

Title: Quasinormal-mode filters and their impacts on future black-hole spectroscopy

Speakers: Sizheng Ma

Series: Strong Gravity

Date: September 22, 2022 - 1:00 PM

URL: <https://pirsa.org/22090090>

Abstract: In this talk, I present a novel framework to do black-hole spectroscopy. This approach is based on a new technique so-called "quasinormal-mode filter", which can be classified into two subsets: rational filter and full filter. On the theoretical level, I explain how to use the rational filter to understand the ringdown of numerical-relativity waveforms. On the observational level, I introduce a way to incorporate the filter into Bayesian analysis. The new Bayesian framework not only allows us to analyze the ringdown of a real gravitational-wave event without Markov chain Monte Carlo, but also yields a natural estimate of the ringdown start time. By applying our method to GW150914, we find strong and self-consistent evidence for the first overtone from multiple perspectives. On the other hand, the relationship between the full filter and the metric reconstruction is discussed. Its connection to a numerical-relativity technique "Cauchy-characteristic Matching" is provided. In the final part of the talk, I also briefly present our recent progress on fully relativistic 3D Cauchy-characteristic Matching.

Zoom link: <https://pitp.zoom.us/j/95593408817?pwd=U2RXZDV2WUtwNDlaRDVBVTJjTEdBQT09>



The QNM filter: a novel framework in black-hole spectroscopy

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In collaboration w/: Ling Sun (OzGrav, ANU), Yanbei Chen (Caltech)

09/22/2022 @Perimeter Institute

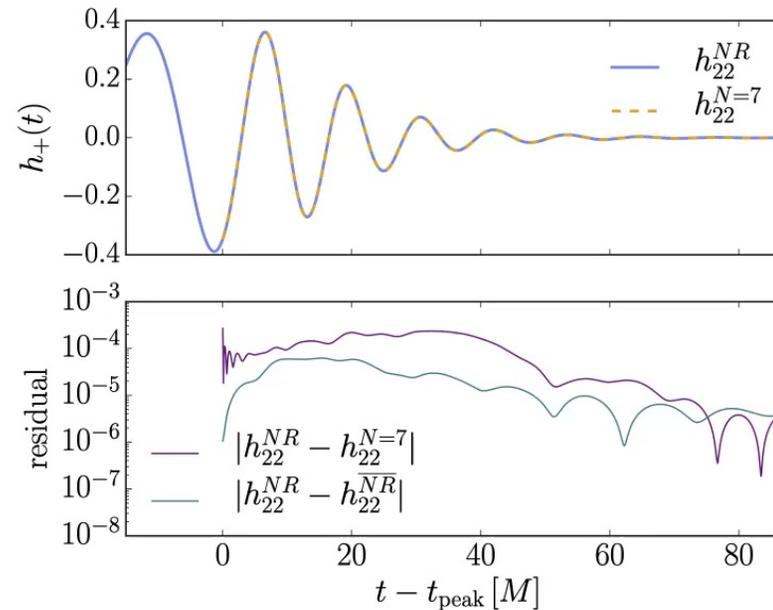
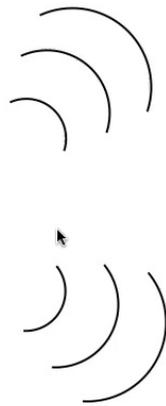
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BH spectroscopy and ringdown modeling



Black Hole



- Overfitting?

$$\omega_{lmn}$$

$$h_{lm} = \sum_n A_n e^{-i\omega_{lmn}t}$$

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Giesler et al., arXiv: 1903.08284

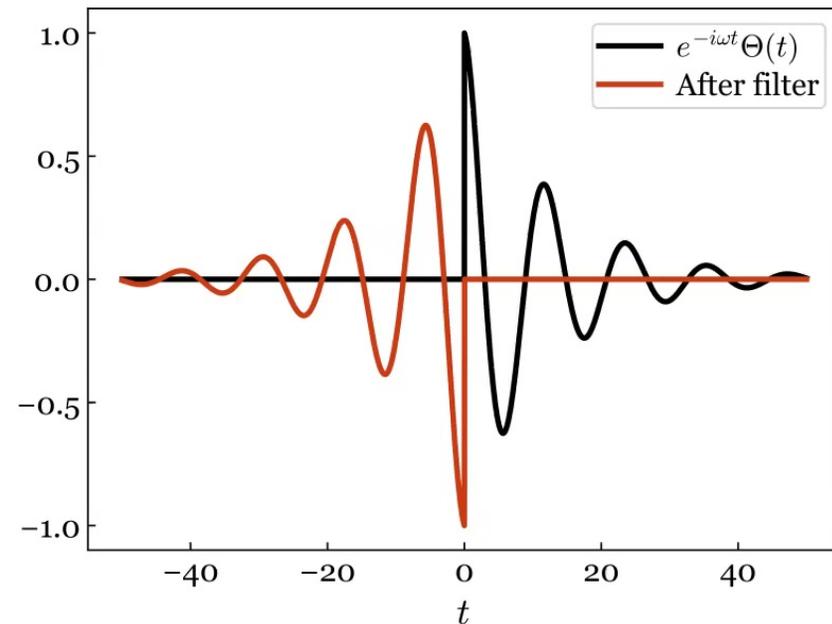


A new approach — a frequency-domain

$$h(t) = e^{-i\omega_{lmn}(t-t_0)} \Theta(t - t_0)$$

Time domain operator: $\left(\frac{d}{dt} + i\omega_{lmn}\right) h(t) = 0$

$$\begin{aligned} &\downarrow \\ &\omega - \omega_{lmn} \\ &\downarrow \\ &\frac{\omega - \omega_{lmn}}{\omega - \omega_{lmn}^*} \end{aligned}$$



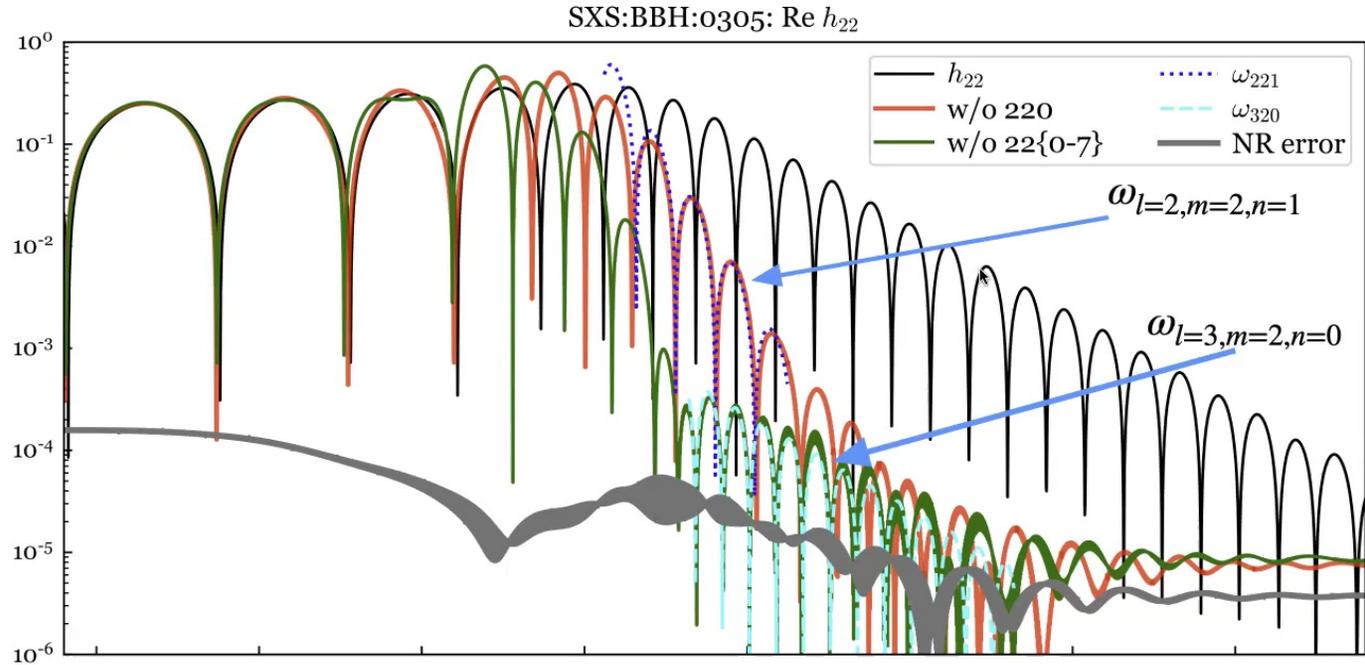
arXiv: 2207.10870

The spherical-spheroidal mixing m



$$\text{Filter} \sim \prod_{lmn} (\omega - \omega_{lmn})$$

GW150914-like
numerical-relativity simulation



arXiv: 2207.10870

Second-order quasinormal mo

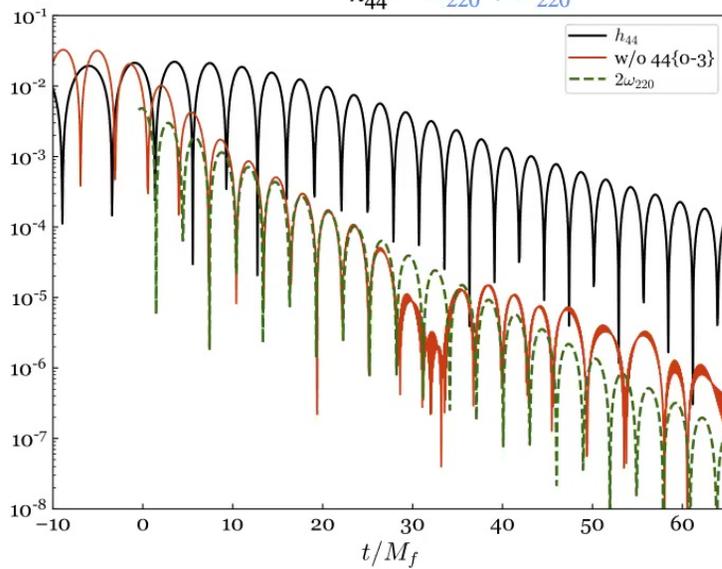


$$h_{lm} = \sum_n A_n e^{-i\omega_{lmn}t} + B e^{-i(\omega_{l_1 m_1 n_1} + \omega_{l_2 m_2 n_2})t}$$

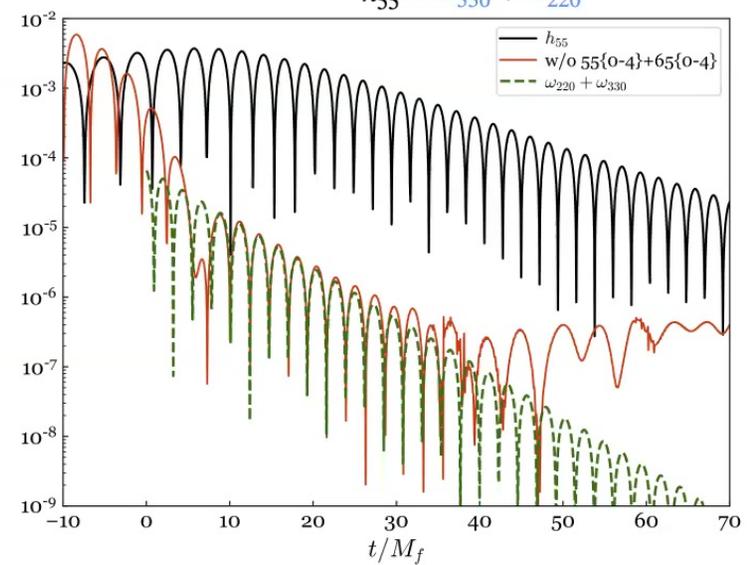
linear effects

second-order effects

$$h_{44} \sim \omega_{220} + \omega_{220}$$



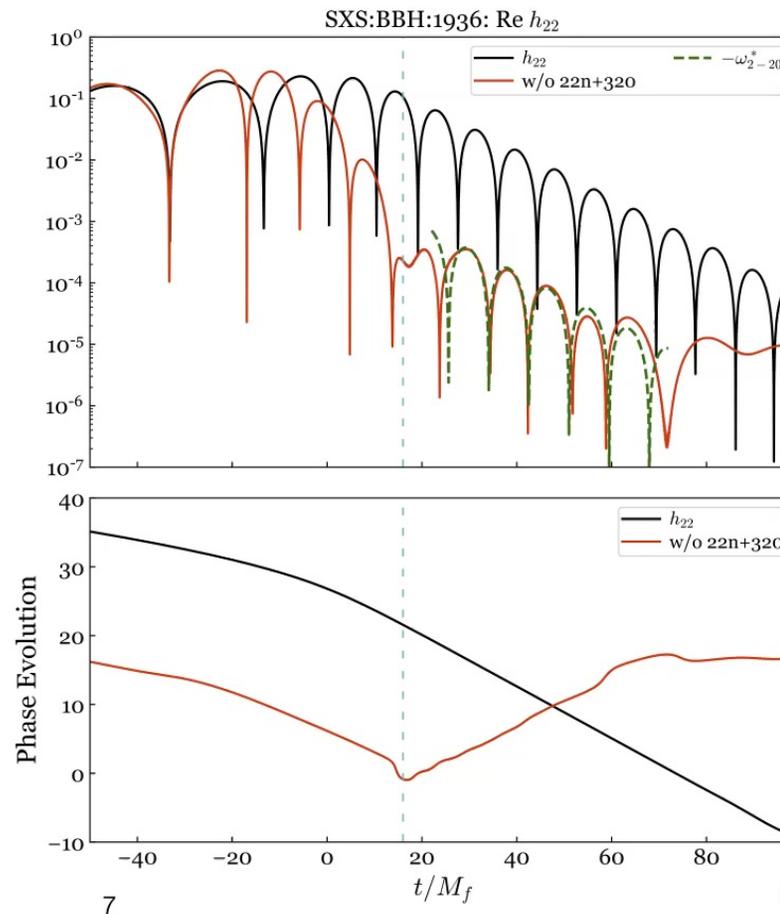
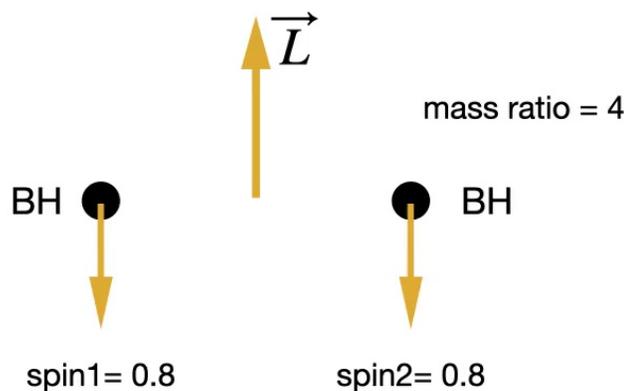
$$h_{55} \sim \omega_{330} + \omega_{220}$$



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arXiv: 2207.10870, 2208.07380, 2208.07374, 2208.07379

Retrograde modes



arXiv: 2207.10870

Takeaway messages



- The goal is to reveal subdominant QNMs by removing the leading one
- The filter is imposed on the entire waveform (avoid spectral leakage)
 1. Its impact on inspiral portion: **a trivial negative time and phase shift** (see 2207.10870 for details)
 2. Its impact on the corresponding QNM: **flipping it** (see 2207.10870 for details)
 3. Its impact on other QNMs: **reducing their amplitudes** (Ma et al. in prep)
- The corresponding QNM can be removed regardless of its amplitude and phase, as well as inclination and polarization angles!

arXiv: 2207.10870

Outline



- Properties of the rational filter in terms of numerical-relativity waveforms
(arXiv: [2207.10870](#))
- Incorporating the rational filter into Bayesian inference (Ma, et al., in prep.)
- The full filter (briefly, arXiv: [2207.10870](#), [2203.03174](#))
- Cauchy-Characteristic Matching (Ma, et al., in prep.)

Incorporating the filter into Bayesian inference



Two additional properties:

- $\mathcal{F}_{lmn}(-f) = [\mathcal{F}_{lmn}(f)]^*$



The filtered signal is still real

- $|\mathcal{F}_{lmn}(f)| = 1$, namely $\mathcal{F}_{lmn}(f) = e^{i\delta_{lmn}(f)}$



The filter has no impact on PSDs!

Incorporating the filter into Bayesian inference



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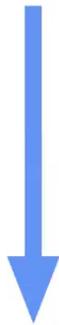
The filter has no impact on PSDs!

Ma, et al., in preparation

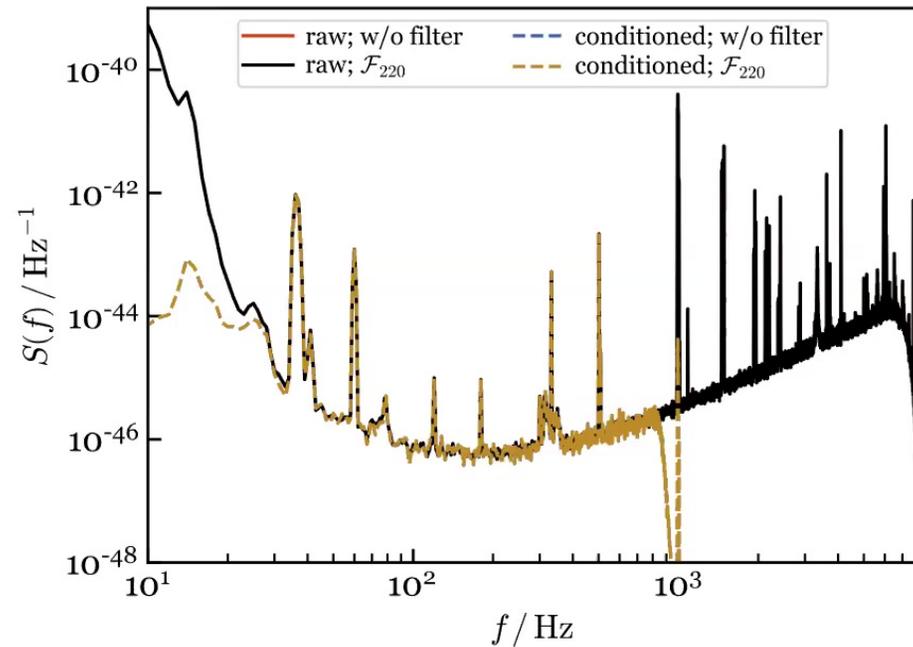
The filter has no impact on PS



$$\tilde{n}_f^F = e^{i\delta_{\text{inn}}(f)} \tilde{n}_f$$



$$E[\tilde{n}_f^F \tilde{n}_{f'}^{F*}] = E[\tilde{n}_f \tilde{n}_{f'}^*] = \frac{1}{2} \delta(f - f') S(f)$$



32 s data of GW150914 Hanford (Welch method)

Isi & Farr (2021), arXiv: 2107.05609 & <https://github.com/maxisi/ringdown>

Ma, et al., in preparation

A new likelihood function

Ma, et al.



Time Domain

Isi et al., arXiv: 2107.05609, 1905.00869

Traditional likelihood function

$$\ln P^{\text{old}}(d_i | A_{lmn}, \phi_{lmn}, M_f, \chi_f, t_0, \text{etc.}) = -\frac{1}{2} \sum_{ij} (d_i - h_i) C_{ij}^{-1} (d_j - h_j)$$

d_i : **GW data** $\in [t_0, t_0 + T]$

h_i : **ringdown template**

C_{ij} : **autocovariance function**

0.2 s

$$h_t = \sum_{lmn} A_{lmn} e^{-(t-t_0)/\tau_{lmn}} \cos [2\pi f_{lmn}(t - t_0) + \phi_{lmn}]$$

New likelihood function

$$\ln P^{\text{new}}(d_i | M_f, \chi_f, t_0) = -\frac{1}{2} \sum_{ij} d_i^F C_{ij}^{-1} d_j^F$$

d_i^F : **filtered GW data** $\in [t_0, t_0 + T]$

C_{ij} : **autocovariance function**



Need only remnant mass and spin



The new likelihood is merely a 2D function

The corresponding QNM can be removed regardless of its amplitude and phase, as well as inclination and polarization angles!

$$\ln P^{\text{old}}(d_i | A_{lmn}, \phi_{lmn}, M_f, \chi_f, t_0, \text{etc.}) \\ = -\frac{1}{2} \sum_{ij} (d_i - h_i) C_{ij}^{-1} (d_j - h_j)$$

versus

$$\ln P^{\text{new}}(d_i | M_f, \chi_f, t_0) \\ = -\frac{1}{2} \sum_{ij} d_i^F C_{ij}^{-1} d_j^F$$



- There is no need to use fancy samplers (e.g. MCMC) to draw samplings
- We can directly plot the likelihood as a function of mass and spin (with t_0 being a hyperparameter)

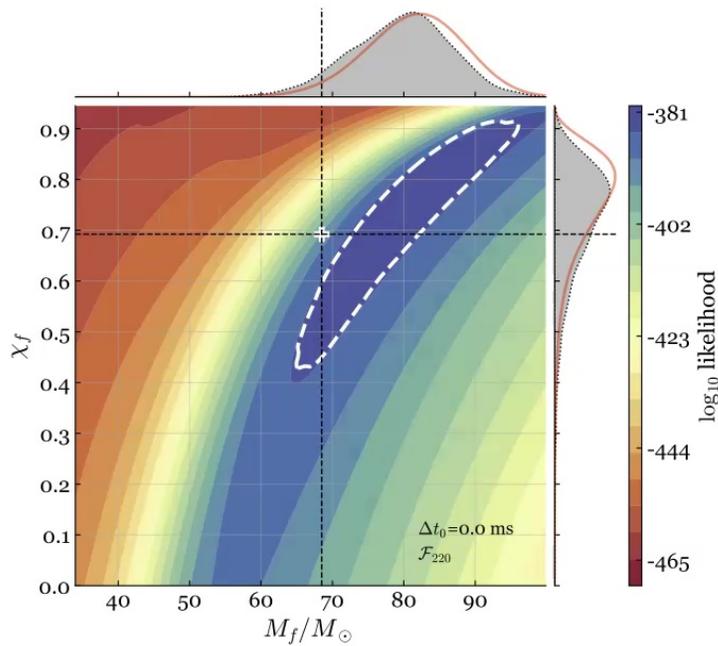
Ma, et al., in preparation

Taking GW150914 as an example



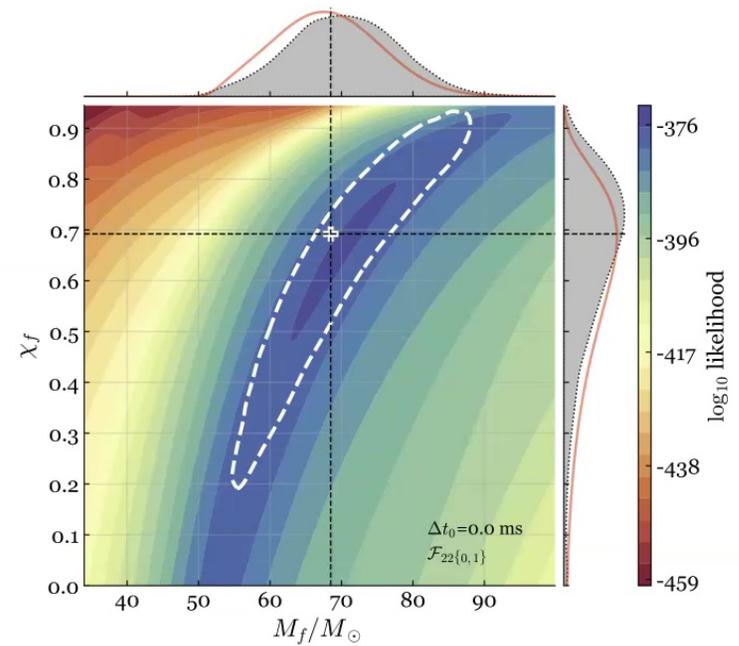
At $t_0 = t_{\text{ref}}$

Fundamental mode only



$t_{\text{ref}} = 1126259462.4083$ GPS at the geocenter
 $T = 0.2$ s

Fundamental mode + the first overtone



- The white dashed contour: the traditional MCMC method, with 90% credible region enclosed
- The plus sign: the IMR result



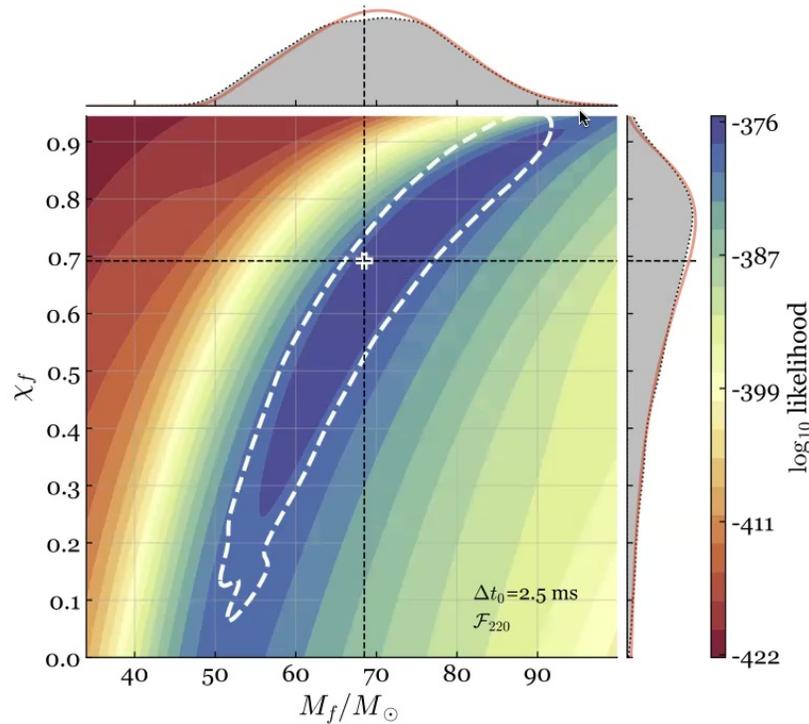
The performance of the new likelihood function

- Fast and efficient: 5.5×10^{-3} s for a single evaluation; 7.8 s for a low-res 2D plot (personal laptop, not fully parallelized); 3 mins for a high-res 2D plot (24-core cluster, not fully parallelized).
- Fully parallelizable: The evaluation of each pair of mass and spin is fully independent.
- The code doesn't slow down after increasing the number of modes because multiple filters can be imposed simultaneously.
- The posteriors are consistent with the traditional MCMC analysis.

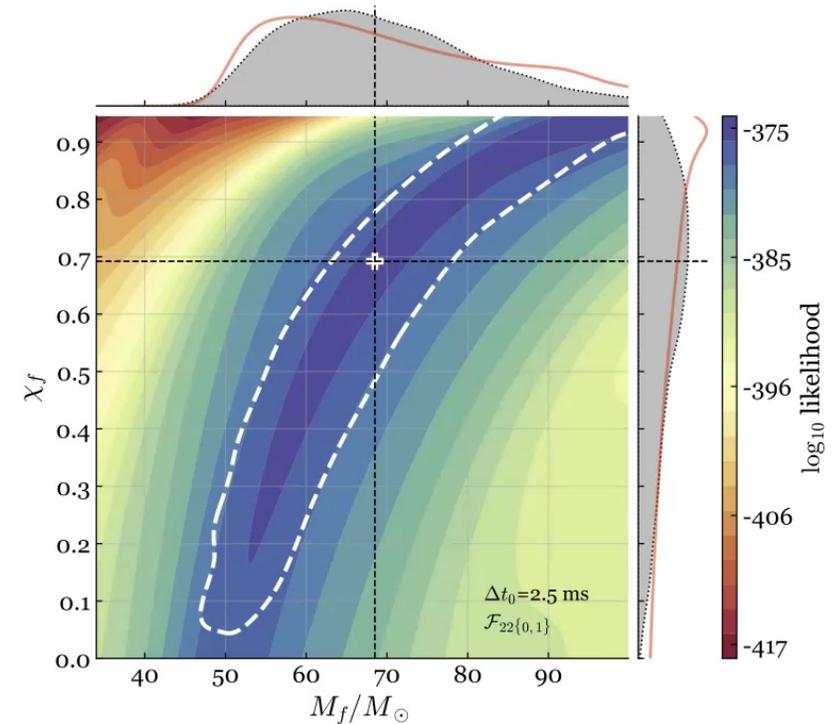
Taking GW150914 as an exam



At $t_0 = t_{\text{ref}} + 2.5 \text{ ms}$ **Fundamental mode only**



Fundamental mode + the first overtone



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Ma, et al., in preparation



The performance of the new likelihood function

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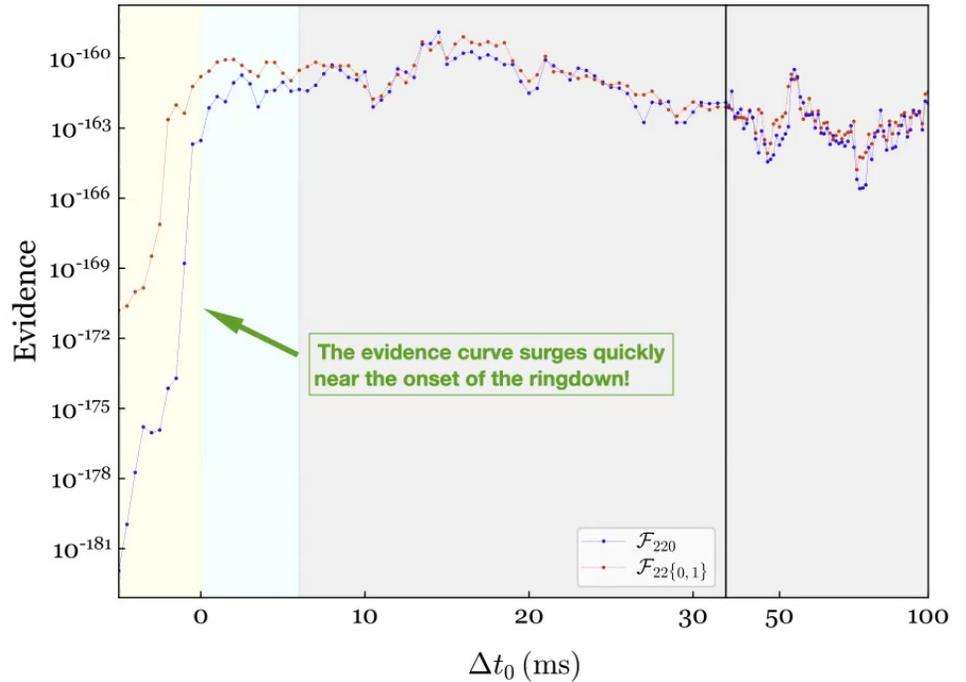
Determining the start time of ringdown (hyperpara



- The model evidence $P_{lmn}(t_0) = \iint dM_f d\chi_f p(d | M_f, \chi_f, t_0) \times \text{prior}$

Red: Fundamental mode + the first overtone
 Blue: Fundamental mode

GW event trigger?

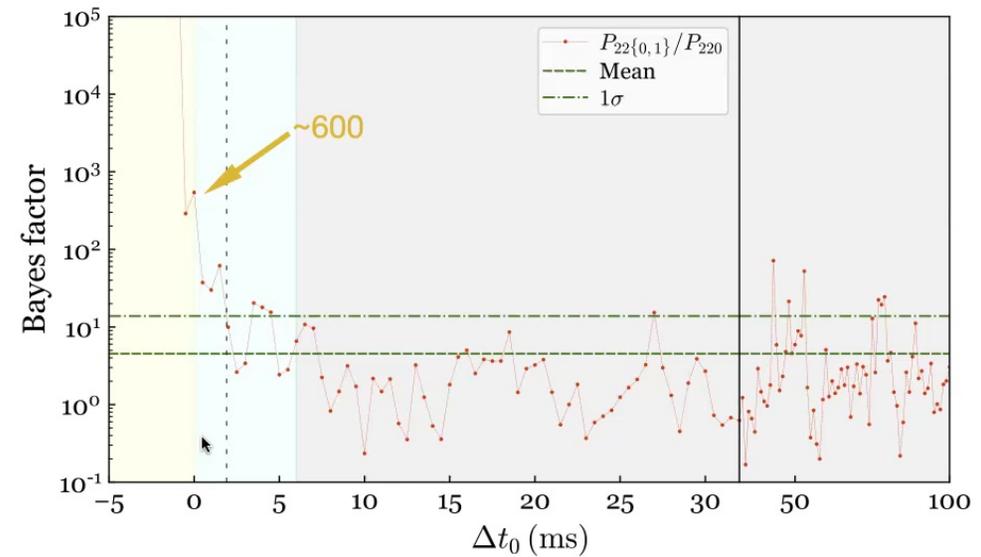


The Bayes factor of the first overtone in the ringdown of GW150914



- The model evidence $P_{lmn}(t_0) = \iint dM_f d\chi_f p(d | M_f, \chi_f, t_0) \times \text{prior}$

- The Bayes factor: $K_{221}(t_0) = \frac{P_{22\{0+1\}}}{P_{220}}$



Ma, et al., in preparation



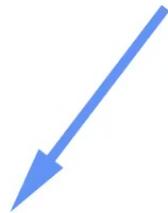
Three approaches to ringdown analysis

1. The traditional MCMC method (building a template)
2. The filter method (removing all ringdown signals)
3. A mixed approach!

The mixed approach



1. Remove a subset of QNM(s) via the filter(s)
2. Fit the filtered data to the remaining QNM(s) model via MCMC



Obtain information of the remaining modes
w/o the impact of other modes (cross-check)

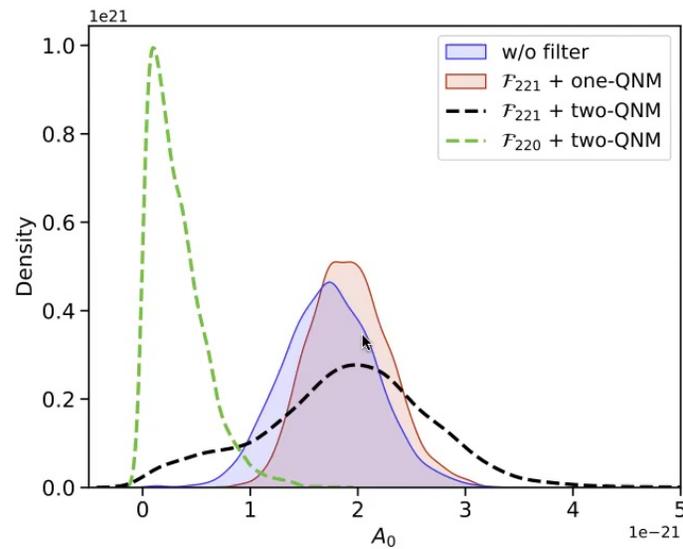


Infer remnant mass and spin from each single QNM
(test the no-hair theorem)

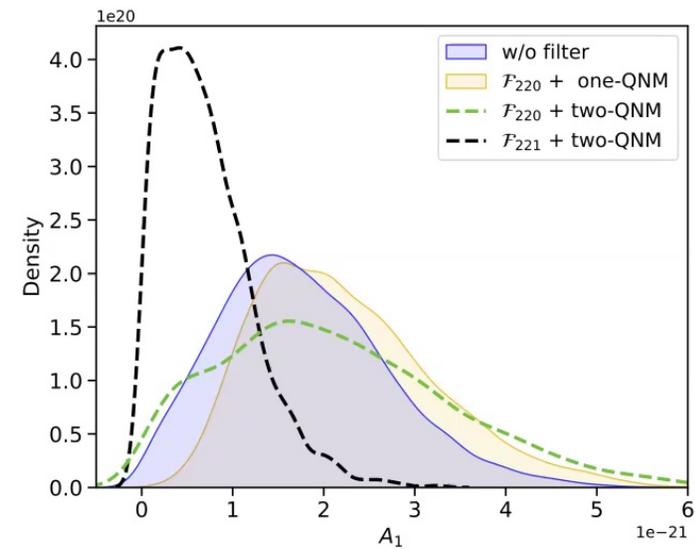
Taking GW150914 as an exam



Obtain information of the remaining modes
w/o the impact of other modes (cross-check)



The amplitude of the fundamental mode



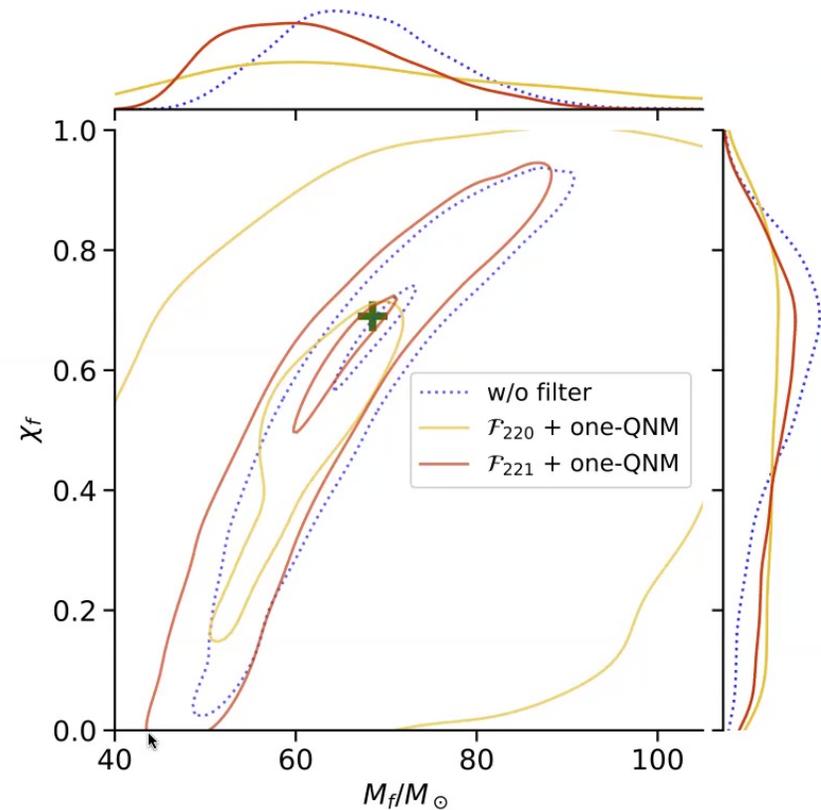
The amplitude of the first overtone

Ma, et al., in preparation

Taking GW150914 as an example



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Ma, et al., in preparation

Outline



- Properties of the rational filter in terms of numerical-relativity waveforms
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QNM filter = rational filter + full filter

The full filter

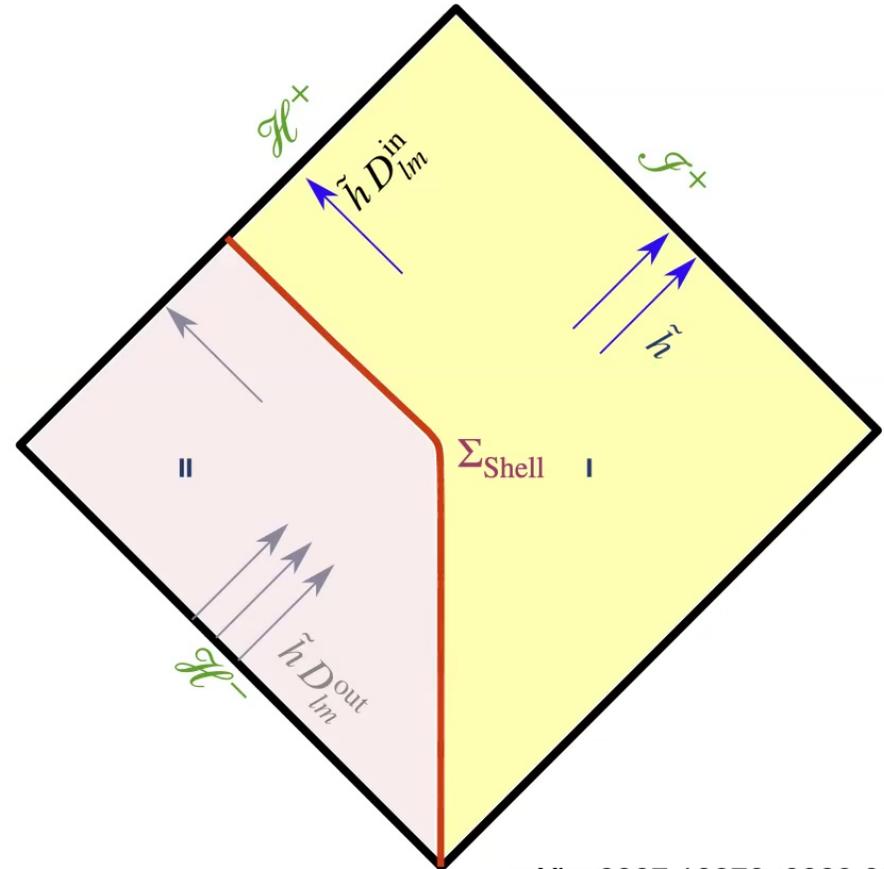


$$R_{lm}^{\text{up}} \sim \begin{cases} r^3 e^{i\omega r_*}, & r_* \rightarrow +\infty, \\ D_{lm}^{\text{out}} e^{i\omega r_*} + \Delta^2 D_{lm}^{\text{in}} e^{-i\omega r_*}, & r_* \rightarrow -\infty, \end{cases}$$

$$D_{lm}^{\text{out}}(\omega_{lmn}) \stackrel{!}{=} 0$$

$$D_{lm}^{\text{out}} \sim \prod_n (\omega - \omega_{lmn})$$

$$\mathcal{F}_{lm}^D = \frac{D_{lm}^{\text{out}}}{D_{lm}^{\text{out}*}}$$

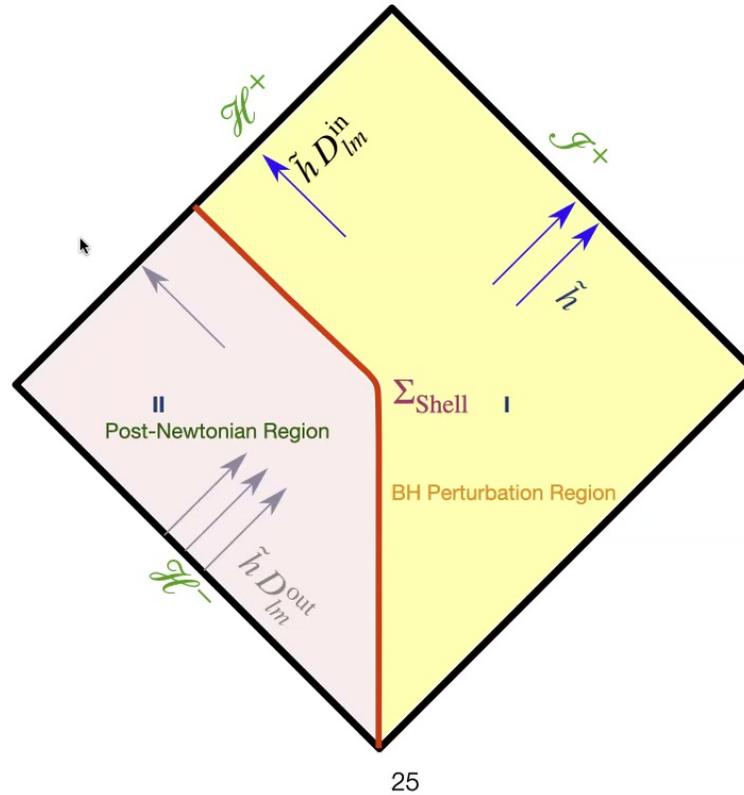


arXiv: 2207.10870, 2203.03174

The hybrid waveform model



Nichols and Chen (2010) & (2012)
arXiv: 1007.2024, 1109.0081



Ma et al., arXiv: 2203.03174

The image wave



$$R_{lm}^{\text{up}} \sim \begin{cases} r^3 e^{i\omega r_*}, & r_* \rightarrow +\infty, \\ D_{lm}^{\text{out}} e^{i\omega r_*} + \Delta^2 D_{lm}^{\text{in}} e^{-i\omega r_*}, & r_* \rightarrow -\infty, \end{cases}$$

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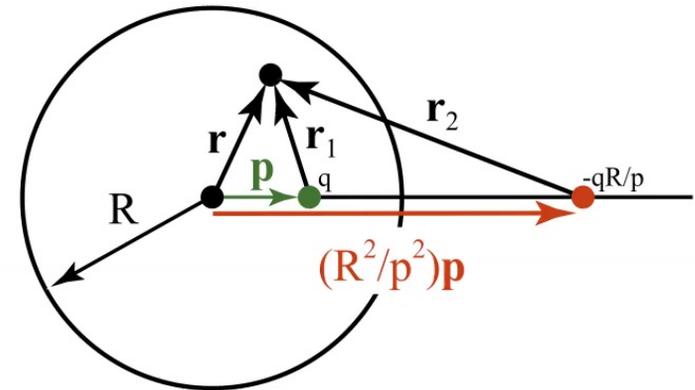
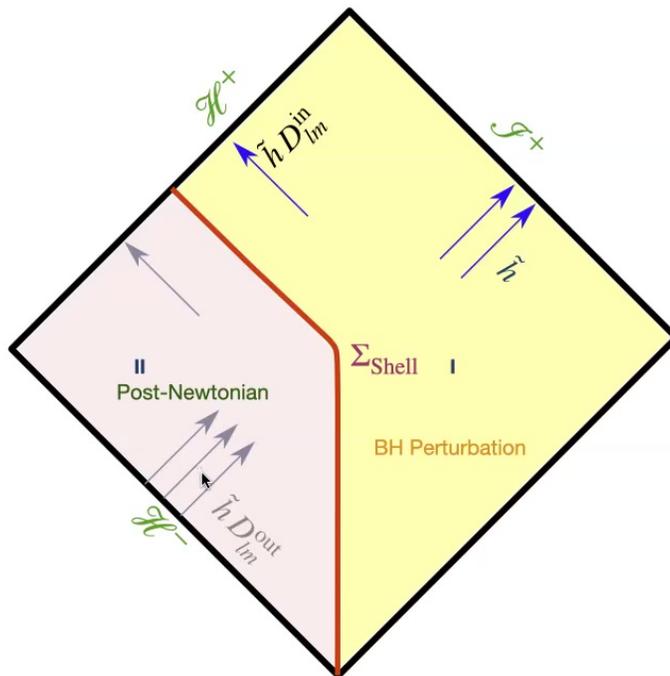
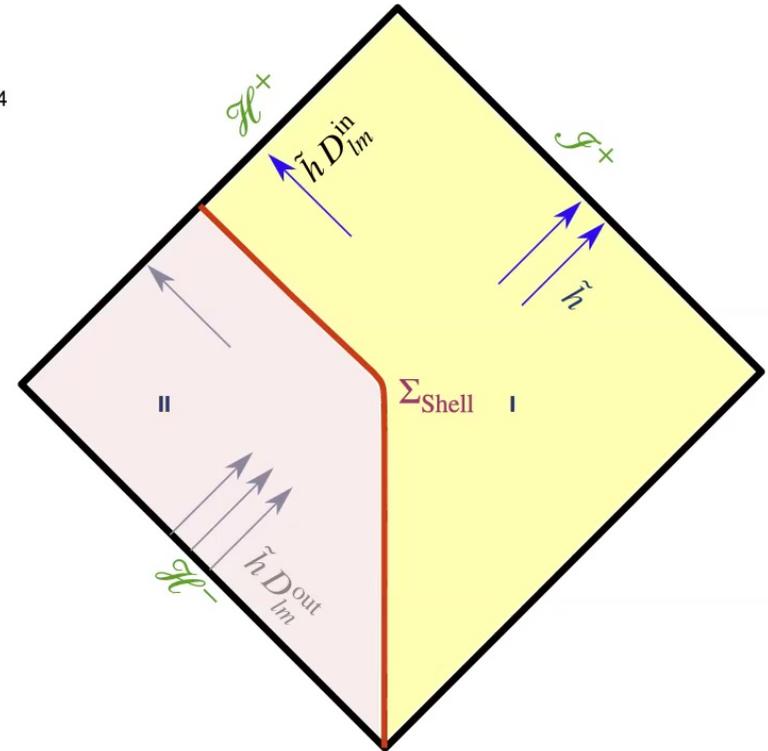


Image Credit: Wikipedia

Metric reconstruction near BBHs and its applications

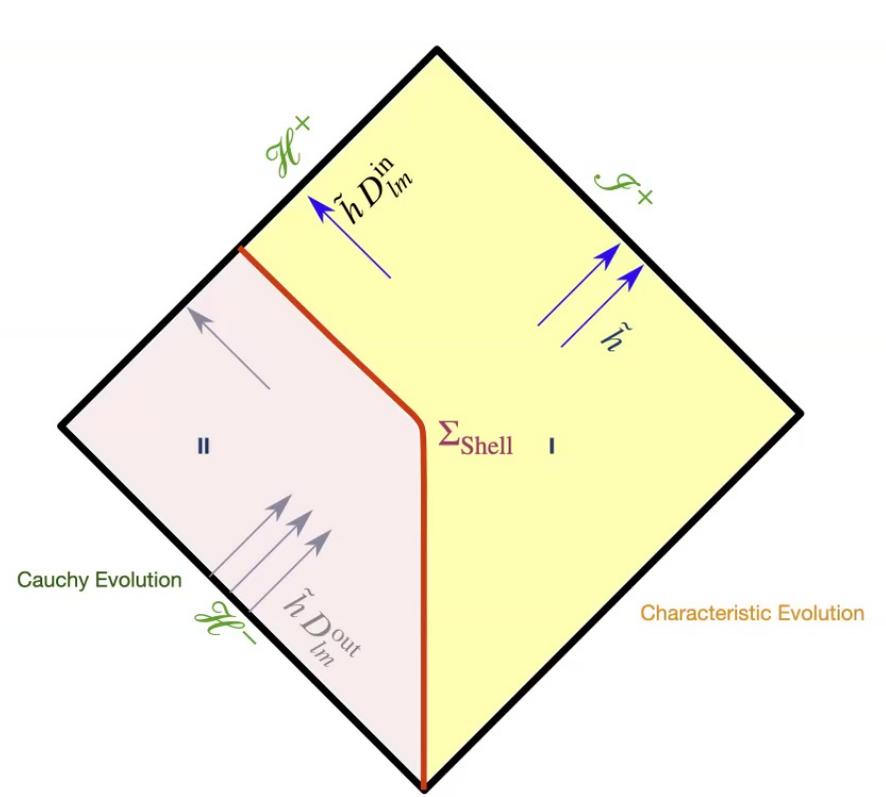
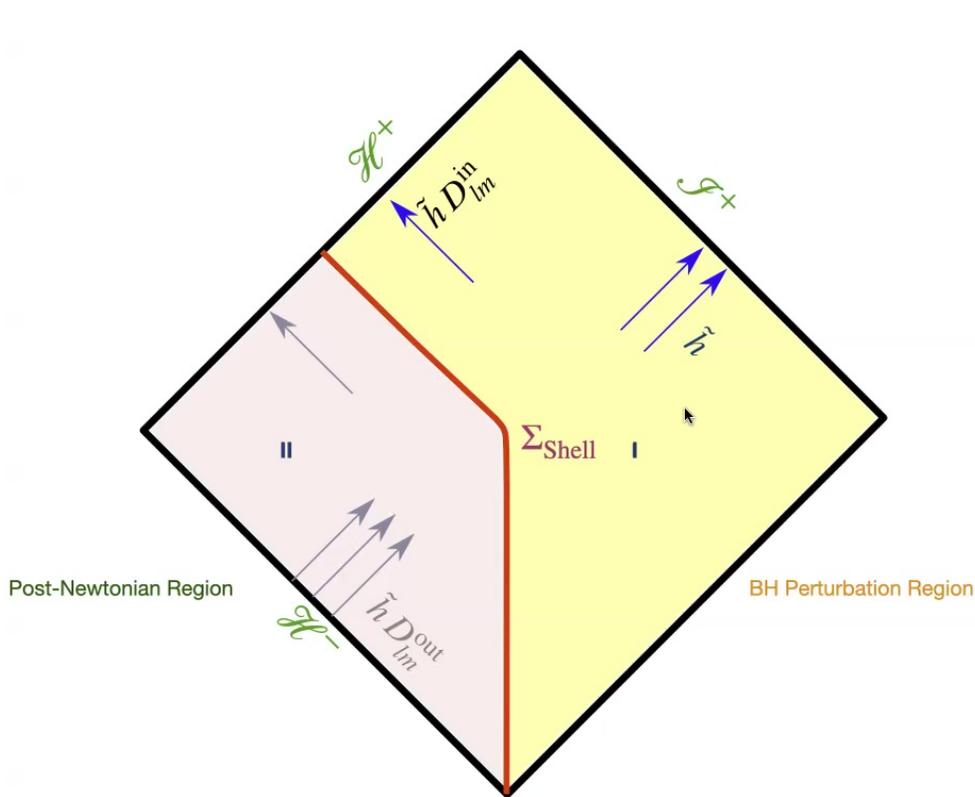


- **Computing gravitational-wave echoes.** Ma et al., arXiv: 2203.03174
- **Second-order BH perturbations?**



Numerical-relativity version of the hybrid method

Cauchy-Characteristic Matching



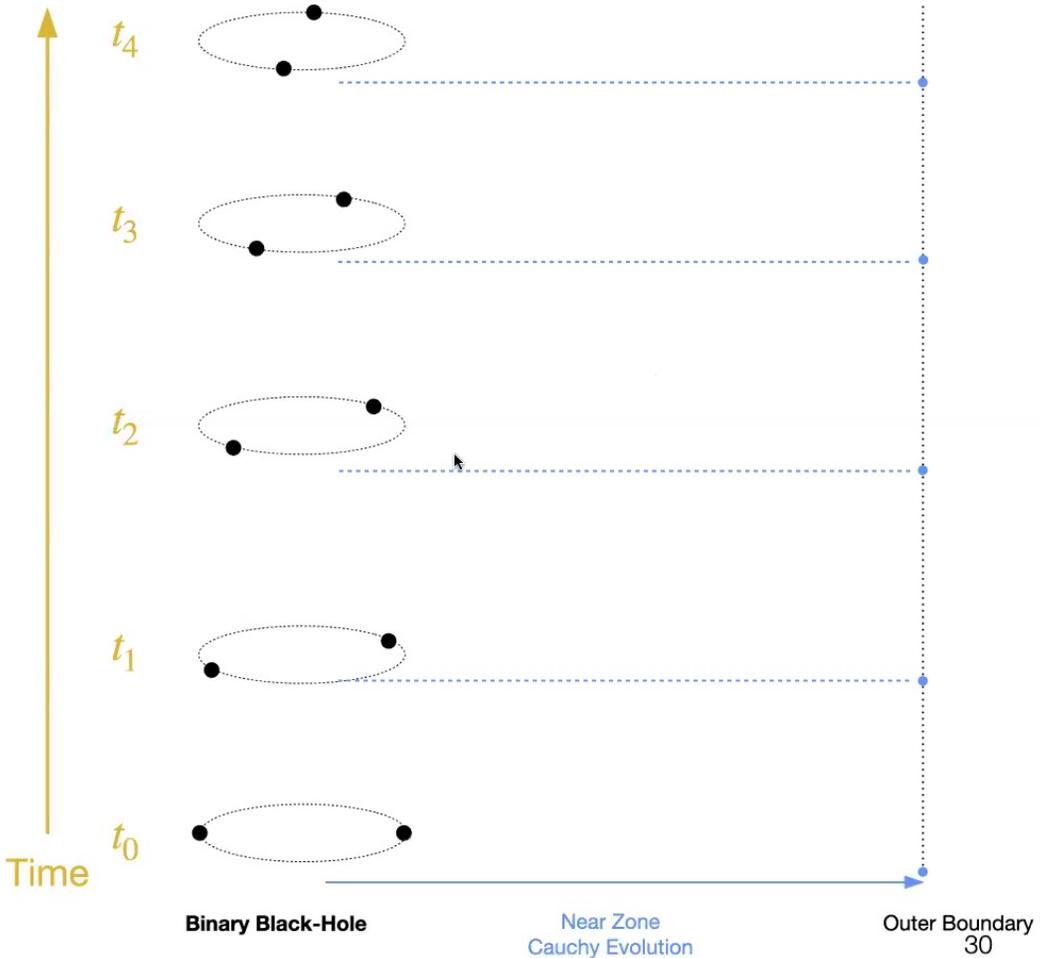
Fully relativistic, 3D CCM in SpECTRE, Ma et al., in preparation

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- **Cauchy-Characteristic Matching** (briefly, Ma, et al., in prep)

Cauchy-Characteristic Extraction

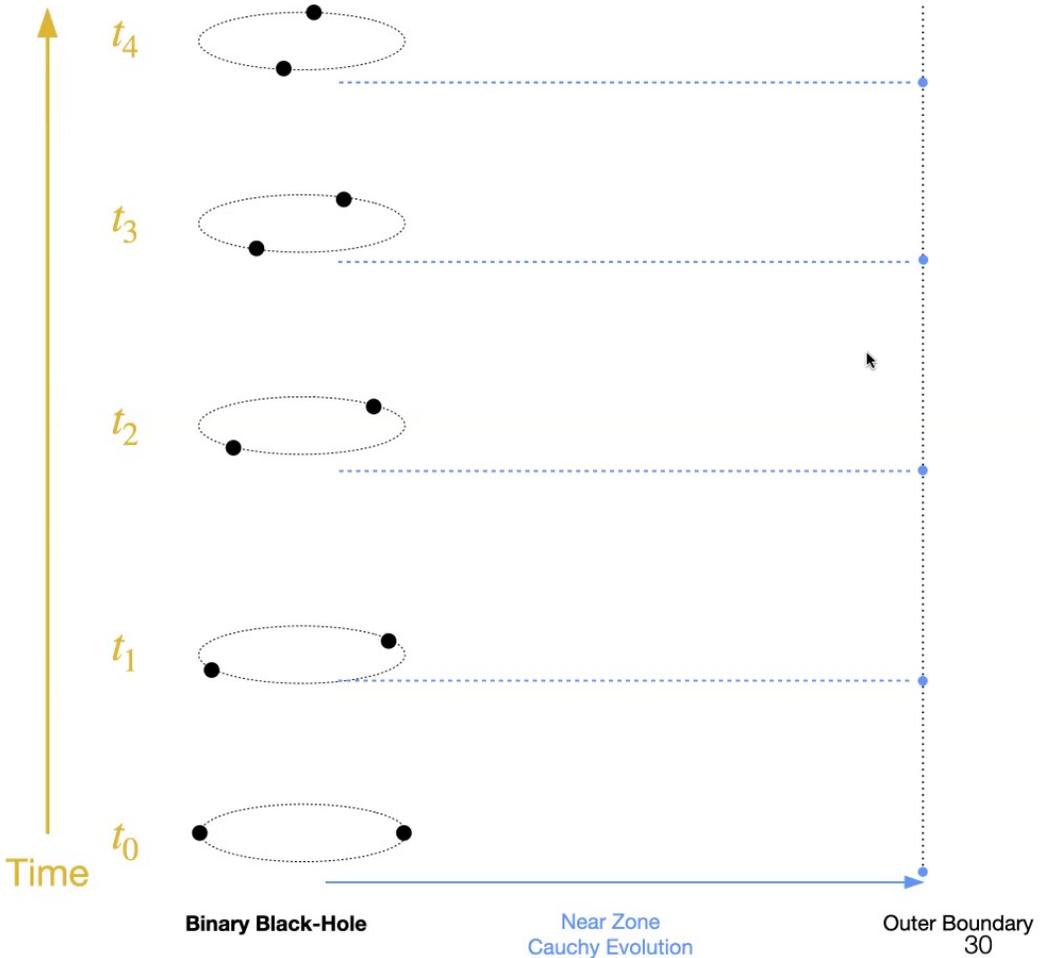


Cauchy-Characteristic Extraction



Sisheng Ma

LIGO



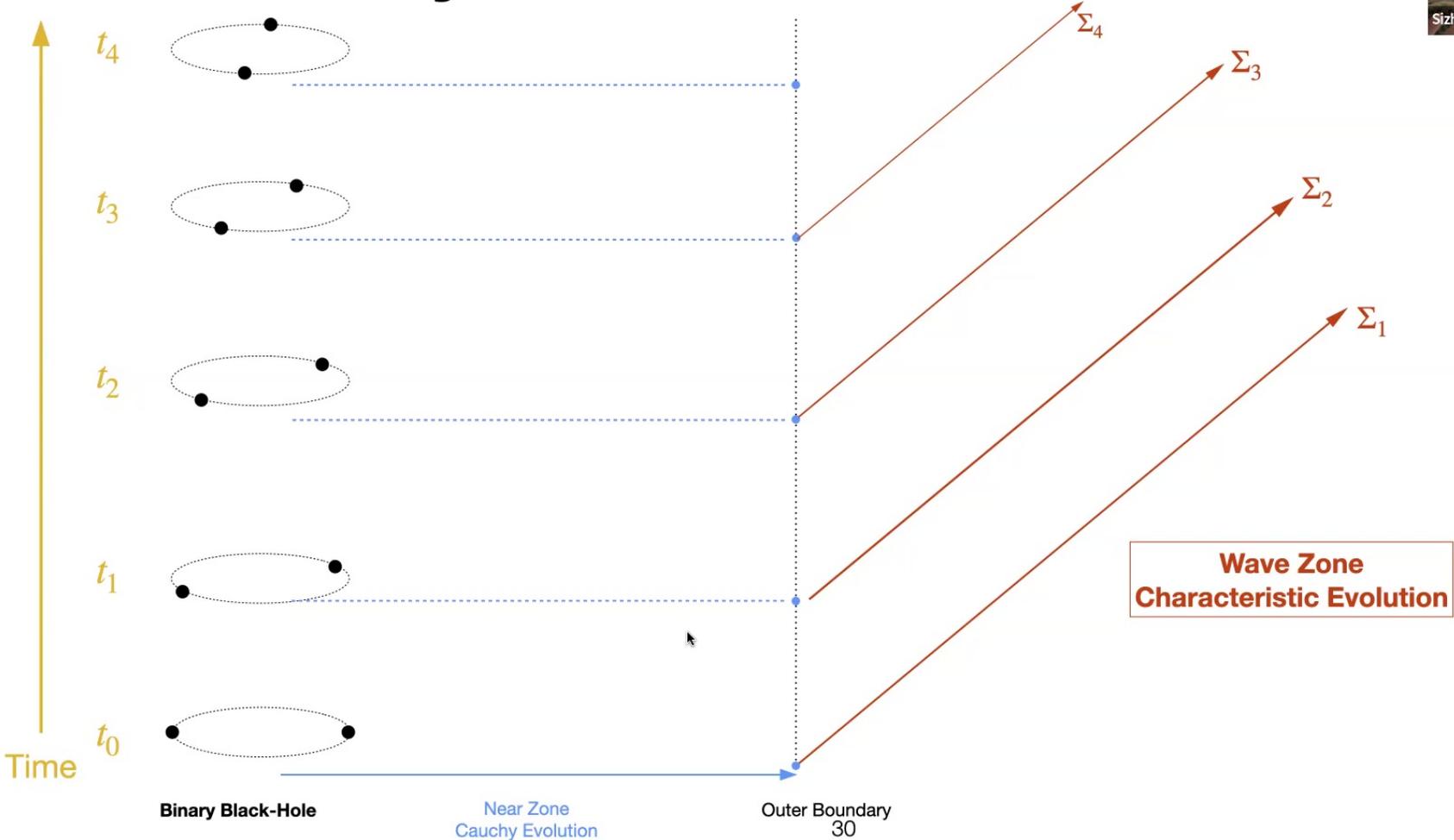
**Wave Zone
Characteristic Evolution**

Cauchy-Characteristic Extraction



Sizheng Ma

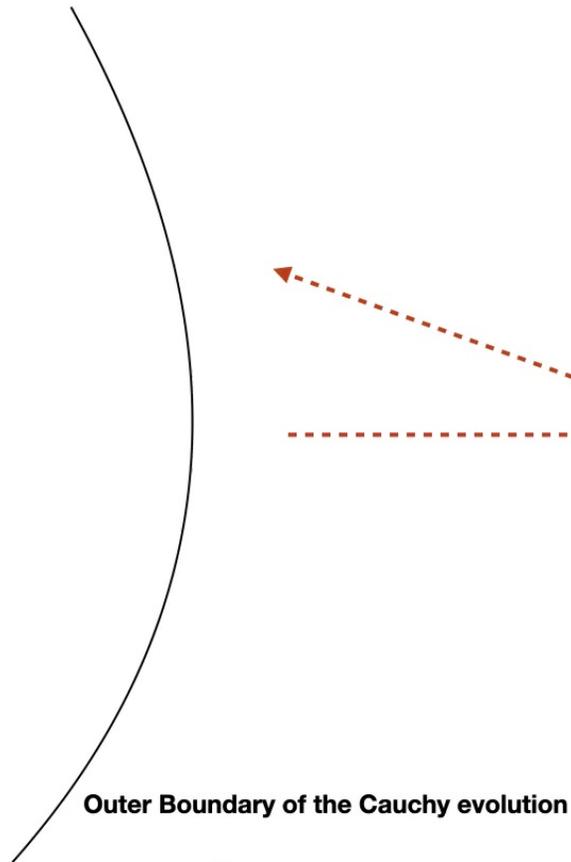
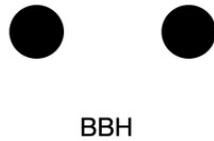
LIGO



The backscattered gravitational wave

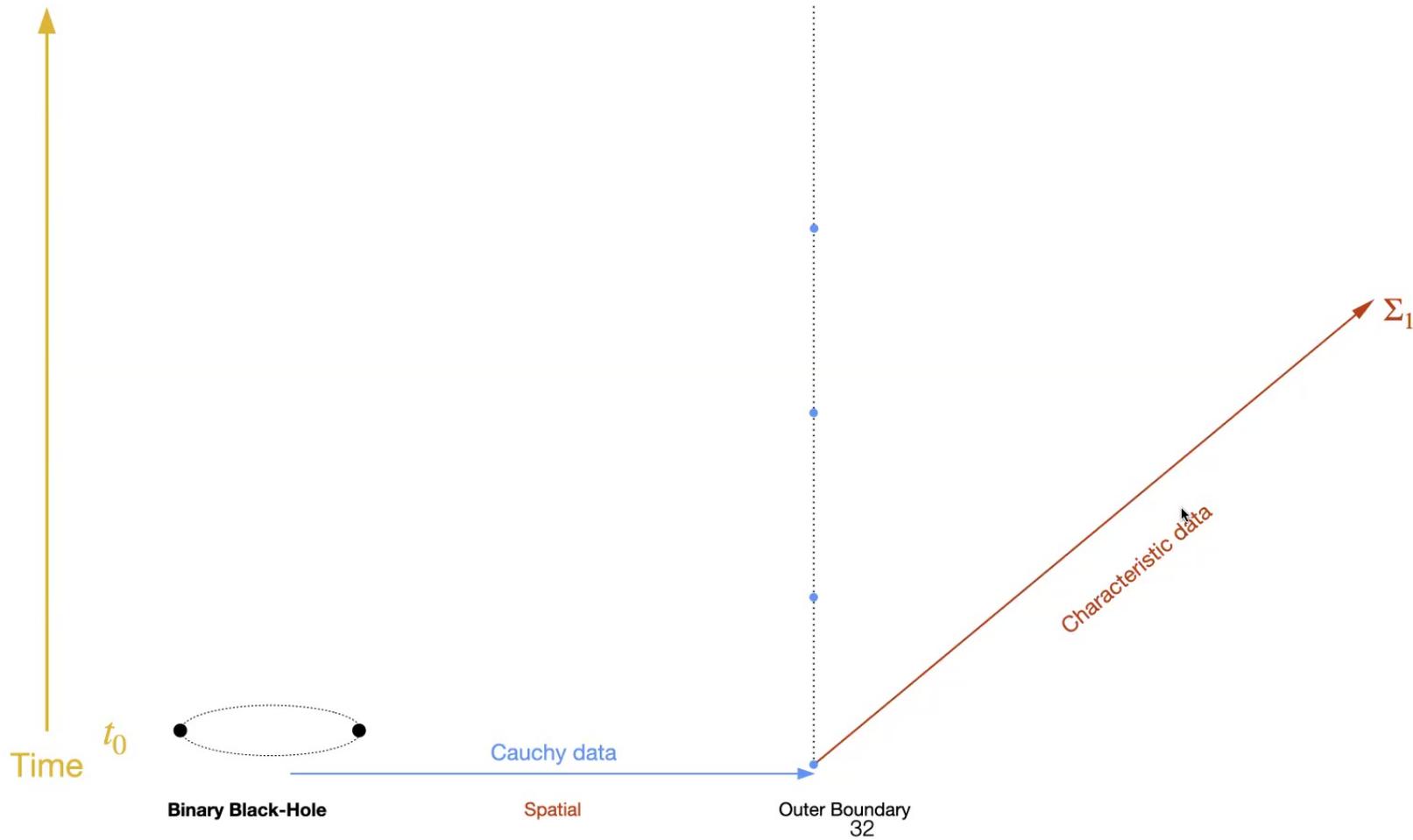


Cauchy computational domain

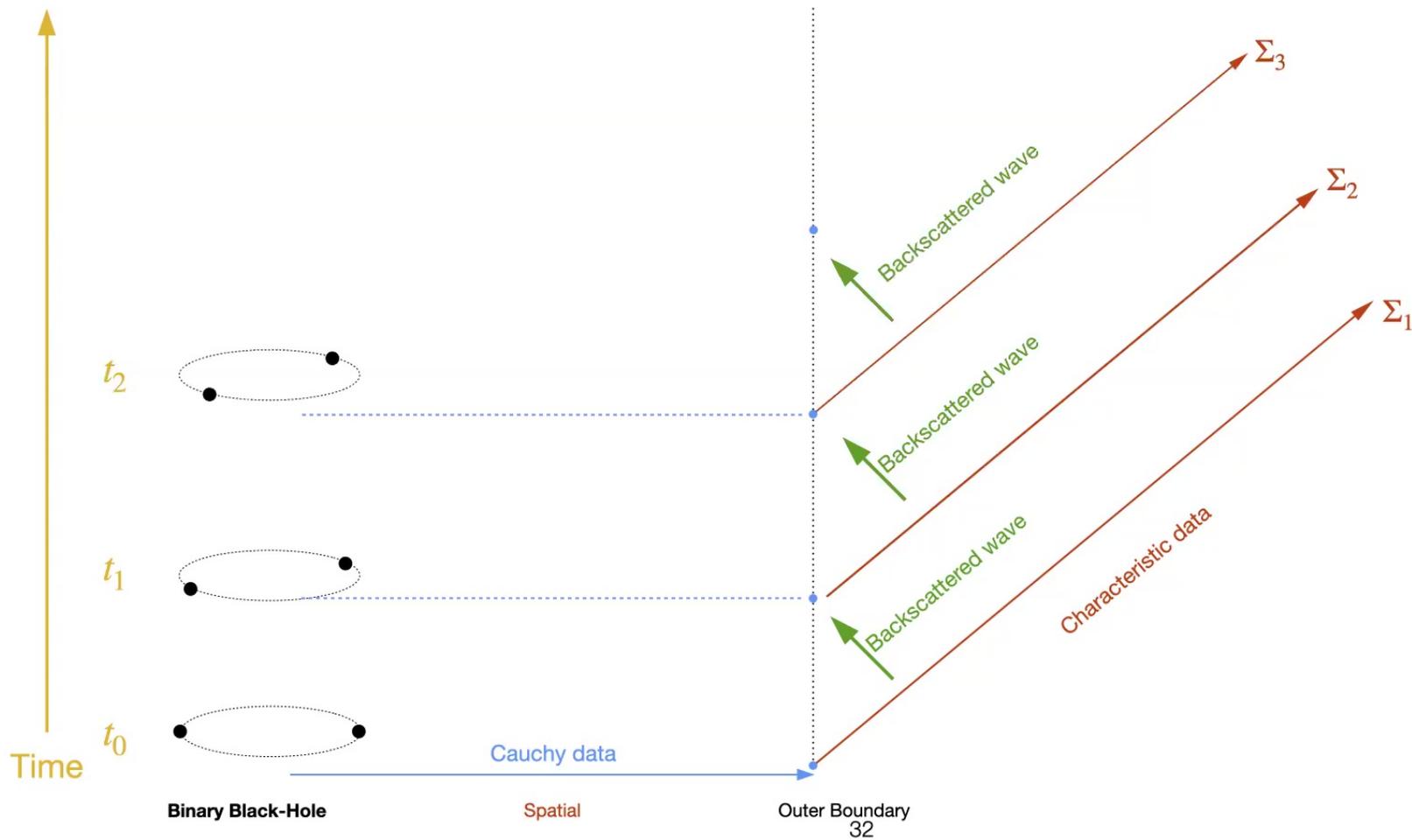


Ψ_0

Cauchy-Characteristic Matching



Cauchy-Characteristic Matching



The Bjorhus boundary conditions



$$\partial_t u^\alpha + A_\beta^{k\alpha} \partial_k u^\beta = F^\alpha$$

Generalized Harmonic
(SpECTRE)

$$d_t u_{ab}^{1-} =$$

Incoming characteristic field
on boundary

4 DOF: Constraint-preserving

4 DOF: Gauge

2 DOF: incoming GWs



The performance of the fully relativistic 3D

- Spurious numerical reflections at the outer boundary are reduced
- The interface between the Cauchy and characteristic system is transparent to physical GWs

Ma et al., in preparation

Summary



1. The rational filter and the corresponding Bayesian framework
 - This novel framework is a powerful tool to do BH spectroscopy (theoretically and observationally)
 - A LIGO O4 pipeline?
2. The full filter; metric reconstruction; echo
3. Cauchy-characteristic matching

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