Title: Quasi-sterile neutrinos from dark sectors

Speakers: Daniele Alves

Series: Particle Physics

Date: March 15, 2022 - 1:00 PM

URL: https://pirsa.org/22030108

Abstract: Quasi-sterile neutrinos are a natural consequence of dark sectors interacting with the Standard Model (SM) sector via neutrino- and vector-portals. Essentially, quasi-sterile neutrinos are light dark sector fermions with two generic properties: (i) they mix with the active neutrinos of the SM, and (ii) they are charged under a vector mediator that couples feebly to SM matter. Various interesting phenomenological consequences result from this class of particles. In this talk, I will discuss one such consequence: new, beyond the SM matter effects that can alter in-medium neutrino oscillations. In particular, for special windows of energy and matter densities, active neutrinos can resonantly oscillate into sterile neutrinos. I will discuss how this feature could be exploited to build a quasi-sterile neutrino model that can explain the MiniBooNE and LSND anomalies, while remaining compatible with observations from long-baseline reactor- and accelerator-based neutrino experiments. I will also discuss the implications to this model from the recent MicroBooNE results, and, time permitting, from solar neutrino measurements.

Zoom Link: https://pitp.zoom.us/j/95362454210?pwd=RzNCeDRkMjJXb2xNbHNkcEc1ZTdkUT09

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# Quasi-Sterile Neutrinos from Dark Sectors

I. BSM Matter Effects in Neutrino Oscillations and the Short-Baseline Anomalies

Daniele S. M. Alves

LANL

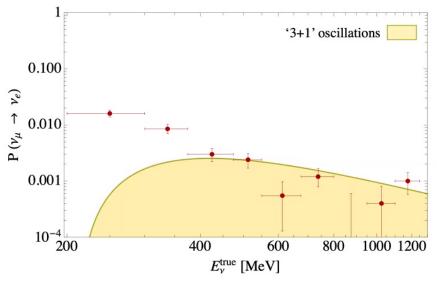
Based on arXiv: 2201.00876 w/ Bill Louis and Patrick deNiverville (and work in progress)

Particle Physics Seminar, Perimeter Institute, March 15th, 2022

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#### The MiniBooNE Low Energy Excess



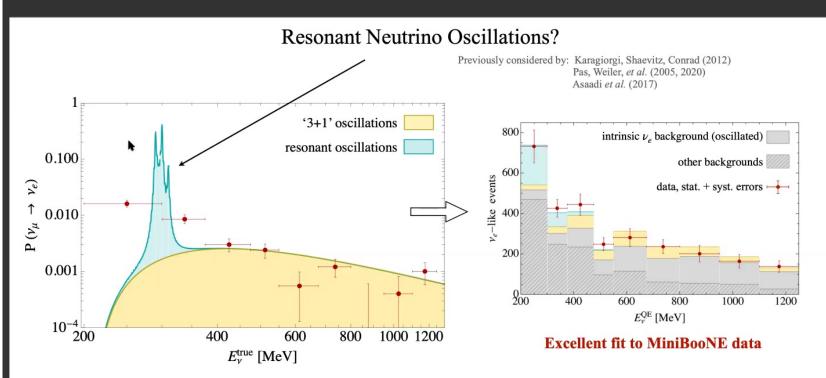


\* data points assume  $E_{\nu}^{true} \approx E_{\nu}^{CCQE}$ , therefore this should be taken as a qualitative depiction of the excess

The MiniBooNE appearance data is compatible, with the '3+1' sterile neutrino interpretation of the LSND excess in the mid-to-high energy range,  $E_{\nu} \gtrsim 400$  MeV, but an excess at lower energies  $E_{\nu} \lesssim 400$  MeV persists

Unaccounted backgrounds or systematics? Or genuine BSM signal?

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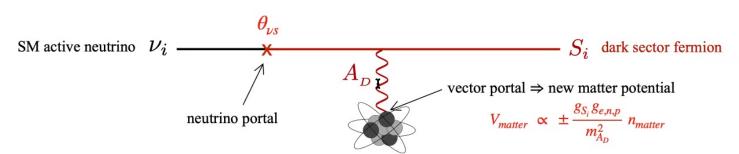
Daniele Spier

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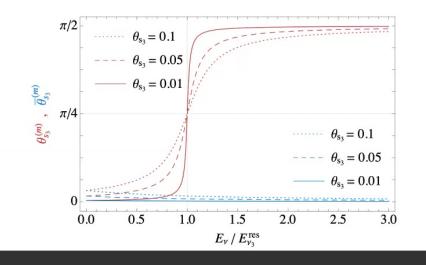
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## Quasi-Sterile Neutrinos and New Matter Effects



New matter potential acting on quasi-sterile neutrino induces resonant active-to-sterile oscillations when the energy and matter potential satisfy the relation:

$$E_{\nu}^{res} = \frac{\delta M_{\nu s}^2 \cos(2\theta_{\nu s})}{2 V_{matter}}$$

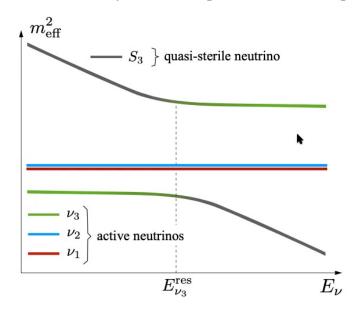


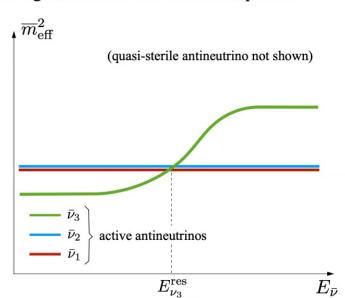


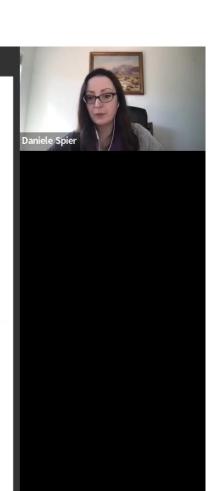
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## Quasi-Sterile Neutrinos and New Matter Effects (toy model)

Effectively, the matter potential alters the propagation eigenstates and their oscillation phases



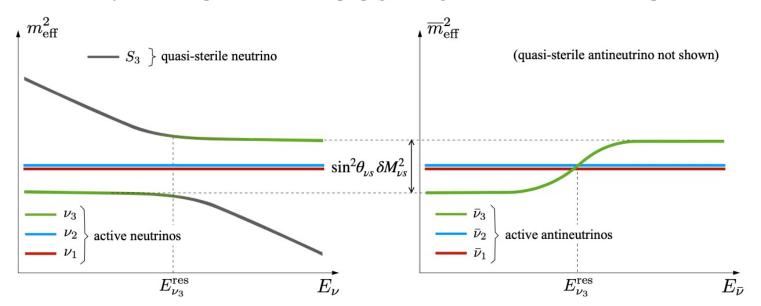




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## Quasi-Sterile Neutrinos and New Matter Effects (toy model)

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In order to preserve the active neutrino spectrum at high energies, we need:

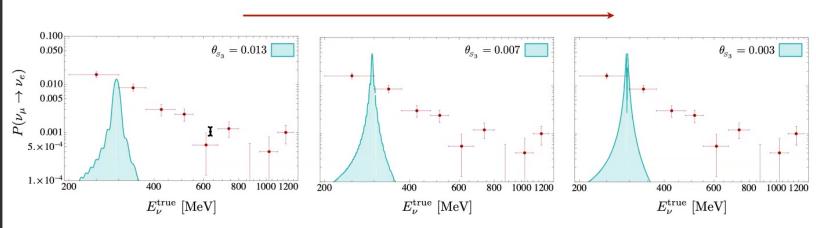
$$\sin^2\!\theta_{\nu s}\,\delta M_{\nu s}^2 \sim 2\,\delta m_{32}^2$$

⇒ this constrains the resonance region to be very narrow



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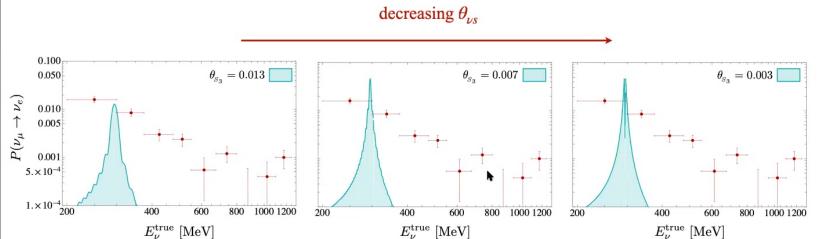




- width of resonance is nearly independent of  $\theta_{\nu s}$
- resonance peak is saturated at  $P(\nu_{\mu} \to \nu_{e})_{\rm max} = 4 |U_{e3}|^2 |U_{\mu3}|^2$  once the active-sterile mixing angle  $\theta_{\nu s}$  is below a critical value  $\theta_{\nu s}^{\rm min}$



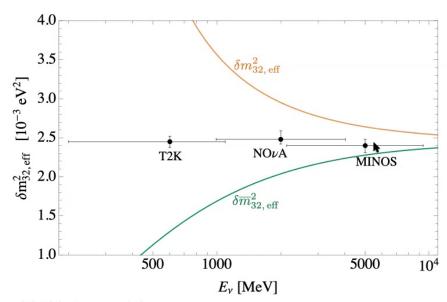
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Three issues with this toy model:

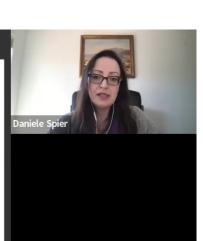
- 1. resonant peak is not large enough to explain the MiniBooNE low energy excess
- 2.  $\nu_{\mu} \rightarrow \nu_{e}$  oscillations vanish outside the resonance region



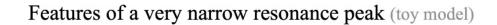


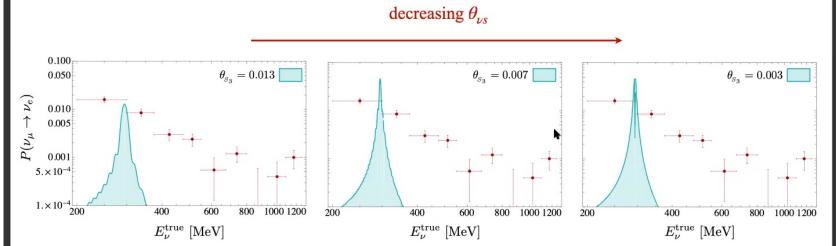
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- 3. the effective  $\delta m_{32}^2$  does not relax fast enough to the SM value at long baseline experiments



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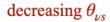


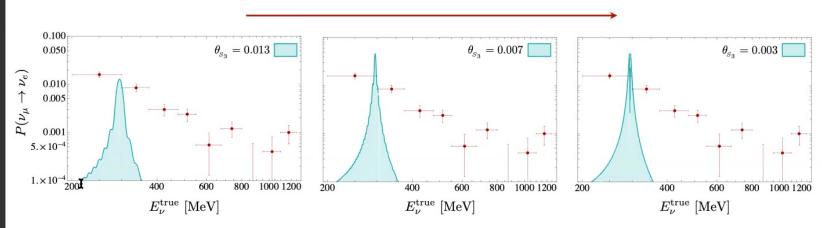


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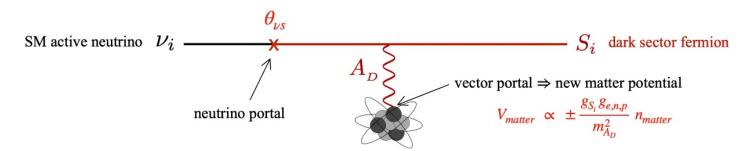


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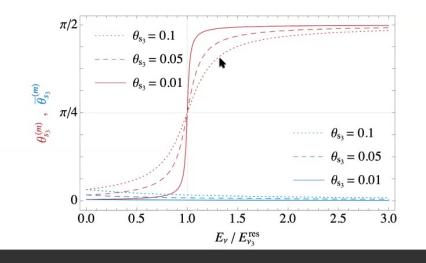
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## Quasi-Sterile Neutrinos and New Matter Effects



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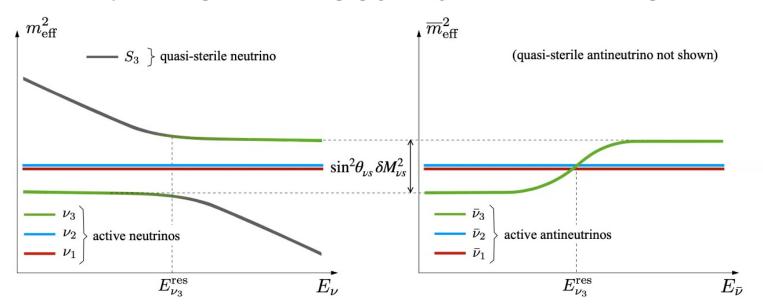
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## Quasi-Sterile Neutrinos and New Matter Effects (toy model)

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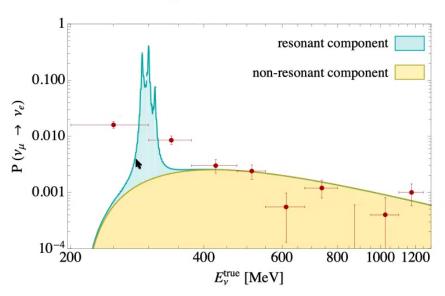


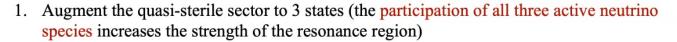
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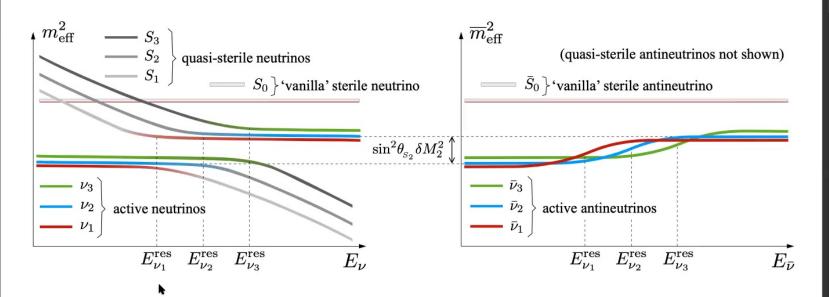


2. Add a 'vanilla' 3+1 sterile neutrino to induce off-resonance  $\nu_{\mu} \rightarrow \nu_{e}$  oscillations (and also explain  $\bar{\nu}_{\mu} \rightarrow \bar{\nu}_{e}$  and  $\bar{\nu}_{e} \nrightarrow \bar{\nu}_{e}$  anomalies)



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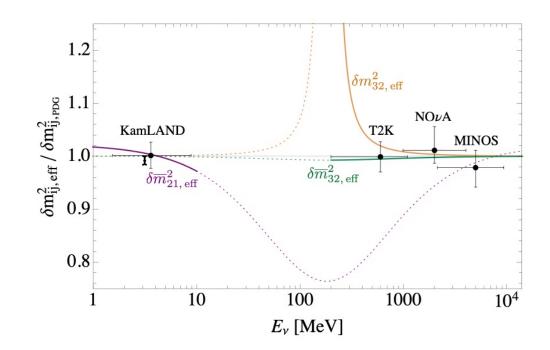




3. Bring the three resonant energies close together to avoid large distortions of the active mass spectrum



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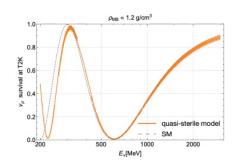
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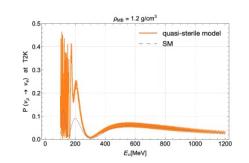
Because the resonant energy depends on the medium density,

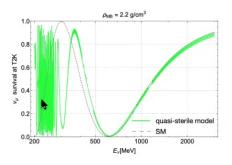
$$E_{\nu}^{res} = \frac{\delta M_{\nu s}^2 \cos 2\theta_{\nu s}}{2 V_{matter}} \sim \frac{1}{\rho_{matter}},$$

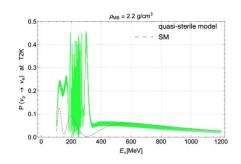
the resonant energy at long-baselines is lower than at short baselines. *e.g.*,

$$egin{array}{ll} E_{
u}^{
m res} ig|_{{}_{
m T2K}} &=& E_{
u}^{
m res} ig|_{{}_{
m MB}} rac{
ho_{{}_{
m MB}}}{
ho_{{}_{
m T2K}}} \\ &pprox & E_{
u}^{
m res} ig|_{{}_{
m MB}} rac{[1.6-2.2] {
m g/cm}^3}{2.6 {
m g/cm}^3} \\ &pprox [180-250] {
m MeV} \end{array}$$



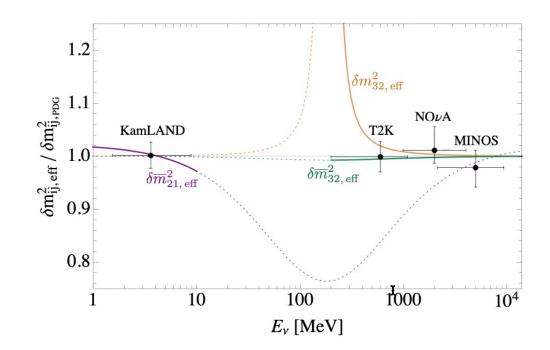








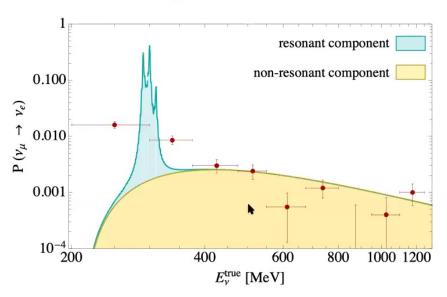
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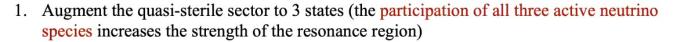


3. Bring the three resonant energies close together to avoid large distortions of the active mass spectrum



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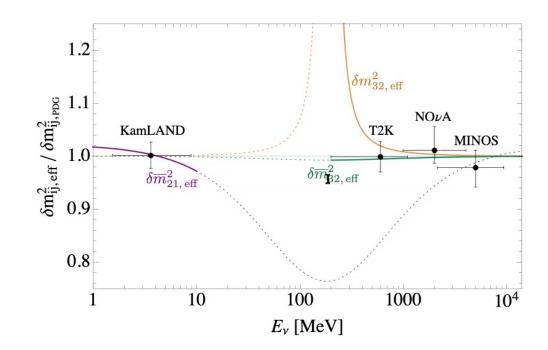




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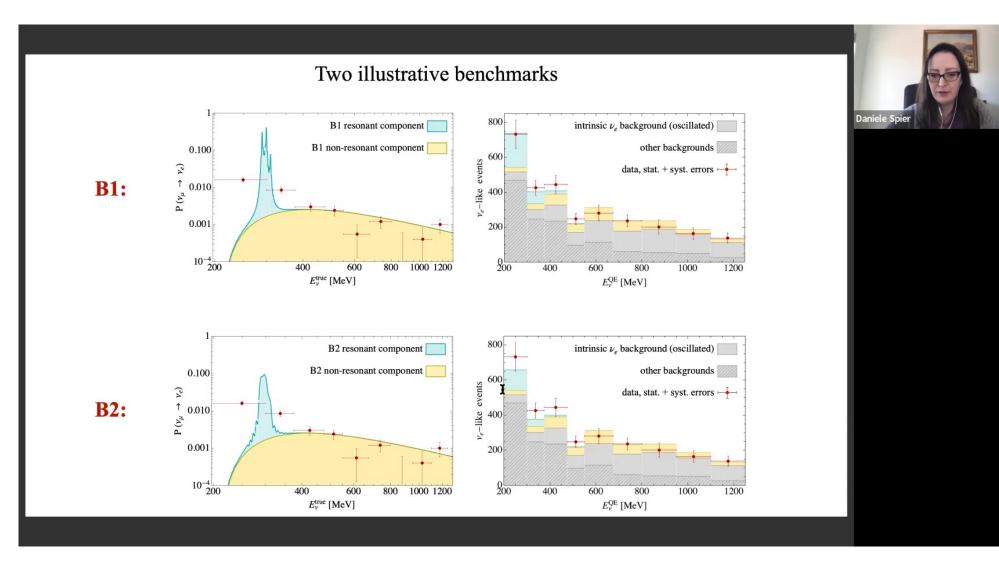
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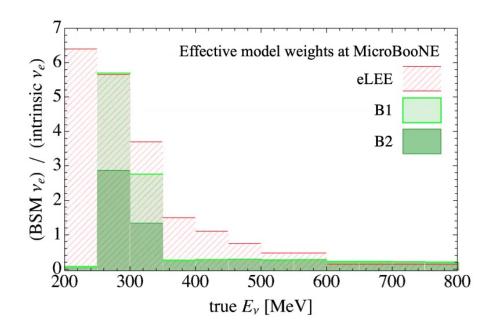


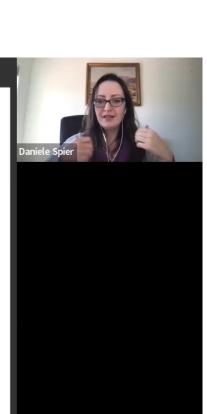
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The flux ratio (BSM  $\nu_e$ )/(intrinsic  $\nu_e$ ) differs significantly between the quasi-sterile models and the "eLEE" empirical model considered in the MicroBooNE analyses

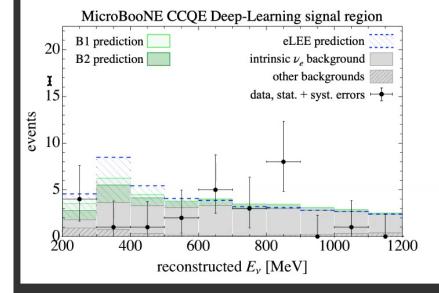


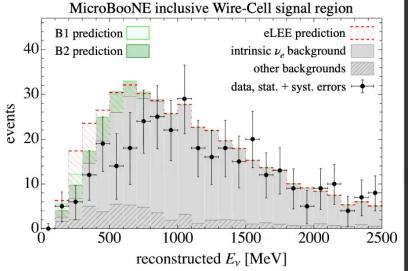


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Expected rates at MicroBooNE from quasi-sterile model are *likely* compatible with (although not favored by) the recent MicroBooNE results

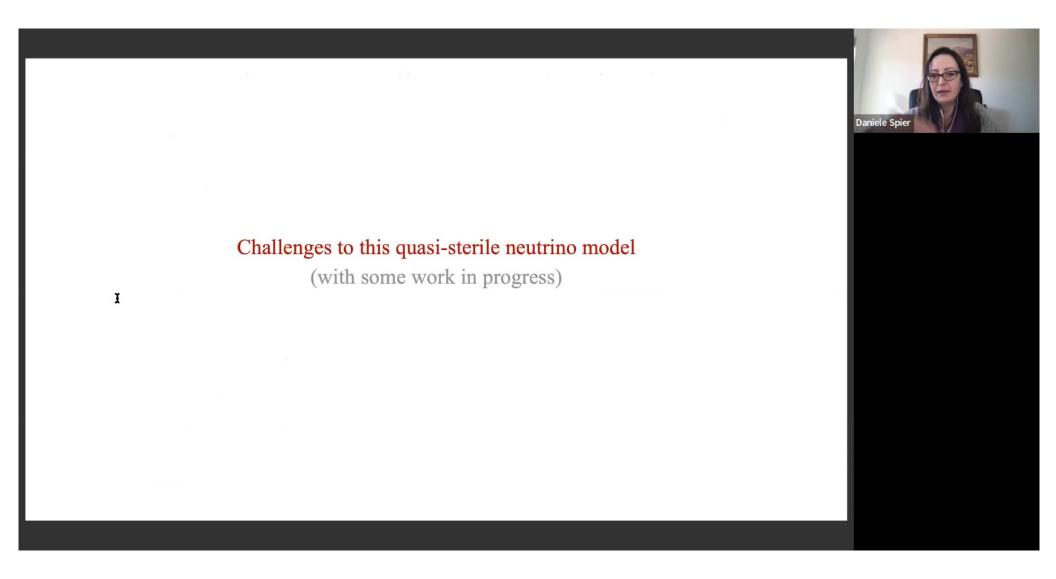
(p-values cannot be calculated without the proper covariance matrix, which cannot be built using current information in the public data releases)







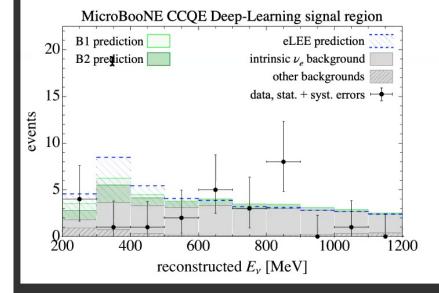
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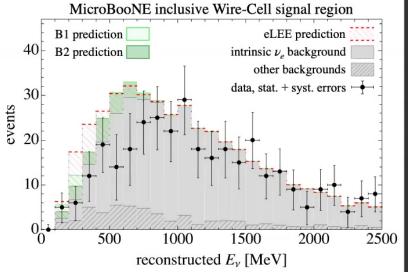


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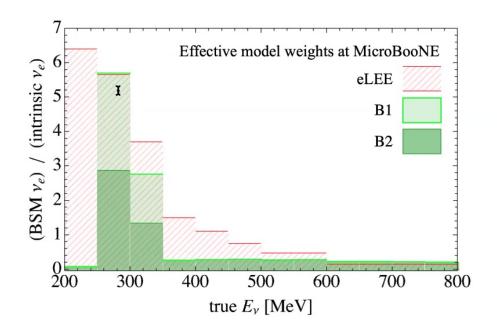
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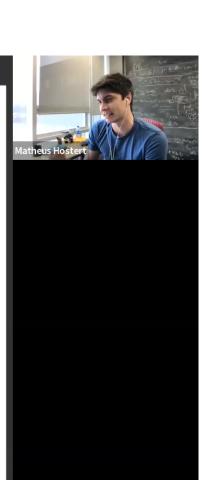




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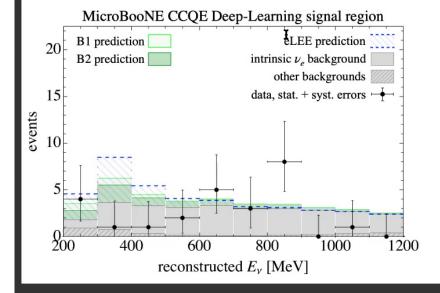


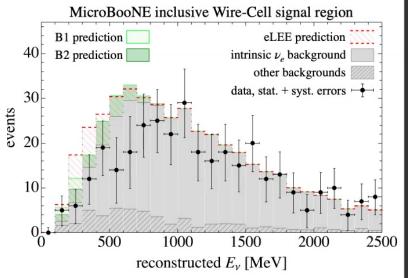


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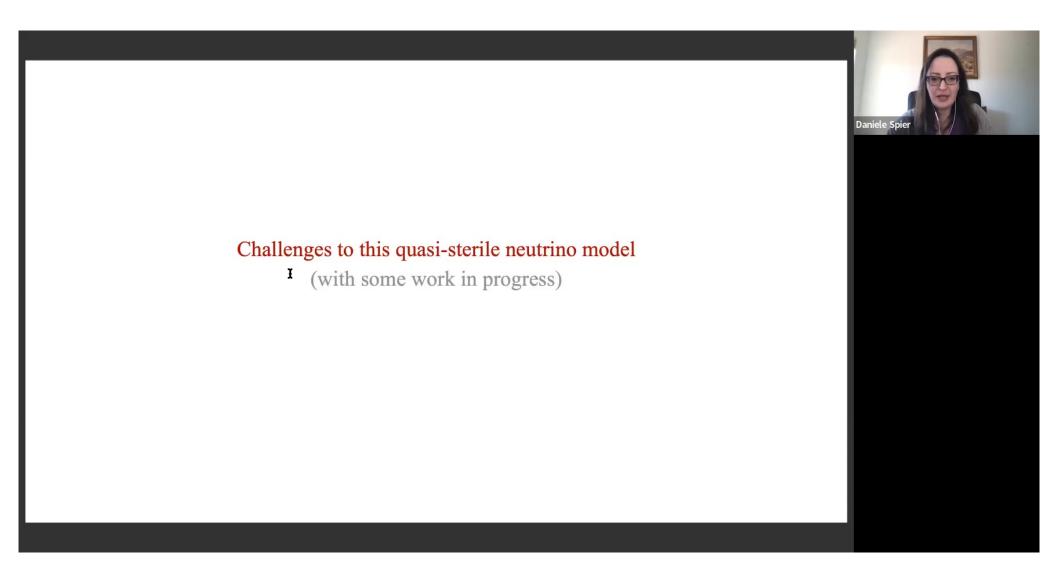
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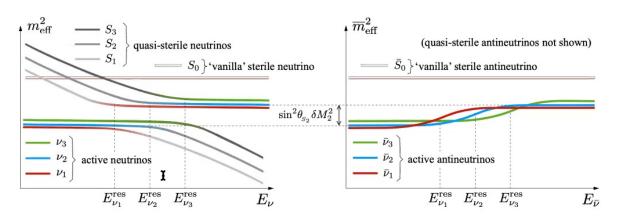


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#### <u>Challenges to this model</u>: 1. tuning of active-quasisterile mixing



The matter-induced distortion of the active neutrino spectrum requires

- (i) alignment with the Yukawa basis, and
- (ii) the same high energy shift for all three species to within ~5%, i.e.,

$$\sin^2 \theta_{_{S_1}} \delta M_1^2 \approx \sin^2 \theta_{_{S_2}} \delta M_2^2 \approx \sin^2 \theta_{_{S_3}} \delta M_3^2 \approx \mathcal{O}(1-10) \times |\delta m_{32}^2|$$

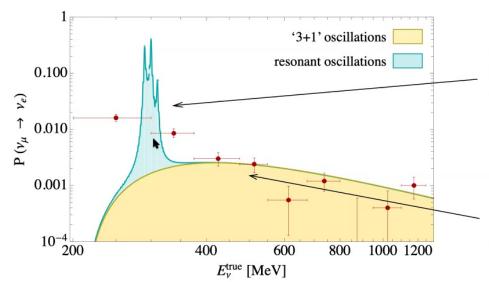
Could be the result of *ad hoc* fine-tuning, or some global symmetry in the Yukawa sector:

$$u_i \longrightarrow S_i$$
 $m_i^{(\text{mix})} \approx \sin \theta_{S_i} M_i$ 



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#### <u>Challenges to this model</u>: 2. the MiniBooNE excess coincidence



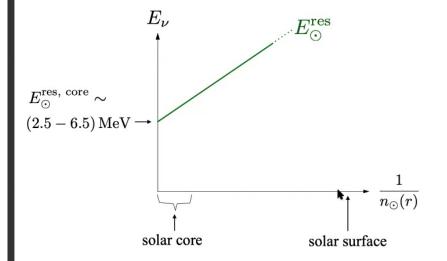
depends on active-quasi-sterile mass splitting and matter potential (mediator's mass and couplings, matter density along MB's beam-line)

depends on active-'vanilla'-sterile mass splitting and mixing angle

MiniBooNE excess coincidence: the resonant and non-resonant mechanisms, which are unrelated, happen to produce  $\nu_{\mu} \rightarrow \nu_{e}$  excesses at MiniBooNE in the same energy window.

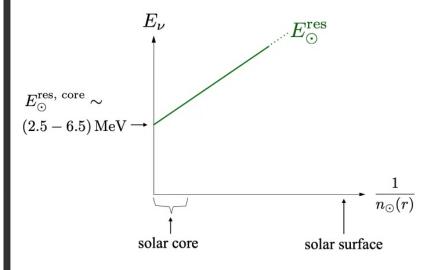


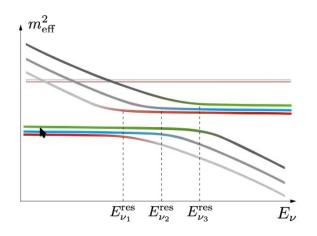
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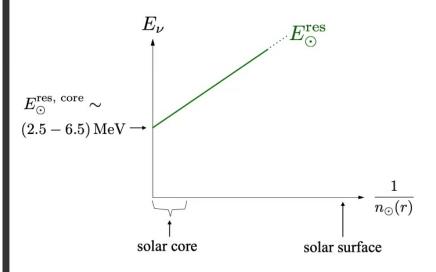
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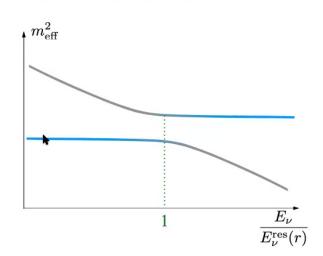


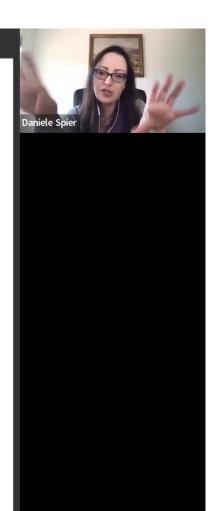




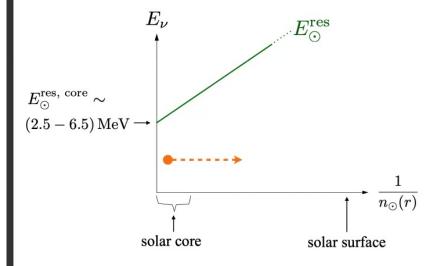
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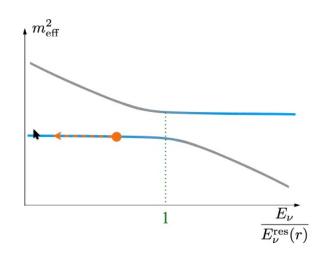


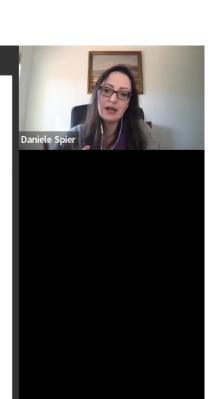




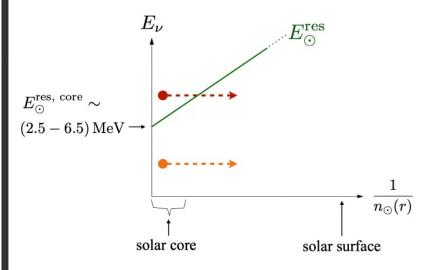
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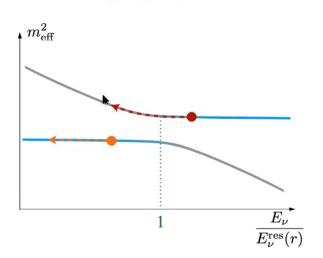






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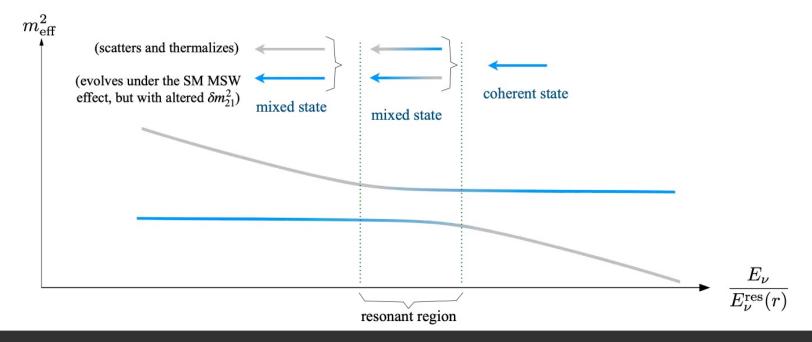
If adiabaticity and coherence are preserved during the solar neutrinos' propagation through the sun, then all active neutrinos with  $E_{\nu} > E_{\odot}^{\rm res,\ core}$  will convert to quasi-sterile neutrinos  $\Rightarrow$  ruled-out



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Possible mechanism to avoid this scenario:

Decoherence of active and quasi-sterile states during resonant region crossing due to the short mean-free-path of quasi-sterile component in the solar medium

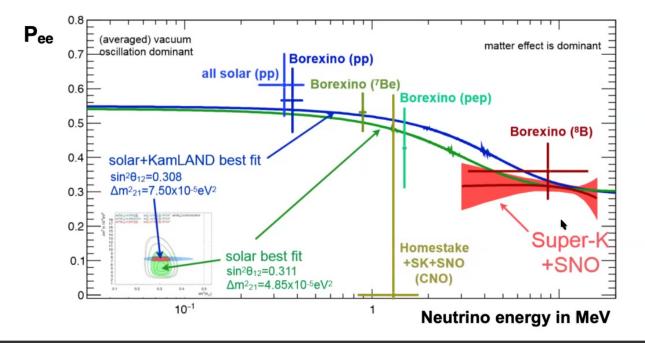


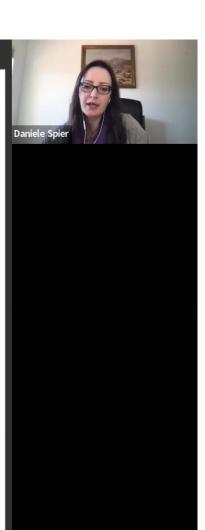


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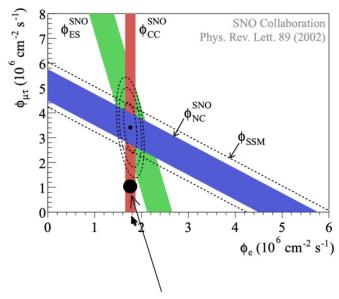
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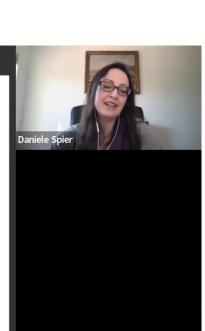


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One more issue: SNO has measured solar neutrino NC scattering via  $\nu + D \rightarrow n + p + \nu$ 



Naively expected from quasi-sterile model due to loss of <sup>8</sup>B solar neutrino flux



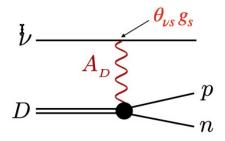
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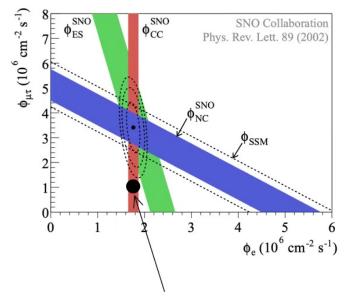
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Potential way out: dark vector mediator contribution to  $\nu + D \rightarrow p + n + \nu$  of order

$$\sigma_{\scriptscriptstyle {
m NC}}^{\scriptscriptstyle {
m BSM}} \sim (0.3-1) \times \left. \sigma_{\scriptscriptstyle {
m NC}}^{\scriptscriptstyle {
m SM}} \right|_{E_{\nu} \sim \mathcal{O}({
m MeV})}$$

$$\sim \left. \mathcal{O}(10^{-45}\,{
m cm}^2) \right.$$





Naively expected from quasi-sterile model due to loss of <sup>8</sup>B solar neutrino flux



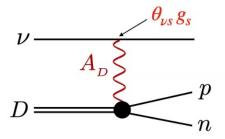
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One more issue: SNO has measured solar neutrino NC scattering via  $\nu + D \rightarrow n + p + \nu$ 

Potential way out: dark vector mediator contribution to  $\nu + D \rightarrow p + n + \nu$  of order

$$\sigma_{\text{\tiny NC}}^{\text{\tiny BSM}} \sim (0.3 - 1) \times \left. \sigma_{\text{\tiny NC}}^{\text{\tiny SM}} \right|_{E_{\nu} \sim \mathcal{O}(\text{\tiny MeV})}$$

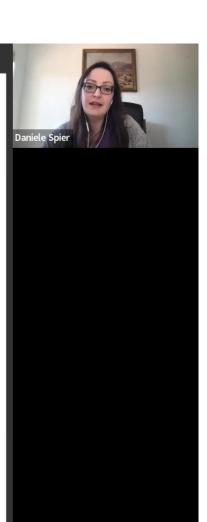
$$\sim \mathcal{O}(10^{-45} \, \text{cm}^2)$$



How viable is this possibility?

1. Need consistent dark mediator completion

k



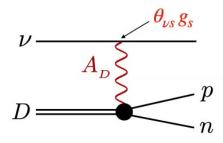
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- 3. In the low energy range\* of  $E_{\nu} \lesssim \mathcal{O}(10\,\text{MeV})$ , the only measurements of this cross section were made by reactor experiments,  $\bar{\nu}_e + D \to p + n + \bar{\nu}_e$ , (Savannah River, ROVNO, Krasnoyarsk, Bugey) with ~20% precision \*CE\nuNS' sensitivity not competitive yet

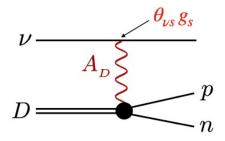


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- ⇒ Significant tension with SNO NC measurements; existence of viable model is not likely.



## Challenges to this model

- 1. tuning of active-quasi-sterile mixing
- 2. the MiniBooNE excess coincidence
- 3. solar neutrinos

These challenges are significant obstacles to the viability of quasi-sterile neutrino models explaining the MiniBooNE low energy excess

Take away:

However, they provide a good illustration of how generic consequences of light dark sectors, such as neutrino+vector portals, can be non-trivially constrained by neutrino data, since generically such models lead to:

- significant distortions of the active neutrino mixing and mass spectrum, even in energy regimes far away from the resonance region
- New matter effects, and therefore non-trivial implications for solar neutrinos



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