

Title: Quantum / optical mechanical sensors

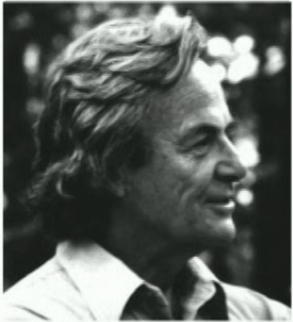
Speakers: Markus Aspelmeyer

Collection: School on Table-Top Experiments for Fundamental Physics

Date: September 22, 2022 - 9:15 AM

URL: <https://pirsa.org/22090020>

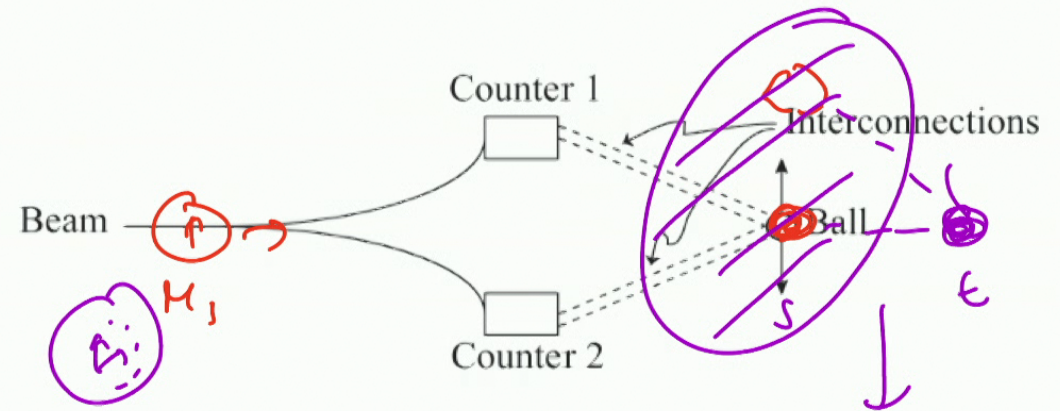
How does gravity react to quantum systems?



R P Feynman

"One should think about designing an experiment which uses a gravitational link and at the same time shows quantum interference"

Chapel Hill Conference 1957



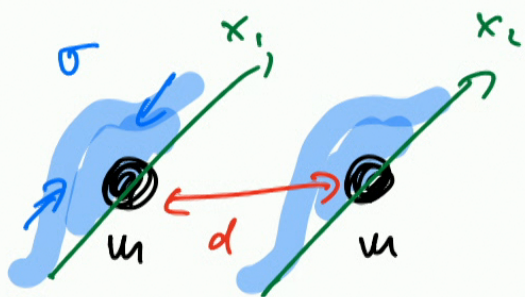
$$\left(|x_{\bar{u}}\rangle_s + |x_d\rangle_s \right) \otimes |x_0\rangle_t$$

$$t \downarrow a_g^{I,\bar{E}} = G \frac{M_s}{d_{I,\bar{E}}^2}$$

$$\left(|x_{\bar{u}}\rangle_s |x_{\bar{u}}\rangle_t + |x_d\rangle_s |x_d\rangle_t \right)$$

$$\langle x_{\bar{u}} | x_d \rangle_t \ll 1$$

ENTANGLED iff

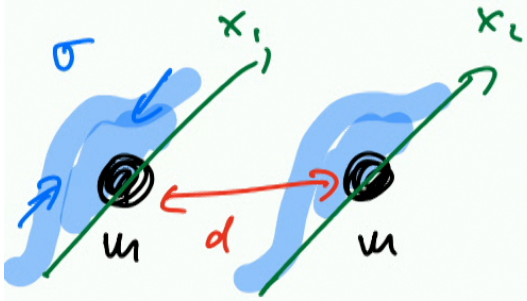


$$W_{\text{int}} \sim G \frac{m^2}{r} = G \frac{m^2}{\sqrt{d^2 + (\kappa_1 - \kappa_2)^2}}$$

$$\approx G \frac{m^2}{d} \left(1 - \frac{(\kappa_1 - \kappa_2)^2}{2d^2} + \dots \right)$$

$$= G \frac{m^2}{d} \left(1 - \frac{\kappa_1^2}{2d^2} - \frac{\kappa_2^2}{2d^2} + \frac{\kappa_1 \kappa_2}{d^2} + \dots \right)$$

$$\rightarrow W_{\text{int}} = G \frac{m^2}{d^3} \kappa_1 \kappa_2$$



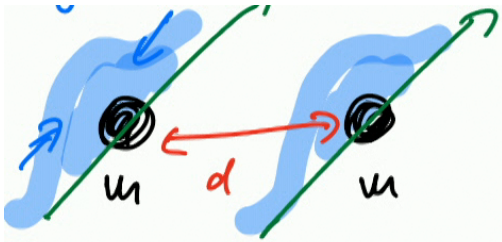
$$V_{\text{int}} \sim G \frac{m^2}{r} = G \frac{m^2}{\sqrt{d^2 + (x_1 - x_2)^2}}$$

$$\approx G \frac{m^2}{d} \left(1 - \frac{(x_1 - x_2)^2}{2d^2} + \dots \right)$$

$$= G \frac{m^2}{d} \left(1 - \frac{x_1^2}{2d^2} - \frac{x_2^2}{2d^2} + \frac{x_1 x_2}{d^2} + \dots \right)$$

$$\rightarrow \overset{1}{V}_{\text{int}} = G \frac{m^2}{d^3} \overset{1}{x}_1 \overset{1}{x}_2 \quad \overset{1}{x} = \sigma (\overset{1}{a} + \overset{1}{a}^+)$$

$$= G \frac{m^2}{d^3} \sigma^2 (\overset{1}{a} + \overset{1}{a}^+) (\overset{1}{j} + \overset{1}{j}^+)$$



$$H_{\text{int}} \sim G \frac{m_1 m_2}{r} = G \frac{m_1 m_2}{\sqrt{d^2 + (\kappa_1 - \kappa_2)^2}}$$

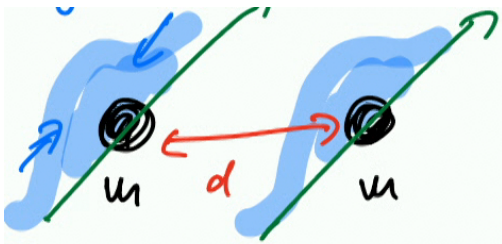
$$\approx G \frac{m_1 m_2}{d} \left(1 - \frac{(\kappa_1 - \kappa_2)^2}{2d^2} + \dots \right)$$

$$= G \frac{m_1 m_2}{d} \left(1 - \frac{\kappa_1^2}{2d^2} - \frac{\kappa_2^2}{2d^2} + \frac{\kappa_1 \kappa_2}{d^2} + \dots \right)$$

$$\rightarrow \overset{1}{H_{\text{int}}} = G \frac{m_1 m_2}{d^3} \overset{1}{\kappa_1 \kappa_2} \quad \overset{1}{\kappa} = \sigma (\overset{1}{a} + \overset{1}{a}^+)$$

$$= G \frac{m_1 m_2}{d^3} \sigma^2 (\overset{1}{a} + \overset{1}{a}^+) (\overset{1}{j} + \overset{1}{j}^+)$$

$$= \frac{1}{2} g (\overset{1}{a} + \overset{1}{a}^+) (\overset{1}{j} + \overset{1}{j}^+)$$



$$H_{\text{int}} \sim G \frac{m}{r} = G \frac{1}{\sqrt{d^2 + (\kappa_1 - \kappa_2)^2}}$$

$$\approx G \frac{m^2}{d} \left(1 - \frac{(\kappa_1 - \kappa_2)^2}{2d^2} + \dots \right)$$

$$= G \frac{m^2}{d} \left(1 - \frac{\kappa_1^2}{2d^2} - \frac{\kappa_2^2}{2d^2} + \frac{\kappa_1 \kappa_2}{d^2} + \dots \right)$$

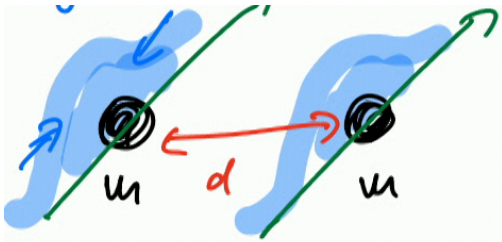
$$\rightarrow H_{\text{int}} = G \frac{m^2}{d^3} \kappa_1 \kappa_2$$

$$\kappa^2 = \sigma^2 (\vec{a}^2 + \vec{a}^{\prime 2})$$

$$= G \frac{m^2}{d^3} \sigma^2 (\vec{a}^2 + \vec{a}^{\prime 2}) (\vec{f}^2 + \vec{f}^{\prime 2})$$

$$= \frac{1}{2} g (\vec{a}^2 + \vec{a}^{\prime 2}) (\vec{f}^2 + \vec{f}^{\prime 2})$$

$$\rightarrow g = \frac{G}{\frac{1}{2}} \frac{m^2}{d} \frac{\sigma^2}{d^2}$$



$$H_{\text{int}} \sim G \frac{u}{r} = G \frac{1}{\sqrt{d^2 + (\kappa_1 - \kappa_2)^2}}$$

$$\approx G \frac{u^2}{d} \left(1 - \frac{(\kappa_1 - \kappa_2)^2}{2d^2} + \dots \right)$$

$$= G \frac{u^2}{d} \left(1 - \frac{\kappa_1^2}{2d^2} - \frac{\kappa_2^2}{2d^2} + \frac{\kappa_1 \kappa_2}{d^2} + \dots \right)$$

$$\rightarrow H_{\text{int}} = G \frac{u^2}{d^3} \kappa_1 \kappa_2 \quad \kappa_i = \sigma (\hat{a}_i + \hat{a}_i^\dagger)$$

$$= G \frac{u^2}{d^3} \sigma^2 (\hat{a}_1 + \hat{a}_1^\dagger) (\hat{a}_2 + \hat{a}_2^\dagger)$$

$$= \frac{1}{\hbar} g (\hat{a}_1 + \hat{a}_1^\dagger) (\hat{a}_2 + \hat{a}_2^\dagger)$$

$$\rightarrow g = \frac{G}{\hbar} \frac{u^2}{d} \frac{\sigma^2}{d^2} \quad \equiv \text{ENTANGLEMENT RATE}$$

→ $g = \frac{G}{\hbar} \frac{m^2}{d} \frac{\sigma^2}{d^2} \equiv \text{ENTANGLEMENT RATE}$

Examples: 1) 2 Atoms separated by 1 μm , $\sigma = \sqrt{\frac{\hbar}{m v}} \sim 10^{-7} \text{ m}$

$$g = 2\pi \cdot 10^{23} \frac{10^{-52}}{10^{-6}} 10^{-2} \sim 10^{-25} \text{ Hz}$$

→ $\tau_{\text{ent}} \sim 10^{25} \text{ sec.} \sim 10^8 \text{ Tian.}$

2) 2 PLANCK MASSES,

→ $g = \frac{G}{\hbar} \frac{m^2}{d} \frac{\sigma^2}{d^2} \equiv \text{ENTANGLEMENT RATE}$

Examples: 1) 2 ATOMS separated by 1 μm , $\sigma = \sqrt{\frac{\hbar}{m\omega}} \sim 10^{-7} \text{ m}$

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→ $\tau_{\text{ent}} \sim 10^{27} \text{ sec.} \sim 10^8 \text{ Tian.}$

2) 2 PLANCK MASSES, $d \sim 100 \mu\text{m}$, $\sigma \sim 1 \mu\text{m}$

$$g = 2\pi \cdot 10^{23} \frac{10^{-16}}{10^{-4}} 10^{-6} \sim$$

→ $g = \frac{G}{\hbar} \frac{m^2}{d} \frac{\sigma^2}{d^2} \equiv \text{ENTANGLEMENT RATE}$

Examples: 1) 2 ATOMS separated by 1 μm , $\sigma = \sqrt{\frac{\hbar}{m v}} \sim 10^{-7} \text{ m}$

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2) 2 PLANCK MASSES, $d \sim 100 \mu\text{m}$, $\sigma \sim 1 \mu\text{m}$

$$g = 2\pi \cdot 10^{23} \frac{10^{-16}}{10^{-4}} 10^{-10} \sim 10 \text{ Hz}$$

→

→ $g = \frac{G}{\hbar} \frac{m^2}{d} \frac{\sigma^2}{d^2} \equiv \text{ENTANGLEMENT RATE}$

Examples: 1) 2 ATOMS separated by 1 μm , $\sigma = \sqrt{\frac{\hbar}{m v}} \sim 10^{-7} \text{ m}$

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2) 2 PLUMBER MASSES, $d \sim 100 \mu\text{m}$, $\sigma \sim 1 \mu\text{m}$

$$g = 2\pi \cdot 10^{23} \frac{10^{-16}}{10^{-4}} 10^{-10} \sim 10 \text{ Hz}$$

→ $\tau_{ent} \sim 100 \mu\text{s}$

→ $g = \frac{G}{\hbar} \frac{m^2}{d} \frac{\sigma^2}{d^2} \equiv \text{ENTANGLEMENT RATE}$

$\frac{\sigma}{d} \sim 10^{-2}, d \sim 10 \mu\text{m}$

Examples: 1) 2 ATOMS separated by $1 \mu\text{m}$, $\sigma = \sqrt{\frac{\hbar}{m v}} \sim 10^{-7} \text{m}$

$g = 2\pi \cdot 10^{23} \frac{10^{-52}}{10^{-6}} 10^{-2} \sim 10^{-25} \text{Hz}$

→ $\tau_{ent} \sim 10^{25} \text{sec.} \sim 10^8 \text{Tyr.}$

g

2) 2 PLUMBER MASSES, $d \sim 100 \mu\text{m}$, $\sigma \sim 1 \mu\text{m}$

$g = 2\pi \cdot 10^{23} \frac{10^{-16}}{10^{-4}} 10^{-10} \sim 10 \text{Hz}$

→ $\tau_{ent} \sim 100 \mu\text{s}$

$$\rightarrow g = \frac{G}{\hbar} \frac{m^2}{d} \frac{\sigma^2}{d^2}$$

$> \Gamma_{dec.}$
 \equiv ENTANGLEMENT RATE

$$\frac{\sigma}{d} \sim 10^{-2}, \quad d \sim 10 \mu\text{m}$$

Examples: 1) 2 ATOMS separated by $1 \mu\text{m}$, $\sigma = \sqrt{\frac{\hbar}{m \omega}} \sim 10^{-7} \text{ m}$

$$g = 2\pi \cdot 10^{23} \frac{10^{-52}}{10^{-6}} 10^{-2} \sim 10^{-27} \text{ Hz}$$

$$\rightarrow \tau_{ent} \sim 10^{27} \text{ sec.} \sim 10^8 \text{ Tian.}$$

$$g \gg 1$$

$$\omega^2 > \frac{\hbar}{G} d \left(\frac{d}{\sigma}\right)^2$$

$$10^{-23} 10^{-5} 10^4$$

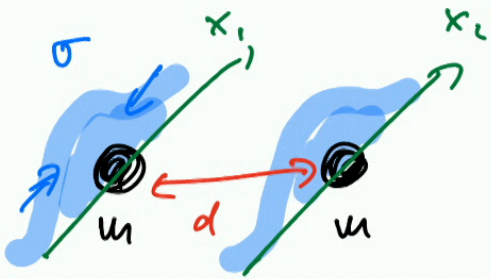
2) 2 PLUMMER MASSES, $d \sim 100 \mu\text{m}$, $\sigma \sim 1 \mu\text{m}$

$$g = 2\pi \cdot 10^{23} \frac{10^{-16}}{10^{-4}} 10^{-10} \sim 10 \text{ Hz}$$

$$\rightarrow \tau_{ent} \sim 100 \mu\text{s}$$

$$\sim (10^{-12})^2$$

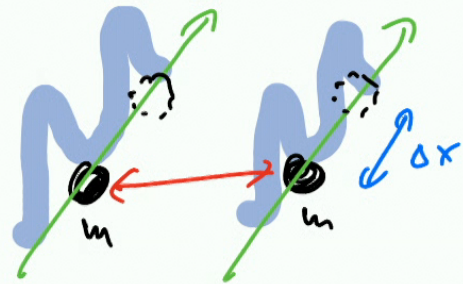
$$\Rightarrow f > \frac{\omega}{d^3} = \frac{10^{-12}}{10^{-18}} \sim 10^3 \frac{\text{kg}}{\text{m}^3}$$



$$H_{\text{int}} \sim G \frac{m^2}{r} = G \frac{m^2}{\sqrt{d^2 + (x_1 - x_2)^2}}$$

$$\approx G \frac{m^2}{d} \left(1 - \frac{(x_1 - x_2)^2}{2d^2} + \dots \right)$$

$$= G \frac{m^2}{d} \left(1 - \frac{x_1^2}{2d^2} - \frac{x_2^2}{2d^2} + \frac{x_1 x_2}{d^2} + \dots \right)$$

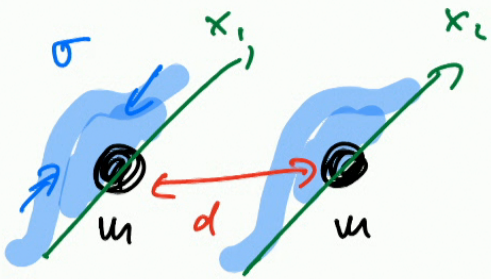


$$\rightarrow H_{\text{int}} = G \frac{m^2}{d^3} x_1 x_2 \quad x_i = \sigma (a_i^{\dagger} + a_i)$$

$$= G \frac{m^2}{d^3} \sigma^2 (a_i^{\dagger} + a_i)(\hat{J}_i + \hat{J}_i^{\dagger})$$

$$= \frac{1}{2} g (a_i^{\dagger} + a_i)(\hat{J}_i + \hat{J}_i^{\dagger})$$

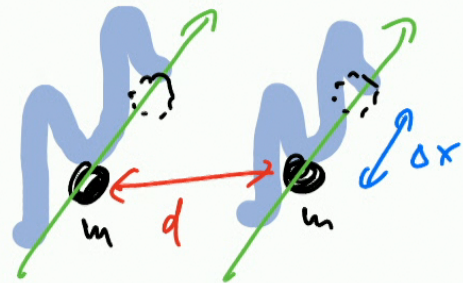
$$\rightarrow g = \frac{1}{2} G \frac{m^2}{d} \frac{\sigma^2}{d^2} \equiv \text{ENTANGLEMENT RATE}$$



$$H_{\text{int}} \sim G \frac{m^2}{r} = G \frac{m^2}{\sqrt{d^2 + (x_1 - x_2)^2}}$$

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$$\rightarrow H_{\text{int}} = G \frac{m^2}{d^3} x_1 x_2 \quad x_i = \sigma (a_i^{\dagger} + a_i)$$

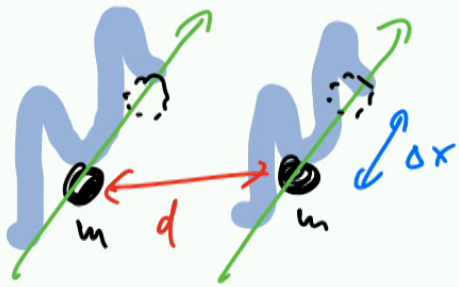
$$\varphi = \frac{1}{\hbar} \int m \phi dt$$

$$\phi = G \frac{m}{r}$$

$$= G \frac{m^2}{d^3} \sigma^2 (a_i^{\dagger} + a_i)(j^{\dagger} + j)$$

$$= \frac{1}{\hbar} g (a_i^{\dagger} + a_i)(j^{\dagger} + j)$$

$$\rightarrow g = \frac{1}{\hbar} G \frac{m^2}{d} \frac{\sigma^2}{d^2} \equiv \text{ENTANGLEMENT RATE}$$



$$|\psi_0\rangle = (|x_L\rangle_1 + |x_R\rangle_1) \otimes (|x_L\rangle_2 + |x_R\rangle_2)$$

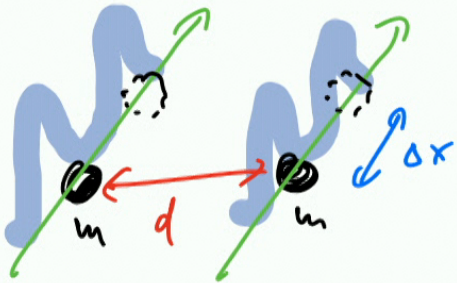
↓
t

$$\varphi = \frac{1}{\hbar} \int m \phi dt$$

$$\phi = G \frac{m}{r}$$

$$r_{\perp} = d$$

$$r_{\parallel} = \sqrt{d^2 + \Delta x^2}$$



$$\varphi = \frac{1}{\hbar} \int m \phi \, dt$$

$$\phi = G \frac{1}{r}$$

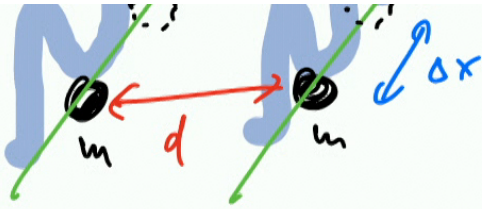
$$r_{\perp} = d$$

$$r_{\parallel} = \sqrt{d^2 + \Delta x^2}$$

$$|\psi_0\rangle = (|x_L\rangle_1 + |x_R\rangle_1) \otimes (|x_L\rangle_2 + |x_R\rangle_2)$$

$$e^{i\varphi_{\perp}} \left(|x_L\rangle_1 \left(|x_L\rangle_2 + e^{i\varphi} |x_R\rangle_2 \right) + |x_R\rangle_1 \left(|x_R\rangle_2 + e^{-i\varphi} |x_L\rangle_2 \right) \right)$$

$$\langle \tilde{x}_L | \tilde{x}_R \rangle = 0 = |e^{i\varphi} + e^{-i\varphi}| \rightarrow \cos \Delta\varphi$$



$$\varphi = \frac{1}{\hbar} \int m \phi dt$$

$$\phi = G \frac{m}{r}$$

$$r_{\pm} = d$$

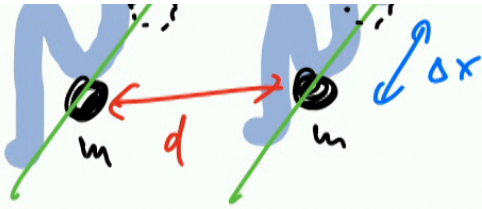
$$r_{\pm} = \sqrt{d^2 + \Delta x^2}$$

$$\Delta \varphi = \frac{1}{\hbar} \left(G \frac{m^2}{d} \left(1 + \frac{\Delta x^2}{d^2} - 1 \right) \right) \cdot t$$

$$g_{\text{end}} = \frac{d}{dt} \Delta \varphi = \frac{G}{\hbar} \frac{m^2}{d} \left(\frac{\Delta x}{d} \right)^2$$

$$e^{i\varphi_1} \left(|x_L\rangle_1 \left(|x_L\rangle_2 + e^{i\varphi} |x_R\rangle_2 \right) + |x_R\rangle_1 \left(|x_R\rangle_2 + e^{i\varphi} |x_L\rangle_2 \right) \right)$$

$$\langle \tilde{x}_L | \tilde{x}_R \rangle = 0 = |e^{i\varphi} + e^{-i\varphi}| \rightarrow \cos \Delta \varphi$$



$$\varphi = \frac{1}{\hbar} \int m \phi dt$$

$$\phi = G \frac{m}{r}$$

$$r_{\pm} = d$$

$$r_{\pm} = \sqrt{d^2 + \Delta x^2}$$

$$\Delta \varphi = \frac{1}{\hbar} \left(G \frac{m^2}{d} \left(1 + \frac{\Delta x^2}{d^2} - 1 \right) \right) \cdot t$$

$$g_{\text{end}} = \frac{d}{dt} \Delta \varphi = \frac{G}{\hbar} \frac{m^2}{d} \left(\frac{\Delta x}{d} \right)^2 \Rightarrow \text{mod'ante } \frac{m^2 \Delta x^2}{d^3} \cdot \epsilon_{\text{oh}}$$

$$e^{i\varphi_1} \left(|x_L\rangle_1 \left(|x_L\rangle_2 + e^{i\varphi} |x_R\rangle_2 \right) + |x_R\rangle_1 \left(|x_R\rangle_2 + e^{i\varphi} |x_L\rangle_2 \right) \right)$$

$$\langle \tilde{x}_L | \tilde{x}_R \rangle = 0 = |e^{i\varphi} + e^{-i\varphi}| \rightarrow \cos \Delta \varphi$$

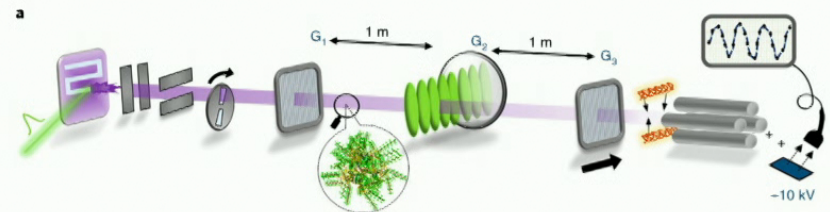
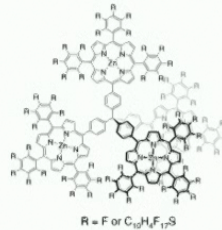
Current experiments involving delocalization of massive objects

Superposition of single systems (beyond atoms & neutrons)

$$\Delta x \cdot M \approx 10^{-29} \text{ [u}\cdot\text{g]}$$

Macro-molecules

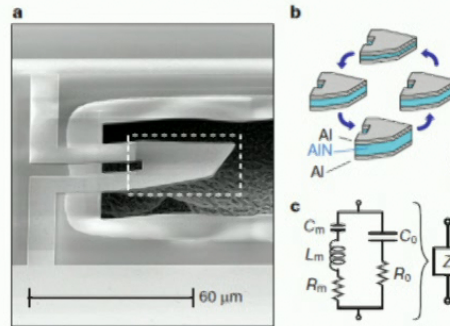
matter-wave interference; 2,000 atoms; $M = 25,000 \text{ a.m.u.} = 4\text{e-}23 \text{ kg}$; particle size $D=5\text{nm}$; superposition size $\Delta x > 500 \text{ nm}$
 Y. Y. Fein *et al.*, *Nat. Phys.* **15**, 1242–1245 (2019)



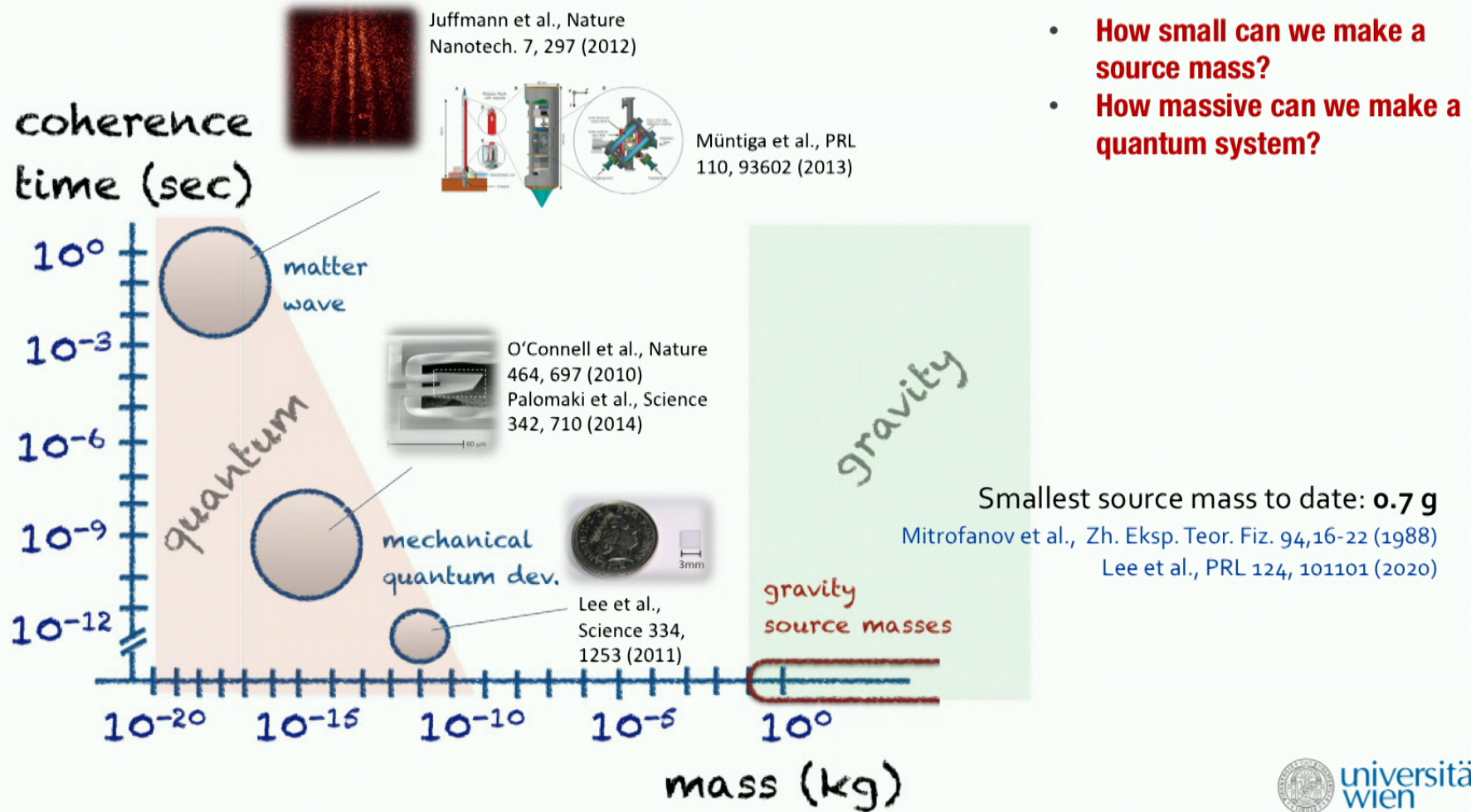
Solid-state mechanical oscillators

Ramsey interference 0+1 between ground state (0) and single-phonon (1) Fock state (6 GHz acoustic mode); $1\text{e}13 \text{ atoms}$, $M = 8\text{e-}13 \text{ kg}$, SiO₂/AlN slab (60umx10umx700nm), superposition size $\Delta x = 2\text{e-}16 \text{ m}$ (thickness oscillation)

A. D. O'Connell *et al.*, *Nature* 464, 697–703 (2010)

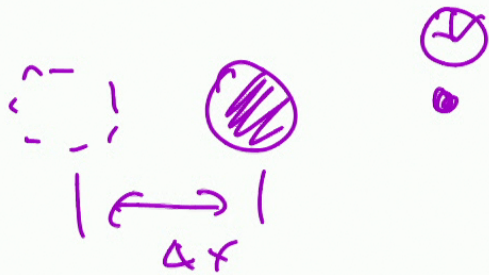


Reality check: quantum systems as gravitational source masses?

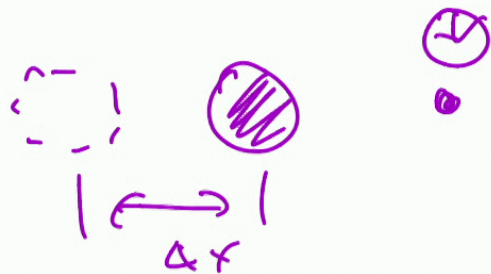


- How small can we make a source mass?
- How massive can we make a quantum system?

$$\left(|x_L\rangle_m |z_L\rangle_g |e_L\rangle_\oplus + |x_R\rangle_m |z_R\rangle_g |e_R\rangle_\oplus \right)$$



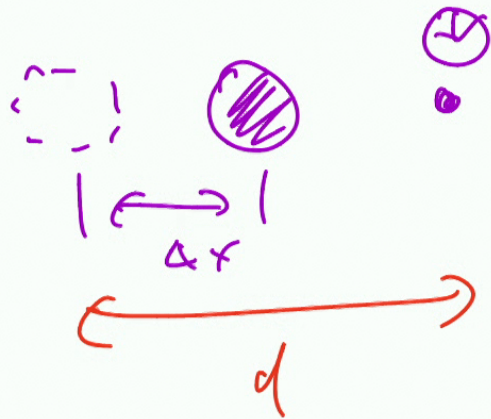
$$\left(|x_L\rangle_m |x_L\rangle_s |\psi_L\rangle_\oplus + |x_R\rangle_m |x_R\rangle_s |\psi_R\rangle_\oplus \right) = |\psi\rangle_+$$



$$\begin{aligned} \rho_m &= \text{Tr}_\oplus |\psi\rangle\langle\psi| \\ &= |\tilde{\psi}_L\rangle\langle\tilde{\psi}_L| + |\tilde{\psi}_R\rangle\langle\tilde{\psi}_R| \end{aligned}$$

$$\left(|x_L\rangle_m |x_L\rangle_s |\tilde{\psi}_L\rangle_\oplus + |x_R\rangle_m |x_R\rangle_s |\tilde{\psi}_R\rangle_\oplus \right) = |\psi\rangle_+$$

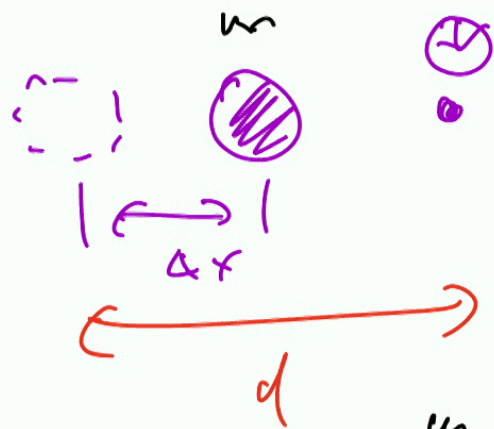
$$m = M_p, \quad d \sim 10^{-7} \text{ m}, \quad \frac{\Delta x}{d} \sim 10^{-2}$$



$$\begin{aligned} S_m &= \text{Tr}_\oplus |\psi\rangle\langle\psi| \\ &= |\tilde{\psi}_L\rangle\langle\tilde{\psi}_L| + |\tilde{\psi}_R\rangle\langle\tilde{\psi}_R| \end{aligned}$$

$$\left(|x_L\rangle_m |x_L\rangle_s |\psi_L\rangle_\oplus + |x_R\rangle_m |x_R\rangle_s |\psi_R\rangle_\oplus \right) = |\psi\rangle_+$$

$$m = M_p, \quad d \sim 10^{-7} \text{ m}, \quad \frac{\Delta x}{d} \sim 10^{-2}, \quad \frac{\Delta v}{v} = \frac{\Delta U}{c^2}$$



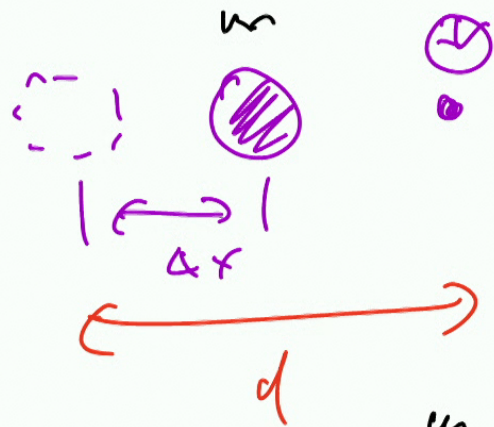
$$S_m = \text{Tr}_\oplus |\psi\rangle\langle\psi|$$

$$= |\tilde{\psi}_L\rangle\langle\tilde{\psi}_L| + |\tilde{\psi}_R\rangle\langle\tilde{\psi}_R|$$

$$\Delta U = -G \frac{m}{d} + G \frac{m}{d - \Delta x}$$

$$\left(|x_L\rangle_m |a_L\rangle_s |\psi_L\rangle_\oplus + |x_R\rangle_m |a_R\rangle_s |\psi_R\rangle_\oplus \right) = |\psi\rangle_+$$

$$m = M_p, \quad d \sim 10^{-7} \text{ m}, \quad \frac{\Delta x}{d} \sim 10^{-2}, \quad \frac{\Delta U}{V} = \frac{\Delta U}{c^2}$$



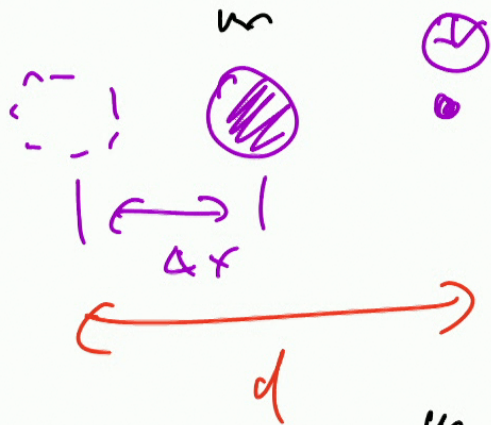
$$S_m = \text{Tr}_\oplus |\psi\rangle\langle\psi|$$

$$= |\tilde{\psi}_L\rangle\langle\tilde{\psi}_L| + |\tilde{\psi}_R\rangle\langle\tilde{\psi}_R|$$

$$\Delta U = -G \frac{m}{d} + G \frac{m}{d - \Delta x} = G \frac{m}{d} \frac{\Delta x}{d}$$

$$\left(|x_L\rangle_m |a_L\rangle_g |\psi_L\rangle_\oplus + |x_R\rangle_m |a_R\rangle_g |\psi_R\rangle_\oplus \right) = |\Psi\rangle_+$$

$$m = M_p, \quad d \sim 10^{-7} \text{ m}, \quad \frac{\Delta x}{d} \sim 10^{-2} \quad \frac{\Delta U}{U} = \frac{\Delta U}{c^2} = \frac{10^{-16}}{10^{16}} \sim \underline{\underline{10^{-32}}}$$

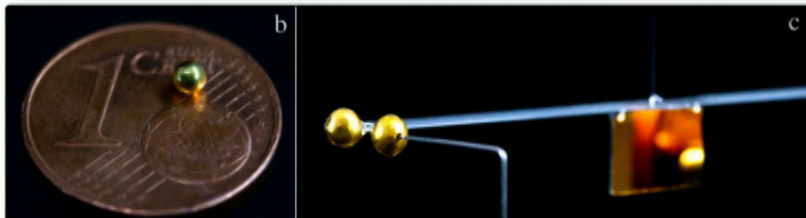
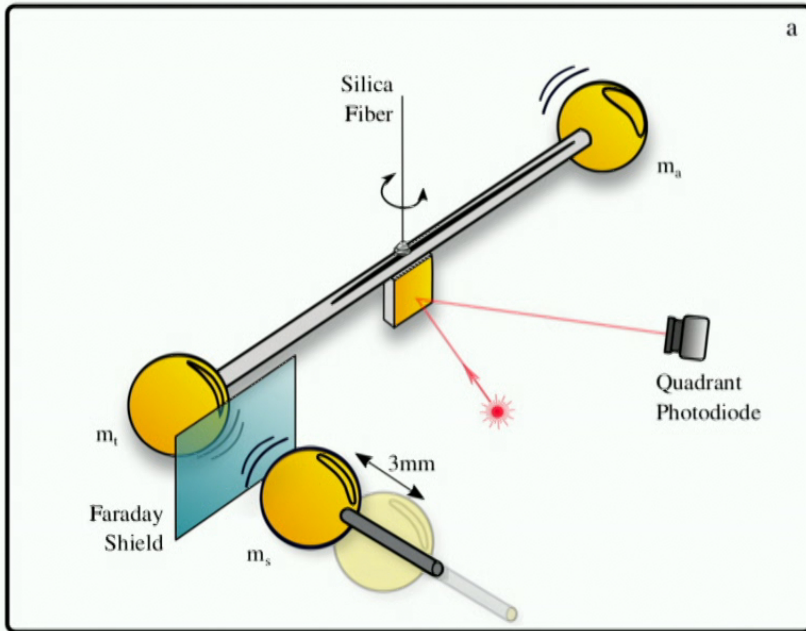


$$S_m = \text{Tr}_\oplus |\Psi\rangle\langle\Psi|$$

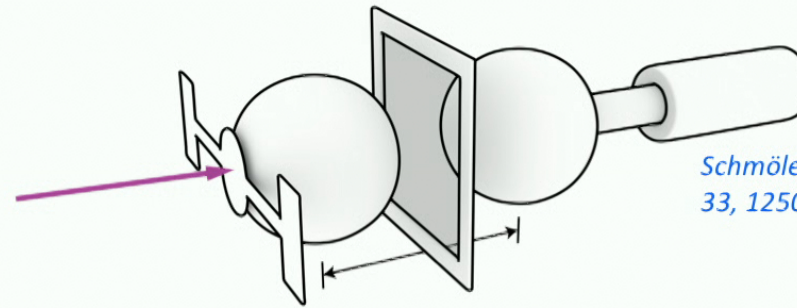
$$= |\tilde{\psi}_L\rangle\langle\tilde{\psi}_L| + |\tilde{\psi}_R\rangle\langle\tilde{\psi}_R|$$

$$\Delta U = -G \frac{m}{d} + G \frac{m}{d - \Delta x} = G \frac{m}{d} \frac{\Delta x}{d} = 6 \cdot 10^{-11} \frac{10^{-31}}{10^{-7}} 10^{-2} \sim 10^{-16}$$

How small can we go? The idea:



Westphal et al., Nature 591, 225 (2021)



Schmölle et al., Class. Quant. Grav. 33, 125031 (2016)

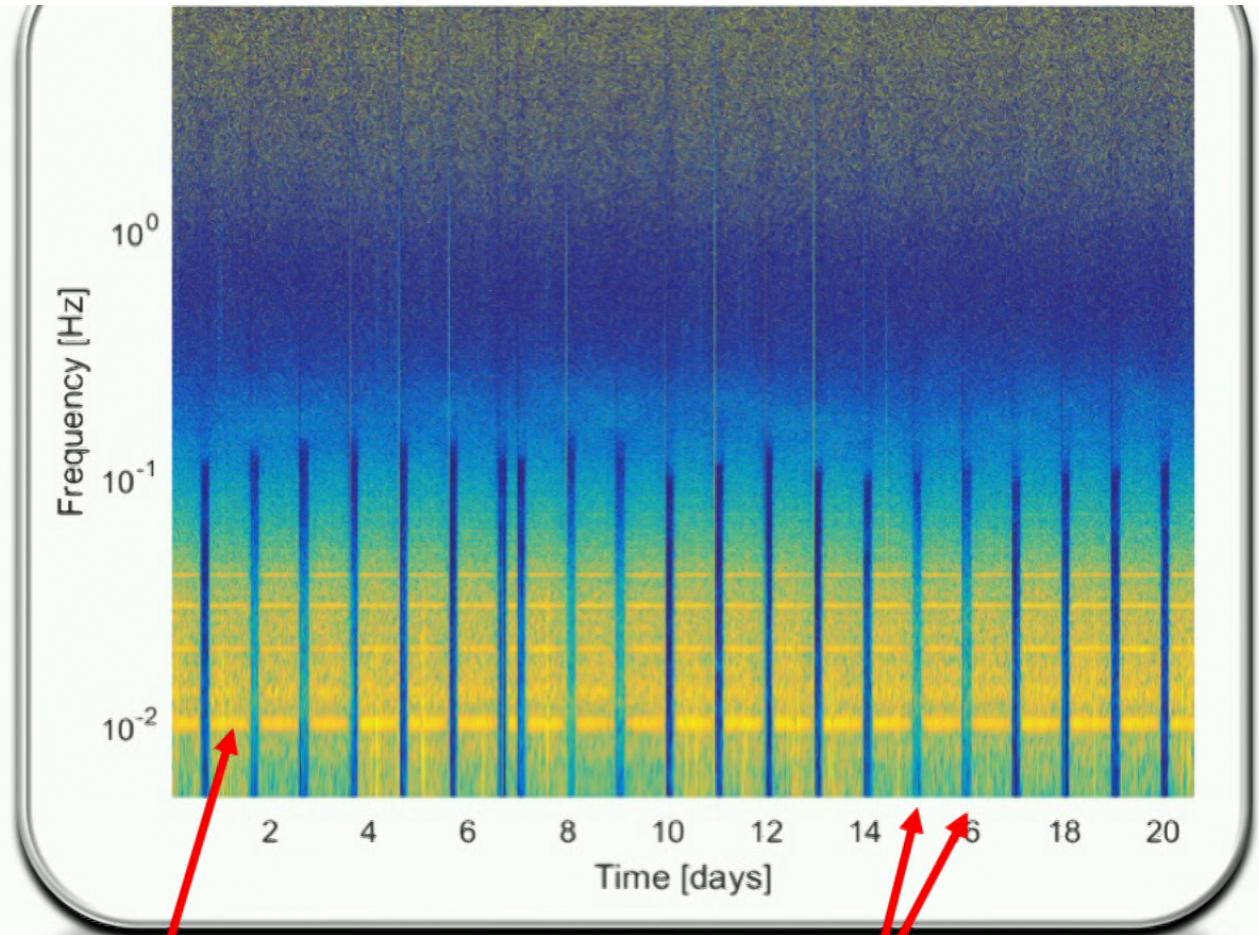
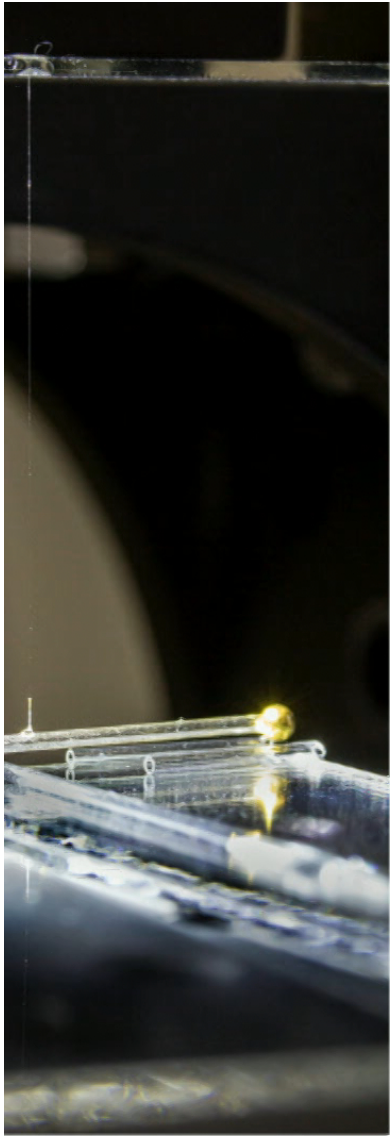
- periodic modulation of a source mass @ f_{mod}
- generates a test mass acceleration @ $n \times f_{\text{mod}}$ ($n=1,2, \dots$)
- fundamental limit: thermal noise of test mass oscillator

$$d(t) = d_0 + d_m \cos \Omega t$$

$$a(t) = G \frac{m}{d(t)^2} = G \frac{m}{d_0^2} \left(1 - 2 \frac{d_m}{d_0} \cos \Omega t + 3 \left(\frac{d_m}{d_0} \right)^2 \cos^2 \Omega t + \dots \right)$$

$$\left. \begin{aligned} m_s &\sim 9 \cdot 10^{-5} \text{ kg} \\ d_0 &\sim 4 \text{ mm} \\ d_m &\sim 1.6 \text{ mm} \end{aligned} \right\} a \sim 3 \cdot 10^{-10} \frac{\text{m}}{\text{s}^2}$$

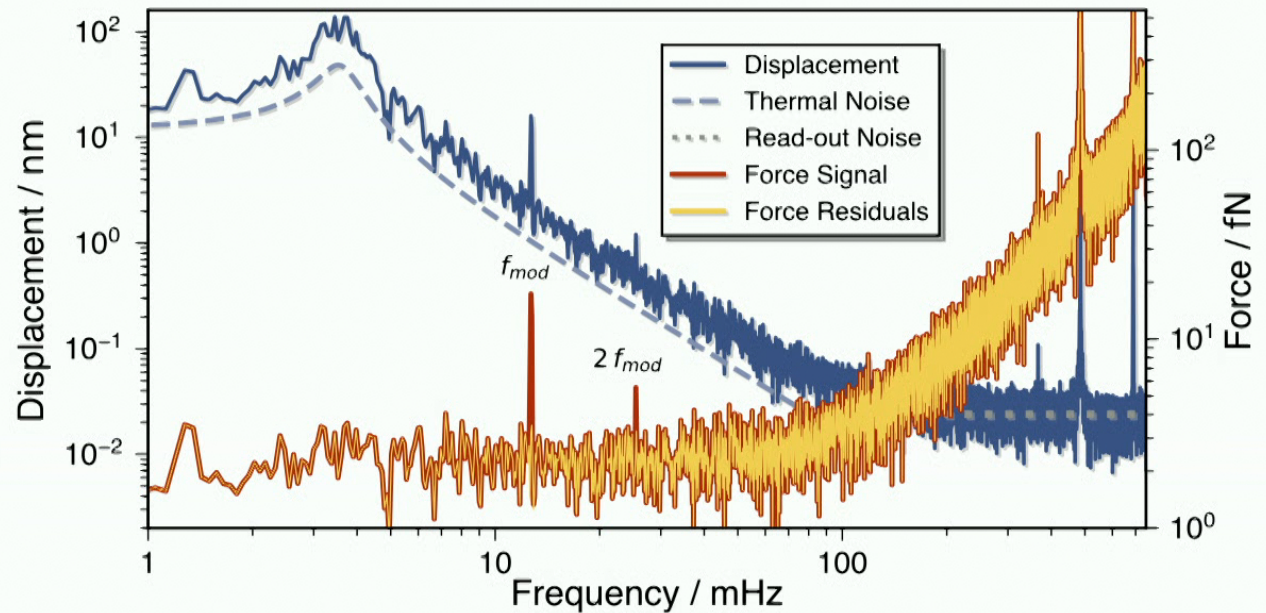
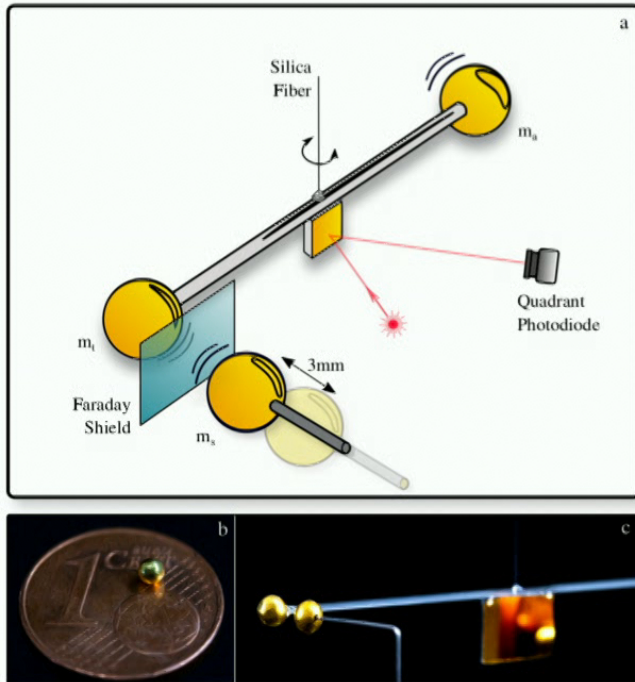




10 mHz = traffic light cycle
Nussdorfer Str @

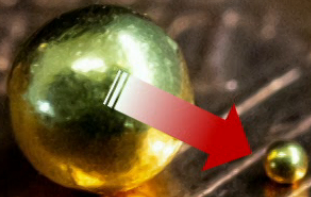
Friday & Saturday night
= party night

Results

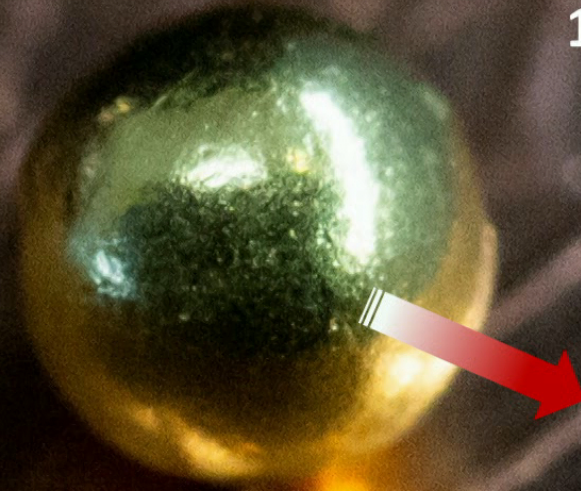


- We observe a linear and quadratic acceleration modulation (at f_{mod} and $2f_{mod}$) produced from a **90mg source mass**
- We resolve an **acceleration modulation of $3e-10 \text{ m/s}^2$** with an accuracy of 10% and a **precision of 1% ($3e-12 \text{ m/s}^2$)**
- The observed coupling deviates from the CODATA value for Newton's constant by 9%, which is covered in the known systematic uncertainties of our experiment (i.e. interaction is >90% gravitational)

Next steps: going smaller in mass...



Planck mass:
 $1e18$ atoms



... by going underground



Quantum controlling levitated solid-state objects

Combining LARGE MASSES with LONG COHERENCE TIMES and FULL MANIPULATION



+



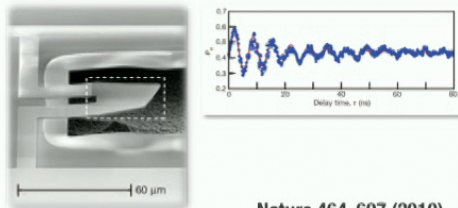
=



Solid-state mechanical quantum devices (clamped):

$10^{10} - 10^{16}$ atoms

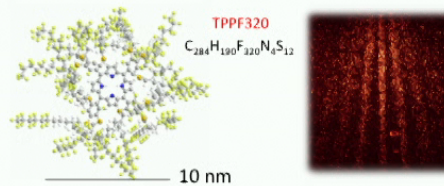
Coherence time
 $10^{-12} - 10^{-8}$ sec



Matter-wave interferometry (free-fall):

$10^0 - 10^4$ atoms

Coherence time
 $10^{-3} - 10^0$ sec



Nature Nanotech. 7, 297 (2012)

Levitated (opto-)mechanics

- Quantum control of a trapped solid state object $\gg 10^{10}$ atoms
- Long coherence times (up to seconds)
- Exceptional force sensitivity
- Externally engineerable (and controllable) arbitrary potential landscape

recent review:

Gonzalez-Ballesterio et al.,
Science **374**, 168 (2021)

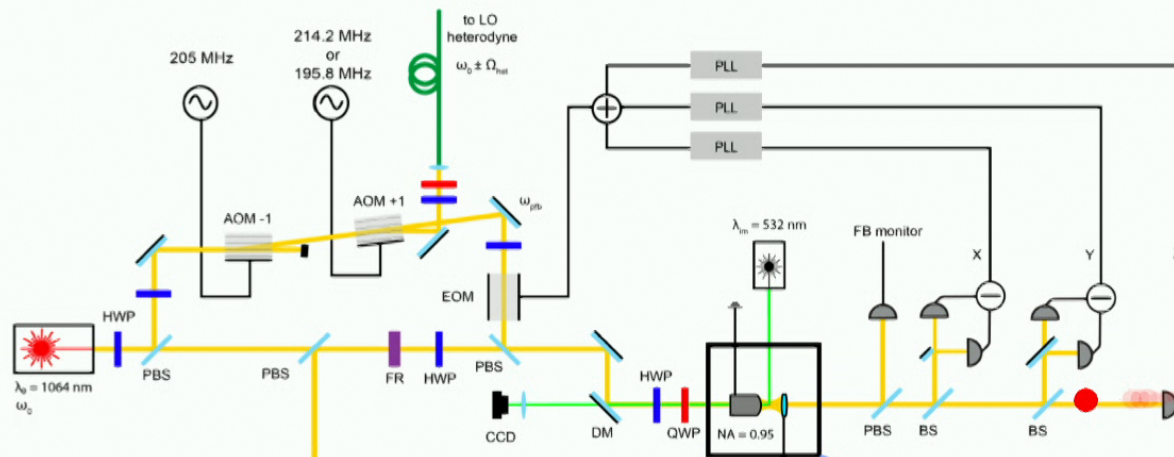
Keeping a particle in high vacuum: feedback control

Magrini et al., Nature 595, 373 (2021)

see also

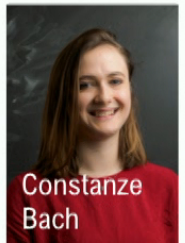
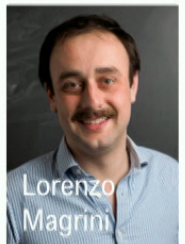
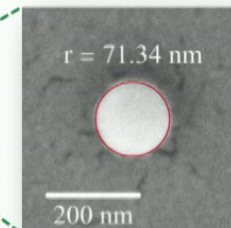
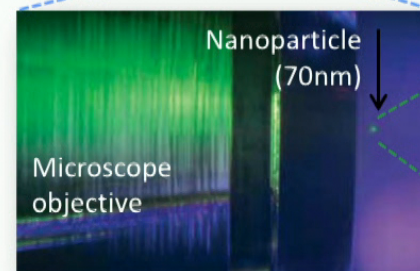
F. Tebbenjohanns et al., PRL 124, 013603 (2020)

F. Tebbenjohanns et al., Nature 595, 378 (2021)



3d parametric feedback stabilizes the particle @ 9e-9 mbar

see also Gieseler et al., Phys. Rev. Lett. 109, 103603 (2012)



Lorenzo Magrini, Constanze Bach
P. Rosenzweig, A. Deutschmann, A. Kugi (TU Wien)

Image of a 150nm glass sphere in its quantum ground state of motion at a room temperature environment

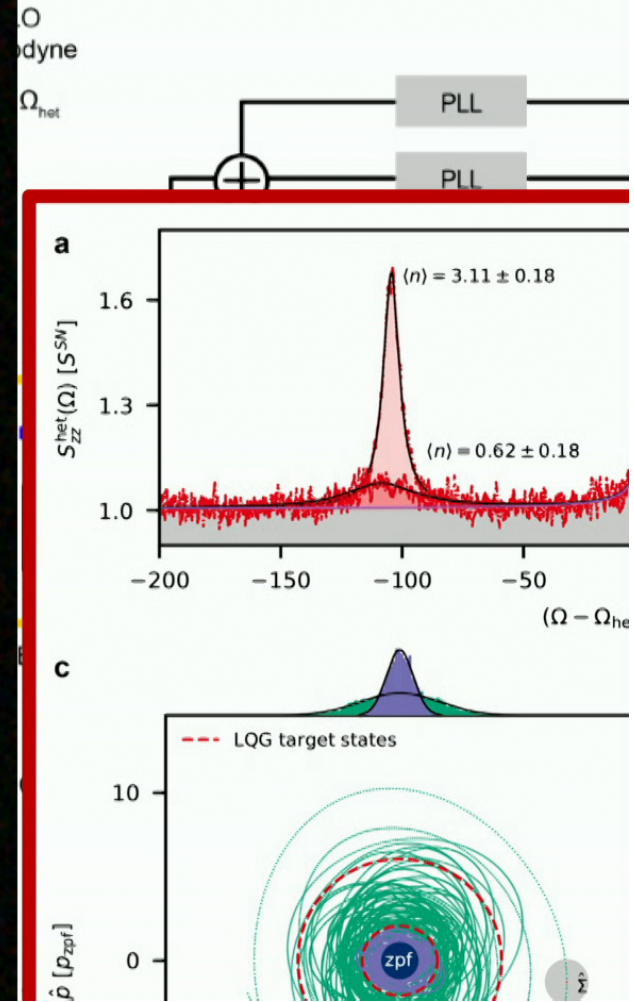


Image of a 150nm glass sphere in its quantum ground state of motion at a room temperature environment



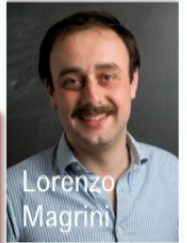
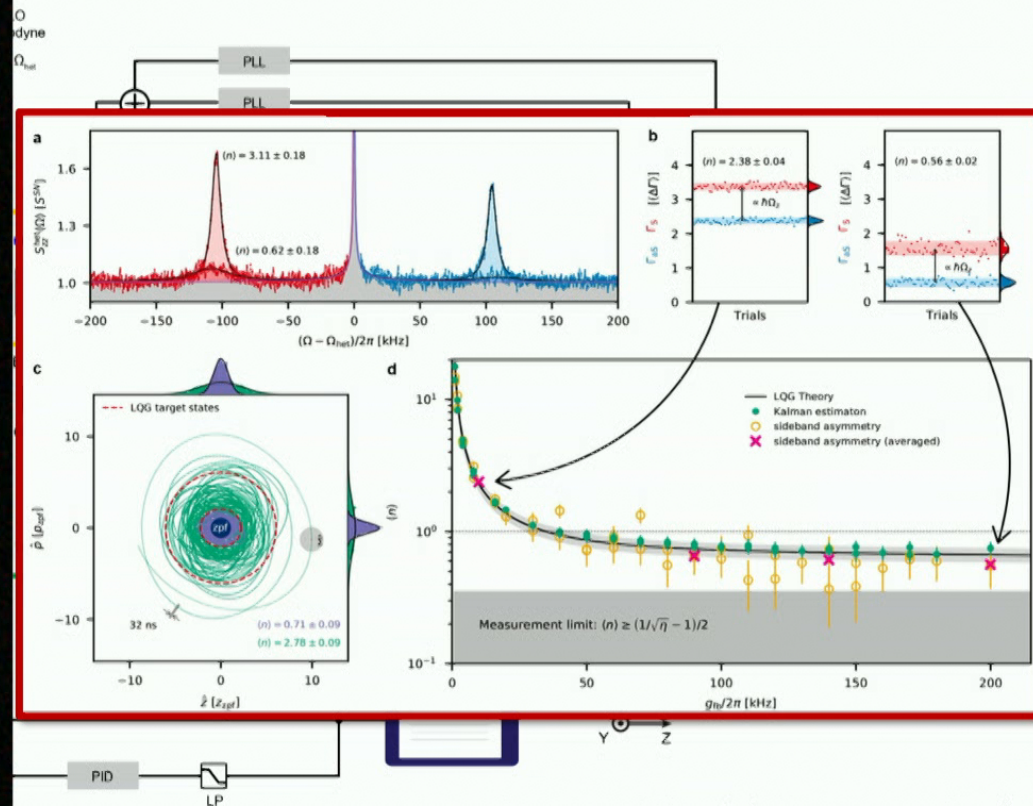
ground-state cooling

Magrini et al., Nature 595, 373 (2021)

see also

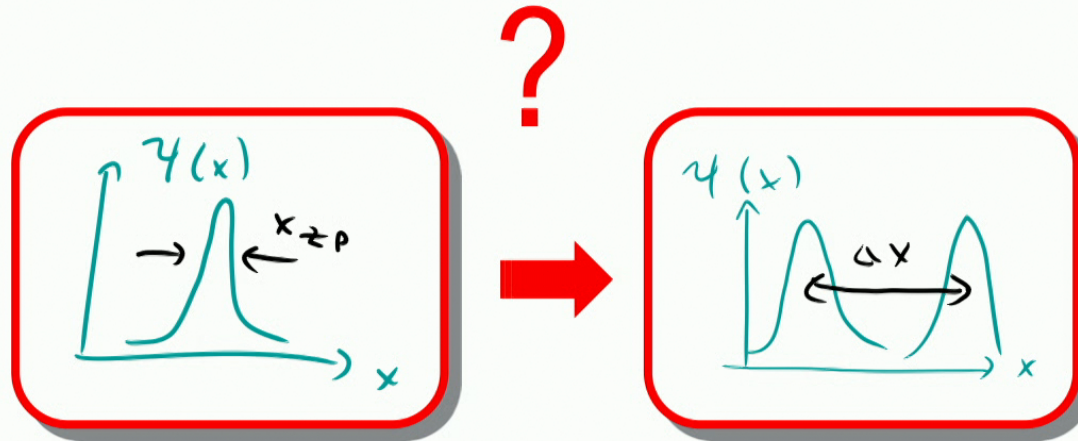
F. Tebbenjohanns et al., PRL 124, 013603 (2020)

F. Tebbenjohanns et al., Nature 595, 378 (2021)



Lorenzo Magrini, Constanze Bach, Nikolai Kiesel
P. Rosenzweig, A. Deutschmann, A. Kugi (TU Wien)

Towards „large“ quantum superposition states?



How to prepare “macroscopic superpositions”?

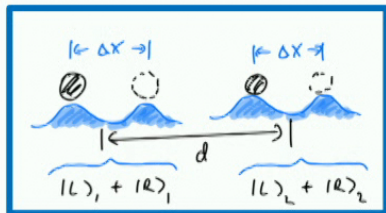
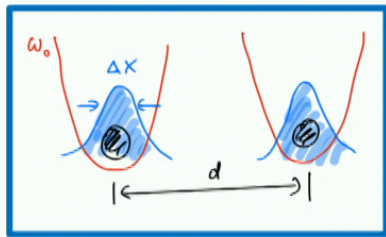
Find at least 2 possible ways of realizing them.

The decoherence challenge...

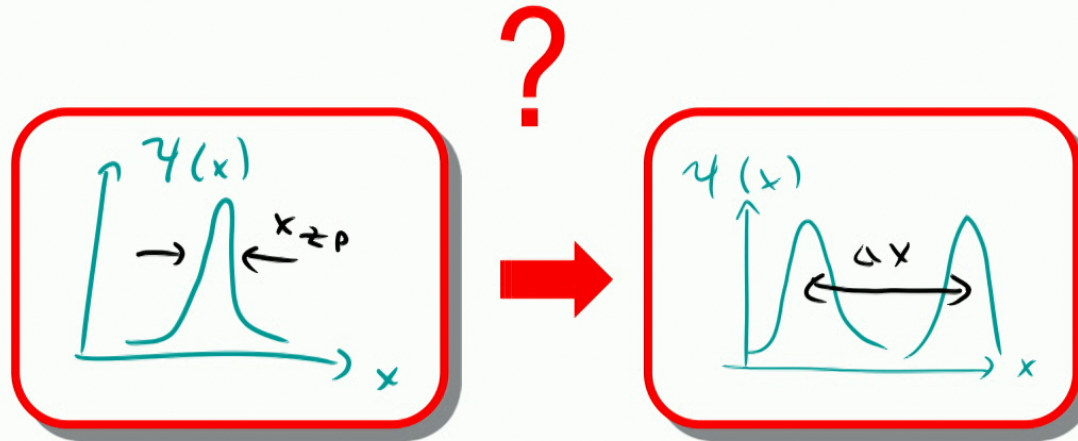
What are the main sources of decoherence?

How to estimate them?

How to avoid them?



Towards „large“ quantum superposition states?



How to prepare “macroscopic superpositions”?

Find at least 2 possible ways of realizing them.