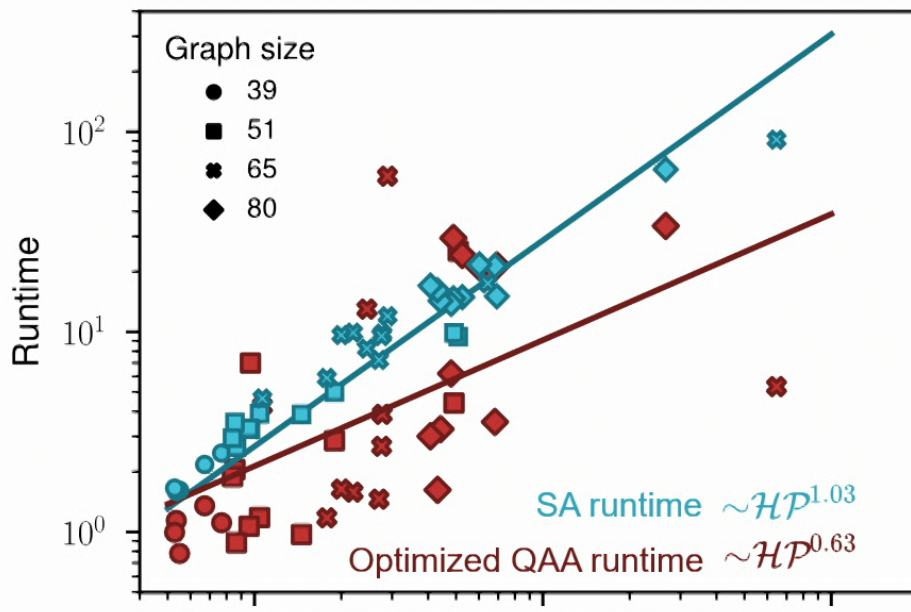


Quantum performance controlled by minimum energy gap

Emergence of superlinear speedup in deep circuit regime

S. Ebadi, A. Keesling*, MC* et al, Science 376, 6598

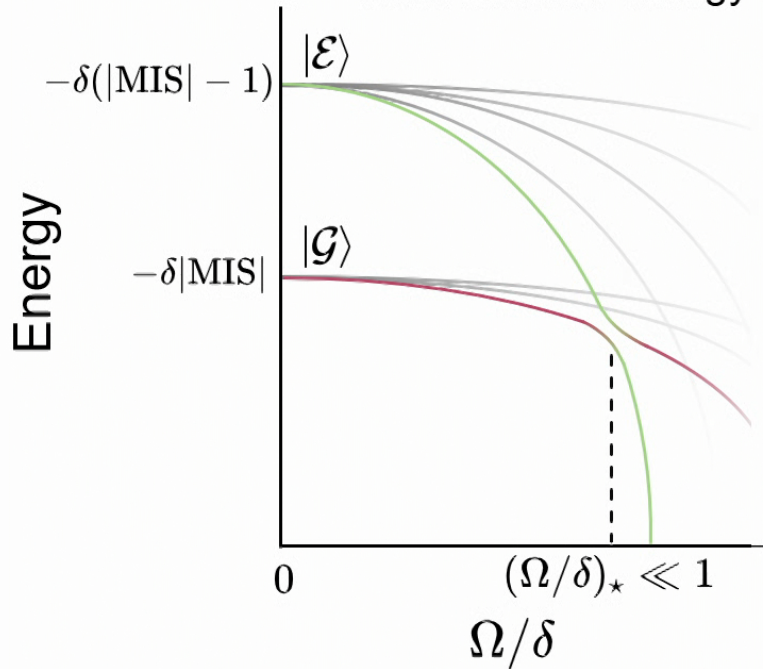
Focus on instances we can optimize (runtime large enough to stay adiabatic)



What controls the minimum energy gap/quantum performance?

Understanding the minimum gap

Ω = Spin-flip energy
 δ = Cost function energy



$$|\mathcal{G}\rangle = \sum_{i \text{ size } |\text{MIS}|} \mathcal{G}_i |i\rangle$$

$$|\mathcal{E}\rangle = \sum_{i \text{ size } |\text{MIS}|-1} \mathcal{E}_i |i\rangle$$

Gap decided by how strongly
 Hamiltonian connects
 $|\mathcal{G}\rangle, |\mathcal{E}\rangle$ via spin-flips

Understand gap from states $|\mathcal{G}\rangle, |\mathcal{E}\rangle$
 Can estimate with perturbation theory for
 $(\Omega/\delta)_* \ll 1$

Understanding the minimum gap

$|\mathcal{G}\rangle, |\mathcal{E}\rangle$ ground states of spin-exchange Hamiltonian by perturbation theory

$$|\mathcal{G}\rangle = \sum_{i \text{ size } |\text{MIS}|} \mathcal{G}_i |i\rangle$$

$$|\mathcal{E}\rangle = \sum_{i \text{ size } |\text{MIS}|-1} \mathcal{E}_i |i\rangle$$

Gap decided by how strongly Hamiltonian connects $|\mathcal{G}\rangle, |\mathcal{E}\rangle$ via spin-flips

Gap behavior in two limits:

1. When $|\mathcal{G}\rangle, |\mathcal{E}\rangle$ are **localized** on certain independent sets,

$$\text{gap} \sim (\Omega/\delta)_{\star}^{\text{hamming distance}}$$

2. When $|\mathcal{G}\rangle, |\mathcal{E}\rangle$ are **delocalized** (uniform superpositions),

$$\text{gap} \sim \sqrt{\frac{D_{|\text{MIS}|}}{D_{|\text{MIS}|-1}}}$$

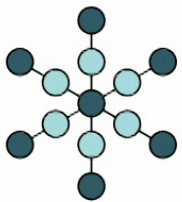
$$\text{Optimized quantum runtime} \sim \sqrt{\frac{D_{|\text{MIS}|-1}}{D_{|\text{MIS}|}}}$$

$$\text{SA runtime} \sim \frac{D_{|\text{MIS}|-1}}{D_{|\text{MIS}|}}$$

Grover-type quadratic speedup over SA!

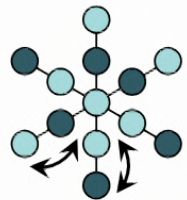
Instance-by-instance performance variation

Star graph has b branches of length ℓ



$|\mathcal{G}\rangle = \text{unique MIS}$

$b = 6, \ell = 2$

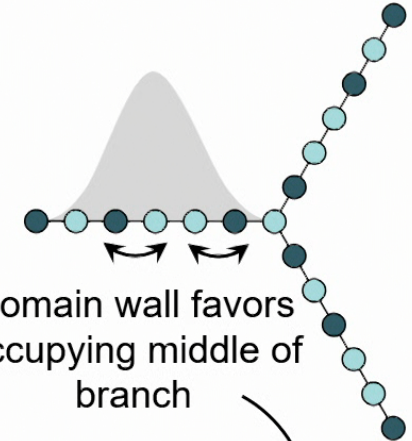


$|\mathcal{E}\rangle = \text{superposition of } \sim \left(\frac{\ell}{2} + 1\right)^b$
 $|\text{MIS}|-1 \text{ states}$

Domain walls

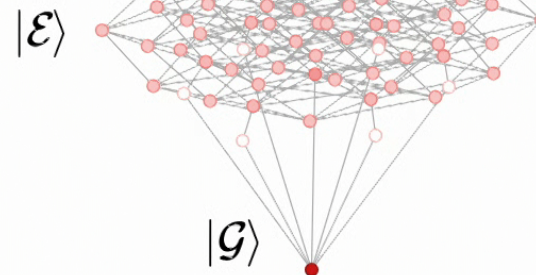
(correspond to the movement of a domain wall on each branch)

Estimate $|\mathcal{E}\rangle$ perturbatively:
 minimize driver Hamiltonian
 at second order in Ω/δ
 (spin-exchange hopping)

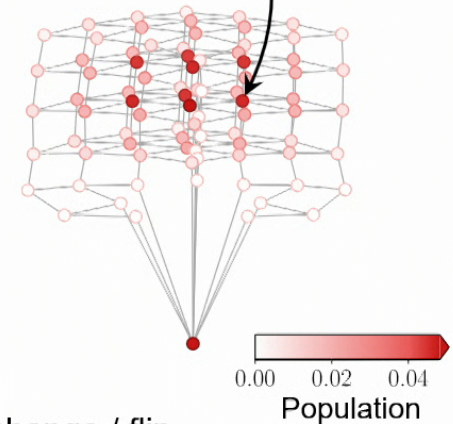


Domain wall favors occupying middle of branch

Delocalized
 $(\ell = 2)$



Localized
 $(\ell > 2)$

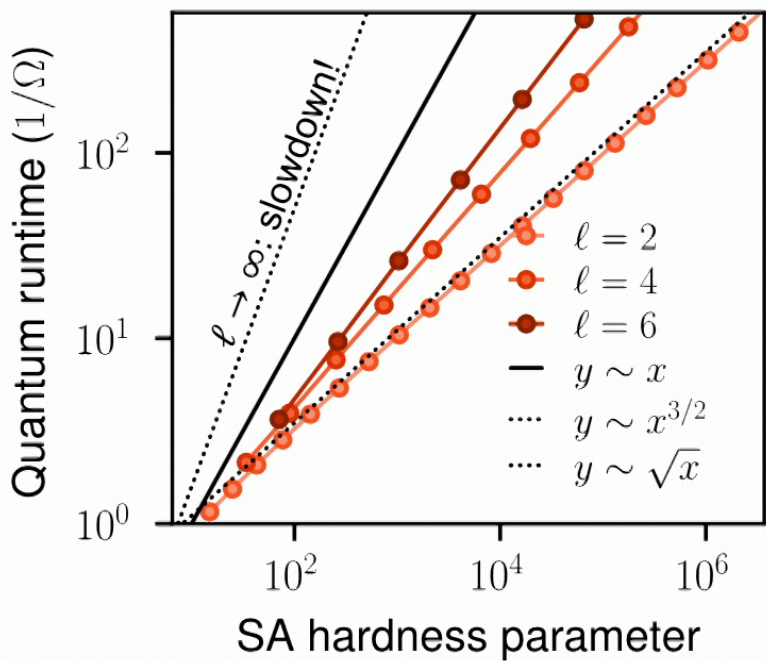


Nodes: independent sets

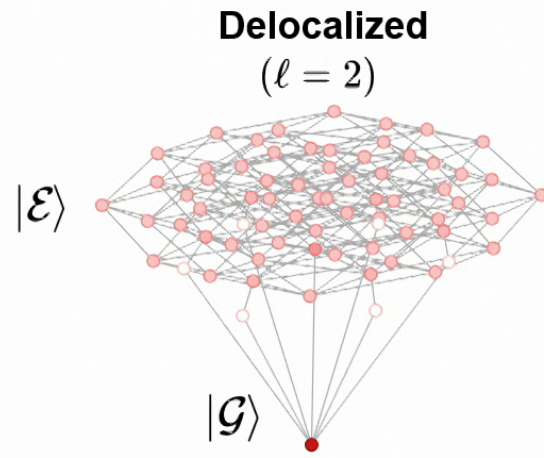
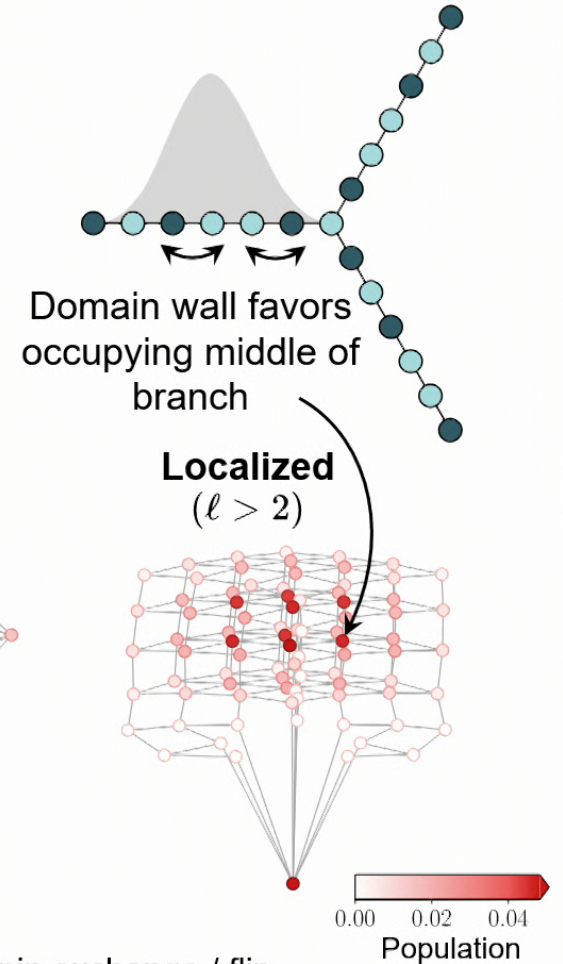
Edges: states connected by spin exchange / flip

Instance-by-instance performance variation

Localization at large Hamming distance
 $\approx b\ell/2$ causes quantum slowdown!



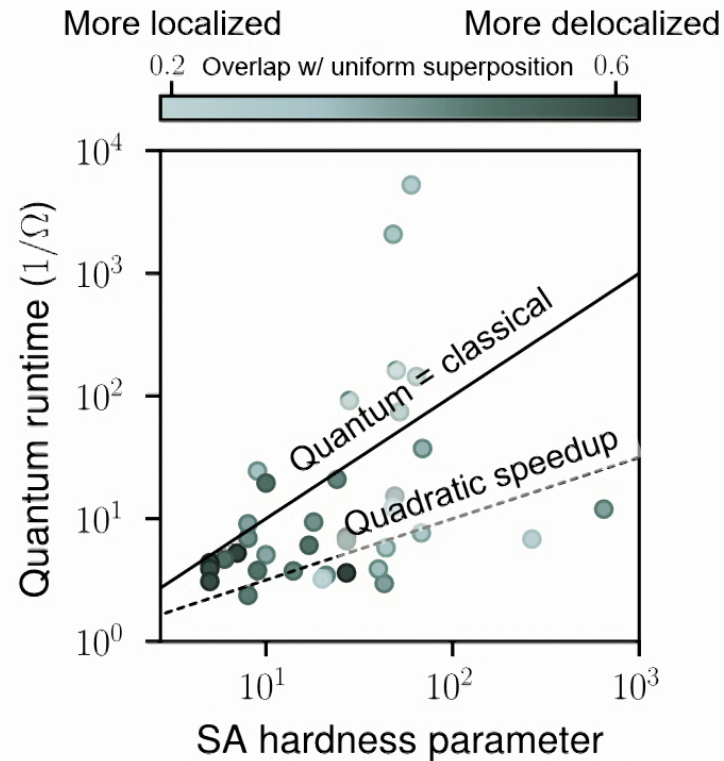
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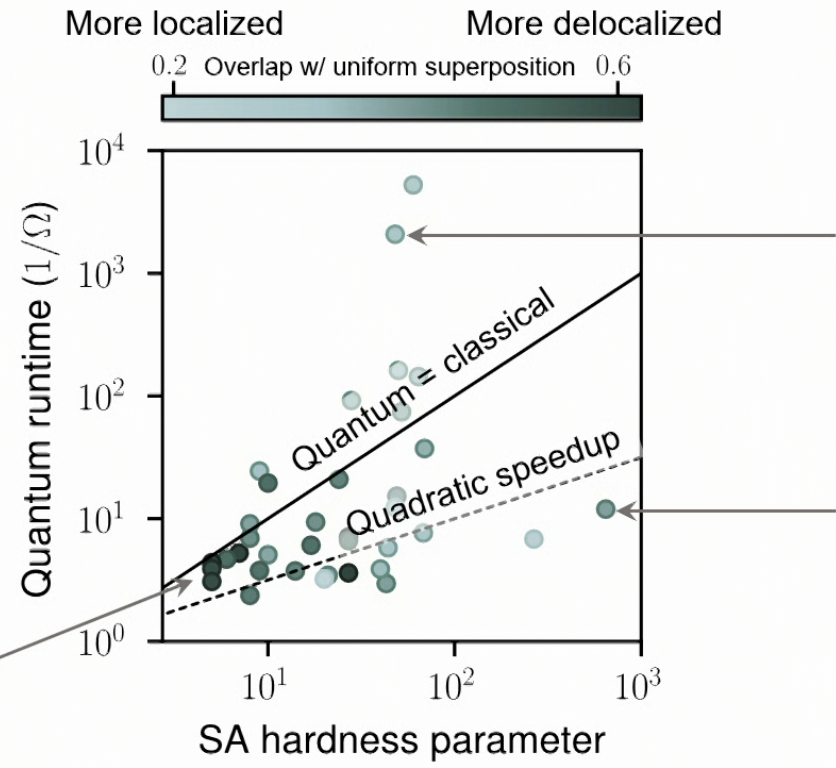
Understanding experimental observations

Study unit-disk graph instances with large SA hardness parameter



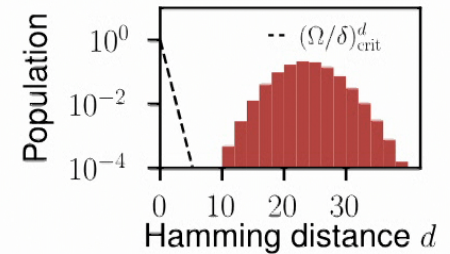
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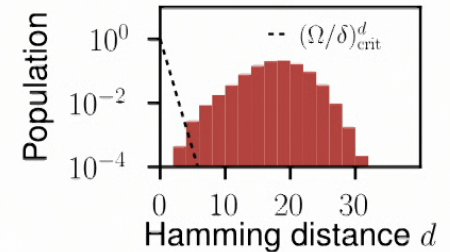


More delocalized:
quadratic speedup

Hard instance for quantum:
large Hamming distance



Easy instance for quantum:
small Hamming distance



We've seen that quantum performance is highly dependent on the problem instance.

Is there a way to generically guarantee a quadratic speedup?

Engineering a quadratic speedup

We want to force $|\mathcal{G}\rangle, |\mathcal{E}\rangle$ to be uniform superpositions.

Add a strong “delocalizing Hamiltonian” λH_ℓ whose ground states are uniform superpositions

$$|\bar{b}\rangle = \frac{1}{\sqrt{D_b}} \sum_{i \text{ size } b} |i\rangle \quad b \in \{1, \dots, |\text{MIS}|\}$$

Inspiration from single-particle quantum mechanics: *kinetic energy* $\sim -\nabla^2$ promotes delocalization



Delocalizing Hamiltonian:
A Laplacian in spin configuration space

Engineering a quadratic speedup

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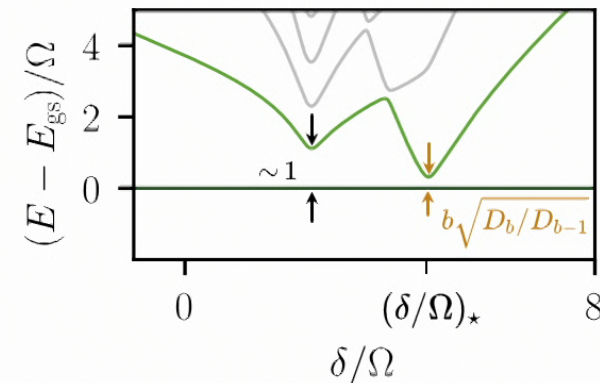
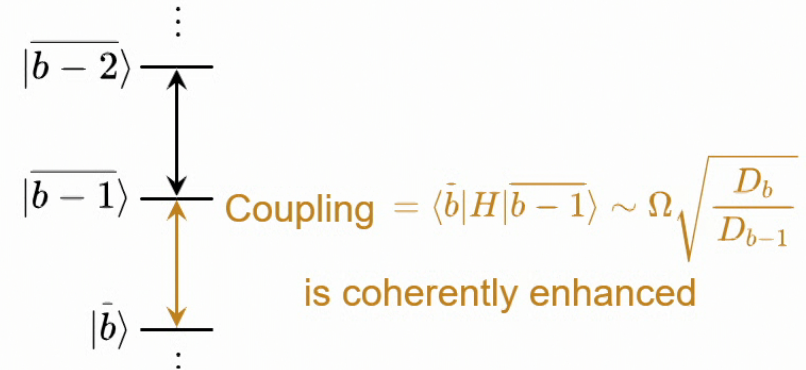


Delocalizing Hamiltonian: *A Laplacian* in spin configuration space

$$H_\ell = -H_{\text{spin-exchange}} + \sum_{(u,v) \in E} n_u(1-n_v) \prod_{(y,v) \in E, y \neq u} (1-n_y)$$

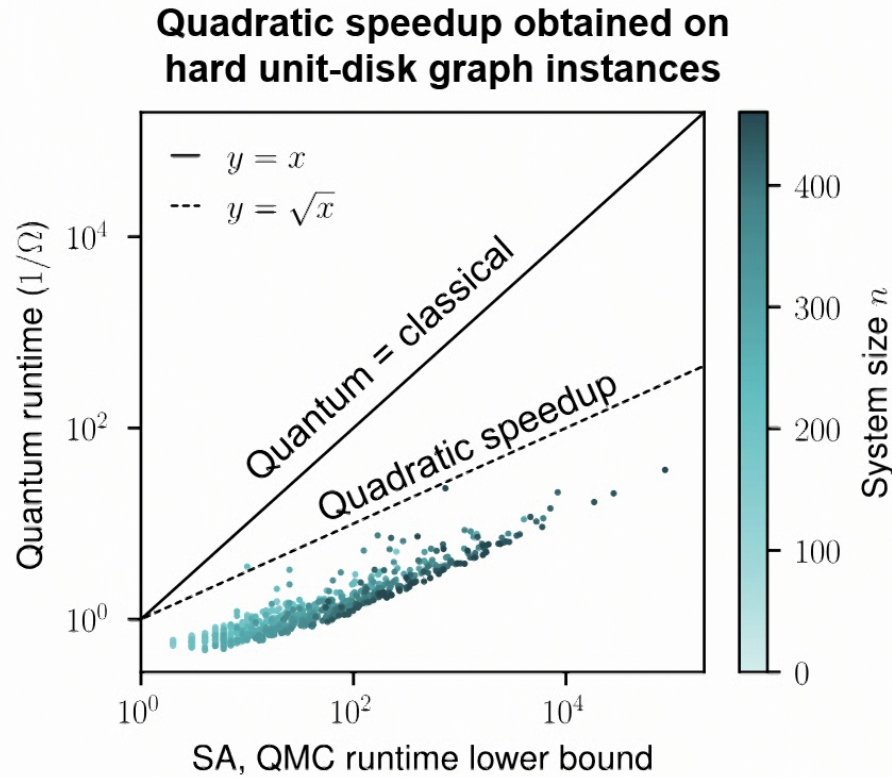
Local terms, no sign problem

λ large: dynamics restricted to H_ℓ ground states

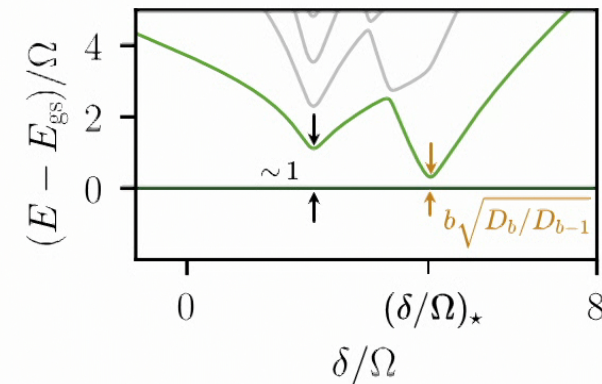
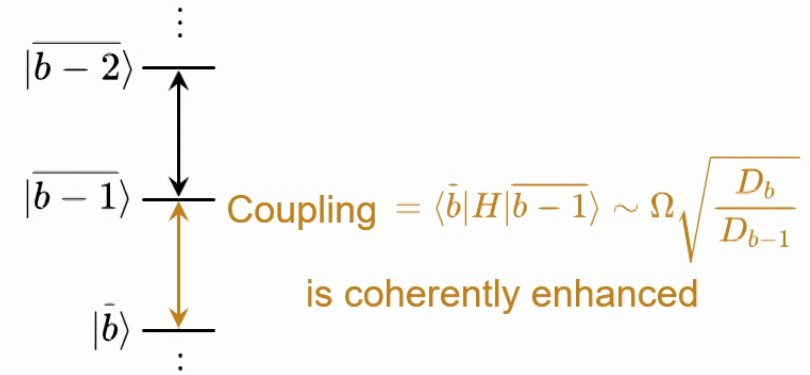


Analytic argument that gap goes as smallest coupling – quadratic speedup!

Engineering a quadratic speedup



λ large: dynamics restricted to H_ℓ ground states



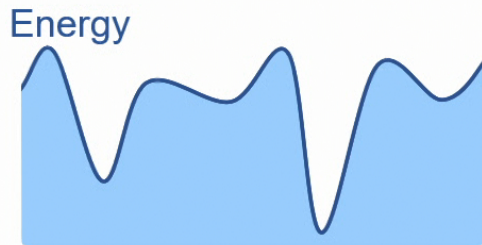
Analytic argument that gap goes as smallest coupling – quadratic speedup!

Summary

- Modification of the QAA provides quadratic speedup over SA
- Framework to understand instance-by-instance performance of the standard QAA
 - Key factor: localization/delocalization of eigenstates at level crossing

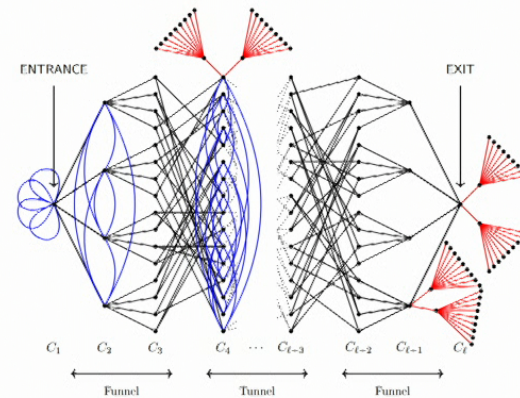
Outlook

Rugged energy landscapes

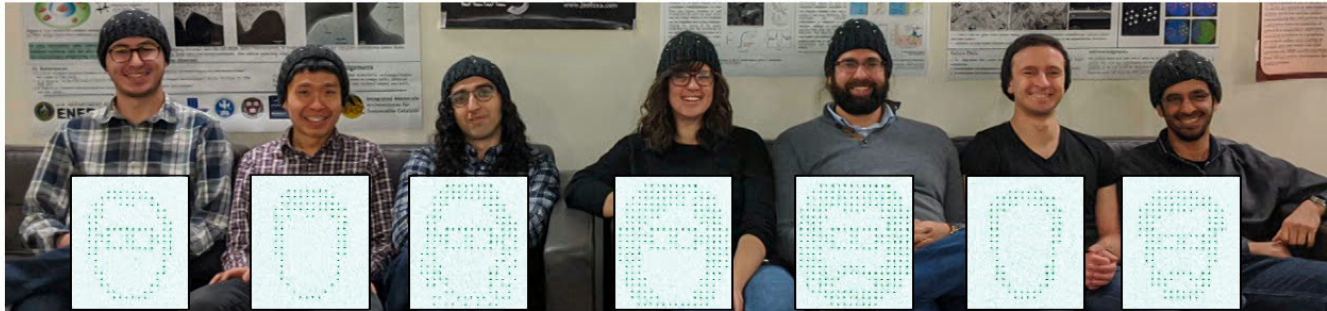


Provably hard to find solution for *all* local classical algorithms
(Overlap Gap Property, Gamarnik 2021)

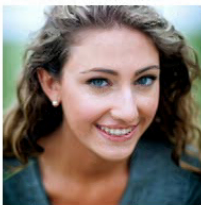
Superpolynomial speedup with unphysical (oracle) Hamiltonians
[Hastings 2020, Gilyen + Vazirani 2020]



Acknowledgements



Harry Levine Tout Wang Sepehr Ebadi Giulia Semeghini Alex Keesling Dolev Bluvstein Ahmed Omran



Beatrice Nash



Leo Zhou



Boaz Barak



Edward Farhi



Hannes Pichler



Rhine Samajdar



Xun Gao



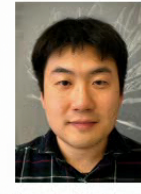
Sambuddha Chattopadhyay



Subir Sachdev



Aram Harrow



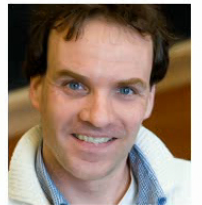
Soonwon Choi



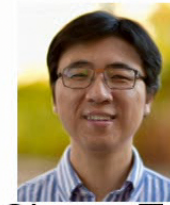
M. Lukin



V. Vuletić



M. Greiner



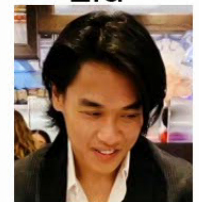
Sheng-Tao Wang



Jin-Guo Liu



Roger Luo



Mao Lin