Title: Superconductivity, charge density wave, and supersolidity in flat bands with tunable quantum metric

Speakers:

Collection: Quantum Criticality: Gauge Fields and Matter

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URL: https://pirsa.org/22050039

Abstract: Predicting the fate of an interacting system in the limit where the electronic bandwidth is quenched is often highly non-trivial. The complex interplay between interactions and quantum fluctuations driven by the band geometry can drive a competition between various ground states, such as charge density wave order and superconductivity. In this work, we study an electronic model of topologically-trivial flat bands with a continuously tunable Fubini-Study metric in the presence of on-site attraction and nearest-neighbor repulsion, using numerically exact quantum Monte Carlo simulations. By varying the electron filling and the spatial extent of the localized flat-band Wannier wavefunctions, we obtain a number of intertwined orders. These include a phase with coexisting charge density wave order and superconductivity, i.e., a supersolid. In spite of the non-perturbative nature of the problem, we identify an analytically tractable limit associated with a 'small' spatial extent of the Wannier functions, and derive a low-energy effective Hamiltonian that can well describe our numerical results.

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Superconductivity, charge density wave and supersolidity in flat bands with tunable quantum metric

ArXiv:2204:02994

Johannes S. Hofmann*, Erez Berg* and Debanjan Chowdhury†



*Department of Condensed Matter Physics, Weizmann Institute of Science, Israel †Department of Physics, Cornell University, USA



May 18th, 2022

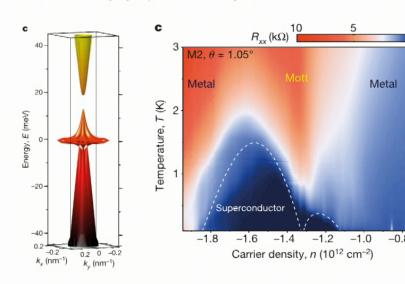
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ARTICLE

doi:10.1038/nature:26160

Unconventional superconductivity in magic-angle graphene superlattices

Yuan Cao¹, Valla Fatemi¹, Shiang Fang², Kenji Watanabe³, Takashi Taniguchi³, Efthimios Kaxiras^{2,4} & Pablo Jarillo-Herrero¹



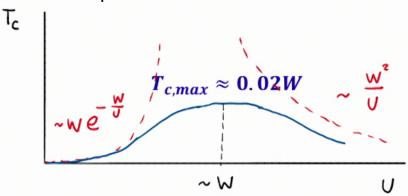
Emery, Kivelson: Nature 1995

Two ingredients:

- pair binding $\sim U$ - phase coherence $T_{ heta}$

Often: $T_c \to 0$ when $W/U \to 0$ as $T_\theta \to 0$

For example: attractive Hubbard model



Paiva, Scalettar, Randeria and Trivedi: PRL 2010

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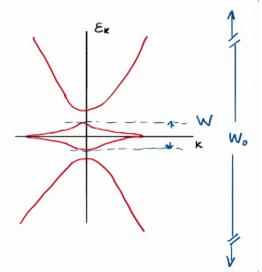
Volovik (2011); Peotta, Torma (2015) Tovmasyan, Huber *et.al.* (2016) Xie, Bernevig *et.al.* (2019) Park, Kim, Lee (2020)

Consider isolated, nearly flat band with W << U << $\Delta_{\rm gap}$:

- Ground state as a function of $|U|/W \rightarrow \infty$?
- Robustness to infinitesimal perturbations (e.g. nearestneighbor interactions)?
 - Superconductor, phase separation, CDW, supersolid, ...



Mean-field theory is not the way to address these questions



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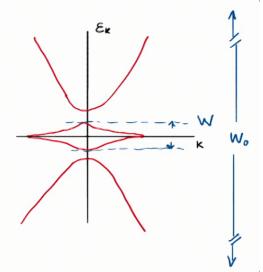
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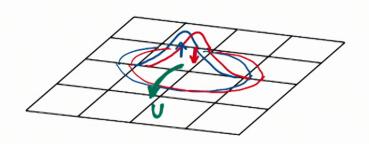
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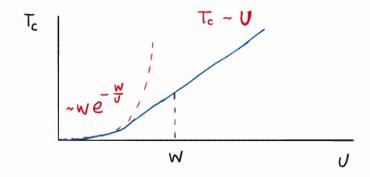
Volovik (2011); Peotta, Torma (2015) Tovmasyan, Huber et.al. (2016) Xie, Bernevig et.al. (2019) Park, Kim, Lee (2020)

Certain flat bands are "special":

Possible for local Cooper-pairs to delocalize even in the flatband limit W → 0

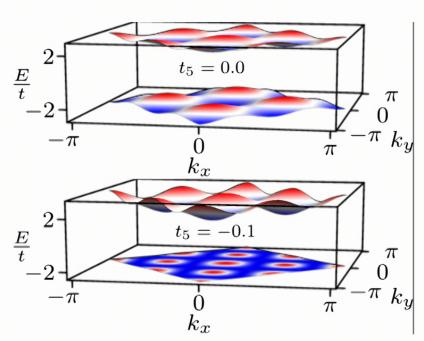
– i.e. $T_c \propto |U|$, instead of the standard $T_c \propto W^2/|U|$





Topological flat bands

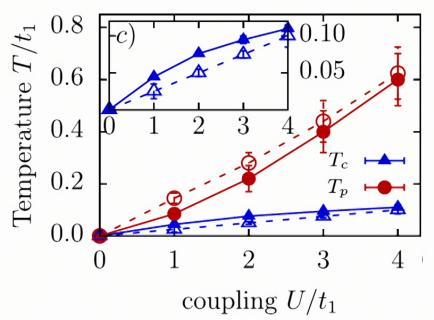
 $C = \pm 1$ Chern bands



Flatness ratio:
$$\mathcal{F} = \frac{W}{\Delta_{\mathrm{gap}}} (=0.2, 0.01)$$

$$H = \sum_{i,j,\sigma} t_{ij}^{\sigma} c_{i\sigma}^{\dagger} c_{j\sigma} - U \sum_{i} n_{i\uparrow} n_{i\downarrow} \left[\text{TRS} : (t_{ij}^{\uparrow})^* = t_{ij}^{\downarrow} \right]$$

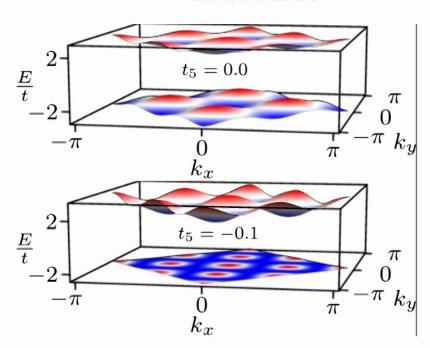
Non-perturbative DQMC results:



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Topological flat bands

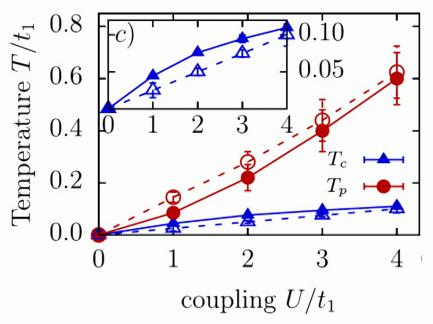
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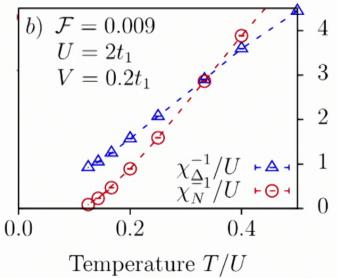
Lessons

 An incoherent metallic normal state can give rise to a superconducting ground state.

(Unlike any "mean-field BCS" transition)

- Strongly fluctuating density and pairing correlations
- Approximate SU(2) invariant limit can be a useful starting point Tovmasyan, Huber et.al. (2016)
- **BUT:** weak nearest neighbor attraction can tip the balance in favor of phase separation!

$$H_V = -V \sum_{\langle i,j \rangle} n_i n_j$$



Key questions

- What is the nature of the competing phases and the associated quantum phase-transitions?
- How does varying the spatial extent of the localized Wannier functions tune the system between different orders?
- How do the interactions induce a finite dispersion for the excitations?
- Is there a theoretical limit in which this band competition can be explored in a controlled fashion, without resorting to BCS mean-field theory?

Hofmann, Berg, Chowdhury; ArXiv:2204:02994

Pirsa: 22050039

Superconductivity in flat bands with tunable metric

$$\hat{H} = \hat{H}_{kin}(\zeta) + \hat{H}_{int}$$

- $\hat{H}_{kin}(\zeta)$ engineered to yield isolated, flat bands with vanishing Berry curvature
 - ζ controls the spatial extent of the Wannier function
 - $-\zeta = 0$ corresponds to a decoupled "atomic limit"

$$\hat{H}_{int} = -\frac{U}{2} \sum_{i,l} \left(\sum_{\sigma} \hat{n}_{i,l,\sigma} - \frac{1}{2} \right)^2 + V \sum_{\langle ij \rangle,l} \left(\sum_{\sigma} \hat{n}_{i,l,\sigma} - \frac{1}{2} \right) \left(\sum_{\sigma} \hat{n}_{j,l,\sigma} - \frac{1}{2} \right)$$

•
$$\hat{H}_{kin}(\zeta) = t \sum_{\mathbf{k},\sigma} \hat{c}_{\mathbf{k},\sigma}^{\dagger} \left(\sin \alpha_{\mathbf{k}} \tau_x + (-1)^{\sigma} \cos \alpha_{\mathbf{k}} \tau_y - \mu \tau_0 \right) \hat{c}_{\mathbf{k},\sigma}$$

 $\alpha_{\mathbf{k}} = \zeta (\cos k_x + \cos k_y)$

Hofmann, Berg, Chowdhury; ArXiv:2204:02994

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Effective theory

For small ζ, we can derive the following effective model:

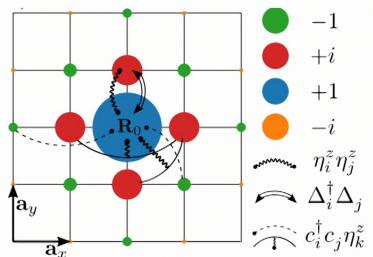
$$\hat{H}_{\text{int}} = -\frac{U_{\text{eff}}}{2} \sum_{\mathbf{r}} (2\hat{\eta}_{\mathbf{r}}^z)^2 - \sum_{\langle \mathbf{r}, \mathbf{r}' \rangle} [J_{\perp}(\hat{\eta}_{\mathbf{r}}^x \hat{\eta}_{\mathbf{r}'}^x + \hat{\eta}_{\mathbf{r}}^y \hat{\eta}_{\mathbf{r}'}^y) + J_z \hat{\eta}_{\mathbf{r}}^z \hat{\eta}_{\mathbf{r}'}^z] + \hat{H}_{\text{hop}}$$

$$\hat{\eta}^{\alpha}_{\mathbf{r}} \equiv (\Psi^{\dagger}_{\mathbf{r}} \eta^{\alpha} \Psi_{\mathbf{r}})/2 \qquad \Psi^{\dagger}_{\mathbf{r}} = (\hat{d}^{\dagger}_{\mathbf{r},\uparrow}, \hat{d}_{\mathbf{r},\downarrow})$$

$$J_{\perp} = \zeta^2 U/4 + \dots, \quad J_z = \zeta^2 U/4 - 2V + \dots$$

$$U_{\text{eff}} = U(2 - \zeta^2)/4$$

$$ilde{H}_{
m hop} \propto U \zeta^2 t_{ij}(k) c_i^\dagger c_j \hat{\eta}_k^z \equiv egin{array}{c} {
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Hofmann, Berg, Chowdhury; ArXiv:2204:02994

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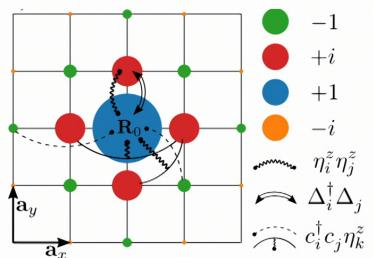
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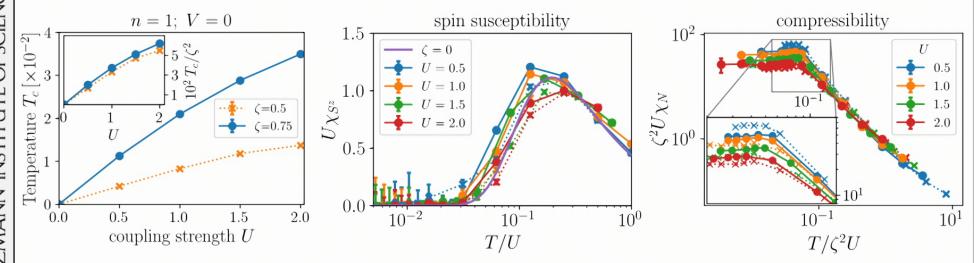
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Superconductivity in flat bands with tunable metric

Non-perturbative DQMC results (V=0.0):



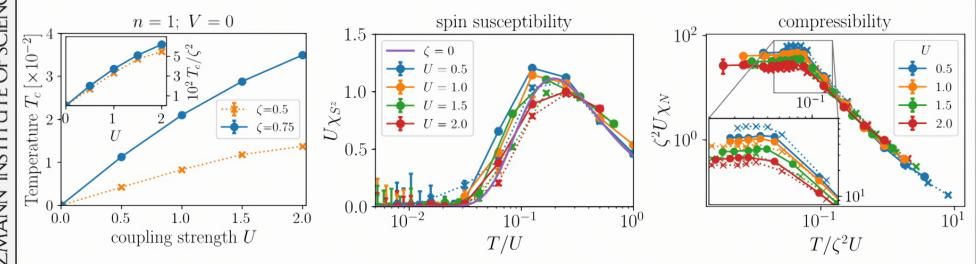
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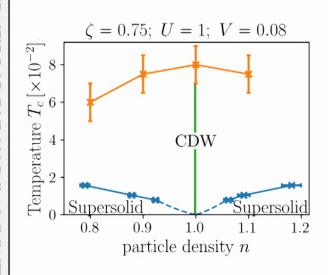
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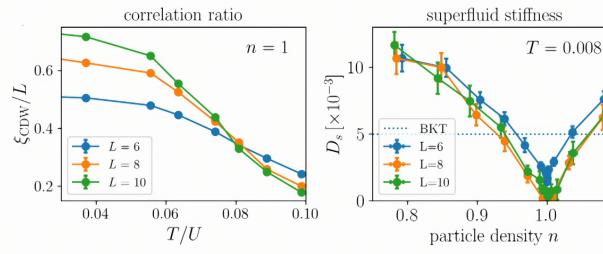
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Supersolid at incommensurate filling

Non-perturbative DQMC results (V=0.08):





Hofmann, Berg, Chowdhury; ArXiv:2204:02994

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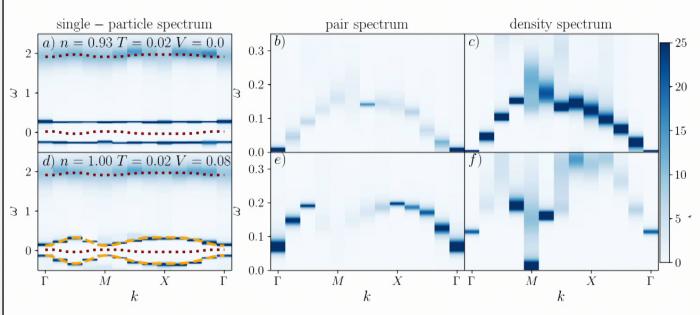
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Excitation spectra & supersolid at quarter filling

Non-perturbative DQMC results:



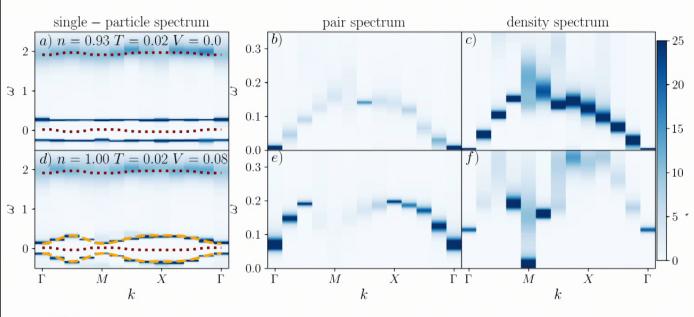
Hofmann, Berg, Chowdhury; ArXiv:2204:02994

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$$\ddot{\Delta} - \frac{U\zeta^2 \Delta_{\text{CDW}}}{4} \left(2 \sum_{a=\pm} \cos k_a + \sum_{a=x,y} \cos 2k_a \right)$$

Hofmann, Berg, Chowdhury; ArXiv:2204:02994

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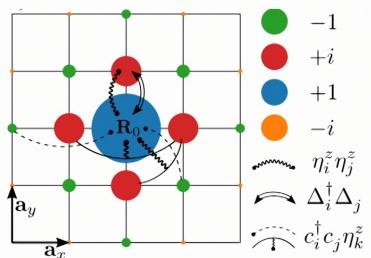
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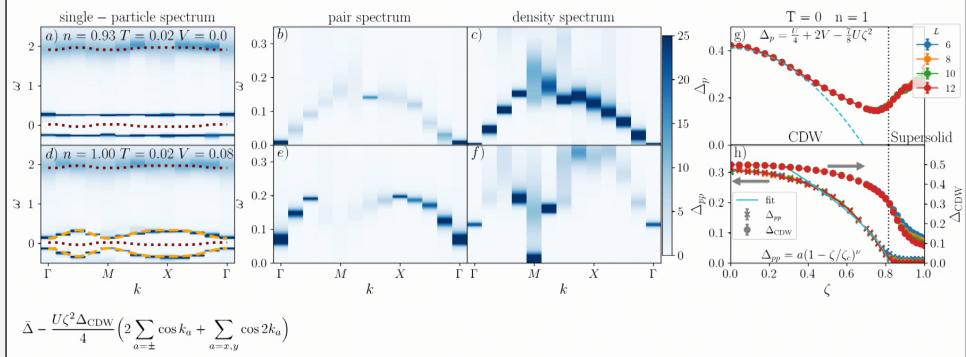
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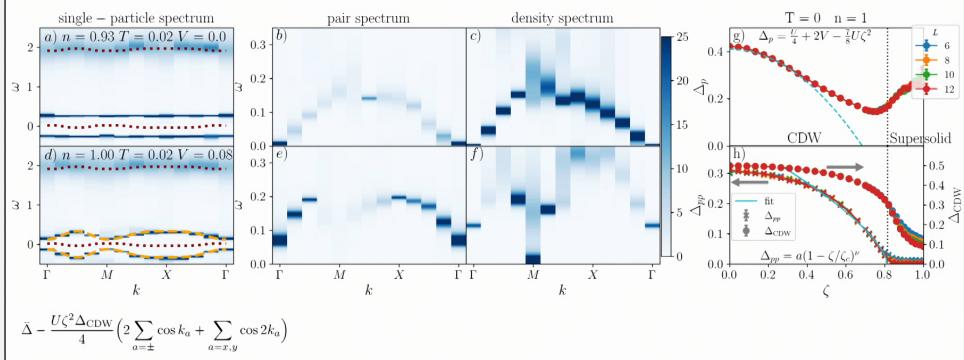
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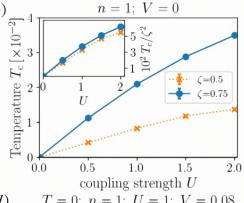
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Conclusion

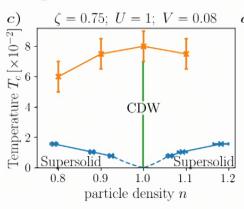
Extended Wannier orbitals → XXZ model + assisted hopping

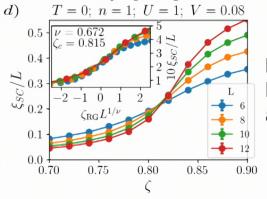
 $\begin{array}{c} \mathbf{a}) \\ -1 \\ +i \\ +1 \\ -i \\ \mathbf{a}_y \\ \mathbf{a}_x \\ \end{array}$



Superconductivity

Charge density wave; supersolid ground state at finite doping





Density-assisted hopping → Supersolid at quarter filling

Hofmann, Berg, Chowdhury; ArXiv:2204:02994

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