Title: Comments on the volume of the black hole interior in 2d gravity

Speakers: Gabor Sarosi

Series: Quantum Fields and Strings

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Abstract: The volume of the interior of a two-sided eternal black hole classically grows forever. I will derive a microscopic formula for this volume in JT gravity by summing the non-perturbative contribution of higher topologies. The non-perturbative corrections lead to the saturation of the volume of the interior at times exponential in the entropy of the black hole. I will connect the microscopic formula for the volume with properties of a "thermo-averaged" density matrix, in particular with its second Renyi entropy, which I argue to measure the number of nearly orthogonal states visited by time evolution. I will discuss various problems with this interpretation.

Zoom Link: https://pitp.zoom.us/j/99880539557?pwd=VWZiMFI5Rk5pdnNZRE5kalRvUk1Zdz09

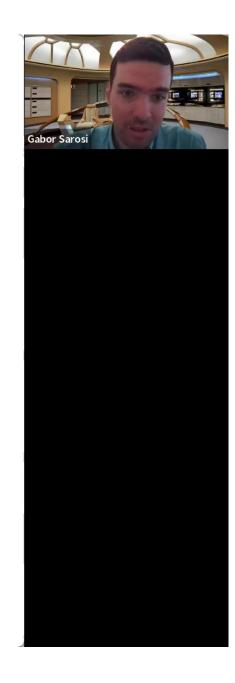
Pirsa: 22040122 Page 1/29

# Comments on the volume of the black hole interior in 2d gravity

Gabor Sarosi (CERN)

Based on:
2107.06286
with Luca Iliesiu & Mark Mezei

and w.i.p. with with Luca Iliesiu & Mark Mezei & Alexey Milekhin



Pirsa: 22040122 Page 2/29

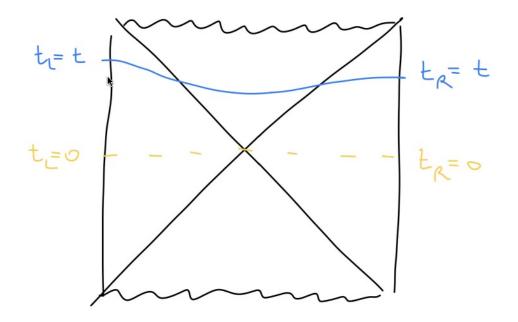
# Plan

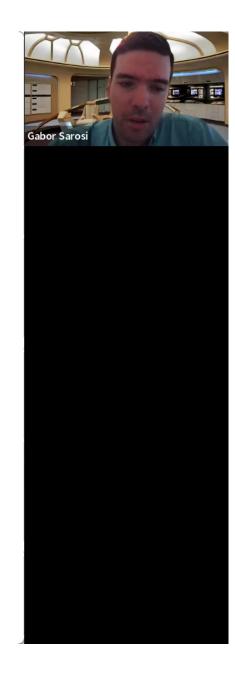
- Motivation
- The volume of the interior in JT gravity
- Microscopic interpretation



Pirsa: 22040122 Page 3/29

Black hole interior: an expanding cosmology

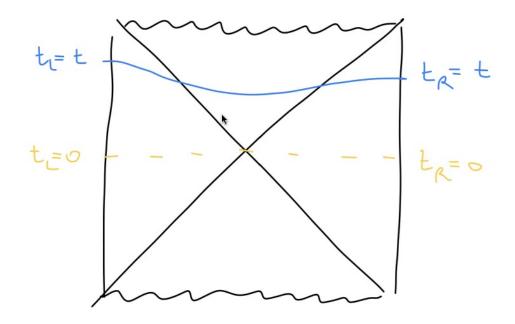




3

Pirsa: 22040122 Page 4/29

Black hole interior: an expanding cosmology



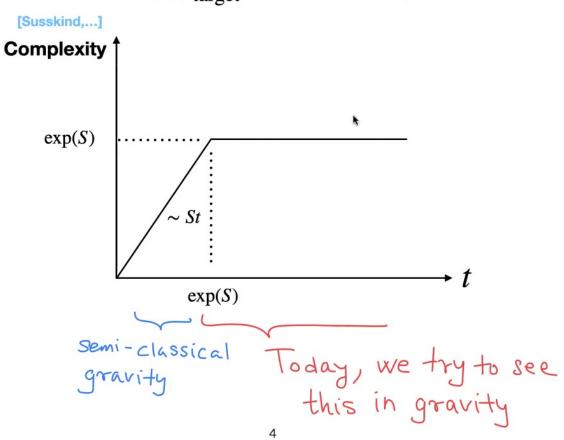
volume of maximal Cauchy slice  $\propto Mt$ 

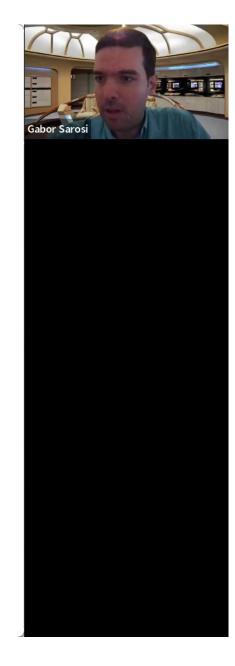
Question: what is the microscopic origin of this "creation of space"?

Pirsa: 22040122 Page 5/29

#### Time dependence

$$|\psi\rangle_{\text{target}} = e^{-iHt}|TFD\rangle$$





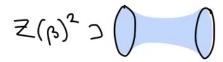
Pirsa: 22040122

Why do we have a chance of seeing something like this in gravity?

Because Euclidean gravity seems to have an unreasonably large regime of validity, e.g.

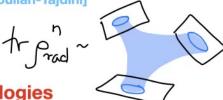
Universal energy level repulsion of chaotic systems

[Saad-Shenker-Stanford,Cotler-Jensen,...]



**Entropy Page curve of an evaporating black hole** 

[Penington-Shenker-Stanford-Yang, Amheiri-Hartman-Maldacena-Shaghoulian-Tajdini]



Both are about corrections coming from summing topologies

Today: how do these effects affect the volume of the interior?

5



# Plan

- Motivation
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Pirsa: 22040122 Page 8/29

### Jackiw-Teitelboim (JT) gravity

Dilaton-gravity in two dimensions:

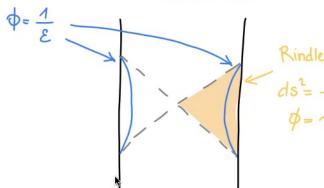
$$I_{JT} = -S_0 \chi(\mathcal{M}) - \frac{1}{2} \int \sqrt{g} \phi(R+2) - \int_{\partial \mathcal{M}} \sqrt{h} \phi(K-1)$$

Arises by dimensional reduction of near-horizon region of near-extremal black holes

[Maldacena-Stanford-Yang,Sarosi,Nayak-Shukla-Soni-Trivedi]

Black hole solution: Global  $AdS_2$  with a cutoff at large constant  $\phi$ 

#### Lorentzian

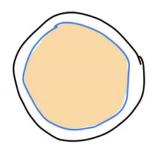


Rindler patch

$$ds^{2} = -(r^{2}r_{s}^{2})dt^{2} + \frac{dr^{2}}{r^{2}r_{s}^{2}}$$

$$d = r$$

#### **Euclidean**



$$ds^2 = d\rho^2 + \sinh^2 \theta dt^2$$

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Pirsa: 22040122 Page 9/29

### Jackiw-Teitelboim (JT) gravity

Higher topologies also contribute, weighted by  $\chi(\mathcal{M}) = 2g + n - 2$ 

[Saad-Shenker-Stanford]

$$Z(\beta) =$$
 +  $+ \cdots$ 

JT gravity is a matrix integral

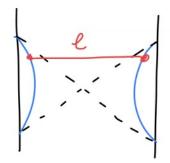
$$[Z(\beta)]_{\text{JT gravity}} \simeq \int dH e^{-V(H)} \text{Tr}[e^{-\beta H}]$$

In the genus expansion and in the double scaling limit (large matrix, zoomed to the bottom of the spectrum)

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Pirsa: 22040122 Page 10/29

#### Classical volume



$$\ell_{\rm ren} = 2 \log \left( 2 \cosh \left[ \frac{2\pi}{\beta} t \right] \right) \approx \frac{4\pi}{\beta} t$$

#### Quantum volume (perturbative): using HH wave function [Yang, Harlow-Jafferis]

$$\psi_{\beta/2}^{\text{Disk}}(\mathcal{E}) =$$

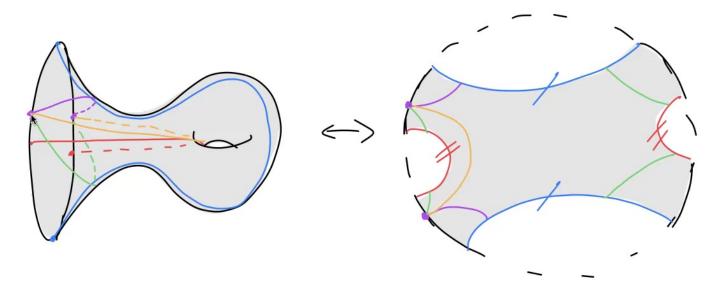
$$= \int_0^\infty dE \rho_0(E) e^{-\beta E/2} [4e^{-\ell/2} K_{i\sqrt{8E}} (4e^{-\ell/2})]$$

$$\langle \ell \rangle =$$

$$= \frac{e^{-S_0}}{Z_{\text{disk}}} \int e^{\ell} d\ell \, |\psi_{\frac{\beta}{2} + it}^{\text{Disk}}(\ell)|^2 \ell$$

Non-perturbative quantum volume:

Challenge: infinite number of extremal geodesics on higher genus surfaces



Taking minimal geodesic on each surface is not an option:
we want to continue to Lorentzian!
Euclidean minimal geodesic changes abruptly:
leads to non-analiticity

10

Pirsa: 22040122 Page 12/29

Non-perturbative quantum volume:

Challenge: infinite number of extremal geodesics on higher genus surfaces

Prescription: average over a well defined set of extremal geodesics

$$\langle \ell \rangle \equiv \sum_{g} e^{S_0(1-2g)} \sum_{\gamma} \langle \ell_{\gamma} \rangle_{\text{wiggles \& moduli space}}$$

**Divergent!** 

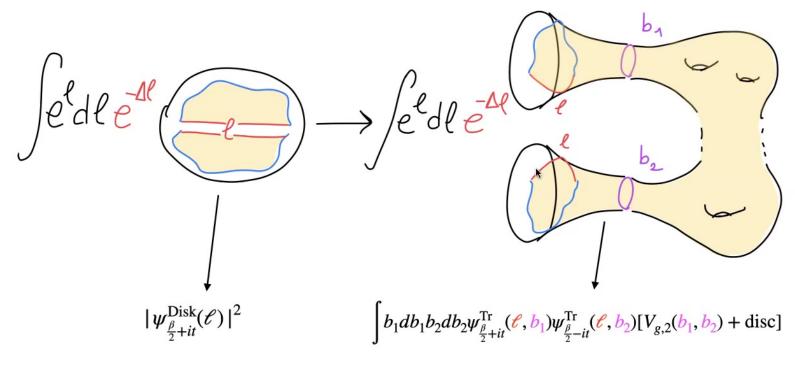
**Natural regularization:** 

$$\langle \ell \rangle \equiv -\lim_{\Delta \to 0} \frac{d}{d\Delta} \sum_{g} e^{S_0(1-Qg)} \sum_{\gamma} \langle e^{-\Delta \ell_{\gamma}} \rangle_{\text{wiggles \& moduli space}}$$

Related to two point function [Yang,Saad,...]

11

Pirsa: 22040122 Page 13/29



(sum over non-int geodesics turns 1bndy moduli space into 2bndy moduli space [Mirzakhani])

12

Pirsa: 22040122 Page 14/29

#### **Result:**

$$\langle \mathcal{E}(t) \rangle = -\frac{e^{-S_0}}{4\pi^2 Z_{\rm disk}(\beta)} \int_0^\infty dE_1 dE_2 \langle \rho(E_1) \rho(E_2) \rangle M(E_1, E_2) e^{-\frac{1}{2}\beta(E_1 + E_2) - i(E_1 - E_2)t}$$

$$M(E_1, E_2) = \frac{16\pi^4}{(E_1 - E_2)[\cosh 2\pi\sqrt{2E_1} - \cosh 2\pi\sqrt{2E_2}]}$$

$$\propto \frac{1}{(E_1 - E_2)^2}$$

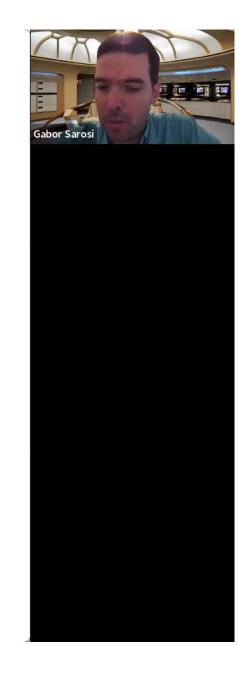
#### In matrix integrals, universally:

$$\langle \rho(E_1)\rho(E_2)\rangle = \langle \rho(E_1)\rangle \langle \rho(E_2)\rangle + \langle \rho(E_1)\rangle \delta(E_1-E_2) - \frac{\sin^2\left[\pi\langle \rho(E_2)\rangle(E_1-E_2)\right]}{\uparrow}$$

$$\frac{\text{disk}}{\text{time indep}}$$

$$\frac{\text{divergence}}{\left(\triangle \text{regulated}\right)}$$

$$\frac{13}{}$$



#### **Result:**

$$\langle \mathcal{E}(t) \rangle = -\frac{e^{-S_0}}{4\pi^2 Z_{\rm disk}(\beta)} \int_0^\infty dE_1 dE_2 \langle \rho(E_1) \rho(E_2) \rangle M(E_1, E_2) e^{-\frac{1}{2}\beta(E_1 + E_2) - i(E_1 - E_2)t}$$

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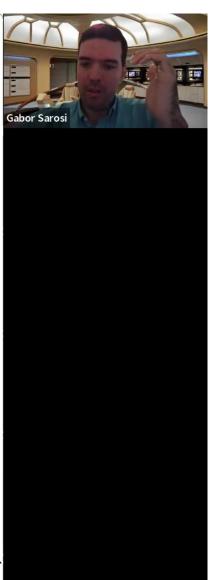
$$\Rightarrow \pi^2(E_1-E_2)^{\frac{1}{2}}$$

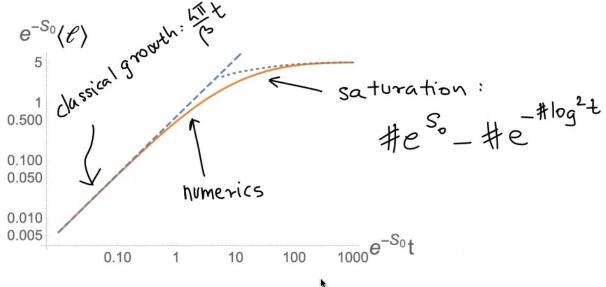
$$\Rightarrow \text{doubly}$$

$$\text{divergence}$$

$$\text{divergence}$$

$$\text{Aregulated}$$







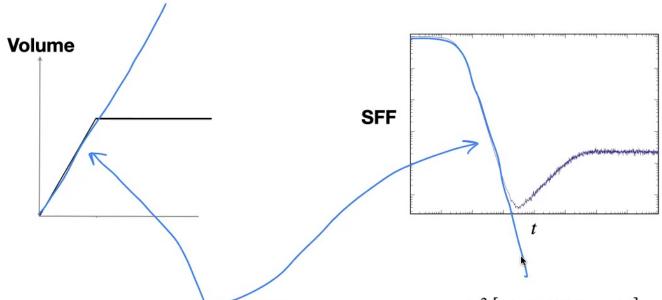
14

Pirsa: 22040122 Page 17/29

Very similar quantity: spectral form factor  $\overline{Z(\beta - it)Z(\beta + it)} \equiv SFF(t)$ 

$$\overline{Z(\beta - it)Z(\beta + it)} \equiv SFF(t)$$

$$\int_0^\infty dE_1 dE_2 \langle \rho(E_1) \rho(E_2) \rangle e^{-\frac{1}{2}\beta(E_1+E_2)-i\left(E_1-E_2\right)t} \times \begin{cases} M(E_1,E_2) & \text{for } \langle \mathscr{E}(t) \rangle \\ 1 & \text{for SFF}(t) \end{cases}$$



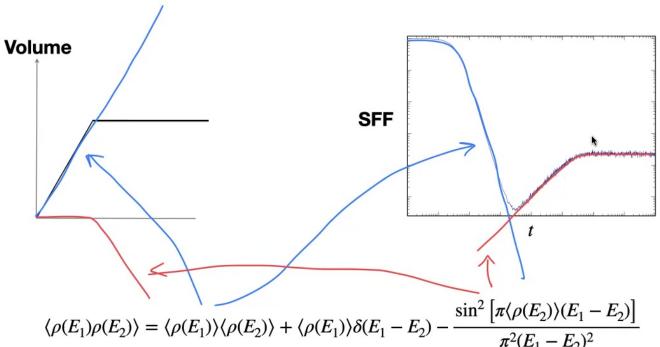
$$\langle \rho(E_1)\rho(E_2)\rangle = \langle \rho(E_1)\rangle\langle \rho(E_2)\rangle + \langle \rho(E_1)\rangle\delta(E_1 - E_2) - \frac{\sin^2\left[\pi\langle \rho(E_2)\rangle(E_1 - E_2)\right]}{\pi^2(E_1 - E_2)^2}$$

15

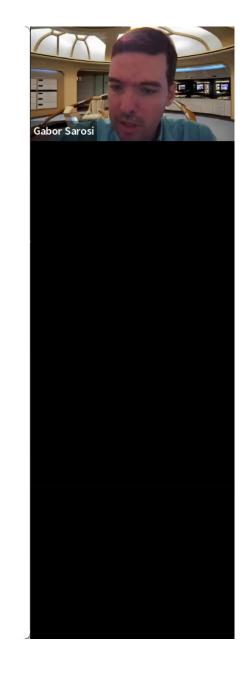
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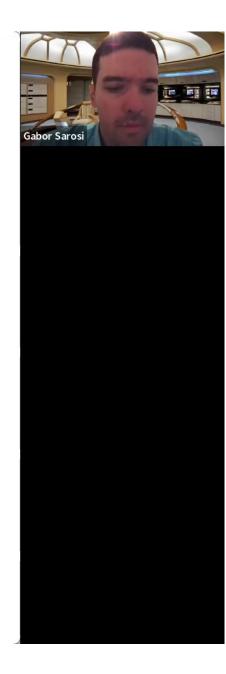


15



# Plan

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Pirsa: 22040122 Page 20/29

### Spectral complexity

The final formula for the volume makes sense in any matrix integral, can also be derived for more general dilaton potential (with methods of [Maxfield-Turiaci,Witten])

Call it spectral complexity

$$\mathscr{C}(t) \equiv \frac{1}{Z(\beta)^2} \sum_{i \neq j} \frac{1}{[E_i - E_j]^2} e^{-\beta(E_i + E_j)} [1 - \cos(E_i - E_j)t]$$

Can we interpret this quantity?

17

Pirsa: 22040122 Page 21/29

### Time-averaged density matrix

An extremely simple proxy for complexity:

How much of the Hilbert space is spanned by time evolved thermofield double states up to a given time?

$$\{e^{-iHt}|TFD\rangle|t < T\}$$

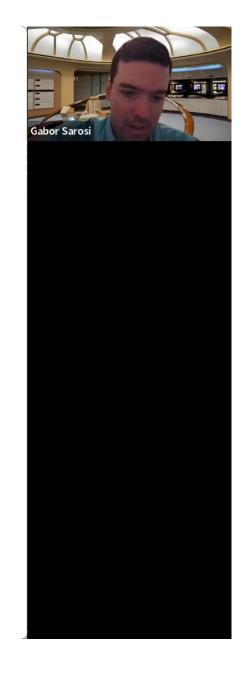
**Consider the density matrix:** 

$$\rho(T) = \frac{1}{T} \int_{0}^{T} dt e^{-iHt} |TFD\rangle \langle TFD| e^{iHt}$$

First guess: rank $\rho$ 

However, instantly full rank due to extremely small eigenvalues Replace rank with number of eigenvalues bigger than some threshold:

$$\mathbf{rank}_{\epsilon}\rho \equiv \sum_{\substack{\lambda_i > \epsilon \\ 18}} 1 = \mathbf{Tr}\Theta(\rho - \epsilon)$$



Pirsa: 22040122 Page 22/29

$${\rm rank}_{\epsilon}\rho \equiv \sum_{\lambda_i > \epsilon} 1 = {\rm Tr}\Theta(\rho - \epsilon)$$

How is it a proxy for complexity?

Consider an ensemble of states  $\mathscr{E}$ ,  $C(\mathscr{E}) \equiv \min$  num. of gates to prep.  $\mathscr{E}$ 

then 
$$C(\mathcal{E}) \ge \frac{\log |\mathcal{E}|}{\log \text{choices}}$$
 [Roberts-Yoshida]

19

Pirsa: 22040122 Page 23/29

$$\mathbf{rank}_{\boldsymbol{\epsilon}}\rho \equiv \sum_{\lambda_{\boldsymbol{i}}>\boldsymbol{\epsilon}} 1 = \mathbf{Tr}\Theta(\rho - \boldsymbol{\epsilon})$$

How is it a proxy for complexity?

Consider an ensemble of states  $\mathscr{E}, \quad C(\mathscr{E}) \equiv \min$  num. of gates to prep.  $\mathscr{E}$ 

then 
$$C(\mathcal{E}) \ge \frac{\log |\mathcal{E}|}{\log \text{choices}}$$
 [Roberts-Yoshida]

 $C(\mathscr{E})$  is also (roughly) the complexity of the most complex state in  $\mathscr{E}$ 

19

Consider an ensemble of states  $\mathscr{E}$ ,  $C(\mathscr{E}) \equiv \min$  num. of gates to prep.  $\mathscr{E}$ 

then 
$$C(\mathscr{C}) \ge \frac{\log |\mathscr{C}|}{\log \text{choices}}$$
 [Roberts-Yoshida]

Take  $\mathscr{E}_T = \text{subspace with } \rho(T) \text{ eigenvalues } > \epsilon$ 

$$|\,\mathscr{C}_T| \propto \operatorname{vol}[\mathbb{C}P^{\mathrm{rank}_{\varepsilon}\rho(T)-1}] \propto \pi^{\mathrm{rank}_{\varepsilon}\rho(T)-1}$$

$$C(\mathscr{E}_T) \ge \#(\operatorname{rank}_{\varepsilon} \rho(T) - 1)$$

Tempting to think that the most complex state in the set is  $U_T | TFD 
angle$ 

Superposing  $U_{\delta t}|\psi\rangle$ ,  $U_{\delta t}^2|\psi\rangle$ ,  $U_{\delta t}^3|\psi\rangle$  is cheap but does have an overhead [Childs-Kothari-Somma]

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20

Pirsa: 22040122 Page 25/29

$$C(\mathcal{E}_T) \ge \#(\operatorname{rank}_{\epsilon} \rho(T) - 1)$$

$${\rm rank}_{\epsilon}\rho \equiv \sum_{\lambda_i > \epsilon} 1 = {\rm Tr}\Theta(\rho - \epsilon)$$

Difficult to compute, but there is a simple lower bound

Cauchy-Schwarz:  $\mathrm{Tr}[\rho\sqrt{\Theta(\rho-\epsilon)}] \leq \sqrt{\mathrm{Tr}\Theta(\rho-\epsilon)\mathrm{Tr}\rho^2}$ 

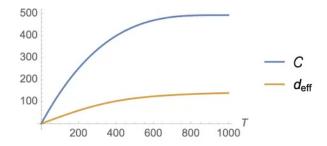
Minimized for  $\rho = \frac{1}{d}I$  (use regularized  $\Theta$  to show)

$$\mathrm{rank}_{\epsilon}\rho \geq \frac{1}{\mathrm{Tr}\rho^2} \equiv d_{\mathrm{eff}} \quad ^{\blacktriangleright} \ \, \mathrm{provided} \,\, \epsilon < d^{-1}$$

21

#### Connection to the volume

$$\rho(T) = \frac{1}{T} \int_{0}^{T} dt e^{-iHt} |TFD\rangle \langle TFD| e^{iHt}$$



#### Pro:

- · Proxy for complexity
- Direct function of the late time volume in gravity

$$d_{\text{eff}} = \frac{1}{\operatorname{Tr}\rho(T)^2} = \frac{1}{\frac{Z(2\beta)}{Z(\beta)^2} + \frac{\mathscr{C}(2T)}{T^2}}$$

$$\frac{1}{Z(\beta)^2} \sum_{i \neq j} \frac{1}{[E_i - E_j]^2} e^{-\beta(E_i + E_j)} [1 - \cos(E_i - E_j)T]$$

$$\propto \mathcal{E}(2T) \text{ for large } T$$

#### Con:

Early time slope different

$$d_{ ext{eff}} pprox rac{T^2}{\mathscr{C}(2T)} pprox rac{eta T}{\gamma^2}$$

 Saturation independent of volume saturating

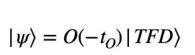


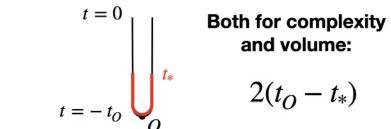
22

Pirsa: 22040122 Page 27/29

#### More cons: Switchback states

We should understand the dual of the volume in other states than the TFD





$$2(t_O - t_*)$$

**Effective dimension** 

$$\begin{split} \rho &= \int_0^{-t_O} U_t |\mathit{TFD}\rangle \langle \mathit{TFD}\,|\, U_{-t} \end{split}$$
 
$$+ \int_0^{t_O} U_t O\Big[U_{-t_O}|\, \mathit{TFD}\rangle \langle \mathit{TFD}\,|\, U_{t_O} \Big] O^\dagger U_{-t} \end{split}$$

Claim:  $Tr \rho^2$  depends only on 2pt func. Does not know about scrambling

Volume in JT

One can derive a microscopic formula involving limits of 6j symbol

Stay tuned...



23

Pirsa: 22040122 Page 28/29

#### Summary

Defined regularized non-perturbative volume in JT gravity Showed that it saturates at  $t \propto e^{S_0}$ , which is expected from complexity Result is similar to SFF, but the origin of "ramp" and plateau are different Volume is dual to a simply calculable quantity, dubbed spectral complexity

#### Questions .

Can spectral complexity match some definition of complexity? Or give a bound?

Is the time-averaged density matrix relevant?

What can we say about other states than the TFD?

24

Pirsa: 22040122 Page 29/29