

Title: Rydberg atoms in Bose-Einstein condensed environments: cold bubble chambers and mesoscopic entanglement

Speakers: Sebastian W ster

Collection: Cold Atom Molecule Interactions (CATMIN)

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Abstract: "S. Tiwari, S. Rammohan, A. Mishra, A. Pendse, A. K. Chauhan, R. Nath, F. Engel, M. Wagner, R. Schmidt, F. Meinert, A. Eisfeld and S. W ster

Indian Institute of Science Education and Research, Bhopal India Palack  University, Olomouc, Czech Republic Indian Institute of Science Education and Research, Pune, India Max Planck Institute for the Physics of Complex Systems, Dresden, Germany Universit t Stuttgart, Germany Max-Planck-Institute of Quantum Optics and MCQST, Garching, Germany

Rydberg Atoms in highly excited electronic states with n=30-200 can be excited within BoseEinstein condensates (BECs), and while lifetimes are shorter than in vacuum [1,2], they live long enough to cause a response of the BEC mean field [3]. During this, thousands of ground-state atoms are present within the Rydberg orbit, allowing the study of atoms moving within atoms [4]. We present beyond-mean field models of the joint Rydberg-BEC dynamics, showing how either can be used to probe the other.

For multiple Rydberg atoms in a single electronic state, we show that the phase coherence of thecondensate allows the tracking of mobile Rydberg impurities akin to the function of bubblechambers in particle physics [5]. For a single Rydberg atom with multiple electronic states, weprovide spectral densities of the BEC as a decohering environment [6], and show that the BECcan image a signature of the entangling evolution that causes Rydberg q-bit decoherence [7] or serve as non-Markovian environment for quantum simulations.

[1] Schlagmu ller et al. PRX 6 (2016) 031020.

[2] Kanungo et al. PRA 102 (2020) 063317.

[3] Balewski et al. Nature 502 (2013) 664.

[4] Tiwari et al. arXiv:2111.05031 (2021)

[5] Tiwari et al. PRA 99 (2019) 043616.

[6] Rammohan et al. PRA 103 (2021) 063307.

[7] Rammohan et al. PRA(Letters) 104 (2021) L060202."

# Rydberg atoms in Bose-Einstein condensed environments



*from bubble chambers to paradigmatic open quantum systems*

Sebastian Wüster

Department of Physics  
Indian Institute of Science Education and Research Bhopal, India

*MPG-Partnergroup with:*

Max-Planck Institute for the Physics of Complex Systems, Dresden, Germany





MAX-PLANCK-GESELLSCHAFT

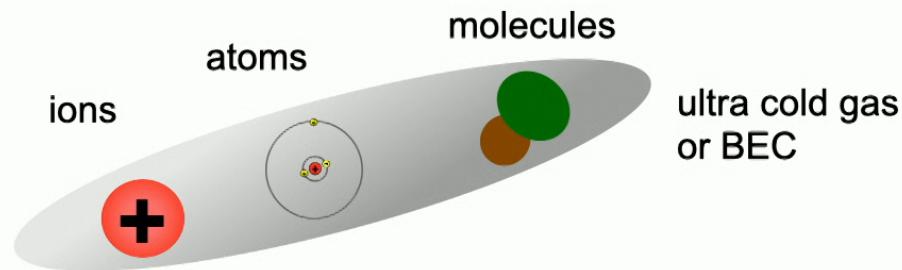
# Geography

**Ultra-cold and Rydberg atoms in Bhopal:  
Max-Planck-IISER Partner Group**





# Cold Atom Molecule Interactions





# Cold Atom Molecule Interactions

## Ion-Rydberg molecules

Zuber et al.  
Nature **5** 453 (2022)

## Rydberg polarons

F. Carmago et al.,  
PRL **120** (2018) 083401.

## Ultra-long-range Mols

C.H. Greene et al.,  
PRL **85** (2000) 2458.

## Cold ion-atom scattering

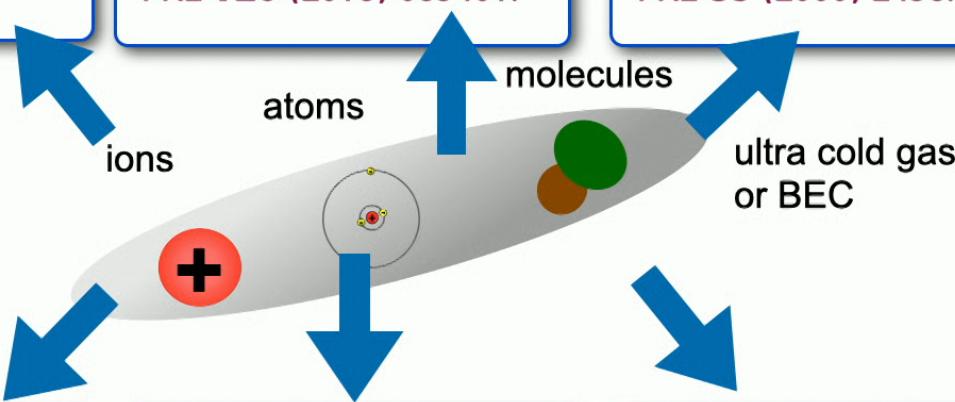
e.g. S. Dutta and S. Rangwala  
PRA(R) **97** (2018) 041401

## Polaron formation

e.g. C. Kohstall et al.  
Nature **485** (2012) 615

## Angular momentum dynamics

e.g. Schmidt and Lemeshko  
PRX **6** (2016) 011012



# Origin of decoherence

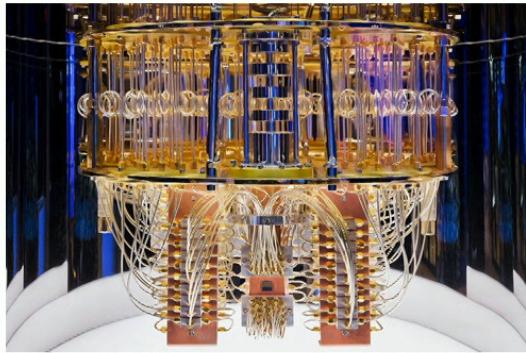
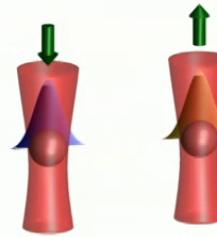


image: **IBM**



$$\text{Qubit: } |\phi\rangle = c_0|0\rangle + c_1|1\rangle$$

# Origin of decoherence

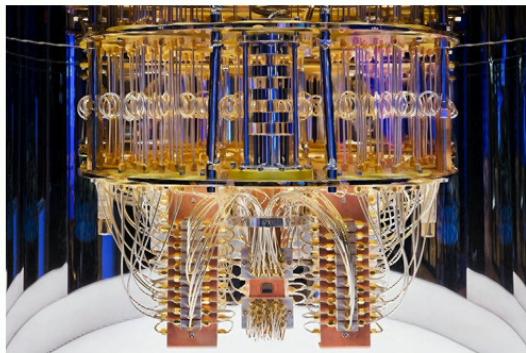
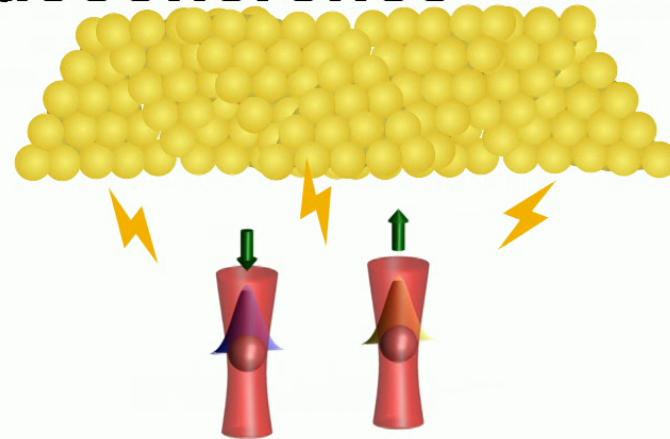


image: IBM



$$\text{Qubit: } |\phi\rangle = c_0|0\rangle + c_1|1\rangle$$

$$\text{Qubit + Env: } |\chi\rangle = (c_0|0\rangle + c_1|1\rangle) \otimes |\Psi_{ini}\rangle$$

$$\text{Interaction: } \hat{H}_{int} = |0\rangle\langle 0| \otimes \hat{h}_0 + |1\rangle\langle 1| \otimes \hat{h}_1$$

# Origin of decoherence

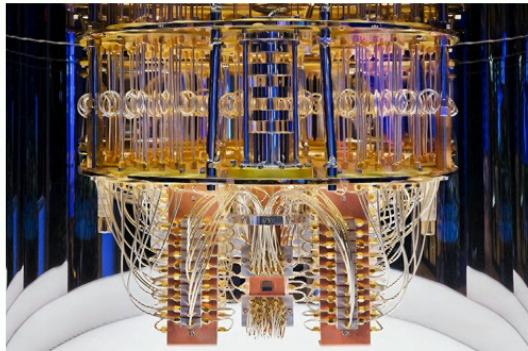
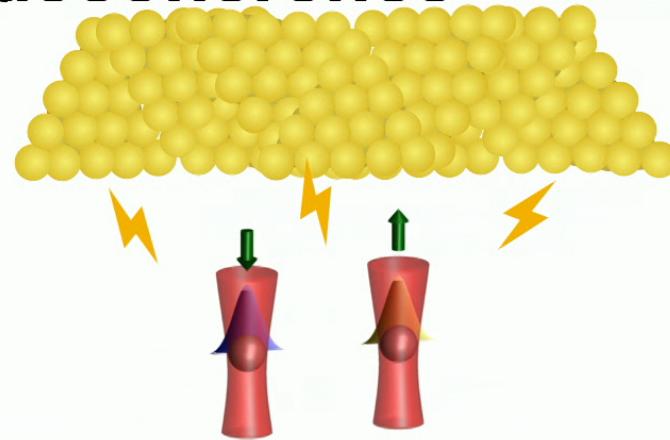


image: IBM



$$\text{Qubit: } |\phi\rangle = c_0|0\rangle + c_1|1\rangle$$

$$\text{Qubit + Env: } |\chi\rangle = (c_0|0\rangle + c_1|1\rangle) \otimes |\Psi_{ini}\rangle$$

$$\text{Interaction: } \hat{H}_{int} = |0\rangle\langle 0| \otimes \hat{h}_0 + |1\rangle\langle 1| \otimes \hat{h}_1$$

$$\text{Entanglement} \\ \Rightarrow \text{decoherence: } |\chi\rangle = c_0|0\rangle \otimes |\Psi_0\rangle + c_1|1\rangle \otimes |\Psi_1\rangle$$

# Origin of decoherence

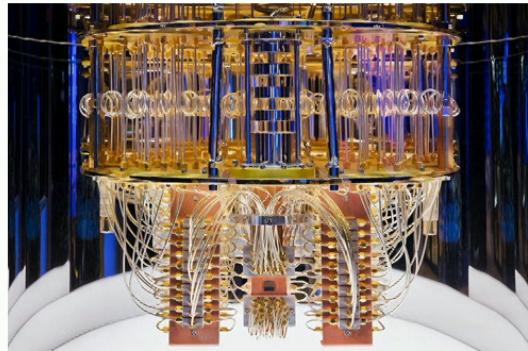
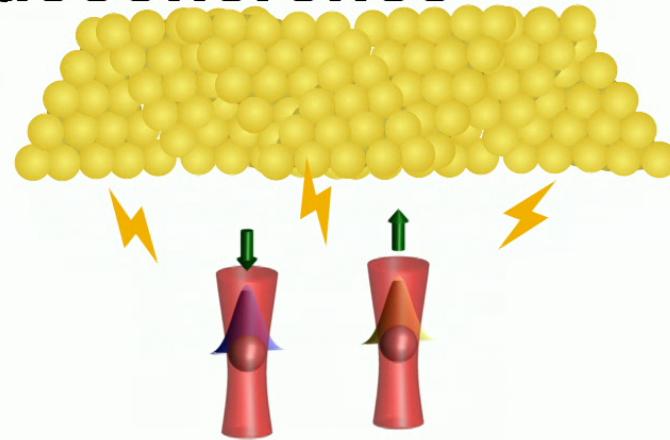


image: **IBM**

Qubit+ Environment:  $|\chi\rangle = c_0|0\rangle \otimes |\Psi_0\rangle + c_1|1\rangle \otimes |\Psi_1\rangle$



# Origin of decoherence

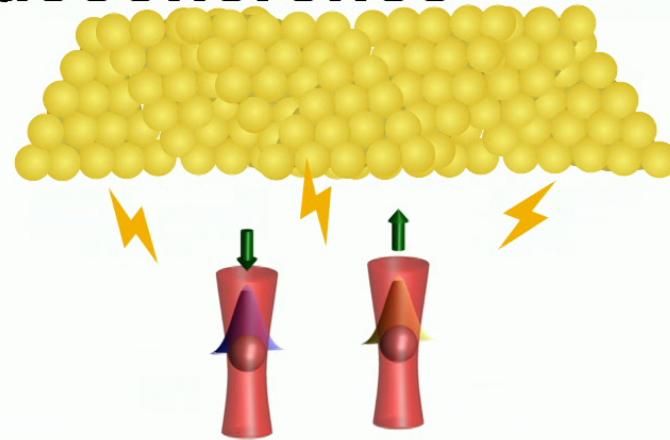
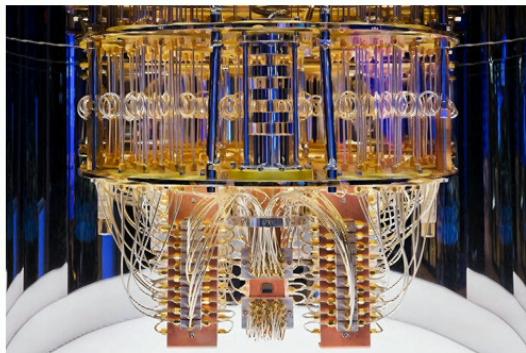


image: IBM

Qubit+ Environment:  $|\chi\rangle = c_0|0\rangle \otimes |\Psi_0\rangle + c_1|1\rangle \otimes |\Psi_1\rangle$

Reduced Qubit:

$$\hat{\rho} = \begin{pmatrix} |c_0|^2 & c_1 c_0^* \langle \Psi_0 | \Psi_1 \rangle \\ c_1^* c_0 \langle \Psi_1 | \Psi_0 \rangle & |c_1|^2 \end{pmatrix}$$



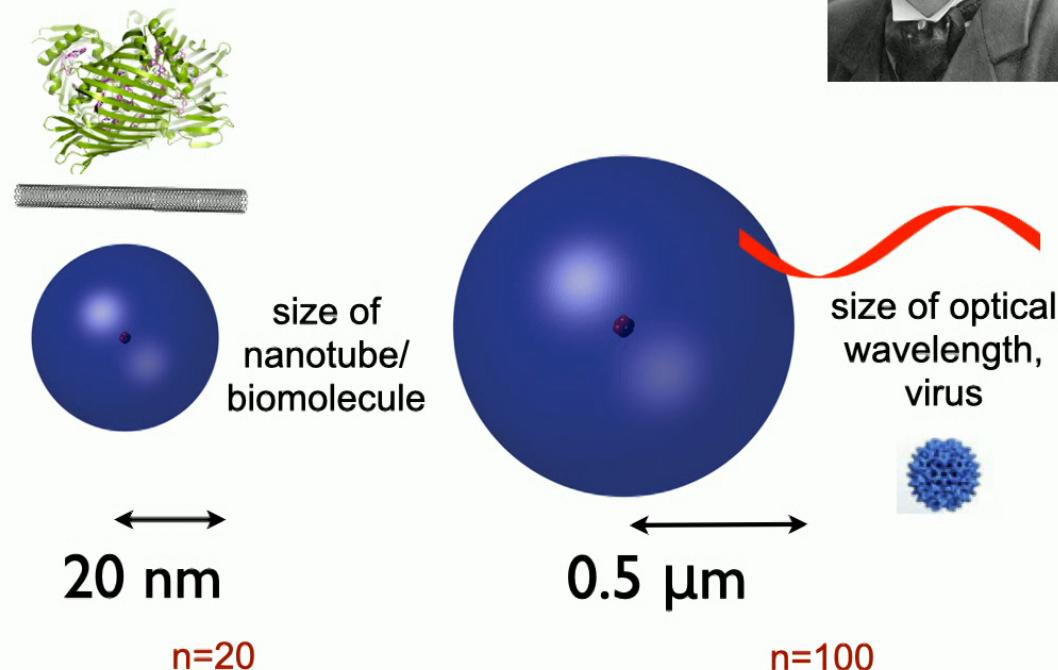
# Rydberg atoms



## What are Rydberg atoms?

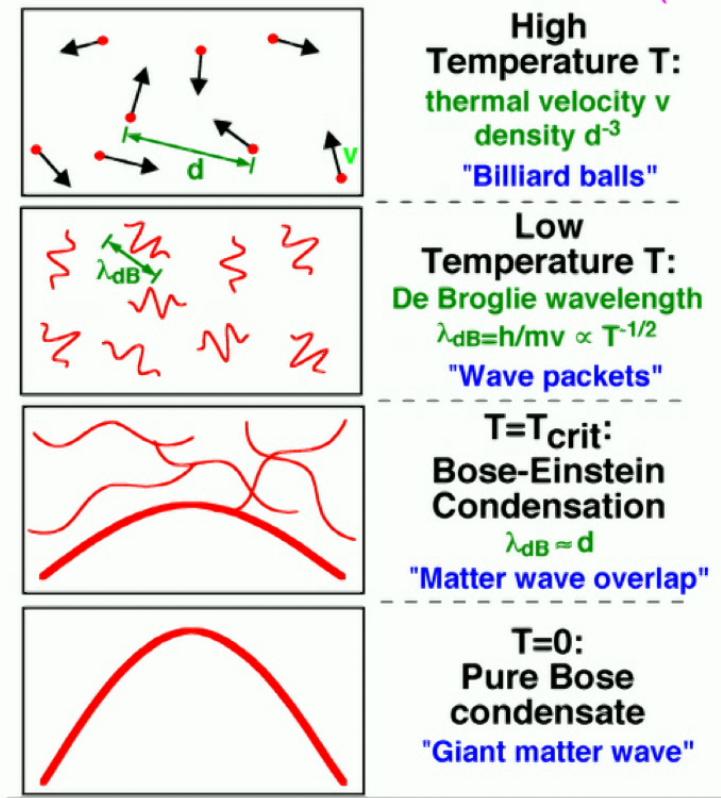
Very high principal quantum number  $n \gg 5$ ,  
these are HUGE atoms...

Johannes Rydberg





# Bose-Einstein Condensates



$$\Psi(\mathbf{R}_1, \dots, \mathbf{R}_N) \rightarrow \otimes_{k=1}^N \phi(\mathbf{R}_k)$$
$$\hat{\Psi}(\mathbf{R}) \rightarrow \langle \hat{\Psi}(\mathbf{R}) \rangle = \phi(\mathbf{R})$$

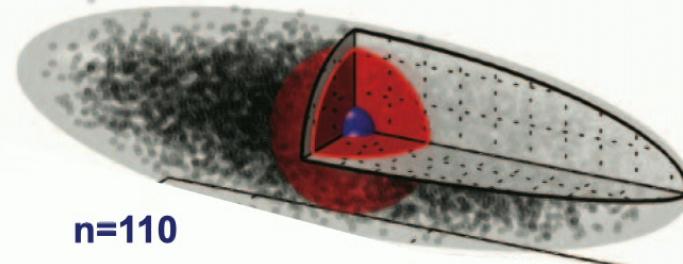
from: Durfee and Ketterle  
Opt. Express **2** (1998) 299.



# Rydberg atoms in BEC



*Single large Rydberg atom in BEC*



$n=110$

$n=202$

J. Balewski *et al.*,  
Nature **502** (2013) 664.

- Naturally excite Rydberg states in an ultracold gas or even BEC
- Extreme atoms in an extreme environment

also:

Celistrino-Texeira *et al.*,  
Phys. Rev. Lett. **115** (2015) 013001.  
F. Carmargo *et al.*,  
Phys. Rev. Lett. **120** (2018) 083401.  
M. Mizoguchi *et al.*,  
Phys. Rev. Lett. **124** (2020) 253201.



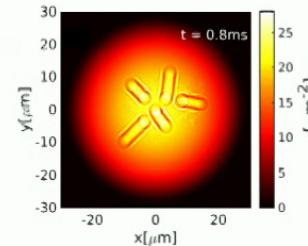
# Outline



## (I) Phase imprinting by Rydberg atoms

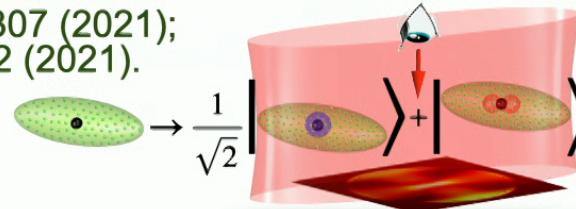
S. Tiwari and S. Wüster,  
PRA **99** 043616 (2019)

S. Tiwari *et al.* NJP **24** 073005 (2022)



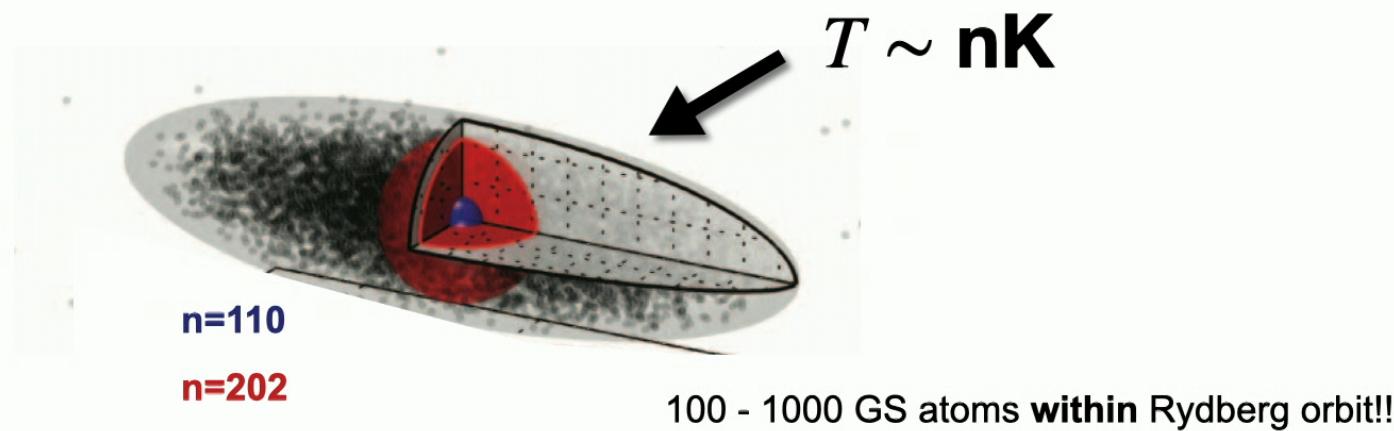
## (II) Decoherence of Rydberg qubits in a BEC

S. Rammohan *et al.* PRA **103**, 063307 (2021);  
Phys. Rev. A (Letters) **104**, L060202 (2021).

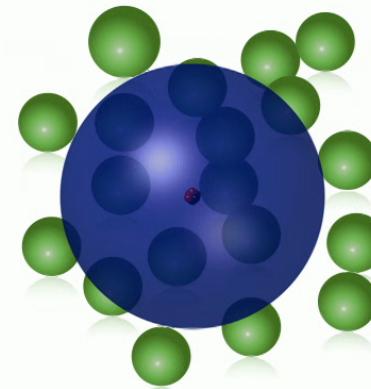




# Rydbergs in BEC

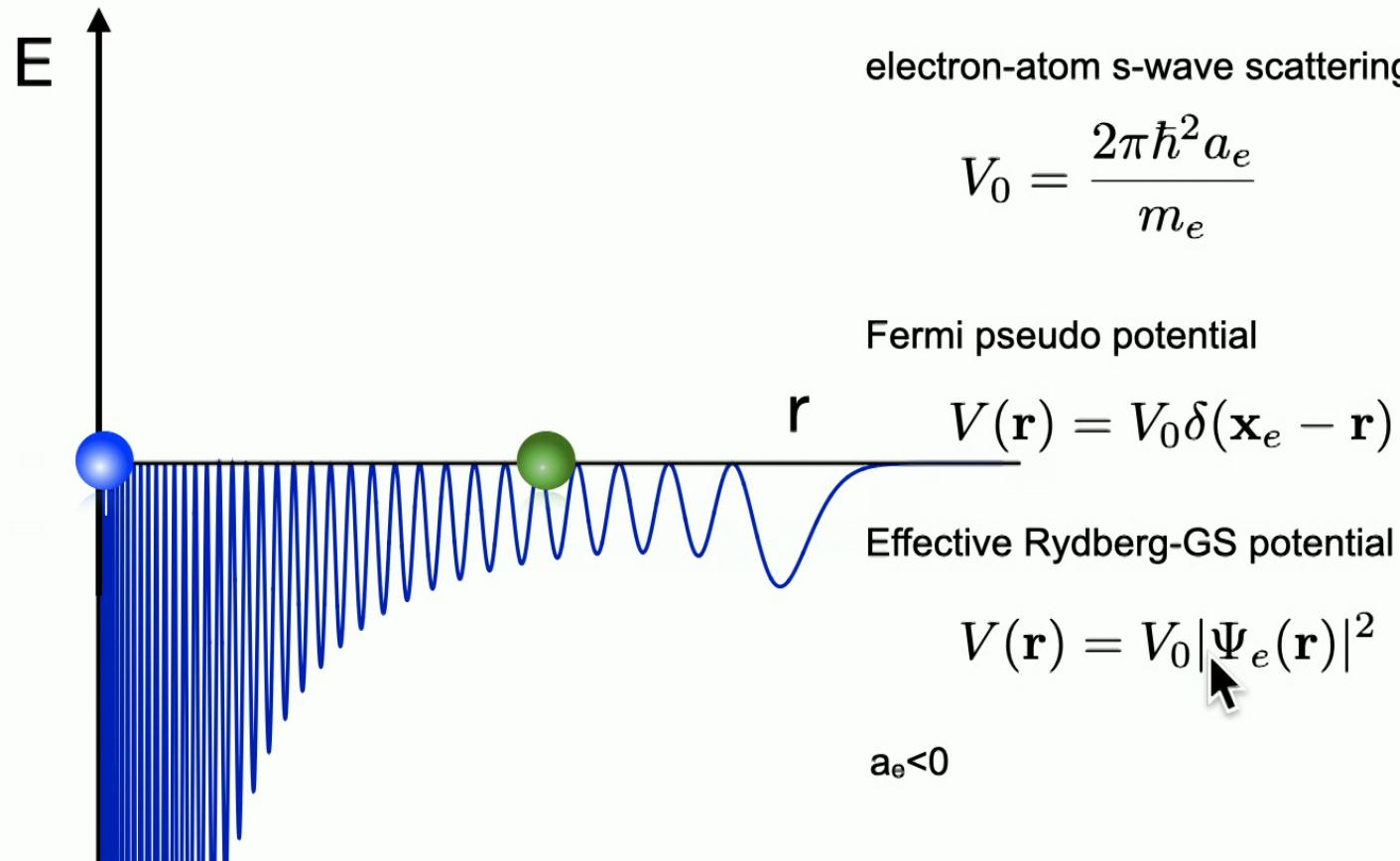


J. Balewski et al.,  
Nature **502** (2013) 664.





# Interactions with ground-state atoms





# Phase imprinting

## GPE

$$i\hbar \frac{\partial}{\partial t} \phi(\mathbf{R}) = \left( -\frac{\hbar^2}{2m} \nabla^2 + g |\phi(\mathbf{R})|^2 + V_0 |\psi(\mathbf{R} - \mathbf{x}_{imp})|^2 \right) \phi(\mathbf{R})$$

Phase imprinting:  $\phi(\mathbf{R}, t) = \phi(\mathbf{R}, 0) e^{-iV_0|\psi(\mathbf{R} - \mathbf{x}_{imp})|^2 t/\hbar}$

see e.g.

Dobrek et al. PRA **60** (1999) R3381,  
R. Mukherjee et al. PRL **115** (2015) 040401.





# Phase imprinting



## GPE

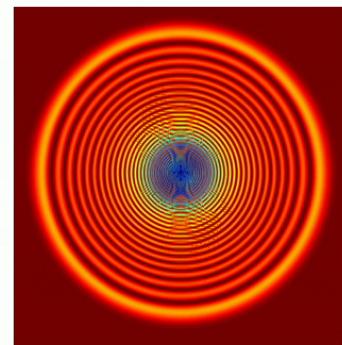
$$i\hbar \frac{\partial}{\partial t} \phi(\mathbf{R}) = \left( -\frac{\hbar^2}{2m} \nabla^2 + g |\phi(\mathbf{R})|^2 + V_0 |\psi(\mathbf{R} - \mathbf{x}_{imp})|^2 \right) \phi(\mathbf{R})$$

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see e.g.

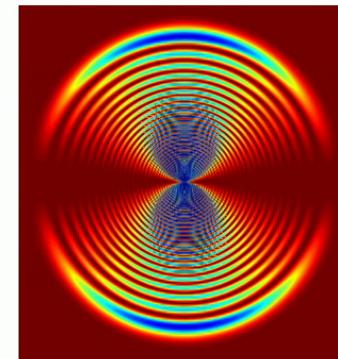
Dobrek et al. PRA **60** (1999) R3381,  
R. Mukherjee et al. PRL **115** (2015) 040401.

$|0\rangle \rightarrow |ns\rangle$



see also: Karpiuk et al.  
NJP **17** (2015) 053046

$|1\rangle \rightarrow |np\rangle$



state-dependent  
Imprinted phase  
pattern

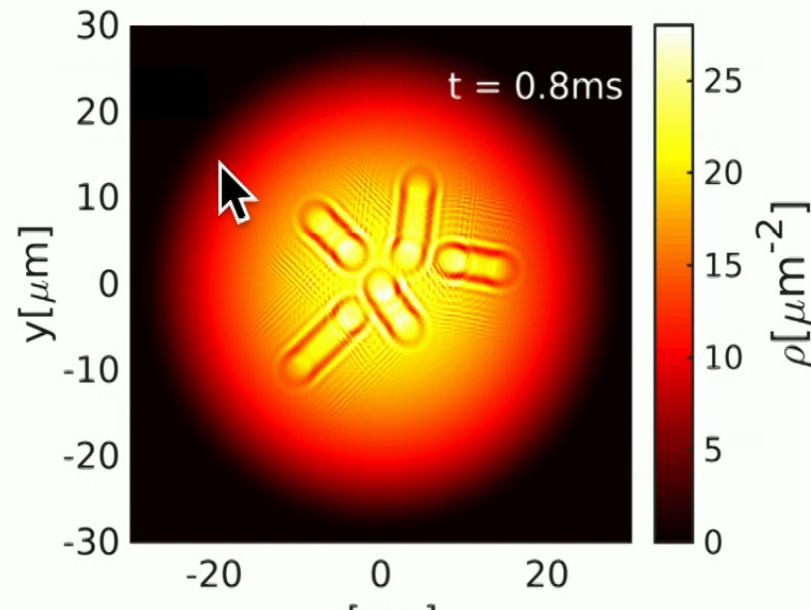
Density perturbation  
weak



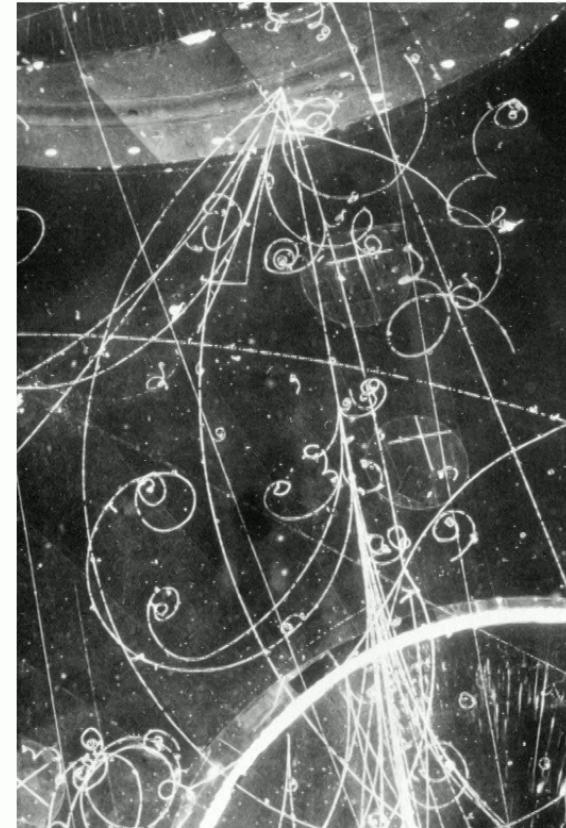
# Rydberg “Bubble chamber” (2D)



S. Tiwari and S. Wüster,  
PRA 99 043616 (2019)



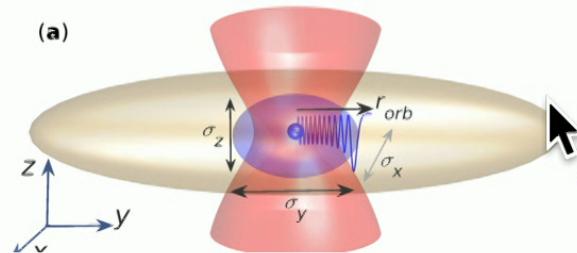
Phase imprint/ momentum kick  
converts into density changes later



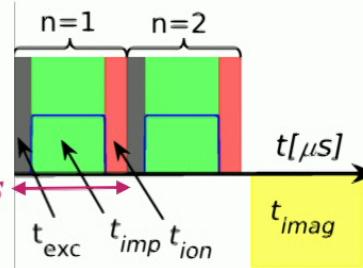


# Repeated excitation

(a)

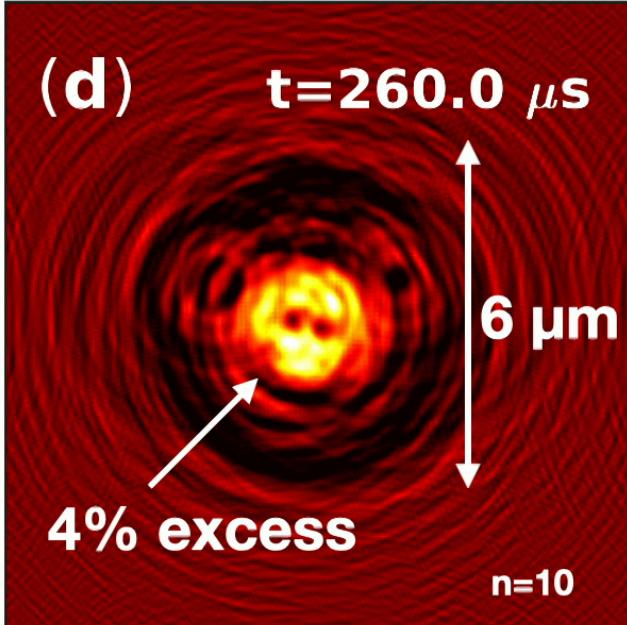


4.8  $\mu$ s



(d)

t = 260.0  $\mu$ s



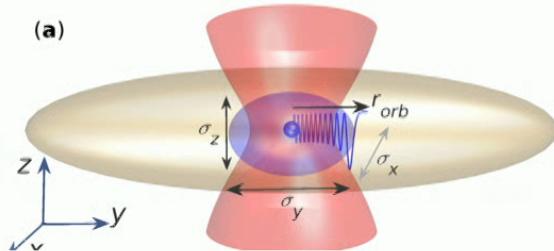
S. Tiwari, F. Engel, M. Wagner, R. Schmidt,  
F. Meinert and S. Wüster NJP **24** 073005 (2022)



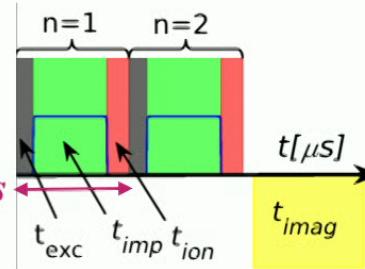
# Repeated excitation



(a)



4.8  $\mu$ s

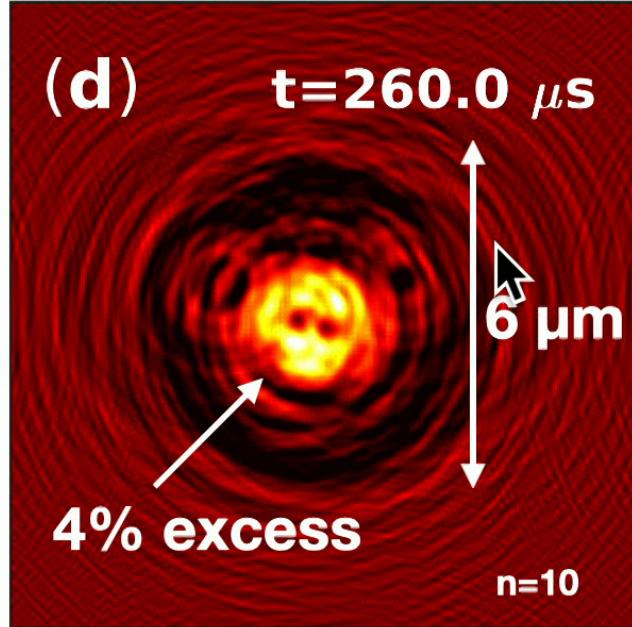


(d)

t=260.0  $\mu$ s

4% excess

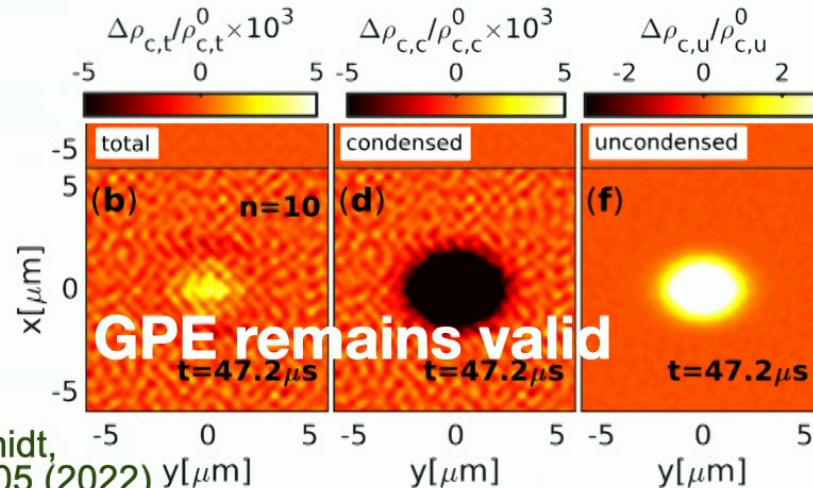
n=10



e.g. Blakie et al., (review)

Advances in Physics **57** 363 (2008)

Truncated Wigner simulation



S. Tiwari, F. Engel, M. Wagner, R. Schmidt,  
F. Meinert and S. Wüster NJP **24** 073005 (2022)

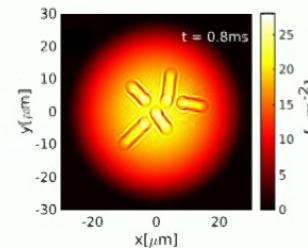


# Outline

## (I) Phase imprinting by Rydberg atoms

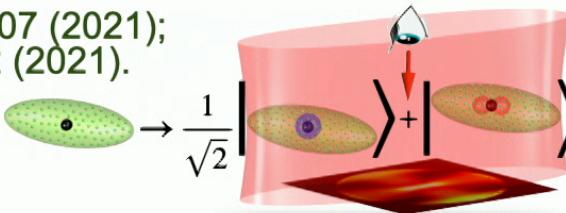
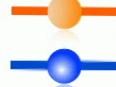
S. Tiwari and S. Wüster,  
PRA **99** 043616 (2019)

S. Tiwari *et al.* NJP **24** 073005 (2022)

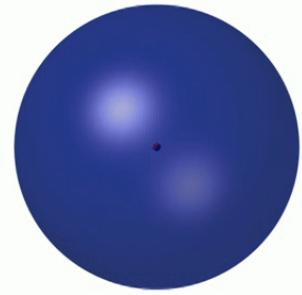
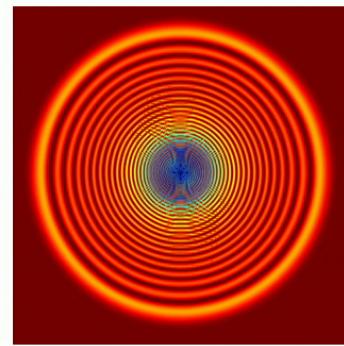


## (II) Decoherence of Rydberg qubits in a BEC

S. Rammohan *et al.* PRA **103**, 063307 (2021);  
Phys. Rev. A (Letters) **104**, L060202 (2021).

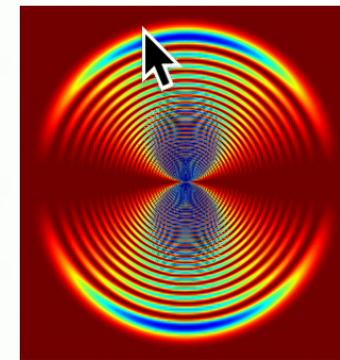


# Rydberg qubit

 $|ns\rangle$  $n=80, l=0$  $|\downarrow\rangle$  $|np\rangle$  $n=80, l=1$  $|\uparrow\rangle$ 

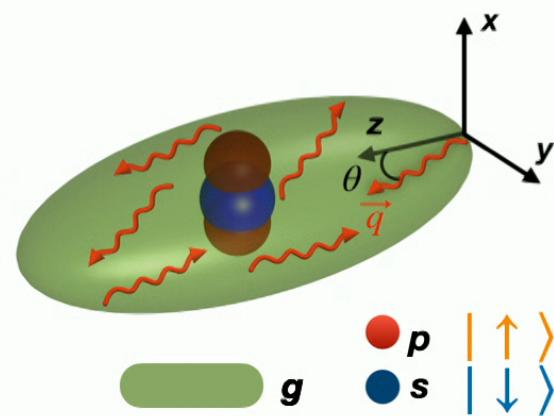
different  
Imprinted  
pattern

see also:  
Karpiuk et al.  
NJP **17** (2015) 053046





# Bogoliubov Spin-Boson model



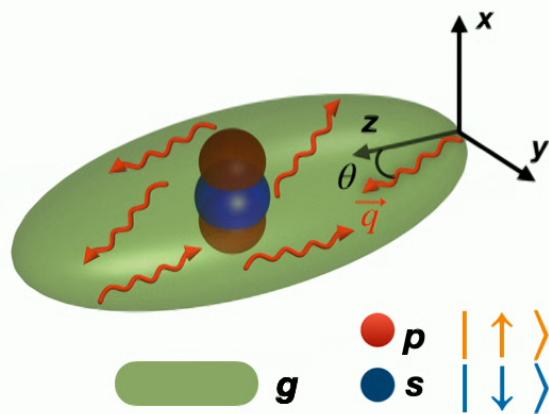
## Many-body Hamiltonian

$$\hat{H} = \sum_k \int d^3\mathbf{x} \left[ \hat{\Psi}_k^\dagger(\mathbf{x}) \left( -\frac{\hbar^2}{2m} \nabla^2 + E_k \right) \hat{\Psi}_k(\mathbf{x}) \right] \quad (1)$$
$$+ \frac{1}{2} \sum_{i,j,s} \int d^3\mathbf{y} \hat{\Psi}_k^\dagger(\mathbf{x}) \hat{\Psi}_i^\dagger(\mathbf{y}) U_{kij s}(\mathbf{x} - \mathbf{y}) \hat{\Psi}_j(\mathbf{y}) \hat{\Psi}_s(\mathbf{x}).$$

see also: Middelkamp et al. PRA **76** (2007) 022507.



# Bogoliubov Spin-Boson model



$g$     $p$     $s$     $\left| \begin{array}{c} \uparrow \\ \downarrow \end{array} \right\rangle$

## Many-body Hamiltonian

$$\hat{H} = \sum_k \int d^3\mathbf{x} \left[ \hat{\Psi}_k^\dagger(\mathbf{x}) \left( -\frac{\hbar^2}{2m} \nabla^2 + E_k \right) \hat{\Psi}_k(\mathbf{x}) \right. \\ \left. + \frac{1}{2} \sum_{i,j,s} \int d^3\mathbf{y} \hat{\Psi}_k^\dagger(\mathbf{x}) \hat{\Psi}_i^\dagger(\mathbf{y}) U_{kij s}(\mathbf{x} - \mathbf{y}) \hat{\Psi}_j(\mathbf{y}) \hat{\Psi}_s(\mathbf{x}) \right]. \quad (1)$$

see also: Middelkamp et al. PRA **76** (2007) 022507.

## Bose gas beyond mean field (Bogoliubov)

$$\hat{\Psi}_g(\mathbf{x}) = \phi_0(\mathbf{x}) + \sum_{\mathbf{q}} \left( u_{\mathbf{q}}(\mathbf{x}) \hat{b}_{\mathbf{q}} - v_{\mathbf{q}}^*(\mathbf{x}) \hat{b}_{\mathbf{q}}^\dagger \right)$$

### Spin-Boson Hamiltonian

$$\hat{H}_{\text{syst}} = \frac{\Delta E(t)}{2} \hat{\sigma}_z$$

$$\hat{\sigma}_z = |p\rangle\langle p| - |s\rangle\langle s|$$

$$\hat{H}_{\text{env}} = \sum_{\mathbf{q}} \hbar \omega_{\mathbf{q}} \tilde{b}_{\mathbf{q}}^\dagger \tilde{b}_{\mathbf{q}}$$

$$\hat{H}_{\text{int}} = \sum_{\mathbf{q}} \frac{\Delta \kappa_{\mathbf{q}}}{2} \left( \tilde{b}_{\mathbf{q}} + \tilde{b}_{\mathbf{q}}^\dagger \right) \hat{\sigma}_z.$$

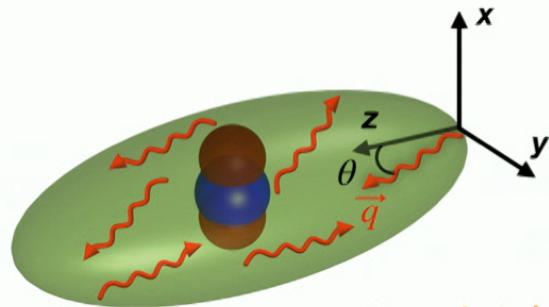


S. Rammohan et al. PRA **103**, 063307 (2021)

**For g-impurities:**  
Jaksch group and Lewenstein group



# Open quantum system



## Spin-Boson Hamiltonian

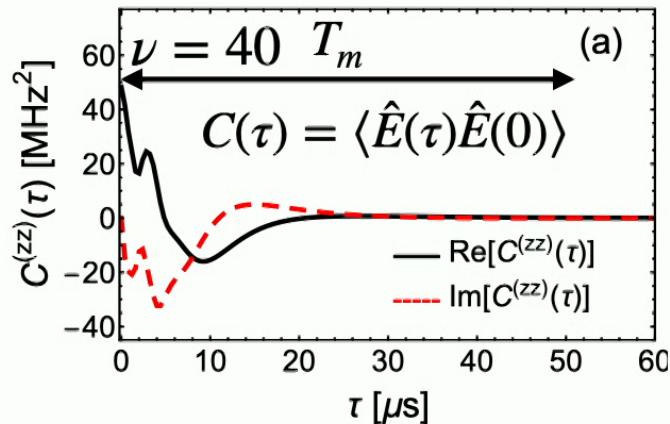
$$\hat{H}_{\text{syst}} = \frac{\Omega_{\text{mw}}}{2} \hat{\sigma}_x + \frac{\Delta E(t)}{2} \hat{\sigma}_z, \quad \hat{H}_{\text{env}} = \sum_{\mathbf{q}} \hbar \omega_{\mathbf{q}} \tilde{b}_{\mathbf{q}}^\dagger \tilde{b}_{\mathbf{q}}$$

**S** — **E**

$$\hat{H}_{\text{int}} = \sum_{\mathbf{q}} \frac{\Delta \kappa_{\mathbf{q}}}{\hbar} (\tilde{b}_{\mathbf{q}} + \tilde{b}_{\mathbf{q}}^\dagger) \hat{\sigma}_z.$$

**E** — **S** +

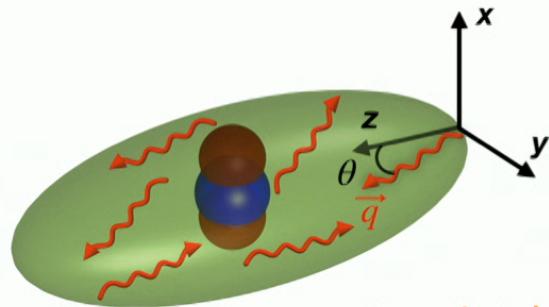
**g**   **p**   **s**    $\hat{H}_{\text{int}} = \hat{E} \otimes \hat{S}$



S. Rammohan et al. PRA 103, 063307 (2021)



# Open quantum system



## Spin-Boson Hamiltonian

$$\hat{H}_{\text{syst}} = \frac{\Omega_{\text{mw}}}{2} \hat{\sigma}_x + \frac{\Delta E(t)}{2} \hat{\sigma}_z, \quad \hat{H}_{\text{env}} = \sum_{\mathbf{q}} \hbar \omega_{\mathbf{q}} \tilde{b}_{\mathbf{q}}^\dagger \tilde{b}_{\mathbf{q}}$$

**S** E

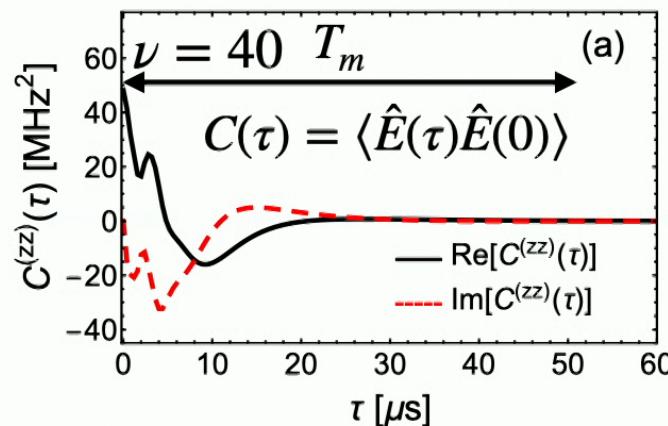
$$\hat{H}_{\text{int}} = \sum_{\mathbf{q}} \frac{\Delta \kappa_{\mathbf{q}}}{2} (\tilde{b}_{\mathbf{q}} + \tilde{b}_{\mathbf{q}}^\dagger) \hat{\sigma}_z.$$

**E** **S** E

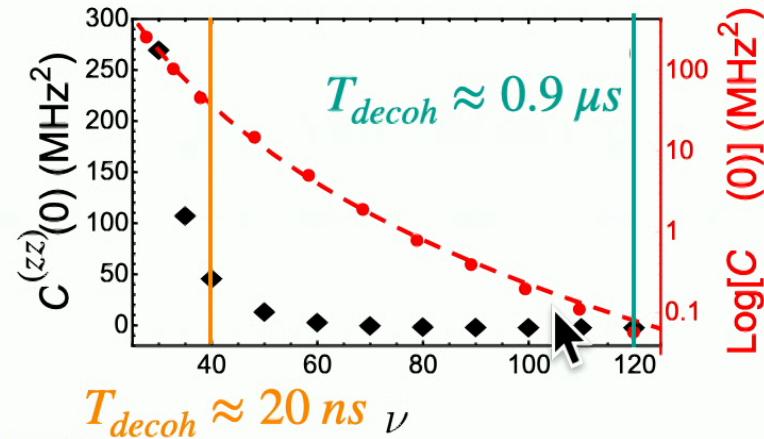


$$\hat{H}_{\text{int}} = \hat{E} \otimes \hat{S}$$

$$T_{\text{decoh}} \approx 1/\sqrt{2C(0)}$$



S. Rammohan et al. PRA 103, 063307 (2021)





# “Image” decohering state

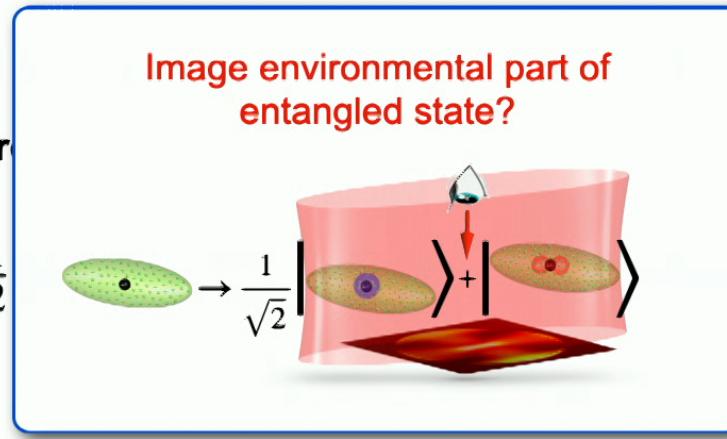
Rydberg state entangled with BEC



$$|\Psi_{\text{tot}}(t)\rangle = \frac{1}{\sqrt{2}}|\uparrow\rangle \otimes |\Psi_{\uparrow}(t)\rangle + \frac{1}{\sqrt{2}}|\downarrow\rangle \otimes |\Psi_{\downarrow}(t)\rangle$$

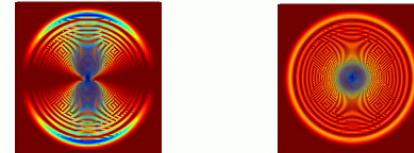
Rabi  $\pi/2$  Micro

$$|\Psi_{\text{tot}}\rangle = \frac{1}{2} \left[ \frac{1}{\sqrt{2}} \left( |\Psi_{\uparrow}\rangle + |\Psi_{\downarrow}\rangle \right) + i \left( |\Psi_{\uparrow}\rangle - |\Psi_{\downarrow}\rangle \right) \right]$$



Measurement of Rydberg state, selection of  $|\uparrow\rangle$

$$|\Psi_{\text{ms}}(t)\rangle \sim A[|\Psi_{\uparrow}(t)\rangle - i|\Psi_{\downarrow}(t)\rangle]$$



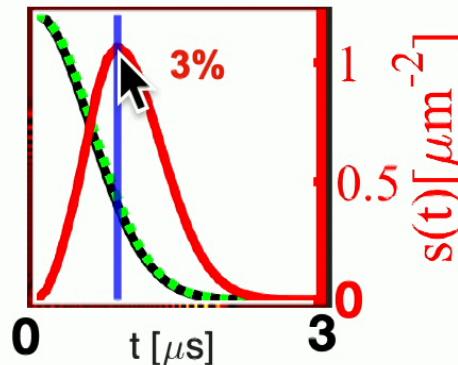
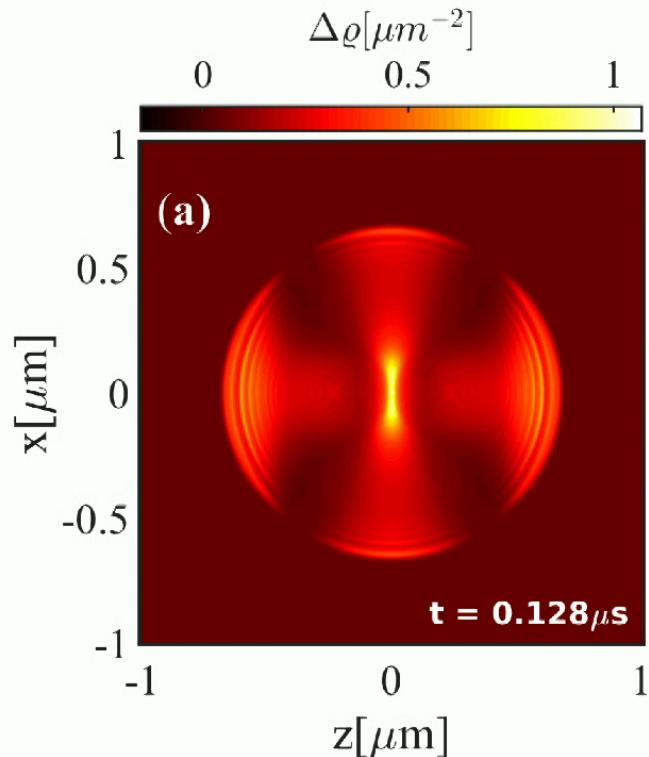


# Imaging decohering environment



BEC column densities in

$$|\Psi_{\text{ms}}(t)\rangle \sim A[|\Psi_{\uparrow}(t)\rangle - i|\Psi_{\downarrow}(t)\rangle] \quad (\text{minus}) \quad \hat{\rho} = (|\Psi_{\uparrow}\rangle\langle\Psi_{\uparrow}| + |\Psi_{\downarrow}\rangle\langle\Psi_{\downarrow}|)/2$$



- Transient glimpse at many-body entanglement at the root of decoherence:

$$|\Psi_{\text{tot}}(t)\rangle = c_{\uparrow}|\uparrow\rangle \otimes |\Psi_{\uparrow}(t)\rangle + c_{\downarrow}|\downarrow\rangle \otimes |\Psi_{\downarrow}(t)\rangle$$

S. Rammohan *et al.*  
Phys. Rev. A (Letters) **104**, L060202 (2021).



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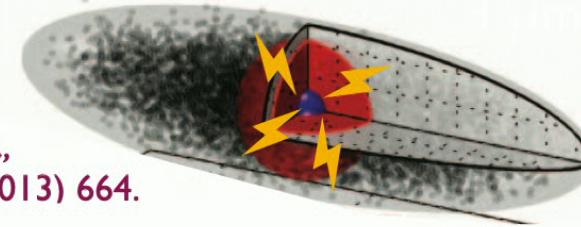


# Summary

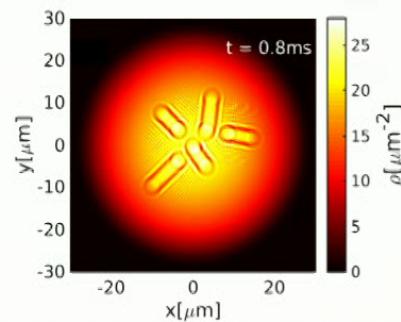


- Combination of two extreme systems:  
Rydberg atoms and BEC:

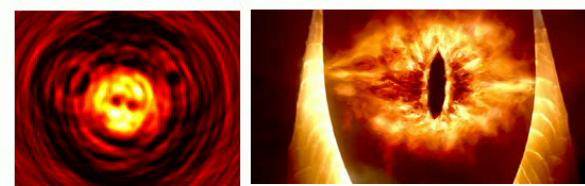
J. Balewski et al.,  
Nature **502** (2013) 664.



- Rydberg bubble chambers



S. Tiwari et al.,  
PRA **99** 043616 (2019);

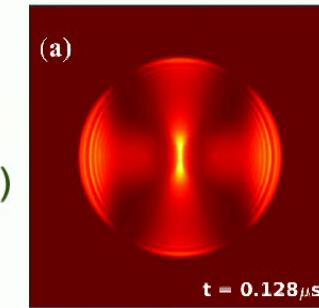


- Attractive BEC response

S. Tiwari et al.,  
NJP **24** 073005 (2022)

- Imaging the decohering interface

S. Rammohan et al.,  
PRA **103**, 063307 (2021);  
PRA (Letters) **104**, L060202 (2021)



Thanks for your attention