

Title: Dirac criticality from field theory beyond the leading order

Speakers: Michael Scherer

Collection: Quantum Criticality: Gauge Fields and Matter

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Abstract: Two-dimensional gapless Dirac fermions emerge in various condensed-matter settings. In the presence of interactions such Dirac systems feature critical points and the precision determination of their exponents is a prime challenge for quantum many-body methods. In a field-theoretical language, these critical points can be described by Gross-Neveu-Yukawa-type models and in my talk I will show some results on Gross-Neveu critical behavior using field theoretical approaches beyond the leading order. To that end, I will first present higher-loop perturbative RG calculations for generic Gross-Neveu-Yukawa models and compare estimates for the exponents with recent corresponding results from Quantum Monte Carlo simulations and the conformal bootstrap. Then, I will discuss a more exotic variant of Gross-Neveu-Yukawa models which describes the interacting fractionalized excitations of two-dimensional frustrated spin-orbital magnets. Here, we have provided field-theoretical estimates for the critical exponents employing higher-order epsilon expansion, large-N calculations, and functional renormalization group.

# Dirac criticality from field theory beyond the leading order



**Michael M. Scherer**  
Ruhr University Bochum

May 18, 2022 @ PI

# Outline

- Mini-review of Dirac criticality from field theory
  - Dirac fermions and quantum critical points #Gross-Neveu-models
  - Critical exponents from perturbative RG at 4-loop order #hep
  - What to expect from these results ... and what not?! #many-body-methods

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- **Fractionalized quantum criticality in spin-orbital liquids** ([→ talk by L. Janssen](#))
  - Gross-Neveu-Yukawa\* model
  - 3-loop perturbative RG, 2<sup>nd</sup>-order large-N expansion, functional RG

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- **Collaborators:**

John Gracey (Liverpool), Igor Herbut (Vancouver), Bernhard Ihrig (Cologne), Lukas Janssen (Dresden), Daniel Krutik (Cologne), Peter Marquard (Berlin), Luminita Mihaila (Heidelberg), Shouryya Ray (Dresden), Nikolai Zerf (Berlin)

## Materials and quasi-relativistic fermions

- emergence of **Dirac, Weyl & Majorana** quasi-particle excitations in materials
- interacting Dirac fermions have quantum critical points
  - **graphene**: antiferromagnetic/Kekulé-bond/charge order/...
  - ...

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# Materials and quasi-relativistic fermions

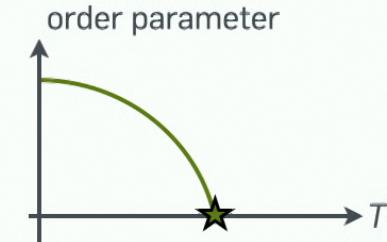
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⇒ **fermionic universality classes** → **What are their critical exponents?**

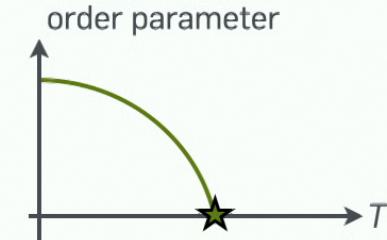
## Three-dimensional Ising universality class

- near critical point of continuous phase transition: **universality**
- **correlation length** diverges for  $T \rightarrow T_c$ :  $\xi(t) \propto |t|^{-\nu}$
- order parameter **correlation function** at  $T_c$ :  $G(\vec{r}) \propto |\vec{r}|^{2-D-\eta}$



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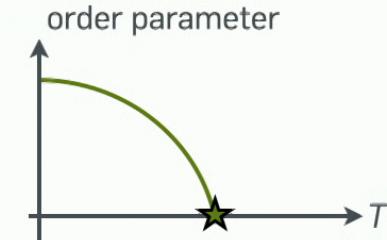


Method	$\nu$	$\eta$
conformal bootstrap	0.629971(4)	0.036298(2)
Monte Carlo	0.63002(10)	0.03627(10)
pRG, 4- $\epsilon$ , 6th order	0.6292(5)	0.0362(2)
functional RGs, DE	0.630(5)	0.034(5)
experiment (fluid mixture)	0.629(3)	0.032(13)

- Kos *et al.* (2016)  
 Hasenbusch (2010)  
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- textbook example and prime benchmark for **development of many-body methods**
- gapless **Dirac fermions** **not** in Ising/O(N) universality classes!

- critical point described by simple **continuum** field theory in  $D = 2+1$  dimensions
  - **Gross-Neveu model:**  $\mathcal{L}_{\text{GN}} = \bar{\psi} \gamma_\mu \partial_\mu \psi + \lambda_\psi (\bar{\psi} \psi)^2$

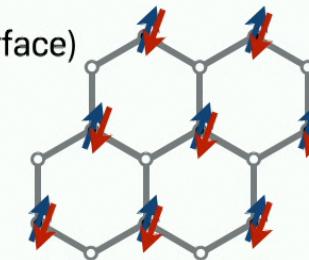
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- simplest **fermionic theory** w/ **critical point** (relativistic, no Fermi surface)

- example: **CDW** order of electrons in graphene Herbut (2006)

- **bosonized version** of model...



# Effective theories for phase transitions in Dirac systems

RUB

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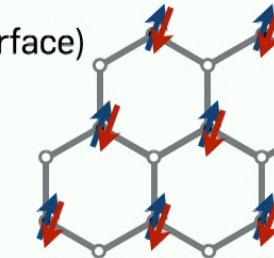
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• **Gross-Neveu-Yukawa model:**  $\mathcal{L}_{\text{GNY}} = \bar{\psi} (\gamma_\mu \partial_\mu + g\phi) \psi + \frac{1}{2} \phi (m^2 - \partial_\mu^2) \phi + \lambda \phi^4$



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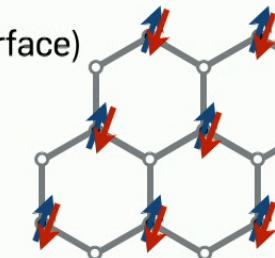
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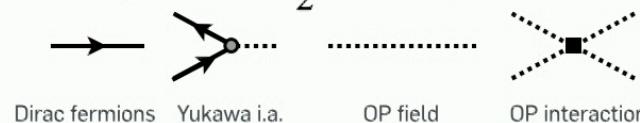
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- both models have **critical point** in  $2 < D < 4$  and lie in **same universality class**

⇒ **chiral Ising universality class**

(it's the Ising model of critical fermions!)

# Fermionic universality classes – developments

RUB

- towards **precision determination** of universality of interacting Dirac fermions

- **Monte Carlo methods**

- microscopic lattice models with 2<sup>nd</sup> order phase transition in GN universality class
- sign-problem free formulations

- Chandrasekharan & Li (2013)
- Wang, Corboz & Troyer (2014)
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- unprecedented precision for O(N) models
- extended to fermionic systems

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- **conformal bootstrap**

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- **renormalization group**

- functional RG
- higher-loop calculations adapted from hep-th...

# GNY model & renormalization group beta functions

- renormalized GNY Lagrangian:

$$\mathcal{L} = Z_\psi \bar{\psi} \gamma_\mu \partial_\mu - \frac{1}{2} Z_\phi (\partial_\mu \phi)^2 + Z_{\phi^2} \frac{m^2}{2} \phi^2 + Z_{\phi \bar{\psi} \psi} g \mu^{\epsilon/2} \phi \bar{\psi} \psi + Z_{\phi^4} \lambda \mu^\epsilon \phi^4$$

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- the  $Z_i$  relate bare and renormalized parameters
- **energy scale**  $\mu$  parametrizes RG flow
- RG  $\beta$  functions:

$$\beta_{g^2} = \frac{dg^2}{d \ln \mu}, \quad \beta_\lambda = \frac{d\lambda}{d \ln \mu}$$

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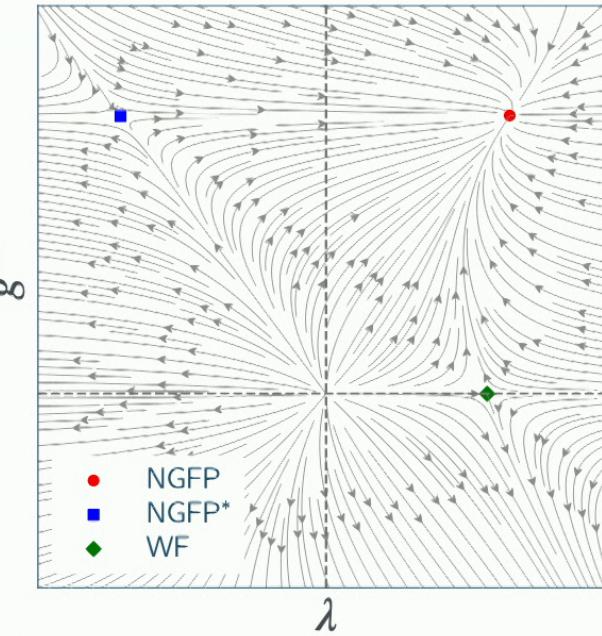
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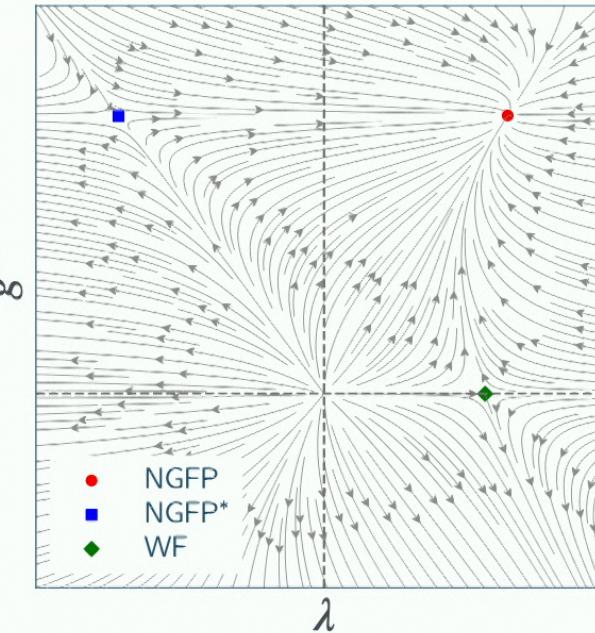
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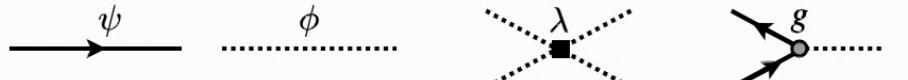
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- **continuous phase transition:**
  - **IR attractive fixed point**
    - ⇒ **critical exponents!**



## Renormalization group constants — tool chain

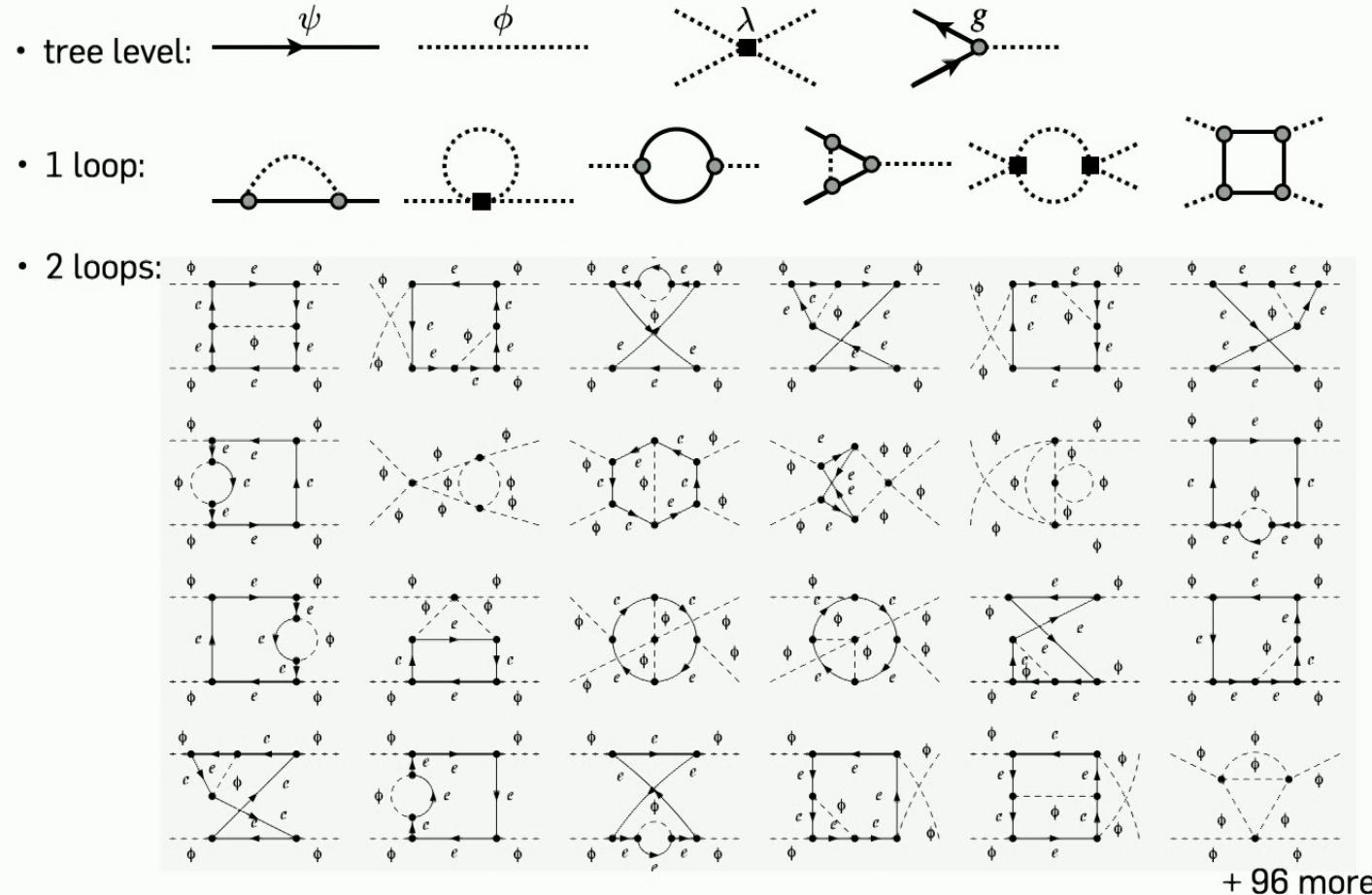
- evaluate renormalization group constants  $Z_i$  up to 4-loop order

• tree level:   $\psi$   $\phi$   $\lambda$   $g$

• 1 loop: 

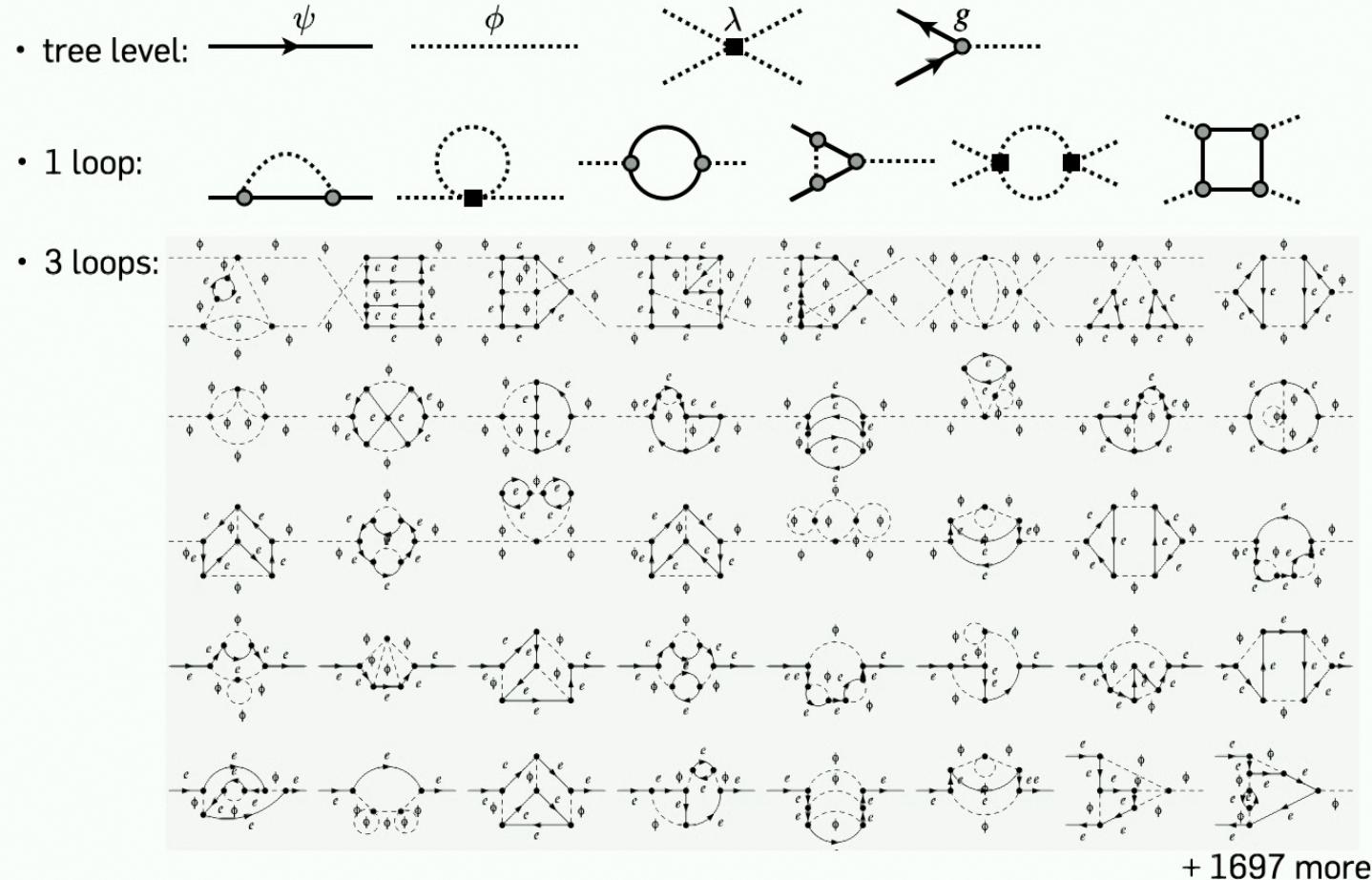
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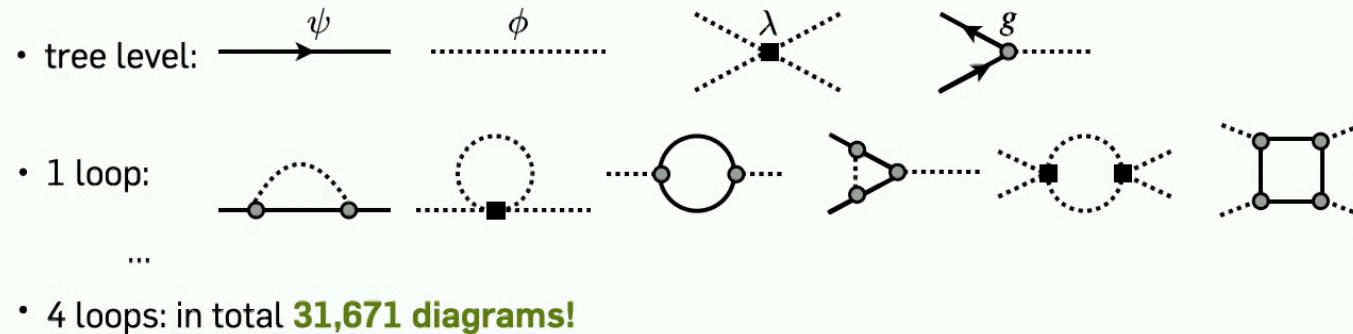
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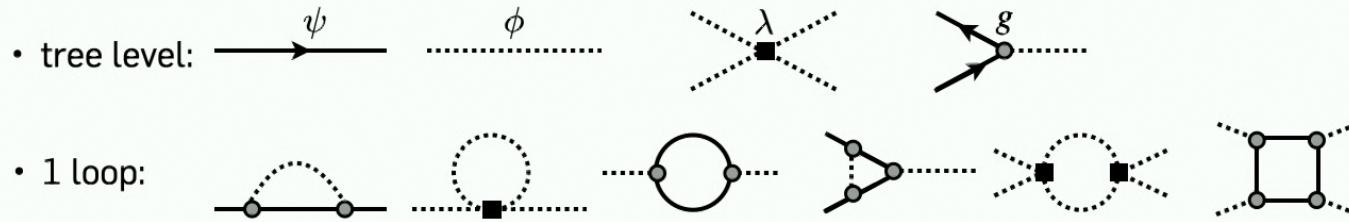
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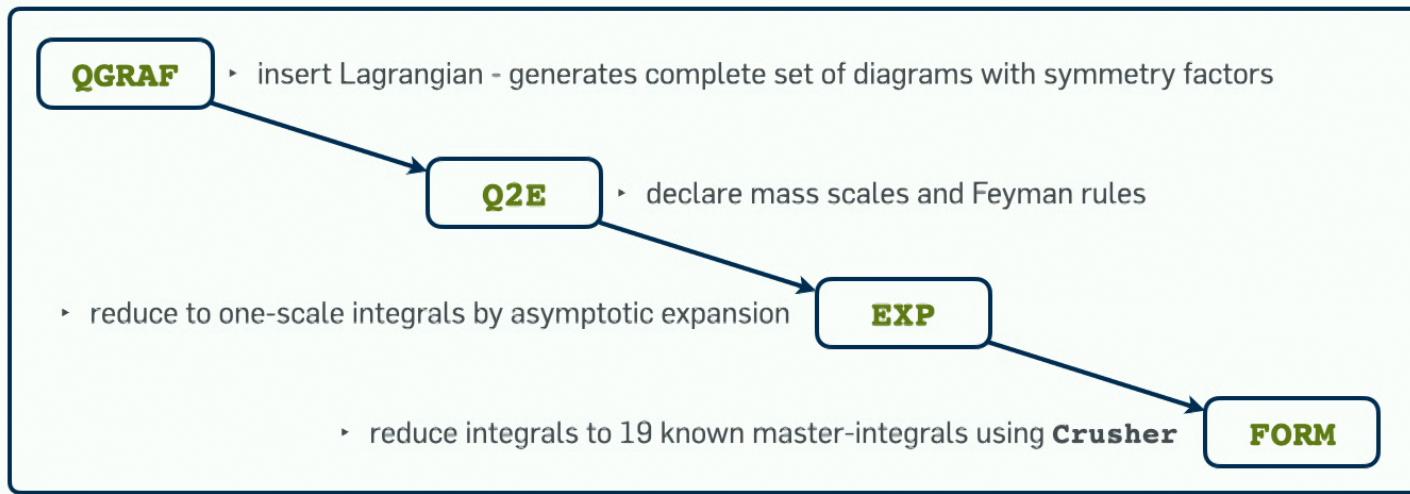
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- 4 loops: in total **31,671 diagrams!**

use tool chain developed for relativistic high-energy physics:



# epsilon expansion for chiral Ising model

- CDW transition in graphene:  $N_f = 2$  flavors of 4-component Dirac fermions

- critical exponents in  $D = 4 - \epsilon$

$$\frac{1}{\nu} = 2 - \frac{20\epsilon}{21} + \frac{325\epsilon^2}{44982} - \frac{(271572144\zeta_3 + 36133009)\epsilon^3}{3821940612} + \frac{(73192843310400\zeta_3 + 179520471709200\zeta_5 - 2472257012904\pi^4 - 86141171013035)\epsilon^4}{4175164363361040} + \mathcal{O}(\epsilon^5),$$

$$\eta_\phi = \frac{4\epsilon}{7} + \frac{109\epsilon^2}{882} + \left( \frac{1170245}{26449416} - \frac{144\zeta_3}{2401} \right) \epsilon^3 + \frac{(162669869280\zeta_3 + 171915696000\zeta_5 - 1203409872\pi^4 + 102456536695)\epsilon^4}{2407822585560} + \mathcal{O}(\epsilon^5),$$

$$\eta_\psi = \frac{\epsilon}{14} - \frac{71\epsilon^2}{10584} - \left( \frac{18\zeta_3}{2401} + \frac{2432695}{158696496} \right) \epsilon^3 + \frac{(1155813964920\zeta_3 + 515747088000\zeta_5 - 3610229616\pi^4 - 556332486445)\epsilon^4}{57787742053440} + \mathcal{O}(\epsilon^5),$$

 Zerf, Mihaila, Marquard, Herbut, MMS (2017)

## Quantum critical behavior of $N_f=2$ chiral Ising model

- GNY exponents from 4-loop RG in  $(4-\epsilon)$ :

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$$\nu^{-1} \approx 2 - 0.95\epsilon + 0.0072\epsilon^2 - 0.095\epsilon^3 - 0.013\epsilon^4$$

$$\eta_B \approx 0.57\epsilon + 0.12\epsilon^2 - 0.028\epsilon^3 + 0.15\epsilon^4$$

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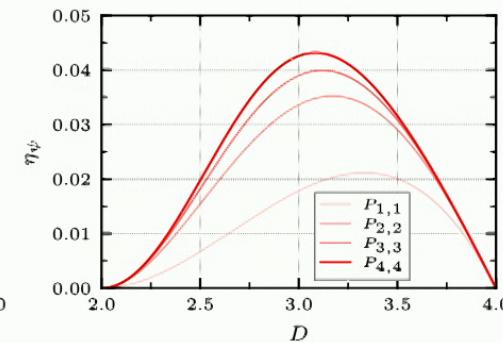
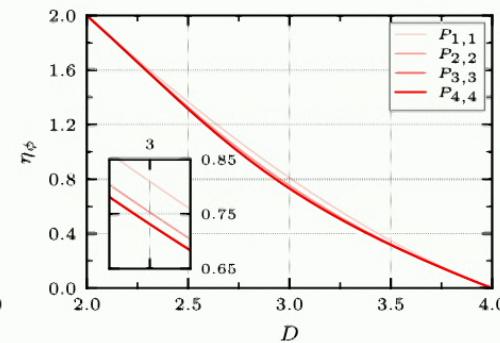
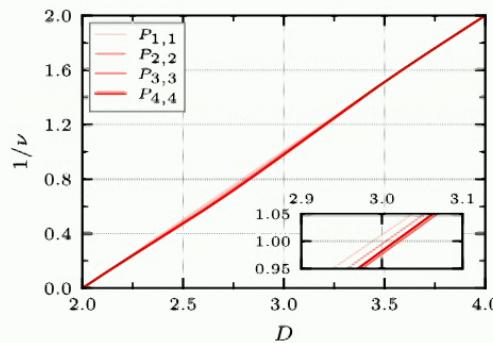
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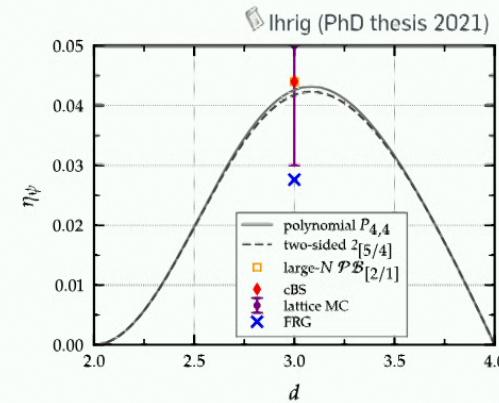
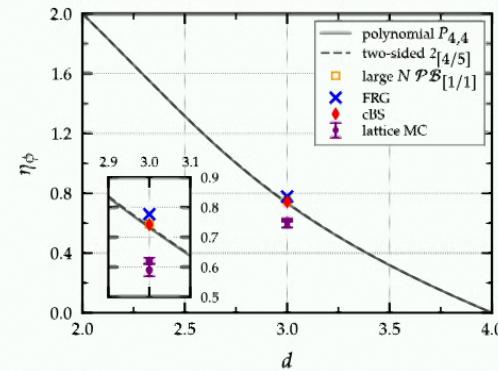
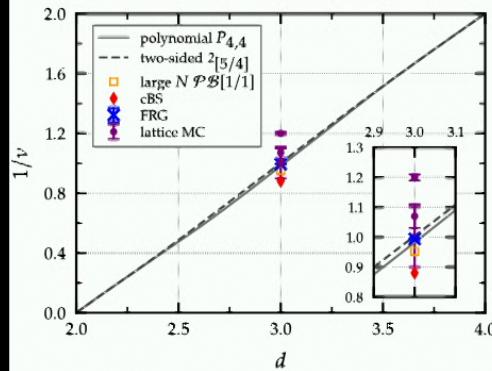
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Method	$1/\nu$	$\eta_B$	$\eta_F$
4-loop pert. RG, $\epsilon^4$	0.99	0.731	0.043
large- $N$ , $1/N^{2/3}$	0.95	0.743	0.044
conformal bootstrap	0.88	0.742	0.044
functional RG	0.994	0.7765	0.0274
Monte Carlo	1.20(1)	0.62(1)	0.38(1)
Monte Carlo	1.1(1)	0.60(2)	0.05(2)

Ihrig, Mihaila, MMS (2018)

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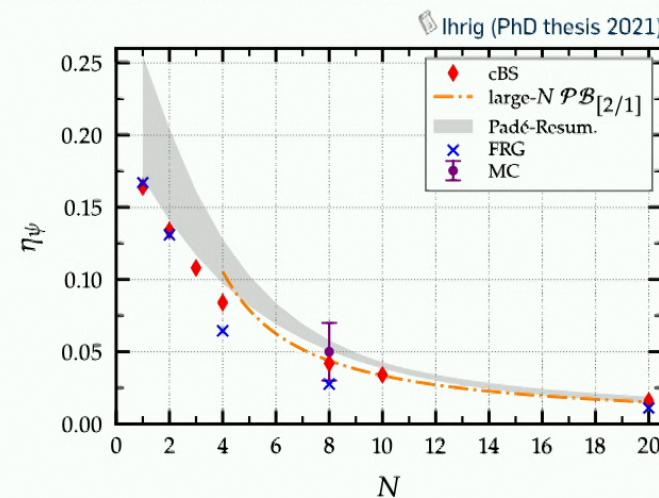
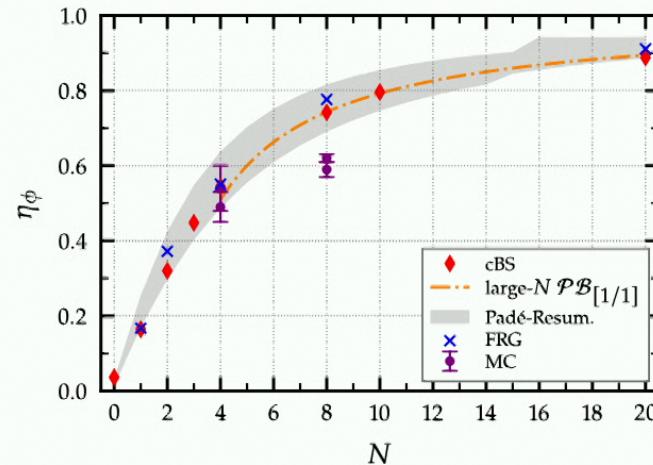
## Perspective for $N_f=2$ chiral Ising criticality

- anomalous dimensions  $\eta_\phi, \eta_\psi$ :
  - very good agreement of perturbative approaches and conformal bootstrap
  - discrepancies with Monte Carlo simulations
- correlation length exponent  $\nu$ :
  - ~10% discrepancies across the methods
  - update on conformal bootstrap ...  Erramilli, Iliesiu, Kravchuk, Liu, Poland, Simmons-Duffin (in progress)
- model with best chance to settle critical behavior across methods, soon #textbook

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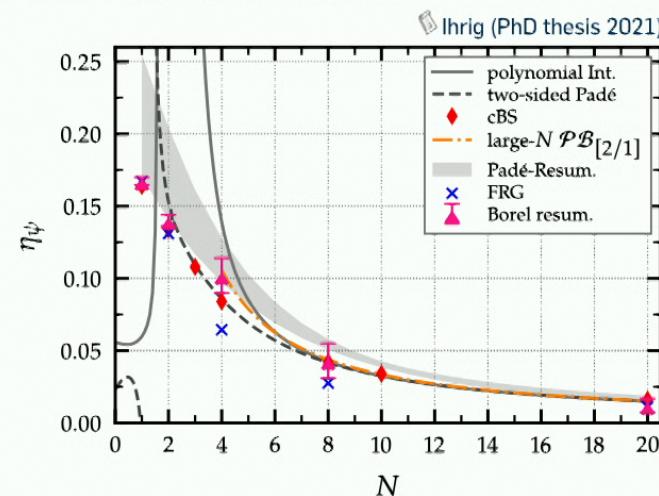
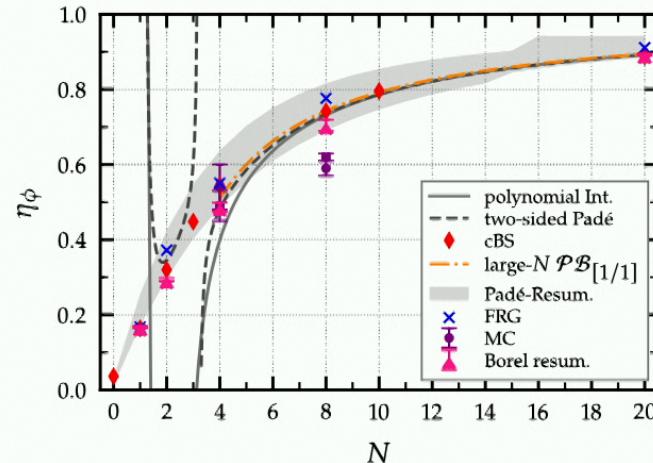
## Related models

- different choices for  $N_f = N/4$ 
  - ▶ results from  $4 - \epsilon$  expansion at 4-loop order with Padé approximants



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  - results from  $4 - \epsilon$  expansion at 4-loop order with Padé approximants



- critical exponents from  $2 + \epsilon$  expansion exhibit pole at  $N = 2$  ( $N_f = 1/2$ )
- impact of pole already shows at  $N = 4$  ( $N_f = 1$ )
- results from 2-sided interpolations not as good at  $N_f = 1$  as they are at  $N_f = 2$ !
- **chiral Ising model at  $N_f = 2$  is a sweet spot for comparisons of different methods!**

## Related models II

- antiferromagnetic transition in graphene ( $N_f = 2$ )

- e.g., Hubbard model on the honeycomb lattice

- **chiral Heisenberg model**

$$\mathcal{L}_{\text{GNY}} = \bar{\psi} \gamma_\mu \partial_\mu \psi + g \vec{\phi} \cdot \bar{\psi} \left( \vec{\sigma} \otimes \mathbb{I}_{2N_f} \right) \psi + \frac{1}{2} \vec{\phi} (m^2 - \partial_\mu^2) \vec{\phi} + \lambda \left( \vec{\phi} \cdot \vec{\phi} \right)^2$$

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- expansion in  $D = 4 - \epsilon$  at 4-loop order

$$\frac{1}{\nu} \approx 2 - 1.527\epsilon + 0.4076\epsilon^2 - 0.8144\epsilon^3 + 2.001\epsilon^4,$$

$$\eta_\phi \approx 0.8\epsilon + 0.1593\epsilon^2 + 0.02381\epsilon^3 + 0.2103\epsilon^4,$$

$$\eta_\psi \approx 0.3\epsilon - 0.05760\epsilon^2 - 0.1184\epsilon^3 + 0.04388\epsilon^4,$$

## Related models III

- 2-dimensional frustrated magnets with spin and orbital degrees of freedom
  - low-energy Dirac fermion excitations from fractionalization
- microscopic model leads to effective **Gross-Neveu-SO(3) model** (→ talk by L. Janssen)

$$\mathcal{L} = \bar{\psi} \gamma_\mu \partial_\mu \psi + \frac{1}{2} \phi_a (-\partial_\mu^2 + m^2) \phi_a + \lambda (\phi_a \phi_a)^2 - g \phi_a \bar{\psi} (\mathbb{I}_{2N/3} \otimes L_a) \psi$$

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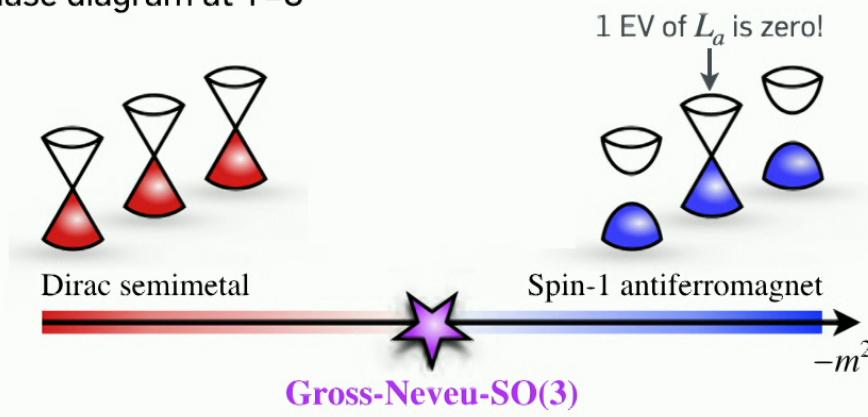


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- mean-field phase diagram at T=0



- only  $2N/3$  of the Dirac cones acquire mass gap and  $1N/3$  remain massless

# Fractionalized quantum criticality in spin-orbital liquids

- critical exponents from field-theory beyond the leading order:

1. expansion in  $D = 4 - \epsilon$  at 3-loop order

$$\frac{1}{\nu} \approx 2 - 0.917\epsilon - 0.0340\epsilon^2 - 0.0735\epsilon^3 + \mathcal{O}(\epsilon^4)$$

$$\eta_\phi \approx 0.333\epsilon + 0.0922\epsilon^2 - 0.0338\epsilon^3 + \mathcal{O}(\epsilon^4)$$

$$\eta_\psi \approx 0.167\epsilon + 0.0360\epsilon^2 - 0.0303\epsilon^3 + \mathcal{O}(\epsilon^4)$$

2. large- $N$  expansion at order  $1/N^2$  for  $\nu, \eta_\phi$  and at order  $1/N^3$  for  $\eta_\psi$

$$\frac{1}{\nu} \approx 1 - \frac{1.621}{N} + \frac{19.922}{N^2} + \mathcal{O}(1/N^3)$$

$$\eta_\phi \approx 1 - \frac{2.026}{N} + \frac{1.564}{N^2} + \mathcal{O}(1/N^3)$$

$$\eta_\psi \approx \frac{0.4053}{N} + \frac{1.0403}{N^2} - \frac{0.7972}{N^3} + \mathcal{O}(1/N^4)$$

3. functional renormalization group...

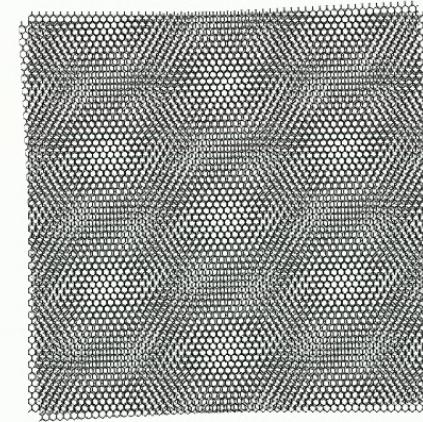
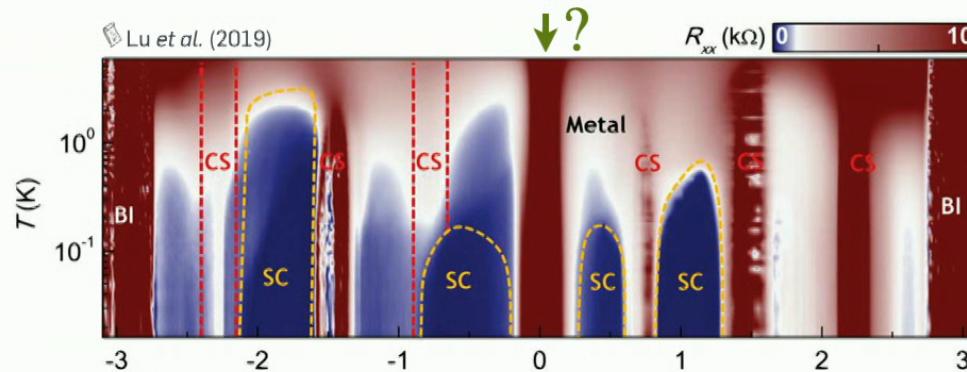
 Ray, Ihrig, Kruti, Gracey, MMS, Janssen (2021)

## Conclusions

- even at high loop order: no guarantee to get quantitative results!
- sometimes it seems to work, though!
  - **chiral Ising model at  $N_f = 2$  is a sweet spot for comparisons of different methods!**
  - **Q:** update from Monte Carlo? **#textbook-example**

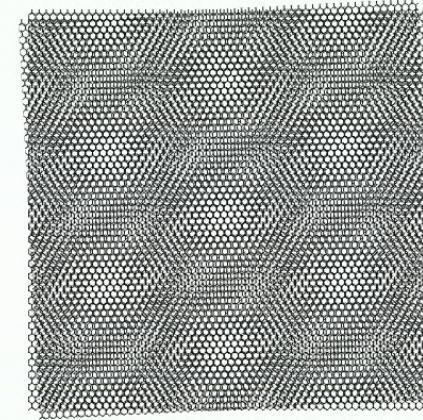
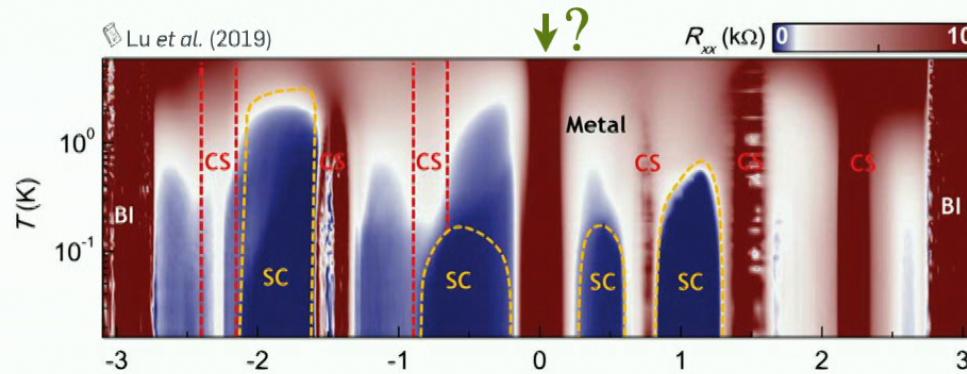
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- chiral Heisenberg model at 4 loops less clear (also no conformal bootstrap data)
- abelian Higgs model  $n_c \approx 183(1 - 1.75\epsilon + 0.80\epsilon^2 + 0.36\epsilon^3)$   $\rightarrow n_c(d=3) \approx 12 \pm 4$

Ihrig, Zerf, Marquard, Herbut, MMS (2019)