

Title: Symplectic formulation of the covariant phase space with boundaries

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Series: Quantum Fields and Strings

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Abstract: There are two different standard ways of endowing a physical theory with a symplectic structure: the canonical and the covariant. The former is derived from the well-known symplectic structure of a certain cotangent bundle. The latter is based on the variational calculus. Including a boundary in the canonical formalism poses no problem, however, in the covariant formalism things break apart. In this talk, I will briefly introduce both formalisms without boundary and explain in detail a new framework that allows us to include boundaries in a straightforward way. To show it in action, time permitting, I will apply it to several theories of gravity. Finally, I will briefly show a new result where we proved that, actually, the canonical and covariant formalisms are equivalent in full generality.

Zoom Link: <https://pitp.zoom.us/j/97641588140?pwd=dFl0SFdHOG5BYko2djNVWk11UkhSUT09>



# Symplectic formulation of the covariant phase space with boundaries

Juan Margalef Bentabol



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# CAN vs CPS

## CAN

Spacetime  $M \rightsquigarrow \mathbb{R} \times \Sigma$

- Dynamical perspective.
- Universal structures.
- Numerical approach.
- Quantization schemes.

- Aparent loss of symmetries.

## CPS

It deals with fields over the whole space-time.

- Explicit symmetries.
- Null infinity.
- Conserved quantities.
- Higher-derivative  $\equiv$  1st-order.

- No canonical structures.

# Differential geometry



## $M$ space-time Finite dimensional manifold

- $TM, V \in \mathfrak{X}(M)$ .
- $T^*M, \Omega^1(M)$ .
- $(TM)^{\otimes r} \otimes (T^*M)^{\otimes s}, \mathfrak{T}(M)$ .

### Operators:

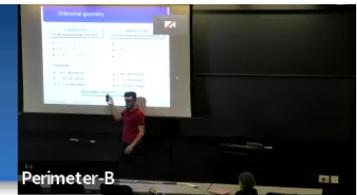
- $d$  ext. derivative.
- $\iota_V$  int. derivative.
- $\mathcal{L}_V$  Lie derivative.

## $\mathcal{F}$ space of fields Infinite-dimensional manifold

- $T\mathcal{F}, \mathbb{V} \in \mathfrak{X}(\mathcal{F})$ .
- $T^*\mathcal{F}, \Omega^1(\mathcal{F})$ .
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- $d\mathbb{I}$  ext. derivative.
- $\iota_{\mathbb{V}}$  int. derivative.
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Bicomplex framework:  $M \times \mathcal{F}$

# Scalar field

$$\mathcal{F} = C^\infty(M)$$
$$L(\phi) = \frac{1}{2} \nabla_\alpha \phi \nabla^\alpha \phi \text{vol}_g \in \Omega^n(M)$$

$$\begin{aligned} \text{d}L &= \nabla_\alpha \text{d}\phi \nabla^\alpha \phi \text{vol}_g = \nabla_\alpha (\text{d}\phi \nabla^\alpha \phi) \text{vol}_g - \text{d}\phi \nabla_\alpha \nabla^\alpha \phi \text{vol}_g = \\ &= E \text{d}\phi + \text{d}\Theta \end{aligned}$$

where  $E = -\square\phi \text{vol}_g$ ,  $\Theta = \iota_W \text{vol}_g$ , and  $W^\alpha = \text{d}\phi \nabla^\alpha \phi$ .

$$\Omega = \int_\Sigma \text{d}\Theta = \int_\Sigma \text{d}\phi \wedge \nabla_{\vec{n}} \text{d}\phi \text{vol}_\Sigma$$

Symplectic form over  $\text{Sol} = \{\phi \in C^\infty(M) / \square\phi = 0\}$

# Action

$$\mathbb{S}^L : \mathcal{F} \rightarrow \mathbb{R} \quad \mathbb{S}^L = \int_M L \quad L \in \Omega^{(n,0)}(M \times \mathcal{F})$$

- If  $L' = L + dY$ , then  $\mathbb{S}^{L'} = \mathbb{S}^L$ .
- If  $\mathbb{S}^{L'} = \mathbb{S}^L$  over contractible bundles, then  $L' = L + dY$ .

$$\text{Sol}(\mathbb{S}^L) = \{\phi \in \mathcal{F} / d_\phi \mathbb{S}^L = 0\}$$

$$d\mathbb{S}^L = \int_M dL \quad dL = E \wedge d\phi + d\Theta$$

$$\text{Sol}(\mathbb{S}^L) = \{\phi \in \mathcal{F} / E(\phi) = 0\} \xrightarrow{\mathcal{J}_{\text{CPS}}} \mathcal{F}$$

$$dL' = E \wedge d\phi + d\Theta' \quad \Theta' = \Theta + dY + dZ$$

# Symplectic structure

Given  $\iota : \Sigma \hookrightarrow M$  Cauchy embedding:

$$\Omega^\iota := \int_{\Sigma} \iota^* d\Theta \quad d\Theta \in \Omega^{(n-1,2)}(M \times \mathcal{F})$$

$$(\Omega')^\iota = \int_{\Sigma} \iota^* d\Theta' = \int_{\Sigma} \iota^* d(\Theta + dY + dZ) = \Omega^\iota$$

Depends on the embedding  $\iota$ .

$$\Omega_{\text{CPS}} := \mathcal{J}_{\text{CPS}}^* \Omega^\iota \text{ symplectic form on } \text{Sol}(\mathbb{S}^L).$$

# Part IV: CAN formalism in a nutshell

J. Margalef Bentabol | Symplectic formulation



# CAN formalism

$Q = Q^1 \times \dots \times Q^M$  configuration space over  $\Sigma$

CAN Lagrangian-action  $\mathcal{L} : TQ \rightarrow \mathbb{R}$

$$\mathcal{L} = \int_{\Sigma} \hat{\mathcal{L}} \quad \hat{\mathcal{L}} \in \Omega^{(n-1,0)}(\Sigma \times TQ)$$

Lagrangian  $(\mathcal{L}, TQ) \rightarrow$  Hamiltonian  $(\mathcal{H}, \text{PCM} \xrightarrow{\mathcal{J}_{\text{CAN}}} T^*Q)$

$$\Omega_{T^*Q} := dq^I \wedge dp_I \quad \longrightarrow \quad \Omega_{\text{CAN}} := \mathcal{J}_{\text{CAN}}^* \Omega_{T^*Q}$$

**Theorem (J.M-B, E.J.S Villaseñor)**

For every theory of any order such that  $\Omega_{\text{CAN}}$  exists, then

$$\Omega_{\text{CPS}} = \Omega_{\text{CAN}}$$

# Part V: Adding boundaries



# Differential geometry on $(M, N)$

$$N \subset \partial M \xrightarrow{j} M \quad \Omega^r(M, N) := \Omega^r(M) \times \Omega^{r-1}(N)$$

- $\underline{d}(\alpha, \beta) := (d\alpha, j^*\alpha - d\beta)$
- $(\alpha_1, \beta_1) \underline{\wedge} (\alpha_2, \beta_2) := \dots$
- $\int_{(M, N)} (a, b) := \int_M a - \int_N b$
- $\underline{\partial}(M, N) := (\partial M \setminus N, \partial N)$

$$\begin{aligned} \int_{(M, N)} \underline{d}(\alpha, \beta) &= \int_M d\alpha - \int_N (j^*\alpha - d\beta) = \\ &= \int_{\partial M \setminus N} j^*\alpha - \int_{\partial N} (j|_{\partial})^*\beta = \int_{(\partial M \setminus N, \partial N)} (j^*\alpha, (j|_{\partial})^*\beta) \end{aligned}$$

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- $\int_{(M, N)} (a, b) := \int_M a - \int_N b$
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- $\underline{F}^*(\alpha, \beta) := (F^*\alpha, (F|_{\partial})^*\beta)$ .

$$\begin{aligned} \int_{(M, N)} \underline{d}(\alpha, \beta) &= \int_M d\alpha - \int_N (j^*\alpha - d\beta) = \\ &= \int_{\partial M \setminus N} j^*\alpha - \int_{\partial N} (j|_{\partial})^*\beta = \int_{\underline{\partial}(M, N)} \underline{j}^*(\alpha, \beta) \end{aligned}$$

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- $\int_{(M, N)} (a, b) := \int_M a - \int_N b$
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$$\int_{(M, N)} \underline{d}(\alpha, \beta) = \int_{\underline{\partial}(M, N)} \underline{j}^*(\alpha, \beta)$$

- $\underline{\mathcal{L}}_V(\alpha, \beta) := \dots \quad \underline{\mathcal{L}}_V(\alpha, \beta) := \dots$

## Differential geometry on $(M, N)$

$$\underline{d}(\alpha, \beta) = (d\alpha, j^*\alpha - d\beta) \quad \underline{\partial}(M, N) := (\partial M \setminus N, \partial N)$$

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$(M, N)$  behaves as a manifold with boundary  $\underline{\partial}(M, N)$

$(M, \partial M)$  behaves like a manifold **without boundary!**

$$\underline{\partial}(M, \partial M) = \emptyset$$

Relative bicomplex framework  
 $(M, N) \times \mathcal{F} := (M \times \mathcal{F}, N \times \mathcal{F})$

# Action

$$\mathbb{S}^{(L, \bar{\ell})} = \int_{(M, \partial M)} (L, \bar{\ell}) \quad L \in \Omega^{(n, 0)}(M \times \mathcal{F})$$



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- If  $(L', \bar{\ell}') = (L, \bar{\ell}) + \underline{d}(Y, \bar{y})$ , then  $\mathbb{S}^{(L', \bar{\ell}')} = \mathbb{S}^{(L, \bar{\ell})}$

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- If  $\mathbb{S}^{L'} = \mathbb{S}^L$  over c.b., then  $L' = L + dY$

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$$\text{Sol}(\mathbb{S}^{(L, \bar{\ell})}) = \{\phi \in \mathcal{F} / \underline{d}_\phi \mathbb{S}^{(L, \bar{\ell})} = 0\}$$

$$\underline{d}\mathbb{S}^L = \int_M \underline{d}L \quad \underline{d}L = E \wedge \underline{d}\phi + \underline{d}\Theta$$



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$$\underline{\mathbf{d}}(L, \bar{\ell}) = (E, \bar{b}) \wedge \mathbf{d}\phi + \underline{\mathbf{d}}(\Theta, \bar{\theta})$$

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$$\text{Sol}(\mathbb{S}^{(L, \bar{\ell})}) = \{\phi \in \mathcal{F} / (E, \bar{b})(\phi) = 0\} \xrightarrow{\mathcal{J}_{\text{CPS}}} \mathcal{F}$$

# Symplectic structure

Given  $\iota : \Sigma \hookrightarrow M$  Cauchy embedding:

$$\Omega^\iota := \int_{(\Sigma, \partial\Sigma)} \underline{\iota}^* \underline{d}(\Theta, \bar{\theta}) = \int_{\Sigma} \iota^* \underline{d}\Theta - \int_{\partial\Sigma} (\iota|_{\partial})^* \underline{d}\bar{\theta}$$

$$(\Omega')^\iota = \int_{(\Sigma, \partial\Sigma)} \underline{\iota}^* \underline{d}(\Theta', \bar{\theta}') = \int_{(\Sigma, \partial\Sigma)} \underline{\iota}^* \underline{d}((\Theta, \bar{\theta}) + \underline{d}(Y, \bar{y}) + \underline{d}(Z, \bar{z})) = \Omega^\iota$$

Depends on the embedding  $\iota$ .

$$\Omega_{\text{CPS}} := j_{\text{CPS}}^* \Omega^\iota \text{ symplectic form on } \text{Sol}(\mathbb{S}^{(L, \bar{\ell})}).$$

# Action

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$$\underline{\mathbf{d}}(L, \bar{\ell}) = (E, \bar{b}) \wedge \mathbf{d}\phi + \underline{\mathbf{d}}(\Theta, \bar{\theta})$$

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Depends on the embedding  $\iota$ .

$$\Omega_{\text{CPS}} := j_{\text{CPS}}^* \Omega^\iota \text{ symplectic form on } \text{Sol}(\mathbb{S}^{(L, \bar{\ell})}).$$

$$\Omega_{\text{CPS}} = \Omega_{\text{CAN}}$$

# Conclusions

- 1 We have developed a new formalism to derive the CPS presymplectic form canonically associated with an action.
- 2 The bicomplex formalism is essential.
- 3 We have proved that the CPS and CAN presymplectic structures are equivalent.
- 4 Also true for higher-order theories and with boundaries.
- 5 We have successfully applied to several theories of gravity solving some open questions.

# Part VII: CPS of GR

J. Margalef Bentabol | Symplectic formulation



# Metric GR

$$L_{\text{EH}}^{(m)}(g) = R_g \text{vol}_g \quad \bar{\ell}_{\text{GHY}}^{(m)}(g) = -2\bar{g}^{\bar{\alpha}\bar{\beta}} \bar{K}_{\bar{\alpha}\bar{\beta}} \text{vol}_{\bar{g}}$$

$$dL_{\text{EH}}^{(m)} = E_{(m)}^{\alpha\beta} dg_{\alpha\beta} + d\Theta_{(m)}^{\text{GR}}$$

$$E_{(m)}^{\alpha\beta} = - \left( \text{Ric}^{\alpha\beta} - \frac{R_g}{2} g^{\alpha\beta} \right) \text{vol}_g$$

$$\Theta_{(m)}^{\text{GR}} = \iota_{\bar{W}} \text{vol}_g$$

$$W_\alpha = \nabla^\beta dg_{\alpha\beta} - g^{\mu\nu} \nabla_\alpha dg_{\mu\nu}$$

$$d\bar{\ell}^{(m)} - j^* \Theta_{(m)}^{\text{GR}} = \bar{b}_{(m)}^{\bar{\alpha}\bar{\beta}} d\bar{g}_{\bar{\alpha}\bar{\beta}} - d\bar{\theta}_{(m)}^{\text{GR}}$$

$$\bar{b}_{(m)}^{\bar{\alpha}\bar{\beta}}(g) = - \left( \bar{K}^{\bar{\alpha}\bar{\beta}} - \bar{K} \bar{g}^{\bar{\alpha}\bar{\beta}} \right) \text{vol}_{\bar{g}}$$

$$\bar{\theta}_{(m)}^{\text{GR}}(g) = \iota_{\bar{V}} \text{vol}_{\bar{g}}$$

# Metric GR

$$L_{\text{EH}}^{(m)}(g) = R_g \text{vol}_g \quad \bar{\ell}_{\text{GHY}}^{(m)}(g) = -2\bar{g}^{\bar{\alpha}\bar{\beta}} \bar{K}_{\bar{\alpha}\bar{\beta}} \text{vol}_{\bar{g}}$$

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$$\Theta_{(m)}^{\text{GR}} = \iota_{\bar{W}} \text{vol}_g$$

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$$\bar{b}_{(m)}^{\bar{\alpha}\bar{\beta}}(g) = - \left( \bar{K}^{\bar{\alpha}\bar{\beta}} - \bar{K} \bar{g}^{\bar{\alpha}\bar{\beta}} \right) \text{vol}_{\bar{g}}$$

$$\bar{\theta}_{(m)}^{\text{GR}}(g) = \iota_{\bar{V}} \text{vol}_{\bar{g}}$$

$$\bar{V}_{\bar{\alpha}} = -j_{\bar{\alpha}}^\alpha \nu^\beta dg_{\alpha\beta}$$

# Tetrad GR

$$\begin{aligned} \Phi_{\text{GR}} : \text{Tetrads} &\longrightarrow \text{Metrics} \\ e_{\alpha}^I &\longmapsto g_{\alpha\beta} = \eta_{IJ} e_{\alpha}^I e_{\beta}^J \end{aligned}$$

# Metric GR

$$L_{\text{EH}}^{(m)}(g) = R_g \text{vol}_g \quad \bar{\ell}_{\text{GHY}}^{(m)}(g) = -2\bar{g}^{\bar{\alpha}\bar{\beta}} \bar{K}_{\bar{\alpha}\bar{\beta}} \text{vol}_{\bar{g}}$$

$$dL_{\text{EH}}^{(m)} = E_{(m)}^{\alpha\beta} dg_{\alpha\beta} + d\Theta_{(m)}^{\text{GR}}$$

$$E_{(m)}^{\alpha\beta} = - \left( \text{Ric}^{\alpha\beta} - \frac{R_g}{2} g^{\alpha\beta} \right) \text{vol}_g$$

$$\Theta_{(m)}^{\text{GR}} = \iota_{\bar{W}} \text{vol}_g$$

$$W_\alpha = \nabla^\beta dg_{\alpha\beta} - g^{\mu\nu} \nabla_\alpha dg_{\mu\nu}$$

$$d\bar{\ell}^{(m)} - j^* \Theta_{(m)}^{\text{GR}} = \bar{b}_{(m)}^{\bar{\alpha}\bar{\beta}} d\bar{g}_{\bar{\alpha}\bar{\beta}} - d\bar{\theta}_{(m)}^{\text{GR}}$$

$$\bar{b}_{(m)}^{\bar{\alpha}\bar{\beta}}(g) = - \left( \bar{K}^{\bar{\alpha}\bar{\beta}} - \bar{K} \bar{g}^{\bar{\alpha}\bar{\beta}} \right) \text{vol}_{\bar{g}}$$

$$\bar{\theta}_{(m)}^{\text{GR}}(g) = \iota_{\bar{V}} \text{vol}_{\bar{g}}$$

$$\bar{V}_{\bar{\alpha}} = -j_{\bar{\alpha}}^\nu dg_{\alpha\nu}$$

# Tetrad GR

$$\begin{aligned} \Phi_{\text{GR}} : \text{Tetrads} &\longrightarrow \text{Metrics} \\ e_{\alpha}^I &\longmapsto g_{\alpha\beta} = \eta_{IJ} e_{\alpha}^I e_{\beta}^J \end{aligned}$$

$$\Phi_{\text{GR}}(e) = \Phi_{\text{GR}}(e') \iff e'_I = \Psi_I^J e_J \text{ with } \Psi \in SO(1, 3)$$

$$L^{(t)}(e) := L_{\text{EH}}^{(m)} \circ \Phi_{\text{GR}}(e) \quad \bar{\ell}^{(t)}(e) := \bar{\ell}_{\text{GHY}}^{(m)} \circ \Phi_{\text{GR}}(e)$$

$$L^{(t)}(e) = \frac{1}{2} \varepsilon_{IJKL} F^{IJ} \wedge e^K \wedge e^L \quad F^{IJ} = d\omega^{IJ} + \omega^I_K \wedge \omega^{KJ}$$

$$\bar{\ell}^{(t)}(e) = -\frac{1}{2} \varepsilon_{IJKL} \left( 2N^I dN^J - \bar{\omega}^{IJ} \right) \bar{e}^K \wedge \bar{e}^L \quad N^I = e_{\alpha}^I \nu^{\alpha}$$

# Tetrad GR

$$dL^{(t)} = E_I^{(t)} \wedge de^I + d\Theta_{(t)}^{\text{GR}}$$

$$E_I^{(t)}(e) = -\varepsilon_{IJKL} F^{JK} \wedge e^L$$

$$\Theta_{(t)}^{\text{GR}} = \frac{1}{2} \varepsilon_{IJKL} e^I \wedge e^J \wedge d\omega^{KL}$$

$$d\bar{\ell}^{(t)} - j^* \Theta_{(t)}^{\text{GR}} = \bar{b}_I^{(t)} \wedge d\bar{e}^I - d\bar{\theta}_{(t)}^{\text{GR}}$$

$$\bar{b}_I^{(t)} = \dots$$

$$\bar{\theta}_{(t)}^{\text{GR}} = \varepsilon_{IJKL} \bar{e}^I \wedge \bar{e}^J \wedge N^K dN^L$$

# Metric vs Tetrad

$$\mathbb{S}^{(t)} = \mathbb{S}^{(m)} \circ \Phi_{\text{GR}} \rightarrow d\mathbb{S}^{(t)} = d\mathbb{S}^{(m)} \cdot d\Phi_{\text{GR}} \rightarrow \text{Sol}_{\text{GR}}^{(t)} = \Phi_{\text{GR}}^{-1} \text{Sol}_{\text{GR}}^{(m)}$$

$$\Theta_{(m)}^{\text{GR}} = \Theta_{(t)}^{\text{GR}} + d\alpha \quad \bar{\theta}_{(m)}^{\text{GR}} = \bar{\theta}_{(t)}^{\text{GR}} + j^* \alpha$$

$$\left( \Theta_{(m)}^{\text{GR}}, \bar{\theta}_{(m)}^{\text{GR}} \right) = \left( \Theta_{(t)}^{\text{GR}}, \bar{\theta}_{(t)}^{\text{GR}} \right) + \underline{d}(\alpha, 0)$$

$$\Omega_{(t)}^{\text{GR}} = \Phi_{\text{GR}}^* \Omega_{(m)}^{\text{GR}}$$

# Conclusions

- 1 We prove the CPS equivalence of metric-GR and tetrad-GR (same for Palatini and Holst).
- 2 The problem is cohomological (relative version), so we need cohomological methods.



# Conclusions

- 1 We prove the CPS equivalence of metric-GR and tetrad-GR (same for Palatini and Holst).
- 2 The problem is cohomological (relative version), so we need cohomological methods.
- 3 Lagrangian boundary terms are essential.
- 4 Symplectic boundary terms are required.



# Thanks for your attention!



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