

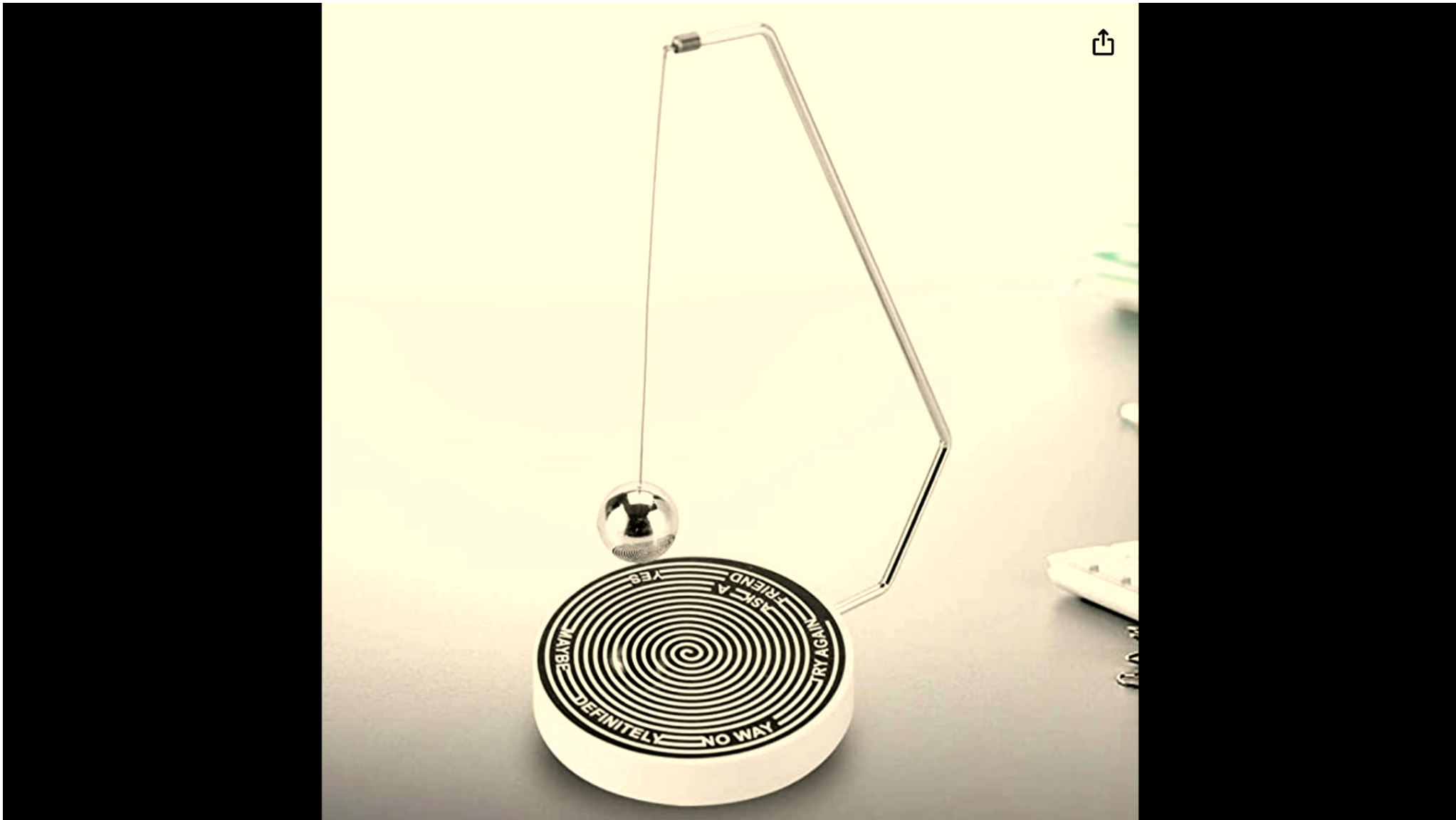
Title: Classical Physics - Lecture 220919

Speakers: Meenu Kumari

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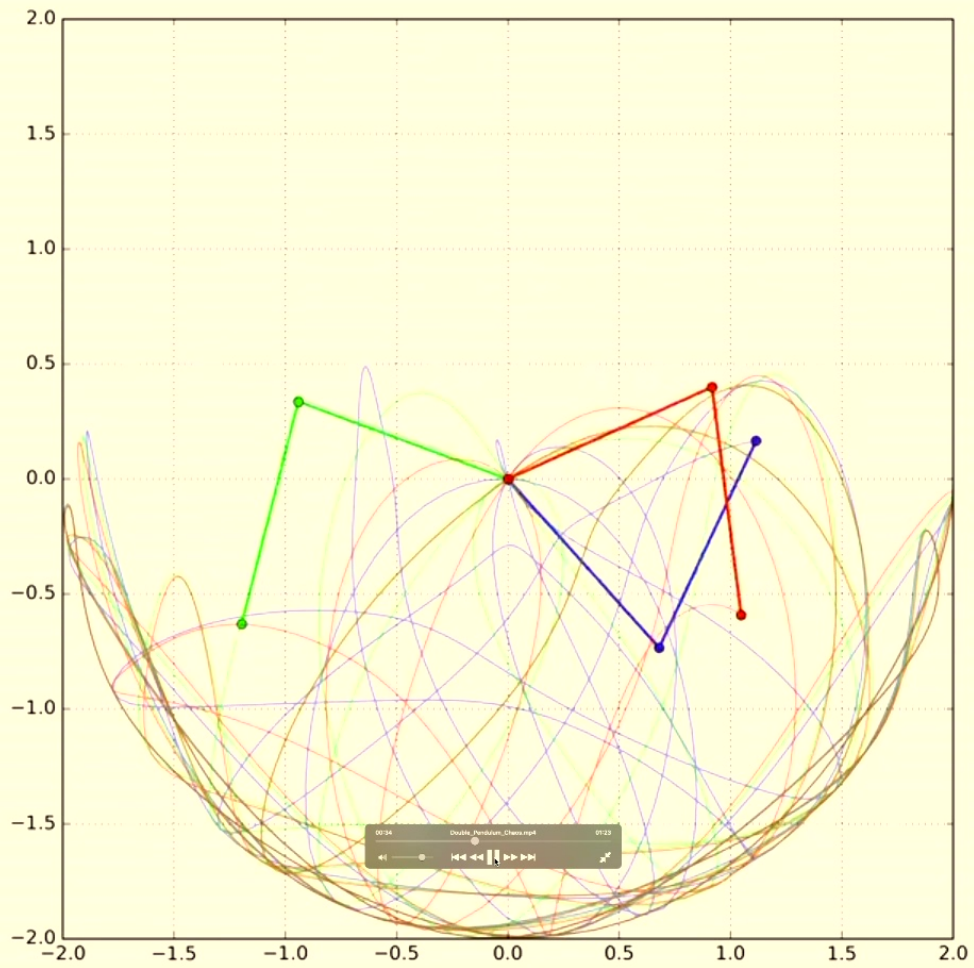
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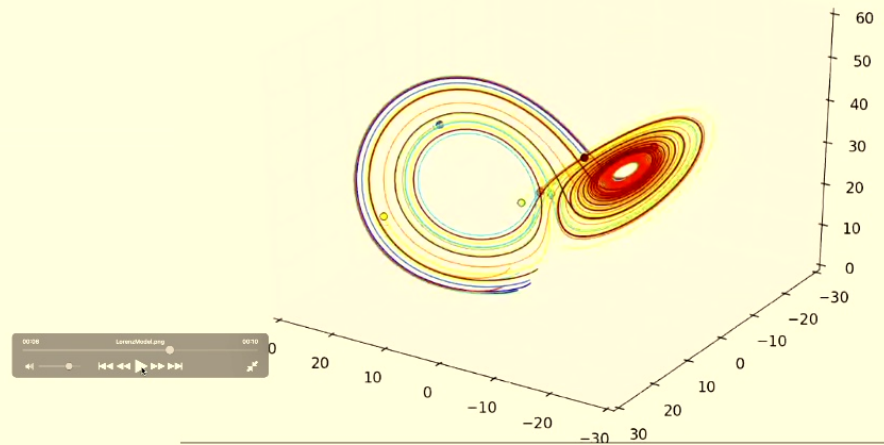


# Classical Chaos

**Lorenz oscillator** : Classic example of chaotic behaviour

- Non-linear deterministic dynamical systems.

$$\begin{aligned}\frac{dx}{dt} &= \sigma(y - x) \\ \frac{dy}{dt} &= x(\rho - z) - y \\ \frac{dz}{dt} &= xy - \beta z\end{aligned}$$



## Chaos

- Unpredictability due to extreme sensitivity to initial conditions
- Nonintegrability - a necessary condition in Hamiltonian system
- Nonlinear deterministic dynamical systems



Heisenberg's uncertainty  
principle

$$\langle x \rangle, \langle p \rangle$$

$$|\psi\rangle$$

$$|\psi'\rangle$$

$$H$$

$$U = e^{-iHt}$$

$$\langle \psi(0) | \psi'(0) \rangle \stackrel{?}{=} \langle \psi(t) | \psi'(t) \rangle$$

Heisenberg's uncertainty  
principle

$$\langle x \rangle, \langle p \rangle$$

$$|\psi(0)\rangle \quad | \quad H$$

$$|\psi'(0)\rangle$$

$$U = e^{-iHt}$$

$$\langle \psi(0) | \psi'(0) \rangle \stackrel{?}{=} \langle \psi(t) | \psi'(t) \rangle$$

$$\underline{UU^\dagger = 1}$$



extreme  
conditions  
necessary condition  
dynamical systems

Lyapunov exponent

$$X_0 \text{ and } X'_0 \in \mathbb{R}^n$$

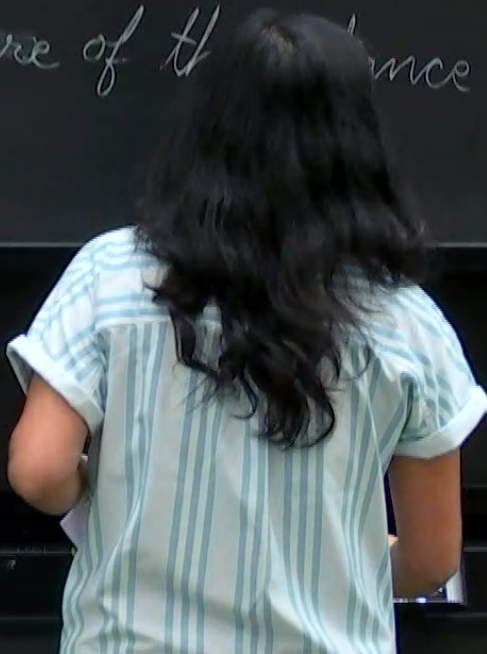
$$\begin{cases} X'_0 = X_0 + \delta X_0 \\ \downarrow \\ X(t) \quad X(t) + \delta X(t) \end{cases}$$

$$d(X_0, t) = \|\delta X(X_0, t)\| = \sqrt{\delta x_1^2 + \delta x_2^2 + \dots + \delta x_n^2}$$

Measure of the distance between two trajectories

Heisenberg's uncertainty principle

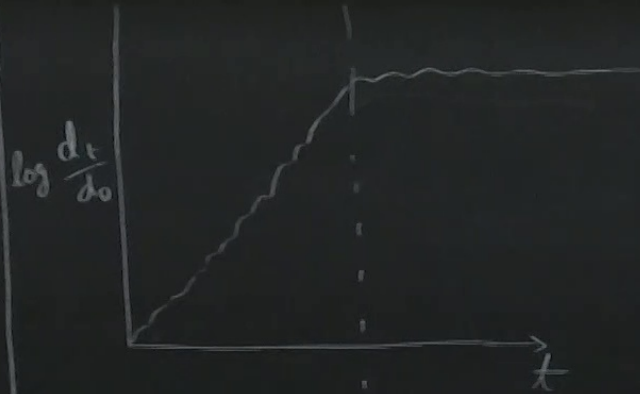
$$\langle x \rangle, \langle p \rangle$$



$$\lambda(x_0, \delta x) = \lim_{t \rightarrow \infty} \lim_{d(x_0, 0) \rightarrow 0} \frac{1}{t} \log \left( \frac{d(x_0, t)}{d(x_0, 0)} \right)$$

$\lambda > 0$  in chaotic systems

$$\log \left( \frac{d(x_0, t)}{d(x_0, 0)} \right)$$



$$\log \frac{d_t}{d_0} = ct$$
$$\Rightarrow d_t = e^{ct} d_0$$

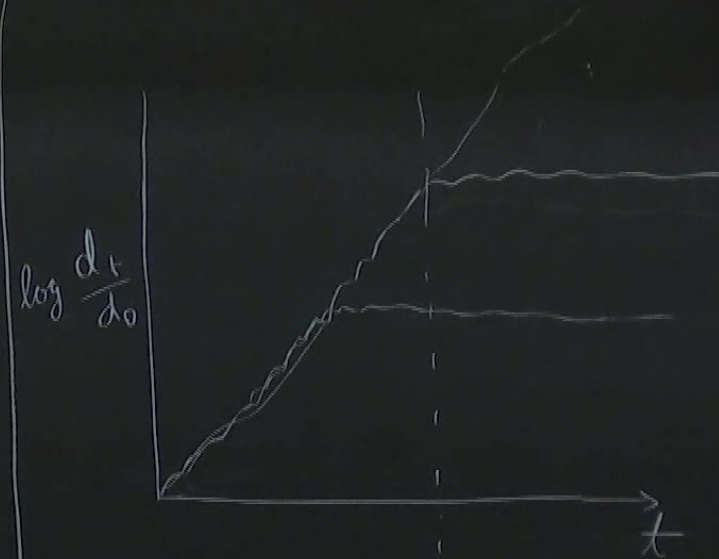
$\vec{F}$  ,  $\vec{\nabla} \cdot \vec{F} < 0$  för dissipative system  
 $\vec{\nabla} \cdot \vec{F} = 0$  för Hamiltonian systems



$$s(X) = \lim_{t \rightarrow \infty} \lim_{d(X_0, 0) \rightarrow 0} \frac{1}{t} \log \left( \frac{d(X_0, t)}{d(X_0, 0)} \right)$$

in chaotic systems

$$X_0 \quad X'_0$$



$$\log \frac{d_t}{d_0} = ct$$

$$\Rightarrow d_t = e^{ct} d_0$$

$$(p_i, q_i)$$

DOF - No. of  $q_i$ 's

Dimension of phase space = 2 DOF

An integrable system:

- No. of DOF = No. of constants of motion  
(COM)

$$\{A, H\} = 0$$

$n$  DOF,  $n$  COM  $\{F_i\}$

$$\rightarrow d_+ = e^- d_0$$



$n$  DOF,  $n$  COM  $\{F_i\}_{i=1}^m$

$$\{F_1, F_2\} = F_3$$

— An involution  $\{F_i, F_j\} =$

— Functionally independent of each other

Integrable systems

① Free

$n$   
 $i=1$   
 $F_j = 0$   
t of

## Integrable systems

- ① Free particle
- ② Simple pendulum
- ③ All 1-d time independent potentials
- ④ Kepler problem, spherically symmetric potentials

## Non-integrable systems

- ① Periodically driven simple pendulum
- ②  $n$ -body problem
- ③ Spherical pendulum
- ④ Double pendulum

$$d_t = \dots$$

$$\text{COM } \{F_i\}_{i=1}^n$$

$$\{F_i, F_j\} = 0$$

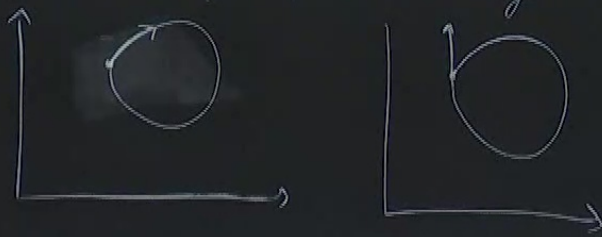
independent of

### Integrable systems

- ① Free particle
- ② Simple pendulum

All 1-d time independent potentials

Kepler problem, spherically symmetric potentials





simple pendulum

(2.74)  
 $F^i$  are the conserved quantities

$$\dot{F}^i = \{H, F^i\} = 0$$

$$\dot{\Psi}^i = \{H, \Psi^i\} = \frac{\partial H}{\partial F^i} = -\Omega_i = \Omega_i(F_j)$$

Solution

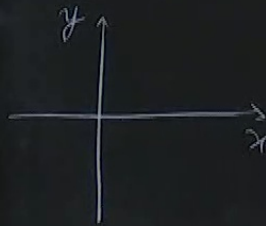
$$F^i(t) = F^i(0)$$

$$\Psi_i(t) = \Psi_i(0) + \underbrace{-\Omega_i}_t t$$



## Non-integrable systems

- ① Periodically driven simple pendulum
- ②  $n$ -body problem
- ③ Spherical pendulum
- ④ Double pendulum



(2.74) —  
 $F^i$  are the

$$\dot{F}^i = \{H, F^i\}$$

$$\dot{\Psi}^i = \{H, \Psi^i\}$$

Solution

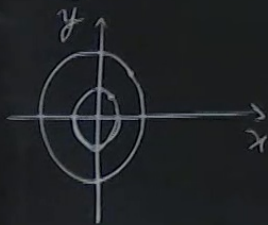
$$F^i(t) =$$

$$\Psi_i(t) =$$



## Non-integrable systems

- ① Periodically driven simple pendulum
- ② n-body problem
- ③ Spherical pendulum
- ④ Double pendulum



n DOF, 2n dim phase space

Energy shell is (2n-1) dimensional

(2.74) —

$F^i$  are the conserved quantities

$$\dot{F}^i = \{H, F^i\} = 0$$

$$\dot{\psi}^i = \{H, \psi^i\} = \frac{\partial H}{\partial F^i} = -\Omega_i = \Omega_i(F_j)$$

Solution

$$F^i(t) = F^i$$

$$\psi^i(t) = \psi^i + t$$

le systems

driven simple pendulum

problem

pendulum

OF, phase space

1) dimensional

(2.74) —

$F^i$  are the conserved quantities

$$\dot{F}^i = \{H, F^i\} = 0$$

$$\dot{\Psi}^i = \{H, \Psi^i\} = \frac{\partial H}{\partial F^i} = -\Omega_i = \Omega_i(F_j)$$

Solution

$$F^i(t) = F^i(0)$$

$$\Psi_i(t) = \Psi_i(0) + \underbrace{\Omega_i}_t$$

