

Title: Quantum Fields and Strings 2021/2022 - Lecture 3

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Collection: Quantum Fields and Strings 2021/2022

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Zero Modes

$$\begin{aligned} i \not{D}(\gamma^5 \phi_n) &= -\gamma^5 i \not{D} \phi_n \\ &= -\lambda_n \gamma^5 \phi_n \end{aligned}$$

$\phi_n + \gamma^5 \phi_n$ are orthogonal if $\lambda_n \neq 0$

$$\text{Tr } \gamma^5 = \int d^4x \sum_n \phi_n^\dagger \gamma^5 \phi_n = \int d^4x \sum_{\text{zero modes}} \phi_n^\dagger \gamma^5 \phi_n$$

$$i \not{D} \phi_n$$

$$\gamma^5 \phi_n$$

are orthogonal if $\lambda_n \neq 0$

$$n = \int d^4x \sum_{\text{zero modes}} \phi^2$$

values ± 1

$$\text{Tr } \gamma^5 = n_+ - n_-$$

$\uparrow$ zero modes with ± 1 eigenvalue of γ^5

Zero Modes

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γ^5 has eigenvalues ± 1

$$\text{Tr } \gamma^5 = n_+ - n_-$$

↑ ↗
zero zero
modes modes
with ± 1 with ± 1
eigenvalue eigenvalue
of γ^5 of γ^5

$$- \gamma^5 : i \not{D} \phi_n \quad 1 = (P_L + P_R)$$

$$- \lambda_n \gamma^5 \phi_n$$

ϕ_n are orthogonal if $\lambda_n \neq 0$

$$\gamma^5 \phi_n = \int d^4x \sum_{\text{zero modes}} \phi_n^\dagger \gamma^5 \phi_n$$

eigenvalues ± 1

$$\text{Tr } \gamma^5 = n_+ - n_-$$

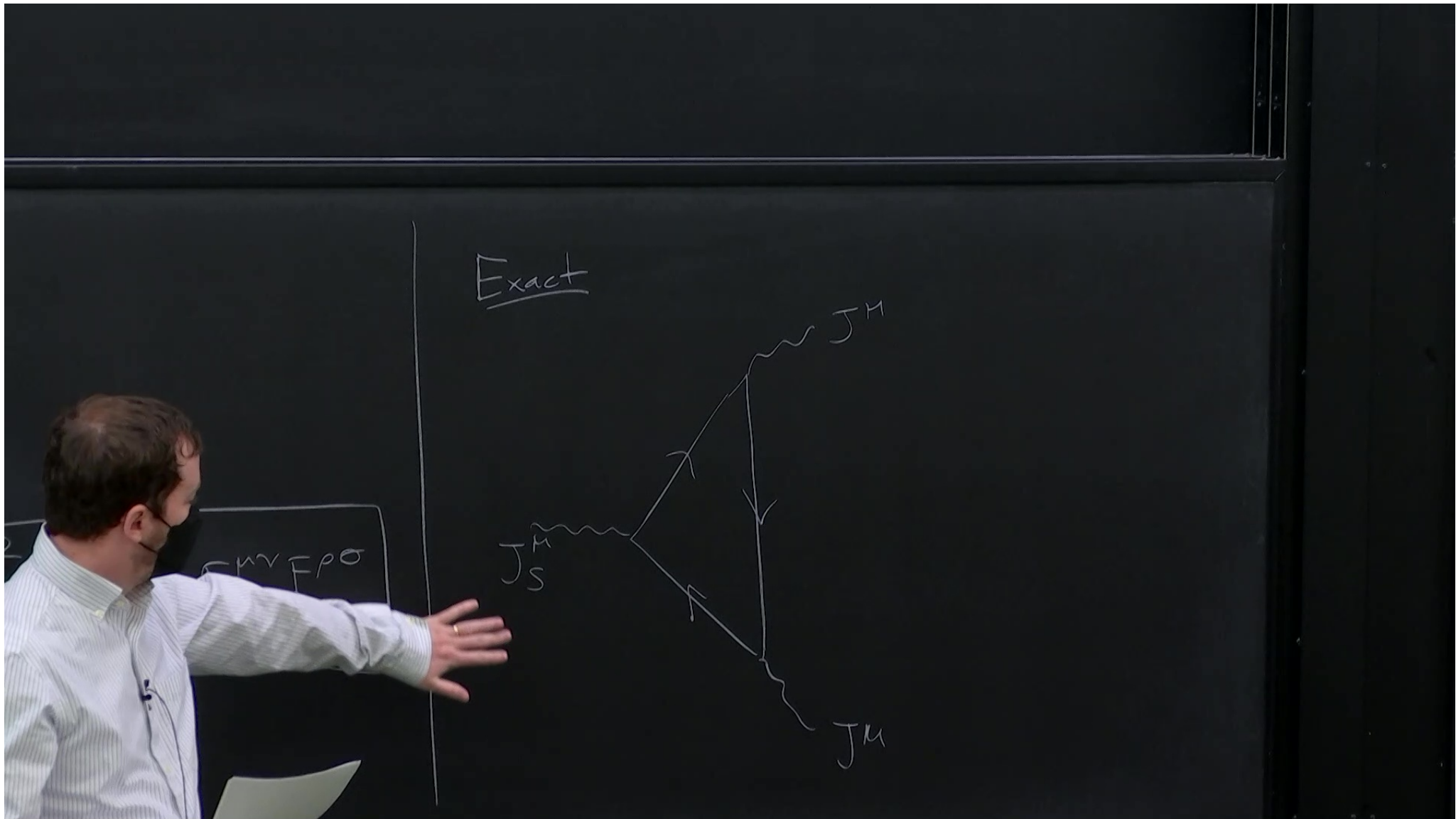
integer!
 $n_+ - n_-$

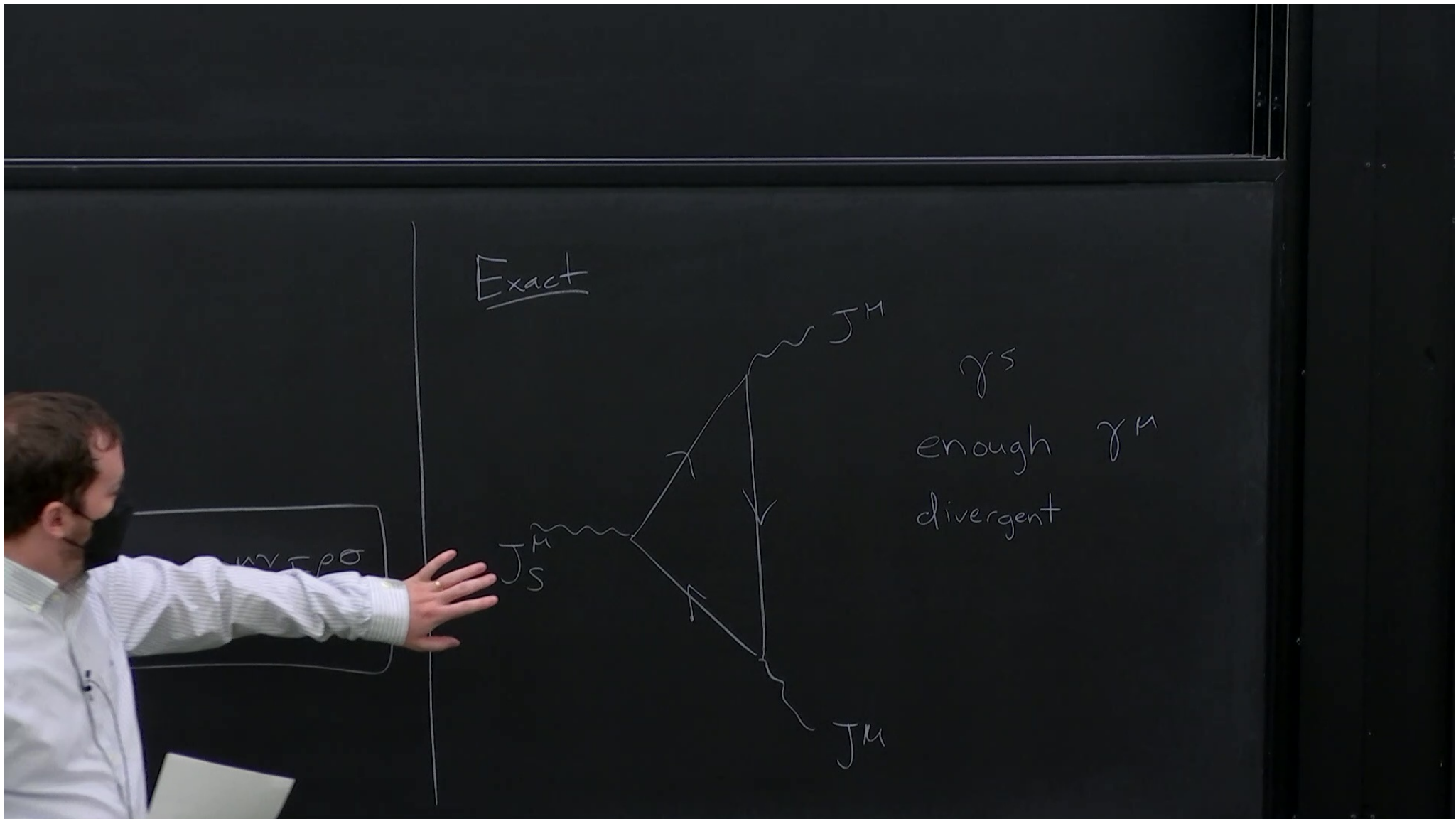
zero modes
with ± 1
eigenvalue
of γ^5

$$\boxed{\text{Index } i \not{D} = \frac{e^2}{32\pi^2} \epsilon_{\mu\nu\rho\sigma} F^{\mu\nu} F^{\rho\sigma}}$$

Atiyah-Singer index theorem

Exact





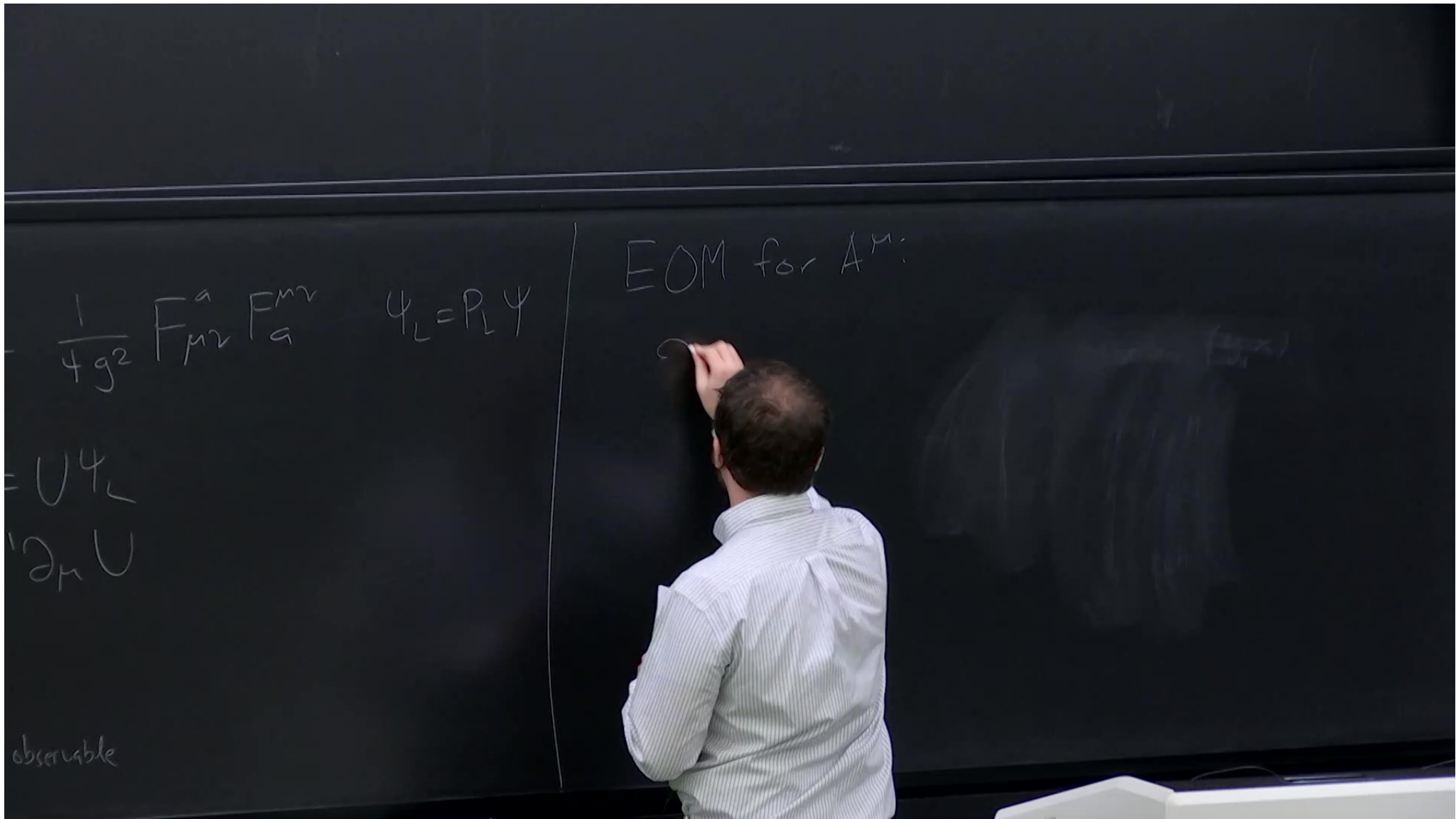
Gauge anomalies

$$\mathcal{L} = \bar{\Psi}_L (i \not{\partial} + i A^a (T_a)_{ij}) \Psi_L^j - \frac{1}{4g^2} F_{\mu\nu}^a F^{\mu\nu}_a \quad \Psi_L = P_L \Psi$$

Gauge invariance

$$\Psi_L \rightarrow e^{i\beta^a(x) T^a} \Psi_L = U \Psi_L$$

$$A_\mu \rightarrow U A_\mu U^{-1} - U^{-1} \partial_\mu U$$

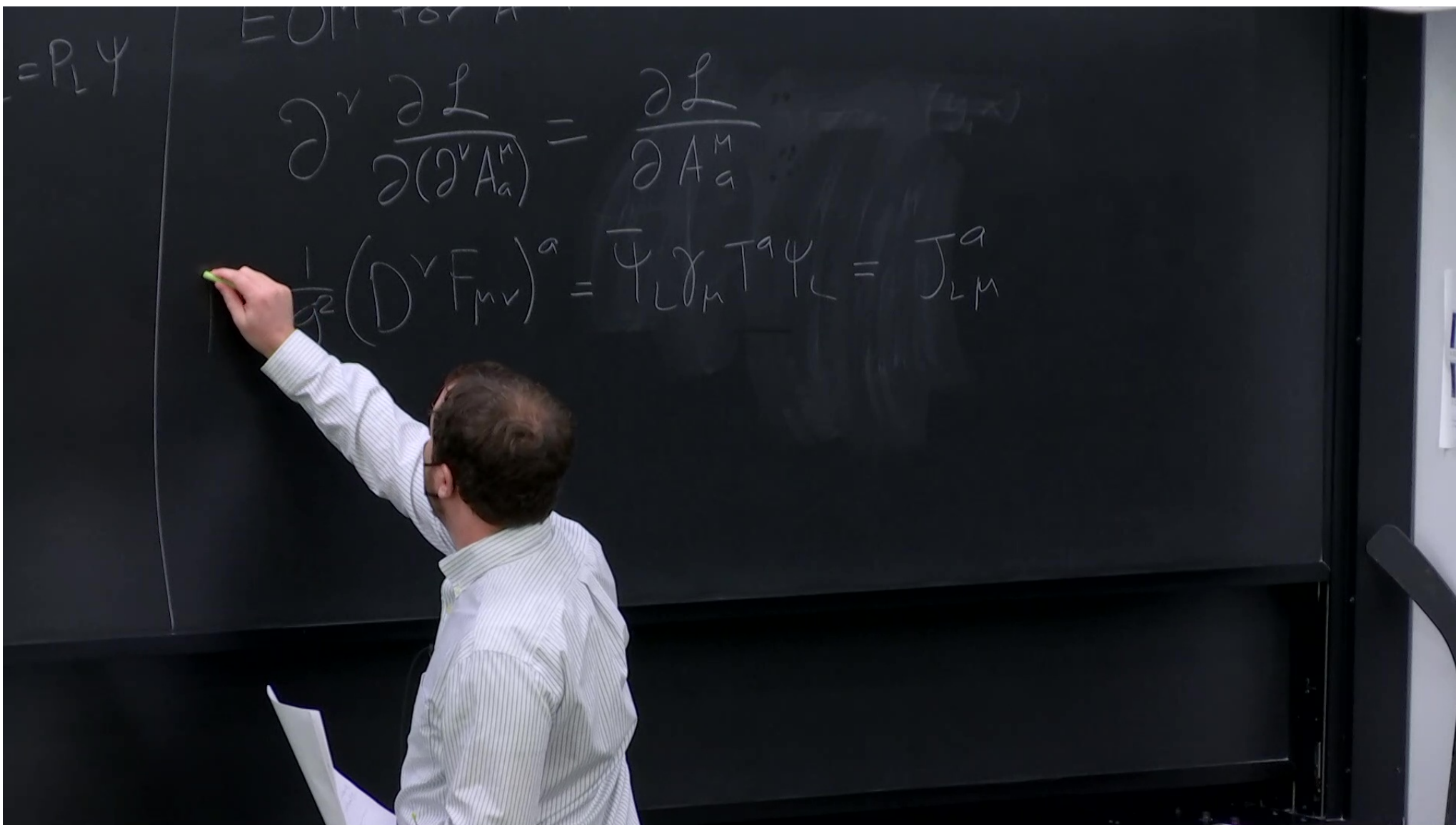


$$\psi_L = P_L \psi$$

EOM for A^μ :

$$\partial^\nu \frac{\partial \mathcal{L}}{\partial (\partial^\nu A_a^\mu)} = \frac{\partial \mathcal{L}}{\partial A_a^\mu}$$

$$\frac{1}{g^2} (D^\nu F_{\mu\nu})^a = \bar{\psi}_L \gamma_\mu T^a \psi_L$$



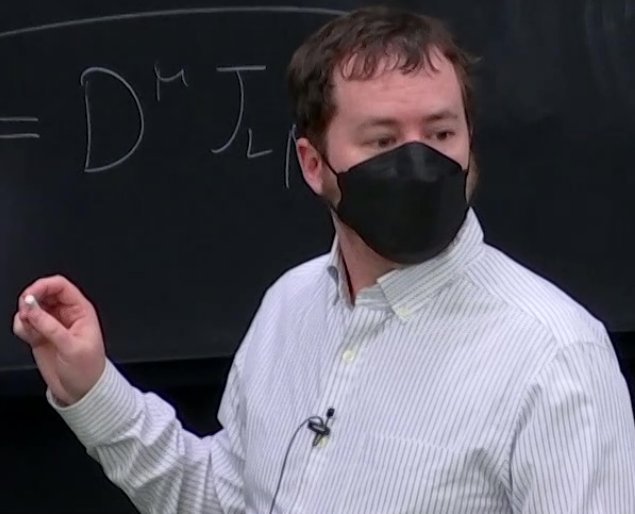
$$= P_L \psi$$

EOM for A^M :

$$\partial^\nu \frac{\partial \mathcal{L}}{\partial (\partial^\nu A_a^M)} = \frac{\partial \mathcal{L}}{\partial A_a^M}$$

$$D^\mu \left(\frac{1}{g^2} (D^\nu F_{\mu\nu})^a \right) = \bar{\psi}_L \gamma_\mu T^a \psi_L = J_{L\mu}^a$$

$$\boxed{0 = D^\mu J_{L\mu}^a}$$



$$= P_L \psi$$

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$$\boxed{0 = D^\mu J_{L\mu}^a}$$

essential for consistency
unlike chiral
current conservation

current $J_{LM} = \Psi_L^\dagger \partial_\mu \Psi_L$
not gauge invariant $\rightarrow$ not observable

$$(\mathcal{D}^\mu J_{LM})^a = \dots \text{tr} [T^a T^b T^c]$$

$\uparrow$ depends on gauge field



current $J_{LM} = \frac{1}{L} \partial_\mu \dots$

not gauge invariant $\rightarrow$ not observable

$$(\partial^\mu J_{LM})^a = \dots \text{tr} [T^a T^b T^c]$$

$\uparrow$ depends on gauge field

consistency of the theory depends on fermion content!

current conservation

Functional Integrals

extra factors of T^a

chiral coupling complicates eigenfunction expansion

$i\not{D}_L$ not Hermitian in Euclidean space

γ has eigenvalues ± 1

Atiyah-Singer index

Euclidian

$$Z[A] = \int \mathcal{D}\psi_L \mathcal{D}\bar{\psi}_L e^{-\int \bar{\psi}_L i \not{D} \psi_L} \mathcal{L} \text{Det } i \not{D} P_L$$

zero ϕ_n & ψ_n
modes

eigenvalues ± 1

Index 12

$32\pi^2$

Atiyah-Singer index theorem

$$\int \bar{\psi}_L i \not{D} \psi_L \quad ? \quad \mathcal{L} \text{ Det}(i \not{D} P_L)$$

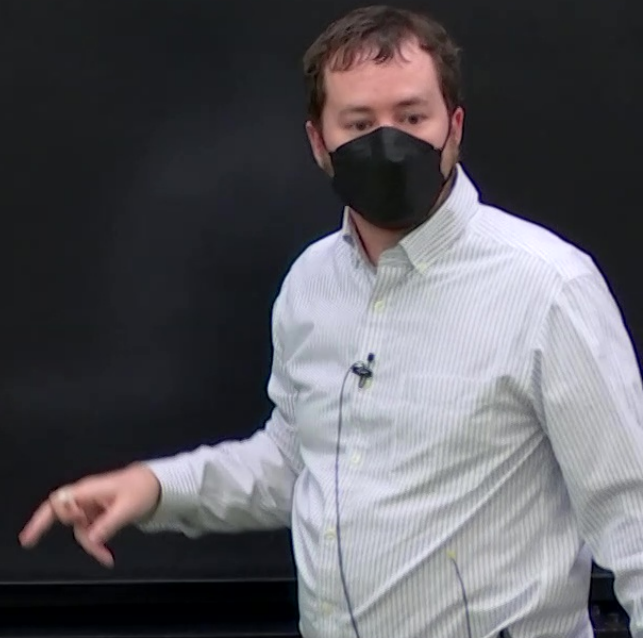
$$P_R i \not{D} P_L = i \not{D} P_L$$

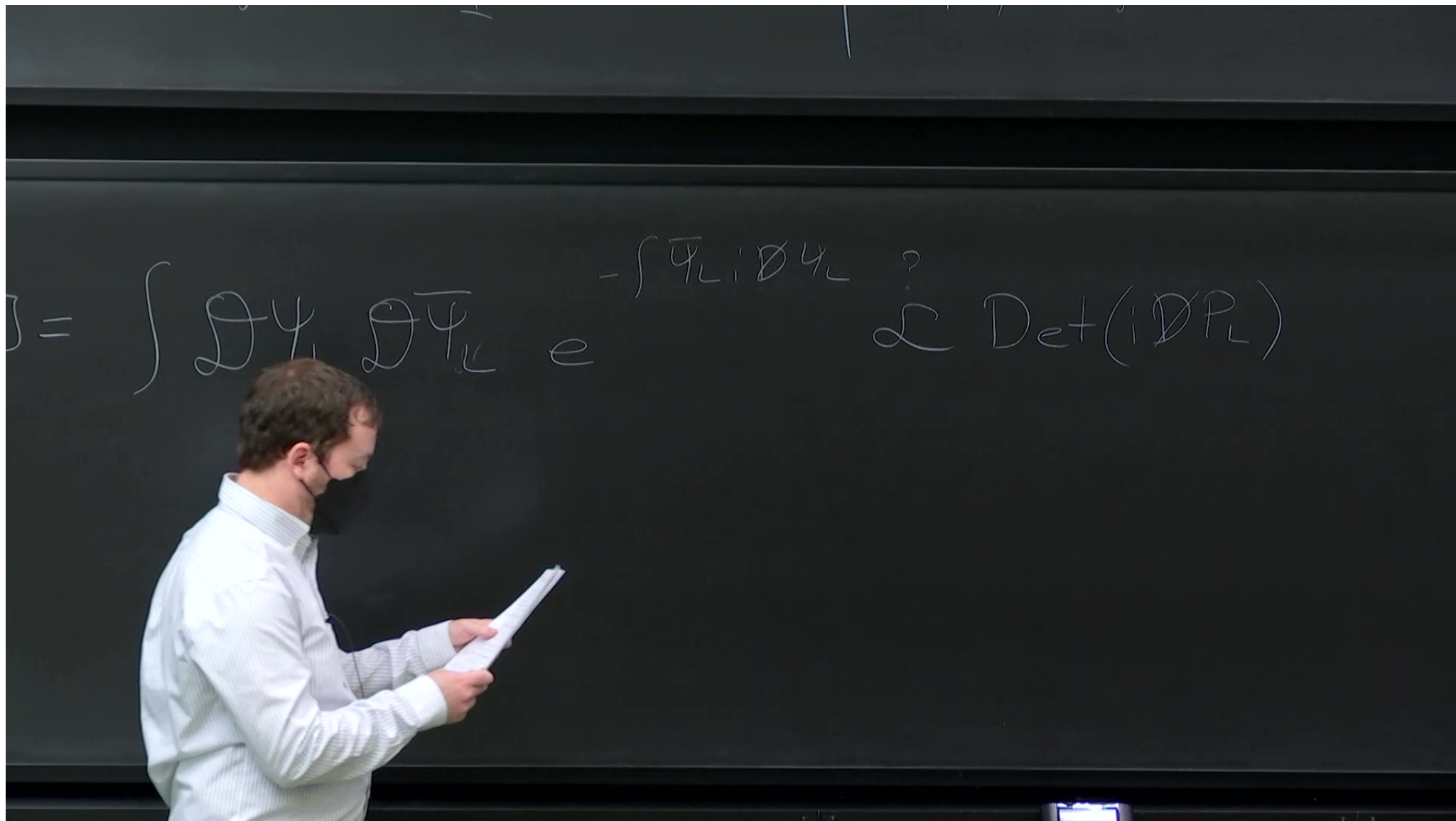
Thyath-Singer index theorem

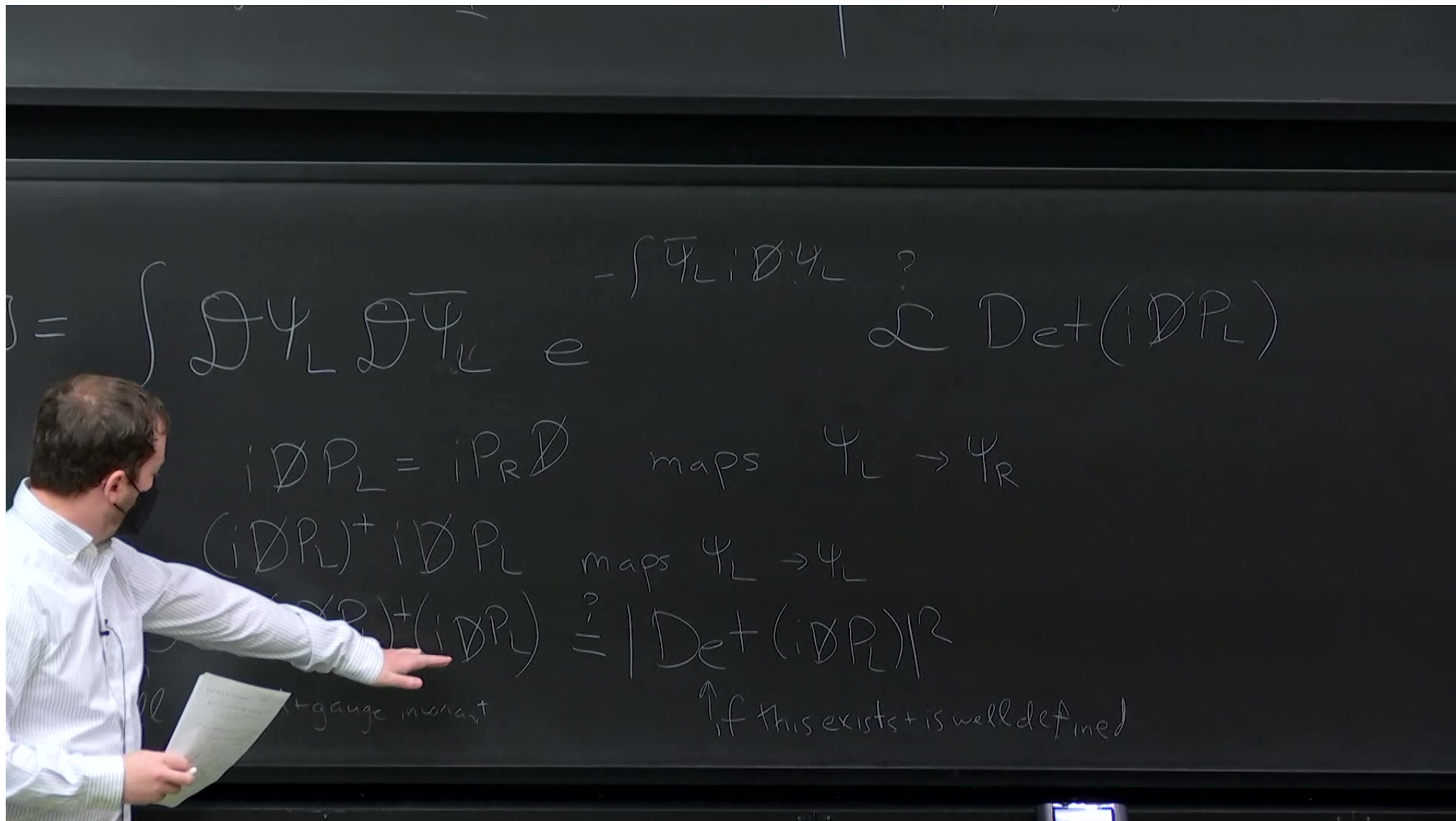
JM

$$? \quad \mathcal{L} \text{ Det}(i\mathcal{D}_L)$$

$$P_R i\mathcal{D}_L = i\mathcal{D}_L$$







$$Z = \int \mathcal{D}\Psi_L \mathcal{D}\bar{\Psi}_L e^{-\int \bar{\Psi}_L i \not{D} \Psi_L} \mathcal{L} \text{Det}(i \not{D} P_L)$$

$$i \not{D} P_L = i P_R \not{D} \quad \text{maps } \Psi_L \rightarrow \Psi_R$$

$$(i \not{D} P_L)^\dagger i \not{D} P_L \quad \text{maps } \Psi_L \rightarrow \Psi_L$$

$$\text{Det}(i \not{D} P_L)^\dagger (i \not{D} P_L) \stackrel{?}{=} |\text{Det}(i \not{D} P_L)|^2$$

defined + gauge invariant

↑ if this exists + is well defined

uclian

$$Z[A] = \int \mathcal{D}\psi_L \mathcal{D}\bar{\psi}_L e^{-\int \bar{\psi}_L i \not{D} \psi_L} \mathcal{L} \text{Det}(i \not{D} P_L)$$

↑ ambiguity

$$i \not{D} P_L = i P_R \not{D}$$

$$(i \not{D} P_L)^+ i \not{D} P_L$$

$$\text{Det}(i \not{D} P_L)^+ (i \not{D} P_L)$$

↑
well defined + gauge invariant

$$\psi_L \rightarrow \psi_R$$

$$\psi_L \rightarrow \psi_L$$

$$\text{Det}(i \not{D} P_L)^2$$

is well defined

Fermion content of theory determines consistency

- fermion in real or pseudoreal reps do not contribute

$$T_L^T = -V T_L V \quad \text{for unitary } V$$

Fermion content of theory determines consistency

- fermion in real or pseudoreal reps do not contribute

$$T_L^T = -V T_L V \quad \text{for unitary } V$$

- massive fermions



