

Title: QFT2 - Quantum Electrodynamics - Afternoon Lecture

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Collection: Special Topics in Physics - QFT2: Quantum Electrodynamics (Cliff Burgess)

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Abstract: This course uses quantum electrodynamics (QED) as a vehicle for covering several more advanced topics within quantum field theory, and so is aimed at graduate students that already have had an introductory course on quantum field theory. Among the topics hoped to be covered are: gauge invariance for massless spin-1 particles from special relativity and quantum mechanics; Ward identities; photon scattering and loops; UV and IR divergences and why they are handled differently; effective theories and the renormalization group; anomalies.

Spin-half fields + Diracology:

(A, B)

$(\frac{1}{2}, 0)$

$(0, \frac{1}{2})$

(left handed)

right-handed

$$A_i = \frac{1}{2} \sigma_i$$

$$B_i = \frac{1}{2} \sigma_i$$

$$e^{i0 \cdot A_i} = D(0)$$

$$J_i = \frac{1}{2} (A_i + B_i) \quad K_i = \frac{1}{2} (A_i - B_i)$$

electron spin- $\frac{1}{2}$ $j = \frac{1}{2}$
charge $-e$
massive
distinct from
antielelectron

$$c_{p\sigma} \neq \bar{c}_{p\sigma}$$

$$\sigma = \pm \frac{1}{2} \quad J_3$$

$$J_i = (A_i + B_i) \quad K_i = (A_i - B_i)$$

choice: $(\frac{1}{2}, 0) \oplus (0, \frac{1}{2})$

$$\Psi = \begin{pmatrix} x \\ x \\ x \\ y \end{pmatrix} \begin{cases} (\frac{1}{2}, 0) \\ (0, \frac{1}{2}) \end{cases}$$

$$U(\lambda) \Psi U^{-1}(\lambda) = \begin{pmatrix} M_L & 0 \\ 0 & M_R \end{pmatrix} \Psi_{(\lambda)}$$

$$A_i = \begin{pmatrix} \frac{1}{2}g & 0 \\ 0 & 0 \end{pmatrix} \quad B_i = \begin{pmatrix} 0 & 0 \\ 0 & \frac{1}{2}g \end{pmatrix}$$

$$J_i = \begin{pmatrix} \frac{1}{2}g & 0 \\ 0 & \frac{1}{2}g \end{pmatrix}$$

$$K_i = \begin{pmatrix} \frac{1}{2}g & 0 \\ 0 & \frac{1}{2}g \end{pmatrix}$$

$$J_i = \frac{1}{2}(A_i + B_i) \quad K_i = \frac{1}{2}(A_i - B_i)$$

$$\sigma = \pm \frac{1}{2} \quad J_3$$

$$J_3 |0, \sigma\rangle = \sigma |0, \sigma\rangle$$

Fiel.

$$(\frac{1}{2}, 0) \oplus (0, \frac{1}{2})$$

$$\Psi = \begin{pmatrix} x \\ x \\ x \\ x \end{pmatrix} \begin{matrix} \left\{ \begin{matrix} (\frac{1}{2}, 0) \\ (\frac{1}{2}, 0) \end{matrix} \right\} \\ \left\{ \begin{matrix} (0, \frac{1}{2}) \\ (0, \frac{1}{2}) \end{matrix} \right\} \end{matrix}$$

$$U(\lambda) \Psi(\lambda) U^{-1}(\lambda) = \begin{pmatrix} M_L & 0 \\ 0 & M_R \end{pmatrix} \Psi(\lambda)$$

$$\begin{pmatrix} \frac{1}{2}g & 0 \\ 0 & 0 \\ 0 & 0 \\ 0 & -\frac{1}{2}g \end{pmatrix}$$

$$J_i = \begin{pmatrix} \frac{1}{2}g & 0 \\ 0 & -\frac{1}{2}g \end{pmatrix}$$

$$K_i = \begin{pmatrix} -\frac{1}{2}g & 0 \\ 0 & \frac{1}{2}g \end{pmatrix}$$

$$J_i = \frac{1}{2}(A_i + B_i) \quad K_i = \frac{1}{2}(A_i - B_i)$$

$$\sigma = \pm 1/2 \quad J_3$$

$$J_3 |0, \sigma\rangle = \sigma |0, \sigma\rangle$$

Field choice. $(1/2, 0) \oplus (0, 1/2)$

$$\Psi = \begin{pmatrix} \chi \\ \bar{\chi} \end{pmatrix} \begin{cases} (1/2, 0) \\ (0, 1/2) \end{cases}$$

$$\sigma_1 = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}$$

$$\sigma_2 = \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix}$$

$$\sigma_3 = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}$$

$$U(\Lambda) \Psi(x) U^{-1}(\Lambda) = \begin{pmatrix} M_L & 0 \\ 0 & M_R \end{pmatrix} \Psi(\Lambda x)$$

$$A_i = \begin{pmatrix} i g & 0 \\ 0 & 0 \end{pmatrix}$$

$$B_i = \begin{pmatrix} 0 & 0 \\ 0 & -i g \end{pmatrix}$$

$$J_i = \begin{pmatrix} i g & 0 \\ 0 & -i g \end{pmatrix}$$

$$K_i = \begin{pmatrix} i g & 0 \\ 0 & i g \end{pmatrix}$$

$$\Psi(x) = \sum_{\sigma=\pm 1/2} \int \frac{d^3p}{(2\pi)^3 2E_p} \left[c_{p\sigma} e^{ipx} u(p, \sigma) + \bar{c}_{p\sigma}^* e^{-ipx} v(p, \sigma) \right]$$

We build $u(p, \sigma)$ by boosting $u(0, \sigma)$
and little group's action is a constraint of $u(0, \sigma)$

$$J_3 u(0, \sigma) = \sigma u(0, \sigma) \quad u(0, \sigma=+1/2) = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 \\ 0 \\ 0 \\ 0 \end{pmatrix} \quad u(0, \sigma=-1/2) = \frac{1}{\sqrt{2}} \begin{pmatrix} 0 \\ 1 \\ 0 \\ 0 \end{pmatrix}$$

and little group's action is a constraint of $u(0, \sigma)$

$$J_3 u(0, \sigma) = \sigma u(0, \sigma) \quad u(0, \sigma = +\frac{1}{2}) = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 \\ 0 \\ 0 \end{pmatrix} \quad u(0, \sigma = -\frac{1}{2}) = \frac{1}{\sqrt{2}} \begin{pmatrix} 0 \\ 1 \\ 0 \end{pmatrix}$$

$$u(0, \sigma \pm \frac{1}{2}) = \beta u(0, \sigma)$$

$$\beta = \begin{pmatrix} 0 & I \\ I & 0 \end{pmatrix}$$

$$\beta = \begin{pmatrix} 0 & I \\ I & 0 \end{pmatrix}$$

In general frame $D(L(\vec{p}))$

$$u(p, \sigma) = \frac{1}{\sqrt{2}} \begin{pmatrix} \sqrt{\epsilon_p + m} - \hat{\vec{p}} \cdot \vec{\sigma} \sqrt{\epsilon_p - m} & 0 \\ 0 & \sqrt{\epsilon_p + m} + \hat{\vec{p}} \cdot \vec{\sigma} \sqrt{\epsilon_p - m} \end{pmatrix} u(0, \sigma)$$

$$u(0, \sigma) = \beta u(0, \sigma)$$

$$\beta = \begin{pmatrix} 0 & I \\ I & 0 \end{pmatrix}$$

$$D(0) \beta u(0, \sigma) = u(0, \sigma)$$

$$D(L) u(0, \sigma) = u(p, \sigma) = \overbrace{D(L) \beta}^{p^\mu} \overbrace{D^{-1}(L)}^{u(p, \sigma)} D(L) u(0, \sigma)$$

In general frame $D(L(p))$

$$u(p, \sigma) = \frac{1}{\sqrt{2}} \begin{pmatrix} \sqrt{\epsilon_p + m} - \hat{p} \cdot \vec{\sigma} \sqrt{\epsilon_p - m} & 0 \\ 0 & \sqrt{\epsilon_p + m} + \hat{p} \cdot \vec{\sigma} \sqrt{\epsilon_p - m} \end{pmatrix} u(0, \sigma)$$

$$u(0, \sigma) = \beta u(0, \sigma)$$

$$\begin{pmatrix} 0 & I \\ I & 0 \end{pmatrix}$$

$$D(L) \beta u(0, \sigma) = u(0, \sigma)$$

$$D(L) u(0, \sigma) = u(p, \sigma) = \overbrace{D(L) \beta}^{p^\mu} \overbrace{D^{-1}(L)}^{u(p, \sigma)} D(L) u(0, \sigma) = \left(\frac{-i p^\mu \gamma_\mu}{m} \right) u(p, \sigma)$$

$$D(L(p))$$

$\sqrt{\epsilon_p + m} - \hat{\vec{p}} \cdot \vec{\sigma} \sqrt{\epsilon_p - m}$	0
0	$\sqrt{\epsilon_p + m} + \hat{\vec{p}} \cdot \vec{\sigma} \sqrt{\epsilon_p - m}$

$$u(0, \sigma)$$

$$-i\gamma_0 = \beta \quad \gamma_0 = i\beta = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \quad u = \left(-i \frac{p^\mu \gamma_\mu}{m} \right) u$$

$$\gamma_k = \begin{pmatrix} 0 & -i\sigma_k \\ i\sigma_k & 0 \end{pmatrix}$$

$$(i p^\mu \gamma_\mu + m) u = 0 \quad \text{Dirac eq}$$

$$(\gamma^\mu \partial_\mu + m) \bar{\Psi}(x) = 0$$

$$v(p, \sigma) = (-)^{2s} (-)^{j-\sigma} u(p, -\sigma)$$

$$v(0, \sigma = +1/2) = \frac{1}{\sqrt{2}} \begin{pmatrix} 0 \\ 1 \\ 0 \\ -1 \end{pmatrix}$$

$$v(0, \sigma = -1/2) = \frac{1}{\sqrt{2}} \begin{pmatrix} 0 \\ 0 \\ 1 \\ 0 \end{pmatrix}$$

$$\gamma_5 = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} \quad [\gamma_5, D(\Lambda)] = 0$$

$$\beta v(0, \sigma) = -v(0, \sigma)$$

$$\beta = \begin{pmatrix} 0 & I \\ I & 0 \end{pmatrix}$$

$$D(L) \beta u(0, \sigma) = u(0, \sigma)$$

$$D(L) u(0, \sigma) = u(p, \sigma) = \overbrace{D(L) \beta}^{p^\mu} \overbrace{D^{-1}(L)}^{u(p, \sigma)} D(L) u(0, \sigma) \\ = \left(\frac{-i p^\mu \gamma_\mu}{m} \right) u(p, \sigma)$$

In general frame $D(L(p))$

$$u(p, \sigma) = \frac{1}{\sqrt{2}} \begin{pmatrix} \sqrt{\epsilon_p + m} - \hat{\vec{p}} \cdot \vec{\sigma} \sqrt{\epsilon_p - m} & 0 \\ 0 & \sqrt{\epsilon_p + m} + \hat{\vec{p}} \cdot \vec{\sigma} \sqrt{\epsilon_p - m} \end{pmatrix} u(0, \sigma)$$

$$\psi(p, \sigma) = D(L) \psi(0, \sigma)$$

$$D(\Lambda) \gamma_\mu D^{-1}(\Lambda) = \Lambda^\nu_\mu \gamma_\nu$$

$$(\frac{1}{2}, 0) \otimes (0, \frac{1}{2}) = (\frac{1}{2}, \frac{1}{2})$$

16 useful basis matrices: $I, \gamma_5, \gamma_\mu, \gamma_\mu \gamma_5, \gamma_{\mu\nu} = \gamma_\mu \gamma_\nu = \frac{1}{2} [\gamma_\mu, \gamma_\nu]$

$$(M, N) = \text{Tr}(MN)$$

$$\left(\frac{-i p^\mu \gamma_\mu}{m} \right) \psi = -\psi$$

$$\{M\} = \begin{pmatrix} 4 \times 4 \end{pmatrix}$$

$$(-i p^\mu \gamma_{\mu+m}) \psi = 0$$

$$\{\gamma_\mu, \gamma_\nu\} = 2\eta_{\mu\nu}$$

$$C = c_1 I + c_5 \gamma_5 + c^\mu \gamma_\mu + c_5^\mu \gamma_\mu \gamma_5 + c^{\mu\nu} \gamma_{\mu\nu}$$

$$\text{Tr}(\gamma_5 C) = 4c_5$$

$$U(p, \sigma) = D(L) U(o, \sigma)$$

$$D(\Lambda) \gamma_\mu D^{-1}(\Lambda) = \Lambda^\nu_\mu \gamma_\nu$$

$$(\frac{1}{2}, 0) \otimes (0, \frac{1}{2}) = (\frac{1}{2}, \frac{1}{2})$$

16 useful basis matrices: $I, \gamma_5, \gamma_\mu, \gamma_\mu \gamma_5, \gamma_{\mu\nu} = \gamma_{\mu\nu} = \frac{1}{2} [\gamma_\mu, \gamma_\nu]$

$$(M, N) = \text{Tr}(MN)$$

$$\left(\frac{-i p^\mu \gamma_\mu}{m} \right) U = -U$$

$$(-i p^\mu \gamma_{\mu+m}) U = 0$$

$$\{M\} = \begin{pmatrix} 4 \times 4 \end{pmatrix}$$

$$\{\gamma_\mu, \gamma_\nu\} = 2\eta_{\mu\nu}$$

$$\psi \rightarrow D(\Lambda) \psi$$

$$\psi^\dagger \rightarrow \psi^\dagger D^\dagger(\Lambda) \neq \psi^\dagger D^{-1}(\Lambda)$$

$$C = c_1 I + c_5 \gamma_5 + c^\mu \gamma_\mu + c_5^\mu \gamma_\mu \gamma_5 + c^{\mu\nu} \gamma_{\mu\nu}$$

$$\text{Tr}(\gamma_5 C) = 4c_5$$

$$\sigma_k^* = \epsilon \sigma_k \epsilon$$

$$\epsilon = \begin{pmatrix} 0 & 1 \\ -1 & 0 \end{pmatrix} = i\sigma_2$$

$$\bar{\psi} := \psi^\dagger \beta$$

$$D^\dagger(\Lambda) \beta = \beta D^{-1}(\Lambda)$$

$$\{\gamma_\mu, \gamma_\nu\} = 2\eta_{\mu\nu}$$

$$\begin{aligned}\psi &\rightarrow D(\Lambda)\psi \\ \psi^\dagger &\rightarrow \psi^\dagger D^\dagger(\Lambda) \neq \psi^\dagger D^{-1}(\Lambda)\end{aligned}$$

$$C = c_1 I + c_5 \gamma_5 + c^\mu \gamma_\mu + c_5^\mu \gamma_\mu \gamma_5 + c^{\mu\nu} \gamma_{\mu\nu}$$

$$\text{Tr}(\gamma_5 C) = 4c_5$$

$$\sigma_k^* = \epsilon \sigma_k \epsilon$$

$$\epsilon = \begin{pmatrix} 0 & 1 \\ -1 & 0 \end{pmatrix} = i\sigma_2$$

$$\begin{aligned}\bar{\Psi} &:= \psi^\dagger \beta \\ &= i\psi^\dagger \gamma^0 = -i\psi^\dagger \gamma_0\end{aligned}$$

$$\begin{aligned}D^\dagger(\Lambda)\beta &= \beta D^{-1}(\Lambda) \\ \bar{\Psi} &\rightarrow \bar{\Psi} D'(\Lambda)\end{aligned}$$

Spin-half fields + Diracology:

$$\bar{\psi}\psi \rightarrow \bar{\psi} \gamma^0 \gamma^0 \psi = \bar{\psi}\psi \text{ scalar}$$

$$\bar{\psi} \gamma_5 \psi \quad (\text{pseudo}) \text{ scalar}$$

$$\left\{ \bar{\psi} \gamma_\mu \psi \quad 4\text{-vector} \right.$$

$$\left\{ \bar{\psi} \gamma_\mu \gamma_5 \psi \quad \parallel \text{ (axial)} \right.$$

$$\bar{\psi} \gamma_{\mu\nu} \psi \quad \text{skew tensor}$$

$$\begin{pmatrix} \psi_L \\ \psi_R \end{pmatrix} \quad \gamma_5 \psi = \psi \quad (\text{LH})$$

$$\gamma_5 \psi = -\psi \quad (\text{RH})$$

$$\psi_L = \frac{1}{2}(1 + \gamma_5)\psi$$

$$\psi_R = \frac{1}{2}(1 - \gamma_5)\psi$$

$$(\frac{1}{2}, 0) \times (\frac{1}{2}, 0) = (0, 0) + \boxed{(1, 0)} \quad (2.11)$$

Spin-half fields + Diracology:

$$\bar{\psi} \psi \rightarrow \bar{\psi} \gamma^0 \gamma^0 \psi = \bar{\psi} \psi \quad \text{scalar}$$

$$\bar{\psi} \gamma_5 \psi \quad (\text{pseudo}) \text{ scalar}$$

$$\left\{ \bar{\psi} \gamma_\mu \psi \right. \quad 4\text{-vector}$$

$$\bar{\psi} \gamma_5 \gamma_\mu \psi \quad \parallel (\text{axial})$$

$$\bar{\psi} \gamma_{\mu\nu} \psi \quad \text{skew tensor}$$

$$(\gamma_L \gamma_{\mu\nu} \gamma_L) \in \epsilon^{\mu\nu\lambda\rho}$$

$$= 0; \gamma_L \gamma^{\lambda\rho} \gamma_L$$

$$t_{\mu\nu} = 0 \in \epsilon_{\mu\nu\lambda\rho} t^{\lambda\rho}$$

$$(\frac{1}{2}, 0) \times (\frac{1}{2}, 0) = (0, 0) + \boxed{(1, 0)} \quad (a.1)$$

Spin-half fields + Dirac algebra

$$\bar{\psi} \psi \rightarrow \bar{\psi} D^\dagger D \psi = \bar{\psi} \psi$$

$$\bar{\psi} \gamma_5 \psi \quad (\text{pseudo}) \text{ scalar}$$

$$\left\{ \begin{array}{l} \bar{\psi} \gamma_\mu \psi \\ \bar{\psi} \gamma_\mu \gamma_5 \psi \end{array} \right. \quad \begin{array}{l} 4\text{-vector} \\ \parallel \text{ (axial)} \end{array}$$

$$\bar{\psi} \gamma_{\mu\nu} \psi \quad \text{skew tensor}$$

$$(\gamma_\mu \gamma_\nu \gamma_\rho) \in \epsilon^{\mu\nu\rho}$$

$$= 0; \gamma_\mu \gamma^\mu \gamma_\nu \gamma^\nu$$

$$t_{\mu\nu} = 0 \in \epsilon_{\mu\nu\lambda\rho} t^{\lambda\rho}$$

$$\gamma_5 = -i \gamma_0 \gamma_1 \gamma_2 \gamma_3$$

$$\epsilon_{\mu\nu\lambda\rho} \gamma^\mu \gamma^\nu \gamma^\lambda \gamma^\rho$$

$\psi \gamma_{\mu\nu} \psi$ (axial)
 skew tensor

$$\frac{i}{4!} \epsilon_{\mu\nu\rho\sigma} \gamma^\mu \gamma^\nu \gamma^\rho \gamma^\sigma$$

$$\psi = \begin{pmatrix} \psi_L \\ \psi_R \end{pmatrix}$$

$$\psi_L = \gamma_L \psi = \frac{1}{2}(1 + \gamma_5) \psi = \begin{pmatrix} 1 & 0 \\ 0 & 0 \end{pmatrix} \psi = \begin{pmatrix} \psi_L \\ 0 \end{pmatrix}$$

$$\psi_R = \gamma_R \psi = \frac{1}{2}(1 - \gamma_5) \psi = \begin{pmatrix} 0 & 0 \\ 0 & 1 \end{pmatrix} \psi = \begin{pmatrix} 0 \\ \psi_R \end{pmatrix}$$

$(\frac{1}{2}, 0) \leftrightarrow (0, \frac{1}{2})$
 under complex conjugation

"Weyl Spinor"

$\psi_{\mu\nu} = \psi_{\nu\mu}$ (axial)
 $\psi_{\mu\nu}$ skew tensor

$$\frac{i}{4!} \epsilon_{\mu\nu\lambda\rho} \gamma^\mu \gamma^\nu \gamma^\lambda \gamma^\rho$$

$$\psi = \begin{pmatrix} \psi_L \\ \psi_R \end{pmatrix}$$

$$\psi_L = \gamma_L \psi = \frac{1}{2}(1 + \gamma_5) \psi = \begin{pmatrix} 1 & 0 \\ 0 & 0 \end{pmatrix} \psi = \begin{pmatrix} \psi_L \\ 0 \end{pmatrix}$$

$$\psi_R = \gamma_R \psi = \frac{1}{2}(1 - \gamma_5) \psi = \begin{pmatrix} 0 & 0 \\ 0 & 1 \end{pmatrix} \psi = \begin{pmatrix} 0 \\ \psi_R \end{pmatrix}$$

$$(1/2, 0) \leftrightarrow (0, 1/2)$$

under complex conjugation

"Weyl Spinor"

$$\text{If } u_L \rightarrow D_L(\Lambda) u_L$$

then $\bar{u}_L^*$ transforms like u_R

"Majorana" spinor: $\psi = \begin{pmatrix} \psi_L \\ -\epsilon \psi_L^* \end{pmatrix} = \underbrace{\begin{pmatrix} 0 & \epsilon \\ -\epsilon & 0 \end{pmatrix}}_C \begin{pmatrix} \psi_L \\ -\epsilon \psi_L^* \end{pmatrix}^*$

$\psi^* = C \psi$
is Lorentz
inv.

$C^2 = 1$

$$\epsilon = \begin{pmatrix} 0 & 1 \\ -1 & 0 \end{pmatrix} = i\sigma_2$$

$\psi \gamma_{\mu} \psi$ (axial)
 $\psi \gamma_{\mu\nu} \psi$ skew tensor

$$\frac{i}{4!} \epsilon_{\mu\nu\rho\sigma} \gamma^{\mu} \gamma^{\nu} \gamma^{\rho} \gamma^{\sigma}$$

$$\psi = \begin{pmatrix} \psi_L \\ \psi_R \end{pmatrix} \quad \begin{matrix} u^+ \\ u^- \end{matrix} \quad \begin{aligned} \psi_L &= \gamma_L \psi = \frac{1}{2}(1 + \gamma_5) \psi = \begin{pmatrix} 1 & 0 \\ 0 & 0 \end{pmatrix} \psi = \begin{pmatrix} \psi_L \\ 0 \end{pmatrix} \\ \psi_R &= \gamma_R \psi = \frac{1}{2}(1 - \gamma_5) \psi = \begin{pmatrix} 0 & 0 \\ 0 & 1 \end{pmatrix} \psi = \begin{pmatrix} 0 \\ \psi_R \end{pmatrix} \end{aligned}$$

$(\frac{1}{2}, 0) \leftrightarrow (0, \frac{1}{2})$
 under complex conjugation

"Weyl Spinor"

If $u_L \rightarrow D_L(\Lambda) u_L$
 then $\bar{u}_L^*$ transforms like u_R

$$\{\gamma_5, \gamma_\mu\} = 0$$

$$(-1)$$

$$\gamma^0(0)$$

$$p_0(0,0,1) = \gamma^0(0)$$

If ψ_1, ψ_2 are Majorana spinors

$$(\bar{\psi}_1 M \psi_2)^* = \chi (\bar{\psi}_1 M \psi_2)$$

$$(\bar{\psi}_1 M \psi_2) = \lambda (\bar{\psi}_2 M \psi_1)$$

$M =$	γ_5	γ_μ	$\gamma_\mu \gamma_5$	$\gamma_{\mu\nu}$
χ	+	-	+	-
λ	+	+	-	+



$$\{\gamma_5, \gamma_\mu\} = 0$$

If ψ_1, ψ_2 are Majorana spinors

$$(\bar{\psi}_1 M \psi_2)^* = \chi (\bar{\psi}_1 M \psi_2)$$

$$(\bar{\psi}_1 M \psi_2) = \lambda (\bar{\psi}_2 M \psi_1)$$

$$\mathcal{L} \Rightarrow c i \bar{\psi} \gamma_5 \psi \quad A^\mu (\bar{\psi}_1 \gamma_\mu \psi_2) = 0$$

$M =$	γ_5	γ_μ	$\gamma_\mu \gamma_5$	$\gamma_{\mu\nu}$
χ	+	-	+	-
χ'	+	+	-	+