

Title: Geometry and topology for physicists 2021/2022 - Lecture 6

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Collection: Geometry and Topology for Physicists 2021/2022

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Formal Properties of \wedge product:

$$1) \quad \omega_1 \wedge (\omega_2 + \omega_3) = \omega_1 \wedge \omega_2 + \omega_1 \wedge \omega_3$$

$$\omega_1 \in \Omega^k(M)$$

$$\omega_2, \omega_3 \in \Omega^l(M)$$

$$2) \quad \omega_1 \wedge (\omega_2 \wedge \omega_3) = (\omega_1 \wedge \omega_2) \wedge \omega_3$$

3) Skew-commutative:

$$\omega_1 \wedge \omega_2 = (-1)^{kl} \omega_2 \wedge \omega_1,$$

$$\omega_1 \in \Omega^k(M)$$

$$\omega_2 \in \Omega^l(M)$$

Re

Reminiscent of fermions

de Rham operator:

This is a map $d: \Omega^k(M) \rightarrow \Omega^{k+1}(M)$

In coordinates

$$d(f^{i_1 \dots i_k} dx_{i_1} \wedge \dots \wedge dx_{i_k})$$
$$= \frac{\partial f^{i_1 \dots i_k}}{\partial x_j} dx_j \wedge dx_{i_1} \wedge \dots \wedge dx_{i_k}$$

$$\omega_1 \wedge \omega_2 = (-1)^{kl} \omega_2 \wedge \omega_1, \quad \omega_1 \in \Omega^k(M)$$

$$\omega \wedge \omega = 0$$

$$\omega_2 \in \Omega^l(M)$$

On \mathbb{R}^2 ,

$$d(f^1 dx_1 + f^2 dx_2)$$
$$= -\frac{\partial f^1}{\partial x_2} dx_1 \wedge dx_2 + \frac{\partial f^2}{\partial x_1} dx_1 \wedge dx_2$$

Properties

1) $d(\omega_1 \wedge \omega_2) = (d\omega_1) \wedge \omega_2 + (-1)^k \omega_1 \wedge d\omega_2$

$$\omega_1 \in \Omega^k$$
$$\omega_2 \in \Omega^l$$

$$\omega_1 \in \Omega^k(M)$$
$$\omega_2 \in \Omega^l(M)$$

$$2) \quad d(d\omega) = 0$$

$$\omega = f dx_{i_1} \wedge \dots \wedge dx_{i_k}$$

$$d\omega = \frac{\partial f}{\partial x_{j_1}} dx_{j_1} \wedge dx_{i_1} \wedge \dots \wedge dx_{i_k}$$

$$d(d\omega) = \frac{\partial^2 f}{\partial x_{j_2} \partial x_{j_1}} \underbrace{dx_{j_2} \wedge dx_{j_1}}_{\text{anti-symmetric}} \wedge dx_{i_1} \wedge \dots \wedge dx_{i_k} = 0$$

Symmetric in
 j_1, j_2

anti-symmetric

In dimension 3

$$\Omega^1(\mathbb{R}^3) \simeq \text{Vect}(\mathbb{R}^3)$$

$$f^i dx_i \mapsto f^i \frac{\partial}{\partial x_j} \delta_{ij}$$

$$\Omega^2(\mathbb{R}^3) \simeq \text{Vect}(\mathbb{R}^3)$$

$$f^i dx_j \wedge dx_k \mapsto f^i \frac{\partial}{\partial x_l} \varepsilon_{ijk}$$

$$\Omega^0 \xrightarrow{d} \Omega^1 \xrightarrow{d} \Omega^2 \xrightarrow{d} \Omega^3$$

$$C^\infty \xrightarrow{\text{Grad}} \text{Vect} \xrightarrow{\text{Curl}} \text{Vect} \xrightarrow{\text{Div}} C^\infty$$

For Curl:

$$f^i dx_j \mapsto \frac{\partial f^i}{\partial x_j} dx_k \wedge dx_l$$

$$f^i dx_i \leftrightarrow f^i \frac{\partial}{\partial x_j} \delta_{ij}$$

$$\frac{\partial f^i}{\partial x_j} dx_j \wedge dx_i \leftrightarrow \epsilon_{kji} \frac{\partial f^i}{\partial x_j} \frac{\partial}{\partial x_k}$$

de Rham sends

$$f^i \frac{\partial}{\partial x_j} \delta_{ij} \rightarrow \epsilon_{kji} \frac{\partial f^i}{\partial x_j} \frac{\partial}{\partial x_k}$$

This is curl.

$\omega_1 \wedge \omega_2 = (-1)^k \omega_2 \wedge \omega_1,$ $\omega_1 \in \Omega^k(M)$
 $\omega \wedge \omega = 0$ $\omega_2 \in \Omega^l(M)$

A first look at cohomology

For this

If M is a manifold,

$$H^k(M) = \left\{ \omega \in \Omega^k(M), d\omega = 0 \right. \\ \left. (\omega \text{ is closed}) \right\}$$

$$\omega \sim \omega' \\ \text{if } \omega - \omega' \\ = d\eta, \\ \eta \in \Omega^{k-1}(M)$$

$\Omega^k(M)$
 $\Omega^l(M)$

$d\eta$

For this to make sense, it's necessary that
 $d(d\eta) = 0$

$\omega \sim \omega'$
 $\omega - \omega'$
 $= d\eta$
 $\in \Omega^{k-1}(M)$

m)
n)

For this to make sense, it's necessary that

$$d(d\gamma) = 0$$

$H^k(M, \mathbb{R})$ are topological invariants

$$b_k = \dim_{\mathbb{R}} H^k(M, \mathbb{R})$$

(typically finite)

ω'
 ω''
 ω'''
 $\omega^{(k-1)}(m)$

Example.

What is $H^*(\mathbb{R})$?

$\Omega^0(\mathbb{R})$ functions $f(x)$

$\Omega^1(\mathbb{R})$ are $g(x) dx$

If $f(x) \in \Omega^0(\mathbb{R})$,

$$df(x) = \frac{df}{dx} dx$$

Closed zero forms are constant

$$H^0(\mathbb{R}) = \mathbb{R}, \text{ the constants}$$

Closed zero forms are constant

$H^0(\mathbb{R}) = \mathbb{R}$, the constant functions

(on any connected manifold this is true)

$H^1(\mathbb{R}) = \left\{ g(x) dx \right\} /$ g which are $g = \frac{df}{dx}$

But $g(x) = \frac{d}{dx} \int_0^x g(y) dy$

So $H^1(\mathbb{R}) = 0$.

$$\text{so } H^1(\mathbb{R}) = 0.$$

S^1

$$\Omega^0 = \{ f(\theta), f(\theta+2\pi) = f(\theta) \}$$

$$\Omega^1 = \{ g(\theta) d\theta, g(\theta+2\pi) = g(\theta) \}$$

If $f(\theta) \in \Omega^0$, $df(\theta) = 0$ iff f is constant

$$H^0(S^1) = \mathbb{R}$$

$(\mathbb{R}^2, \mathbb{Z})$

H^1 : If $\int g(\theta) d\theta$, g is equivalent to 0

in cohomology if it has a periodic anti-derivative

$$g(\theta) d\theta = d\left(\int_{\varphi=0}^{\theta} g(\varphi) d\varphi\right)$$

But $\int_{\varphi=0}^{\theta} g(\varphi) d\varphi$ might not be periodic!

It is periodic iff $\int_0^{2\pi} g(\varphi) d\varphi = 0$

(typically finite)

Exact 1-forms are $g(\theta)d\theta$ with 0 integral

Since we can add a constant $c d\theta$ to
make anything have 0 integral,

$$H^1(S^1) = \mathbb{R}$$

The isomorphism is given by \int_{S^1}

Intuition This is true because
 $H^1(S^1)$ "measures" closed 1d manifolds

live

+ be

$$-g(\varphi)d\varphi = 0$$

Lemma $H^1(\mathbb{R}^n) = 0$ if $n > 0$
 \mathbb{R} if $n = 0$

For $n = 2$

Let's check that $H^1(\mathbb{R}^2) = 0$

A one

$\int_0^{2\pi} -g(\varphi) d\varphi = 0$

Intuition This is true because $H^1(S^1)$ "measures" closed 1d manifolds

$\omega = f^1 dx_1 + f^2 dx_2$
 It's closed if

$$\frac{\partial f^1}{\partial x_2} = \frac{\partial f^2}{\partial x_1}$$

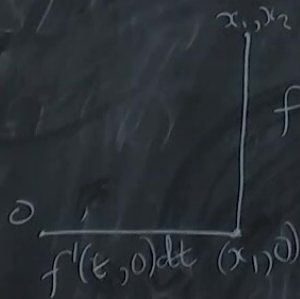
We want to show that there is $F(x_1, x_2)$ with

$$\frac{\partial F}{\partial x_1} = f^1$$

$$\frac{\partial F}{\partial x_2} = f^2$$

Define

$$F(x_1, x_2) = \int_0^{x_1} f'(t, 0) dt + \int_0^{x_2} f^2(x_1, s) ds$$



$$\frac{\partial F}{\partial x_2} = f^2(x_1, x_2) \checkmark$$

$$\frac{\partial F}{\partial x_1} = f'(x_1, 0) + \int_0^{x_2} \frac{\partial f^2}{\partial x_1}(x_1, s) ds$$

$$= f^1(x_1, 0) + \int_0^{x_2} \frac{\partial f^1}{\partial s}(x_1, s) ds$$

$$= f^1(x_1, x_2)$$

$$= f^2(x_1, x_2)$$

$$f^1(x_1, 0) + \int_0^{x_2} \frac{\partial f^2}{\partial x_1}(x_1, s) ds$$

Integration

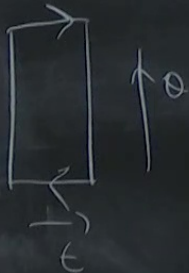
Recall, x_i, y_i are coords

$$dy_1 \wedge \dots \wedge dy_n = \det\left(\frac{\partial y_i}{\partial x_j}\right) dx_1 \wedge \dots \wedge dx_n$$

But, when we change variables in integrations,
we get $\left| \det\left(\frac{\partial y_i}{\partial x_j}\right) \right|$

Defn

An ori



$$t \rightarrow -t$$

$$0 \rightarrow 0 + 2\pi$$

Jacobian
is -ve

Defn

An orientation on a manifold is a coordinate atlas so that $\det \left(\frac{\partial y_i}{\partial x_j} \right) > 0$ between any two patches

Theorem On any compact oriented manifold, we can integrate top forms in a coordinate independent way.

$\frac{d\theta}{dt}$ Integrate top forms in a coordinate independent
 $\rightarrow -d\theta \wedge dt$



$$\begin{aligned}\theta_1 &\sim \theta_1 + 2\pi \\ \theta_2 &\sim \theta_2 + 2\pi \\ \theta_2 &\rightarrow -\theta_2\end{aligned}$$