

Title: Toward realization of novel superconductivity based on twisted van der Waals Josephson junction in Cuprates

Speakers: Philip Kim

Series: Quantum Matter

Date: February 28, 2022 - 12:00 PM

URL: <https://pirsa.org/22020071>

Abstract: Twisted interfaces between stacked van der Waals Cuprate crystals enable tunable Josephson coupling, utilizing anisotropic superconducting order parameters. Employing a novel cryogenic assembly technique, we fabricate high-temperature Josephson junctions with an atomically sharp twisted interface between $\text{Bi}_2\text{Sr}_2\text{CaCu}_2\text{O}_{8+x}$ crystals. The critical current density J_c sensitively depends on the twist angle. While near 0 degree twist, J_c nearly matches that of intrinsic junctions, it is suppressed almost 2-orders of magnitude but remained finite near 45 degree. J_c also exhibits non-monotonic behavior versus temperature due to competition between two supercurrent contributions from nodal and anti-nodal regions of the Fermi surface. Near 45 degree twist angle, we observe two-period Fraunhofer interference patterns and fractional Shapiro steps at half integer values, a signature of co-tunneling Cooper pairs necessary for high temperature topological superconductivity.

Zoom Link: https://pitp.zoom.us/join/9tJcqc-ihqzMvHdW-YBm7mYd_XP9Amhypv5vO



Toward realization of novel superconductivity based on twisted van der Waals Josephson junction in cuprates

Philip Kim

Department of Physics, Harvard University

Acknowledgment



Experiment

Frank Zhao, Nicolar Poccia, Xiaomeng Cui, Hyobin Yoo, Joon Young Park, Yuval Ronen, Rebecca Engelke



TEM

Joules Gardener, Austin Akey (CNS); Sangmin Lee, Miyoung Kim (SNU)

Samples

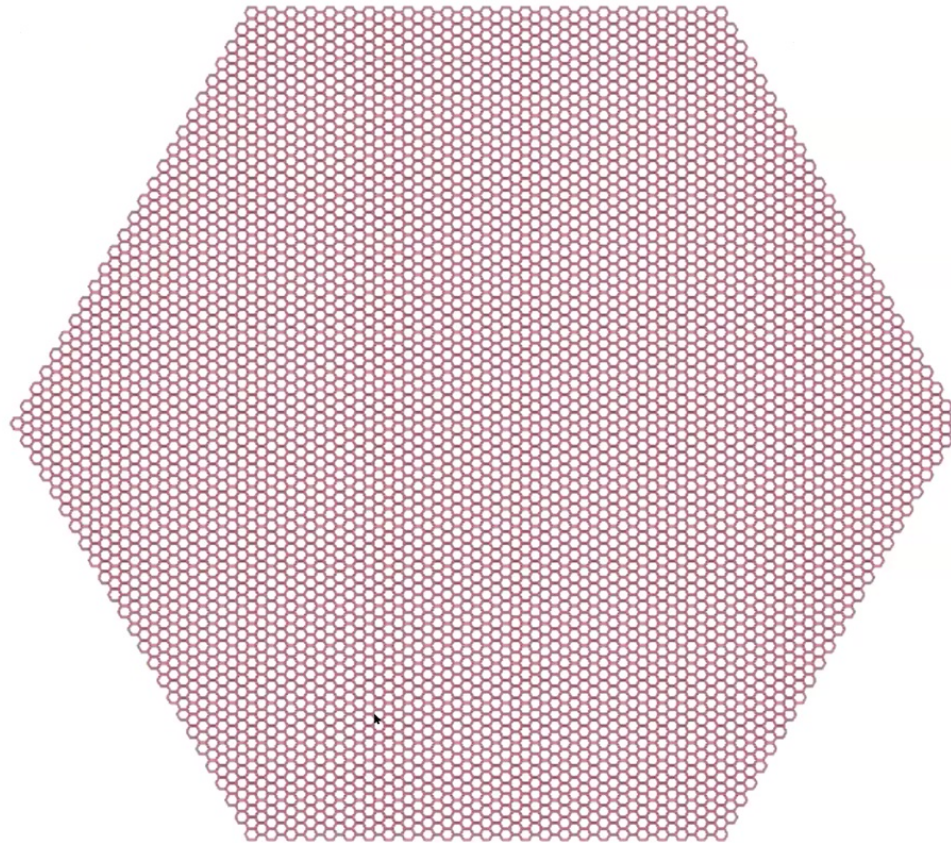
Rudan Zhong, Genda Gu (BNL); Kenji Watanabe, Takashi Taniguchi (NIMS)

Theory

Pavel Volkov, Jed Pixley (Rutgers); Tarun Tummuru, Stephan Plugge, Marcel Franz (UBC)



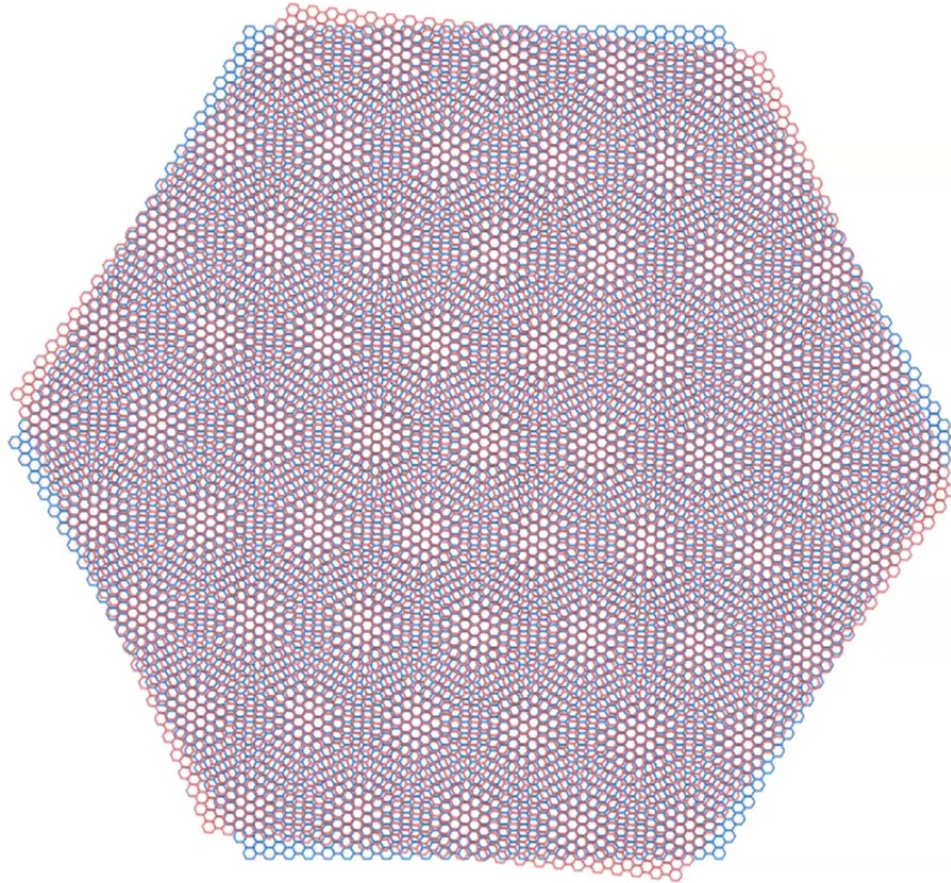
Moire Superlattice Engineering in van der Waals Interface



Twisted van der Waal interface of homo & hetero structure

Animation from Jarillo-Herrero group

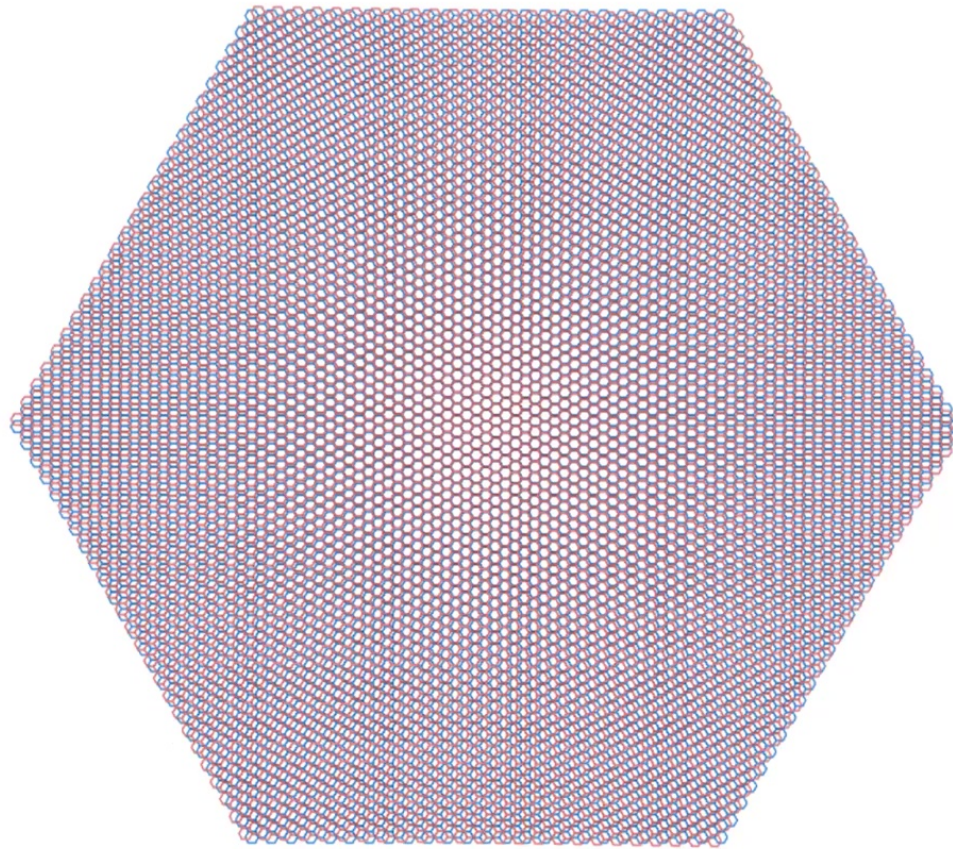
Moire Superlattice Engineering in van der Waals Interface



Twisted van der Waal interface of homo & hetero structure

Animation from Jarillo-Herrero group

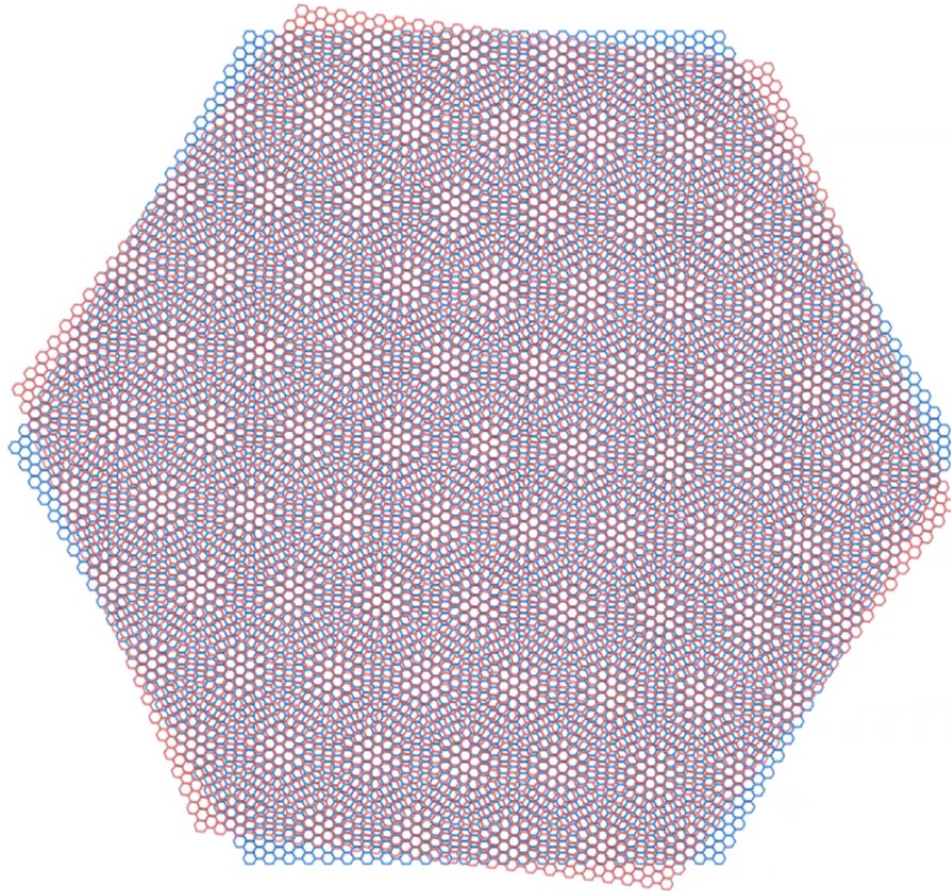
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Twisted van der Waal interface of homo & hetero structure

Animation from Jarillo-Herrero group

Correlated Quantum State in Twisted Graphene Bilayer



ARTICLE

5 APRIL 2018 | VOL 556 | NATURE | 43

doi:10.1038/nature26160

Unconventional superconductivity in magic-angle graphene superlattices

Yuan Cao¹, Valla Fatemi¹, Shiang Fang², Kenji Watanabe³, Takashi Taniguchi³, Efthimios Kaxiras^{2,4} & Pablo Jarillo-Herrero¹

Correlated insulator

Cao et al., Nature 556, 80–84 (2018).

Burg et al., PRL 123, 197702 (2019).

Unconventional superconductivity

Cao et al., Nature 556, 43–50 (2018).

Yankowitz et al., Science 363, 1059–1064 (2019).

Lu et al., Nature 574, 653–657 (2019).

Magnetism and Quantum Anomalous Hall Effect

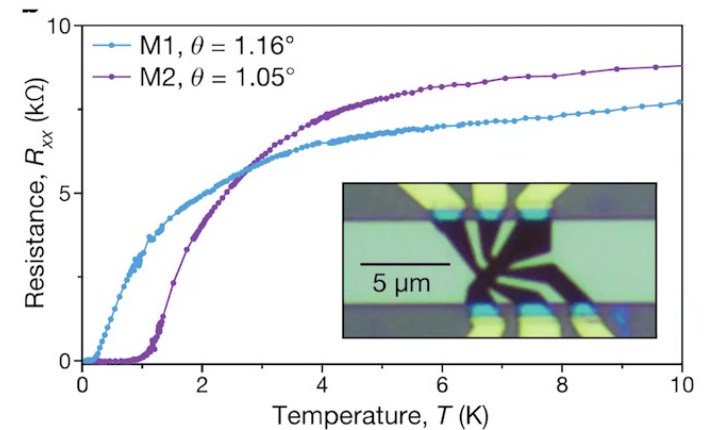
Sharpe et al., Science 365, 605–608 (2019).

Serlin et al., arXiv:1907.00261 (2019).

Trilayers double bilayers

Chen et al. Nature Physics 15, 237–241 (2019).

Chen et al., Nature 572, 215219 (2019). She et al., ArXiv 2019



Double bilayers

Shen et al., arXiv:1903.06952 (2019).

Liu et al., arXiv:1903.08130 (2019).

Cao et al., arXiv:1903.08596 (2019).

Adak et al., arXiv:2001.09916 (2020)

TMDC/TMDC; TMDC/TBG

Tang et al. arXiv:1910.08673 (2019).

Regan et al., arXiv:1910.09047 (2019).

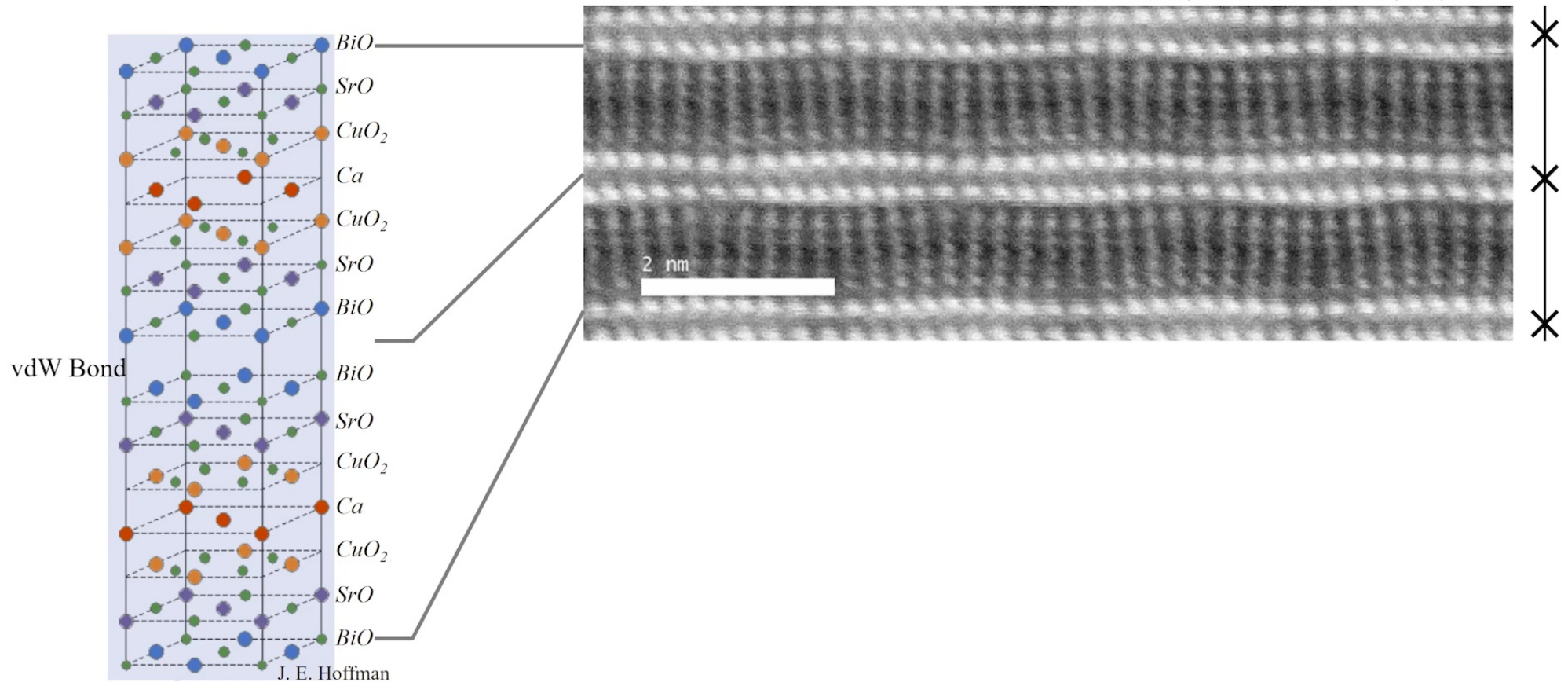
Wang et al., arXiv:1910.12147 (2019).

Arora et al., Nature 583, 379 (2020)

$\text{Bi}_2\text{Sr}_2\text{CaCu}_2\text{O}_{8+x}$: A Cuprate van der Waals Material



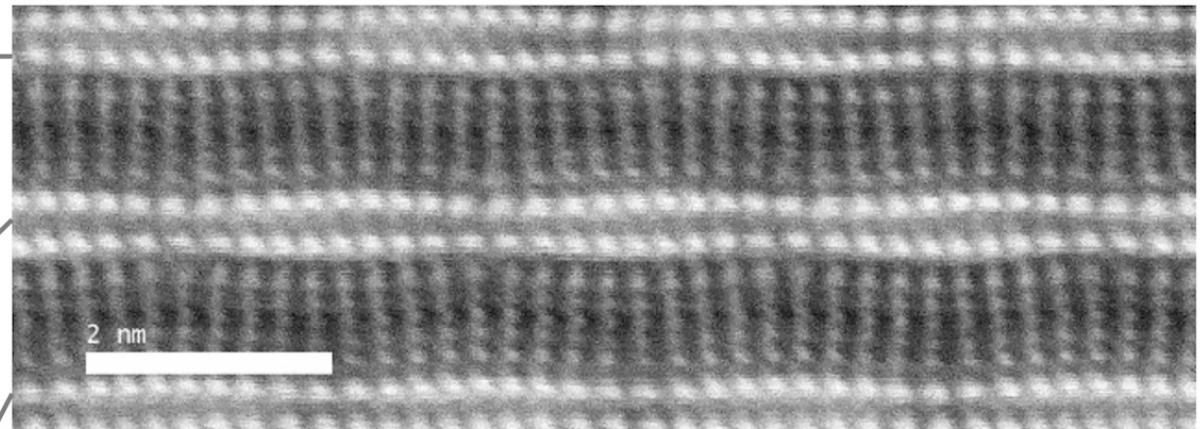
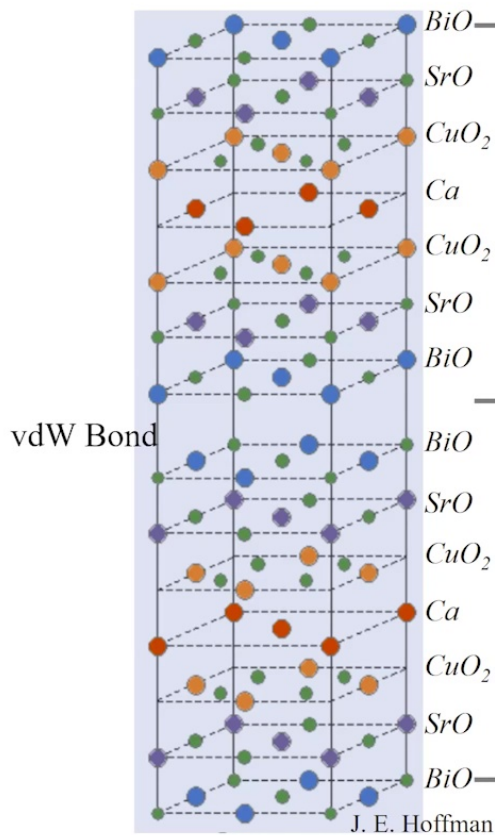
N. Poccia et al. Phys. Rev. Mater. **4**, 114007 (2020)



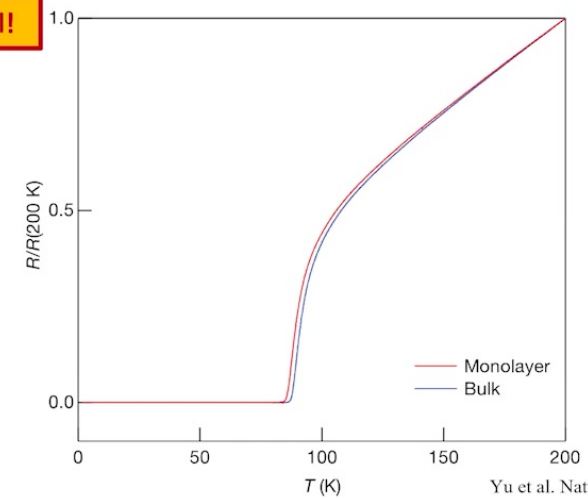
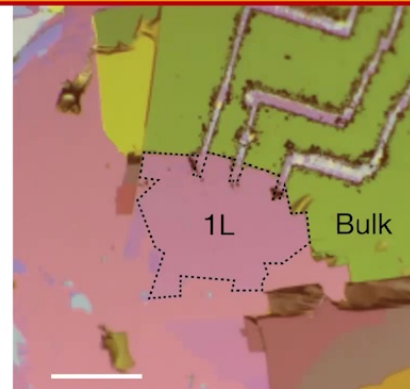
$\text{Bi}_2\text{Sr}_2\text{CaCu}_2\text{O}_{8+x}$: A Cuprate van der Waals Material



N. Poccia et al. Phys. Rev. Mater. **4**, 114007 (2020)



Superconducting down to a half unit cell!



Twisted Cuprate Junctions



nature
physics

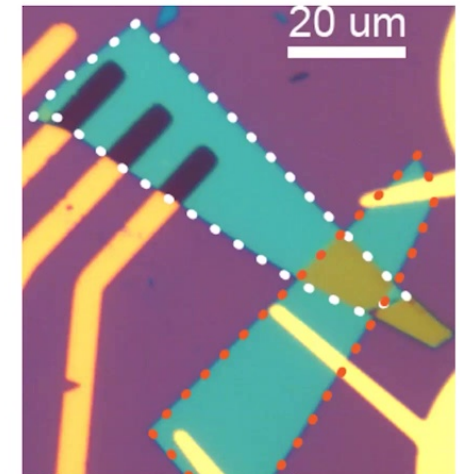
ARTICLES

<https://doi.org/10.1038/s41567-020-01142-7>



High-temperature topological superconductivity in twisted double-layer copper oxides

Oguzhan Can^{1,2}, Tarun Tummuru^{1,2}, Ryan P. Day^{1,2}, Ilya Elfimov^{1,2}, Andrea Damascelli^{1,2} and Marcel Franz^{1,2} ✉



S. Y. F Zhao et al., arXiv 2108.13455

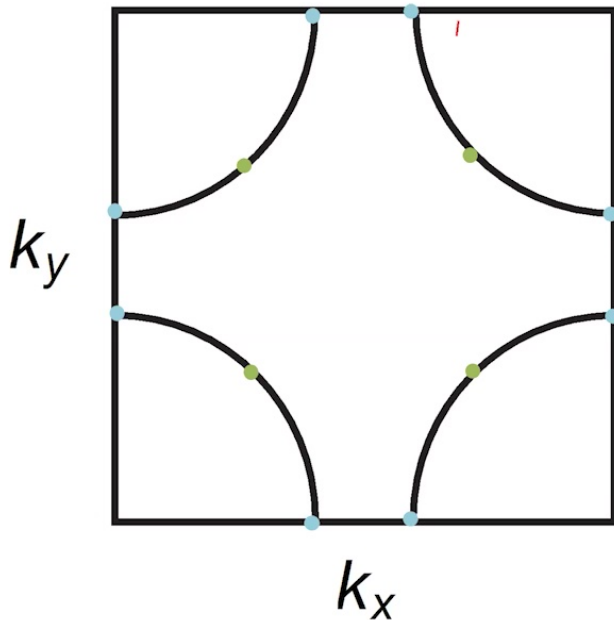
Magic angles and current-induced topology in twisted nodal superconductors

Pavel A. Volkov,^{*} Justin H. Wilson, and J. H. Pixley
*Department of Physics and Astronomy, Center for Materials Theory,
Rutgers University, Piscataway, NJ 08854, USA*
(Dated: December 16, 2020)

Nodal d-wave Superconductor



Cuprate Brillouin Zone



D-wave order parameter

$$\Delta_{\mathbf{k}} = \Delta_0 (\cos k_x - \cos k_y)$$

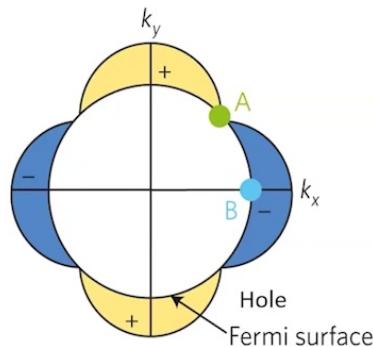
BdG Hamiltonian:

$$\hat{H} = \sum_{\mathbf{k}} \begin{pmatrix} c_{\uparrow}(\mathbf{k}) \\ c_{\downarrow}^{\dagger}(\mathbf{k}) \end{pmatrix}^{\dagger} \begin{pmatrix} \varepsilon(\mathbf{k}) & \Delta(\mathbf{k}) \\ \Delta^*(\mathbf{k}) & -\varepsilon(\mathbf{k}) \end{pmatrix} \begin{pmatrix} c_{\uparrow}(\mathbf{k}) \\ c_{\downarrow}^{\dagger}(\mathbf{k}) \end{pmatrix}$$

$$\approx \sum_{\mathbf{k}} \begin{pmatrix} c_{\uparrow}(\mathbf{k}) \\ c_{\downarrow}^{\dagger}(\mathbf{k}) \end{pmatrix}^{\dagger} \begin{pmatrix} \mathbf{v}_F \cdot \mathbf{p} & \mathbf{v}_{\Delta} \cdot \mathbf{p} \\ \mathbf{v}_{\Delta} \cdot \mathbf{p} & -\mathbf{v}_F \cdot \mathbf{p} \end{pmatrix} \begin{pmatrix} c_{\uparrow}(\mathbf{k}) \\ c_{\downarrow}^{\dagger}(\mathbf{k}) \end{pmatrix}$$

near the nodal points

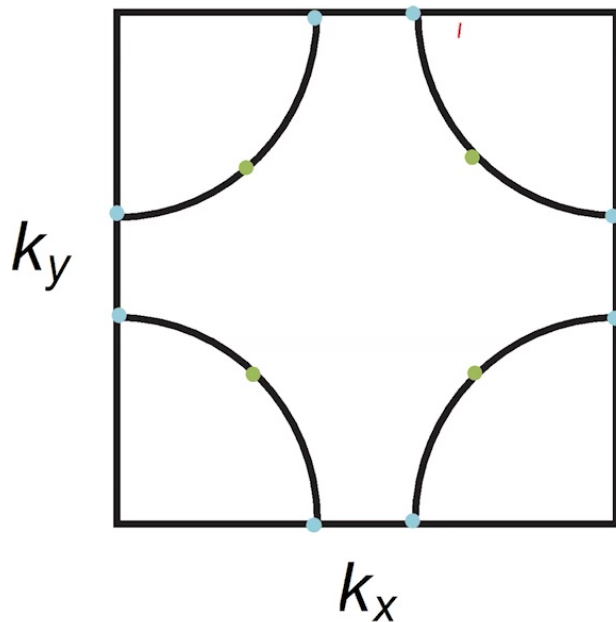
Gap vanishes at the nodal points



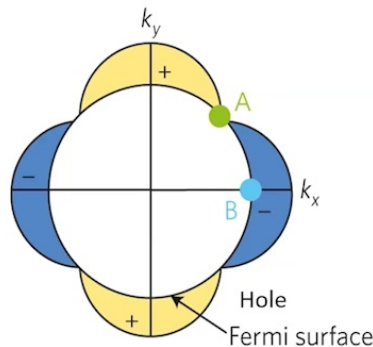
Nodal d-wave Superconductor



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D-wave order parameter

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BdG Hamiltonian:

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$$\approx \sum_{\mathbf{k}} \begin{pmatrix} c_{\uparrow}(\mathbf{k}) \\ c_{\downarrow}(\mathbf{k}) \end{pmatrix}^{\dagger} \begin{pmatrix} \mathbf{v}_F \cdot \mathbf{p} & \mathbf{v}_{\Delta} \cdot \mathbf{p} \\ \mathbf{v}_{\Delta} \cdot \mathbf{p} & -\mathbf{v}_F \cdot \mathbf{p} \end{pmatrix} \begin{pmatrix} c_{\uparrow}(\mathbf{k}) \\ c_{\downarrow}(\mathbf{k}) \end{pmatrix}$$

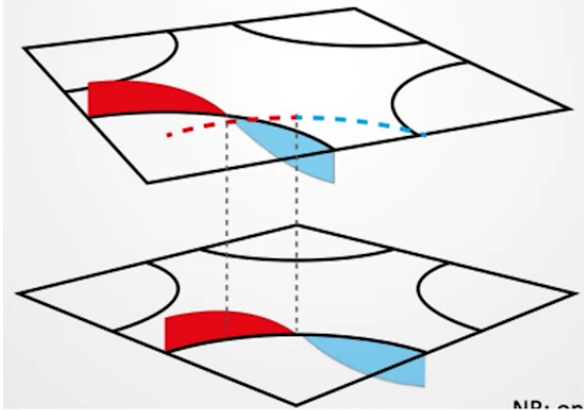
near the nodal points

$$E_{0,\pm}(\mathbf{k}) = \pm \sqrt{v_F^2 k_{\parallel}^2 + v_{\Delta}^2 k_{\perp}^2}$$

Twisted Cuprate Junction



Stacked twist layer



Volkov et al. arXiv: 2012.07860

Double layer Nambu basis

$$\Psi_{\mathbf{k}} = (c_{\mathbf{k}\uparrow 1}, c_{-\mathbf{k}\downarrow 1}^\dagger, c_{\mathbf{k}\uparrow 2}, c_{-\mathbf{k}\downarrow 2}^\dagger)^T$$

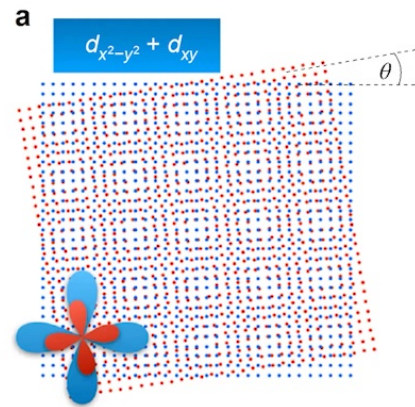
BdG Hamiltonian

$$\mathcal{H} = \sum_{\mathbf{k}} \Psi_{\mathbf{k}}^\dagger h_{\mathbf{k}} \Psi_{\mathbf{k}} + E_0.$$

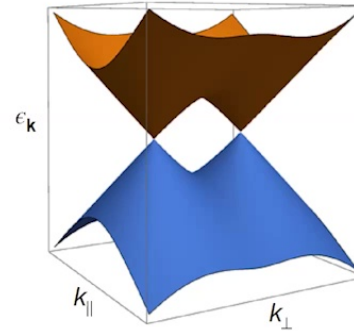
$$h_{\mathbf{k}} = \begin{pmatrix} \xi_{\mathbf{k}} & \Delta_{\mathbf{k}1} & g & 0 \\ \Delta_{\mathbf{k}1}^* & -\xi_{\mathbf{k}} & 0 & -g \\ g & 0 & \xi_{\mathbf{k}} & \Delta_{\mathbf{k}2} \\ 0 & -g & \Delta_{\mathbf{k}2}^* & -\xi_{\mathbf{k}} \end{pmatrix}$$

layer1 Interlayer coupling layer2

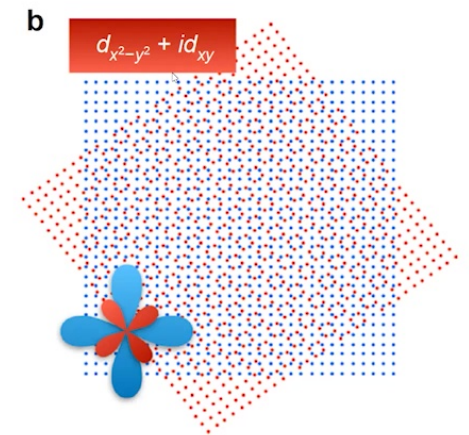
Small twist angle



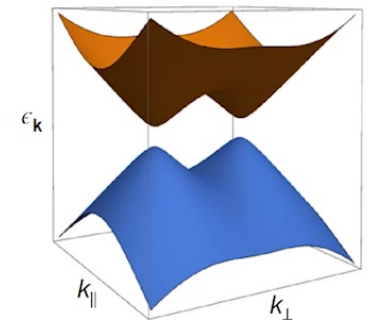
$$\theta = 9.53^\circ < \theta_c^-$$



Large twist angle



$$\theta = 36.87^\circ > \theta_c^-$$



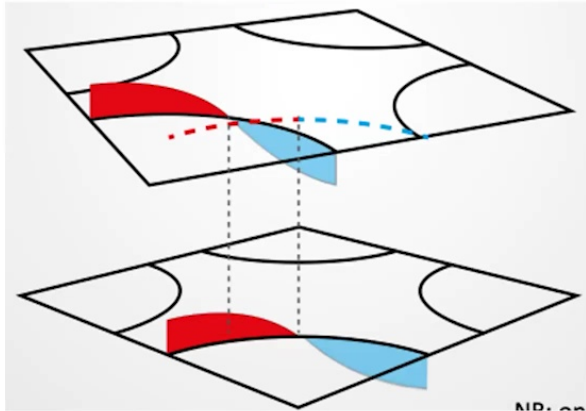
Can et al. Nature Phys (2021)

Twisted Cuprate Junction



SC with broken TRS

Stacked twist layer



Volkov et al. arXiv: 2012.07860

Double layer Nambu basis

$$\Psi_{\mathbf{k}} = (c_{\mathbf{k}\uparrow 1}, c_{-\mathbf{k}\downarrow 1}^\dagger, c_{\mathbf{k}\uparrow 2}, c_{-\mathbf{k}\downarrow 2}^\dagger)^T$$

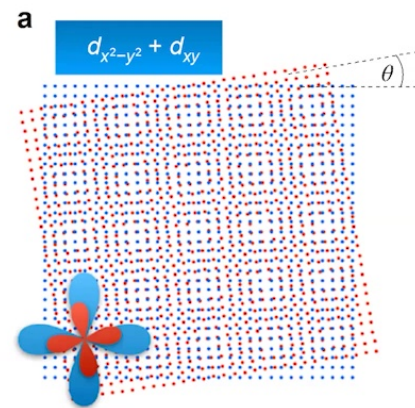
BdG Hamiltonian

$$\mathcal{H} = \sum_{\mathbf{k}} \Psi_{\mathbf{k}}^\dagger h_{\mathbf{k}} \Psi_{\mathbf{k}} + E_0.$$

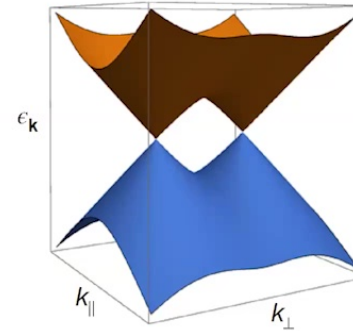
$$h_{\mathbf{k}} = \begin{pmatrix} \xi_{\mathbf{k}} & \Delta_{\mathbf{k}1} & g & 0 \\ \Delta_{\mathbf{k}1}^* & -\xi_{\mathbf{k}} & 0 & -g \\ g & 0 & \xi_{\mathbf{k}} & \Delta_{\mathbf{k}2} \\ 0 & -g & \Delta_{\mathbf{k}2}^* & -\xi_{\mathbf{k}} \end{pmatrix}$$

layer1 Interlayer coupling layer2

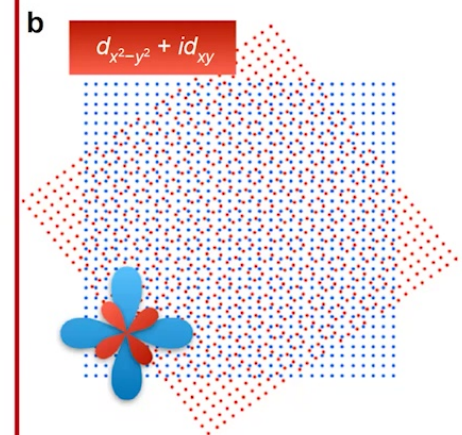
Small twist angle



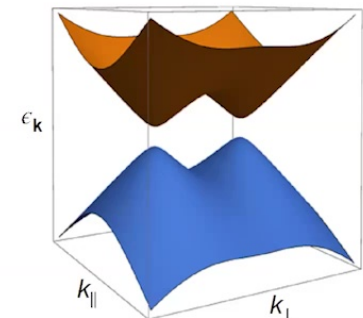
$$\theta = 9.53^\circ < \theta_c^-$$



Large twist angle

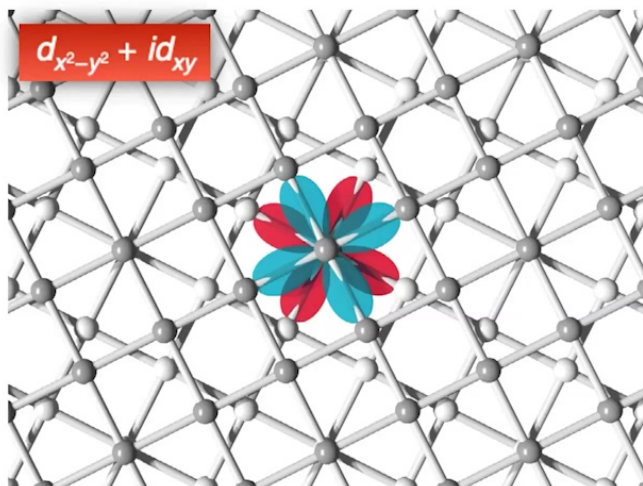


$$\theta = 36.87^\circ > \theta_c^-$$



Can et al. Nature Phys (2021)

45° Twisted Junction: Superconductivity with Broken TRS



BdG Free Energy

$$\mathcal{F}(\varphi) = E_0 - 2k_B T \sum_{k,a=\pm} \ln[2\cosh[E_{ka}(\varphi)/2k_B T]]$$

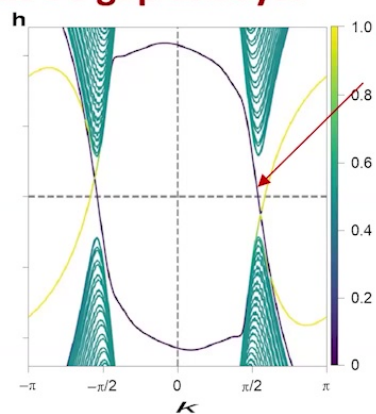
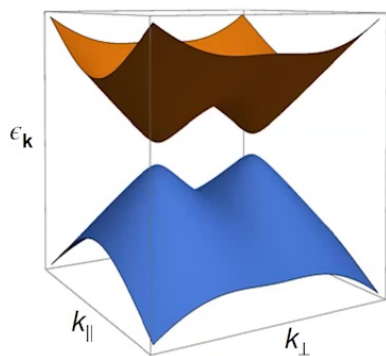
Ginzburg Landau Free Energy

$$\mathcal{F}(\varphi) = \mathcal{F}_0 + 2B_0\psi^2[-\cos(2\theta)\cos\varphi + \mathcal{K}\cos(2\varphi)]$$

Phase diff. between SCs

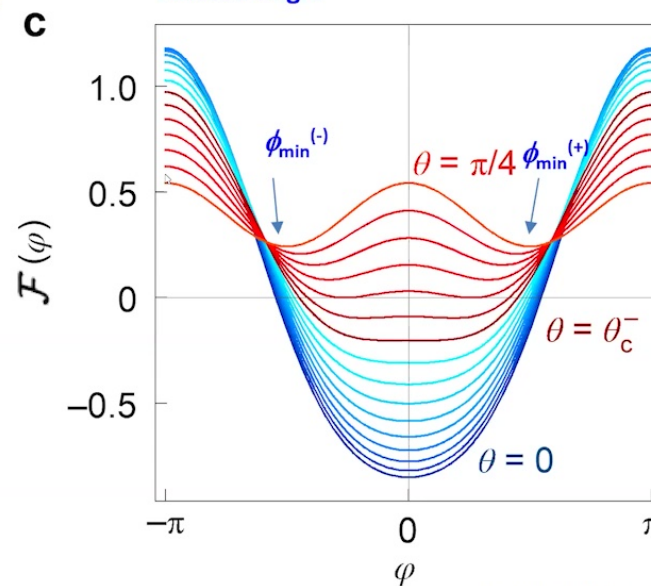
Twisted angle

No nodal points: bulk has SC gap always!



Topological edge states in the gap may appear depending the coupling parameters

BdG spectrum of stripe SC

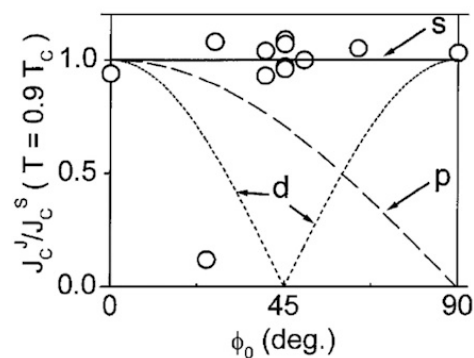
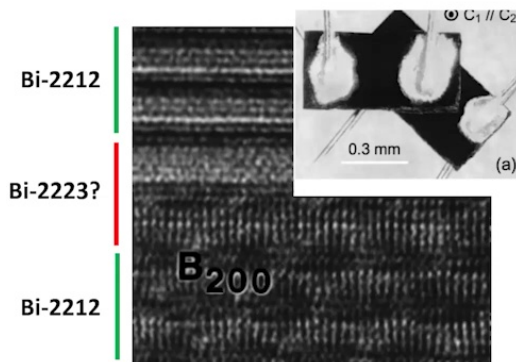


Can et al. Nature Phys (2021)

Experiments on Twist Cuprate Josephson Junctions

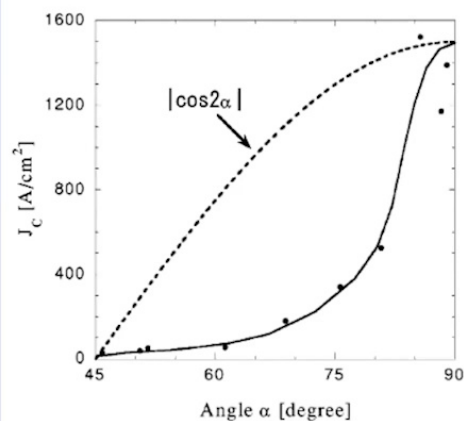
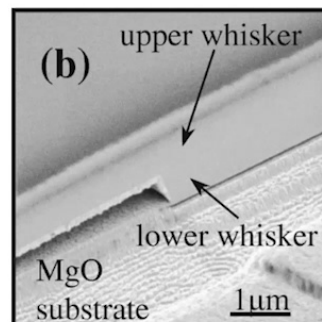


Bulk Crystals
30 hrs
'below melting temperature'



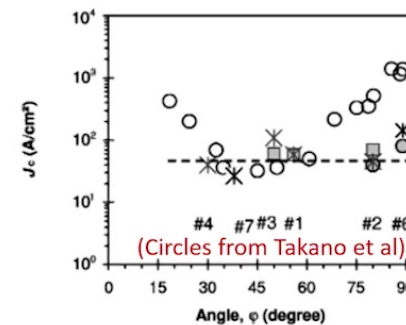
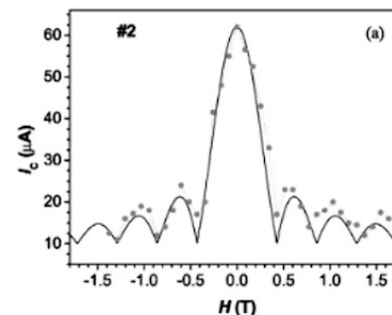
Li et al, PRL 83, 4160 (1999); Zhu et al PRB 57 8601 (1998)

Artificial Whisker Junctions
30 min
850 °C



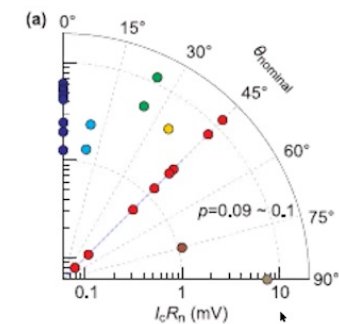
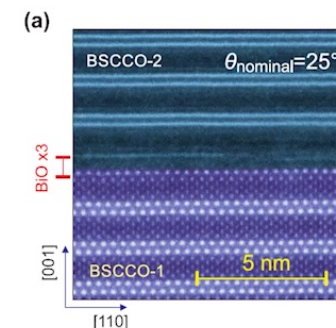
Takano et al, PRB 65, 140513 (2002)

Natural Whisker Junctions
20 min
845 °C



Latyshev et al, PRB 70, 094517 (2004)

Exfoliated Flakes
10 min
450-530 °C



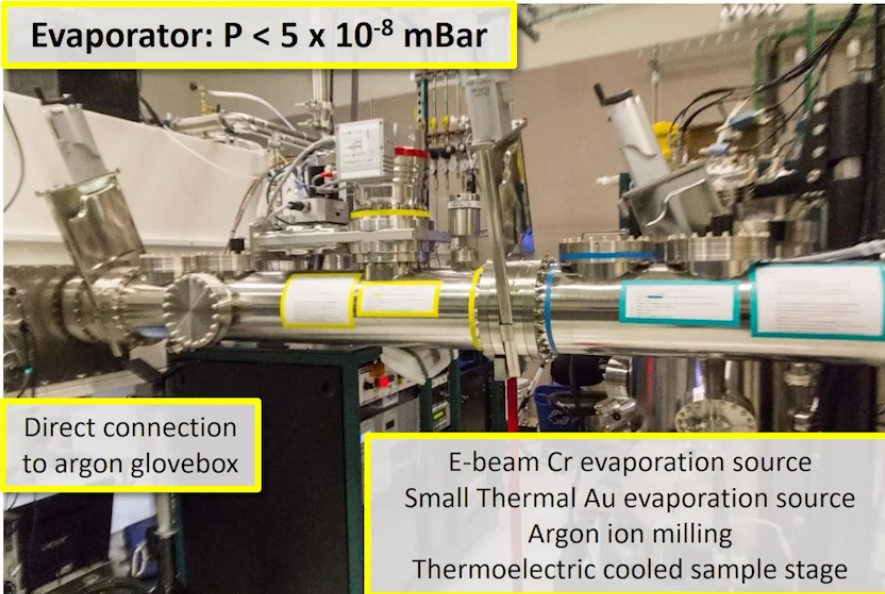
Zhu et al, ArXiv 1903.07965 (2021)

Fab in a Glovebox UHV



- *Stacking*
- *Selection*
- ✓ *Deposition*
- ✓ *Etch*
- ✓ *Packaging*

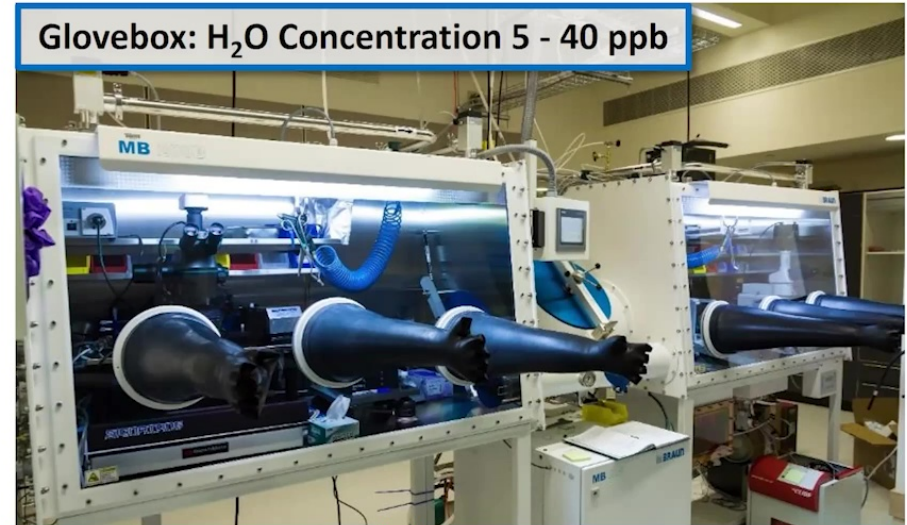
Evaporator: $P < 5 \times 10^{-8}$ mBar



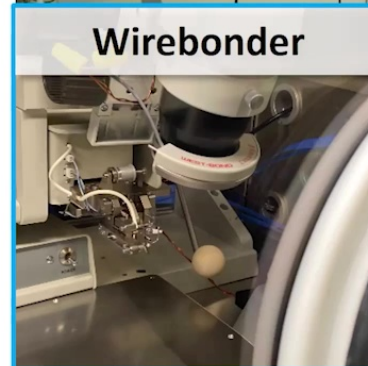
Direct connection
to argon glovebox

E-beam Cr evaporation source
Small Thermal Au evaporation source
Argon ion milling
Thermoelectric cooled sample stage

Glovebox: H_2O Concentration 5 - 40 ppb



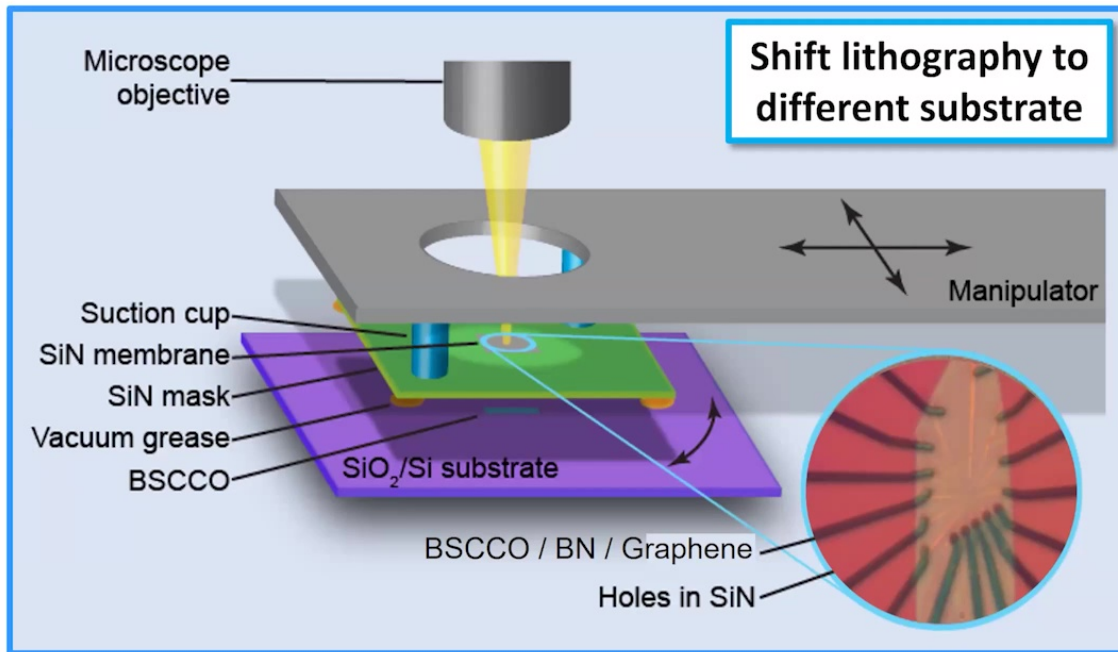
Wirebonder



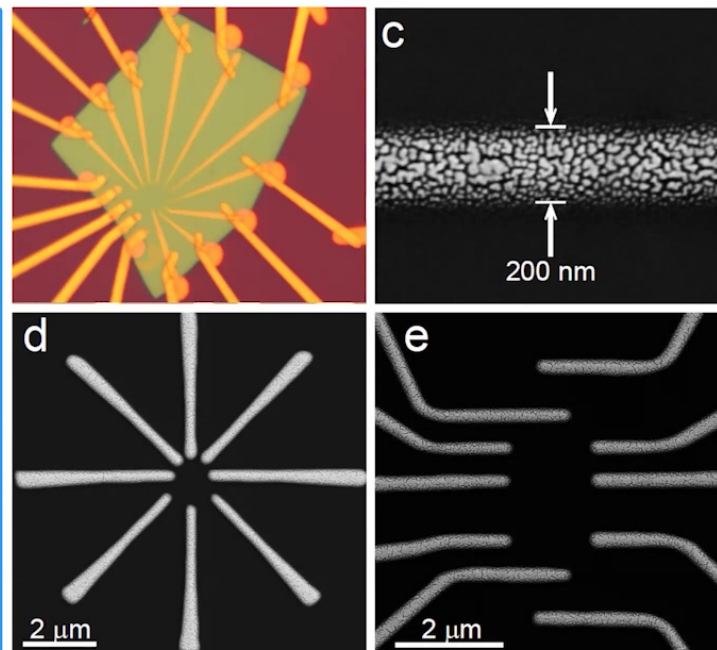
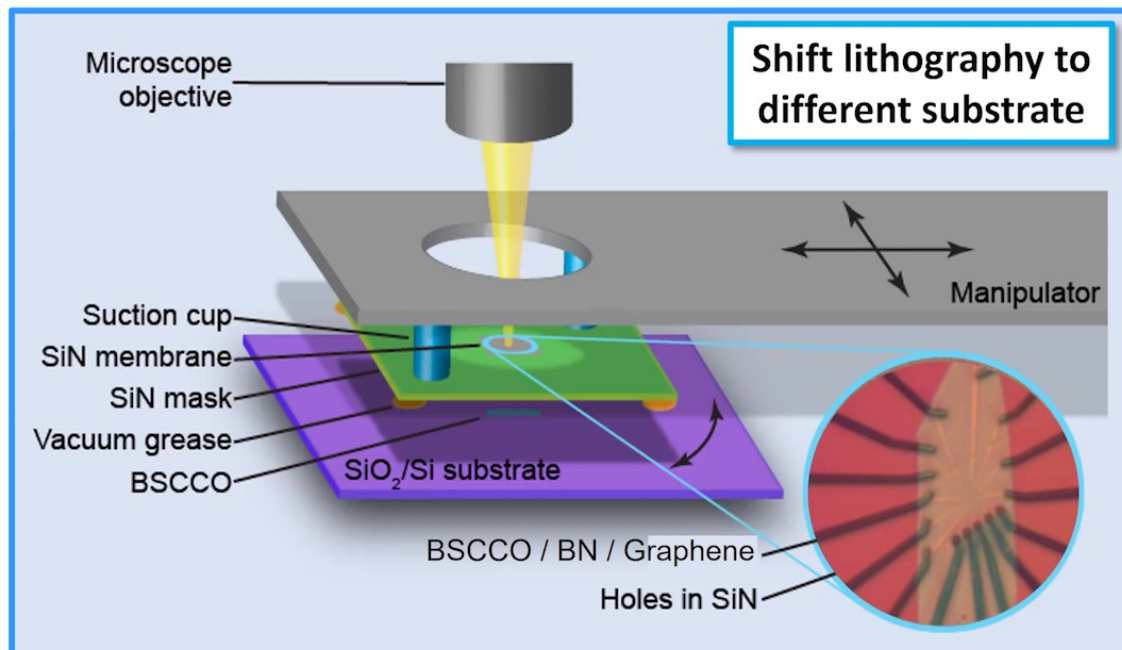
Transfer Stage



Stencil Mask Lithography



Stencil Mask Lithography



Technique advantages:

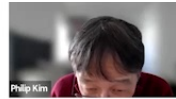
- No air exposure
- No heating
- No solvent, no residue.

Technique capabilities

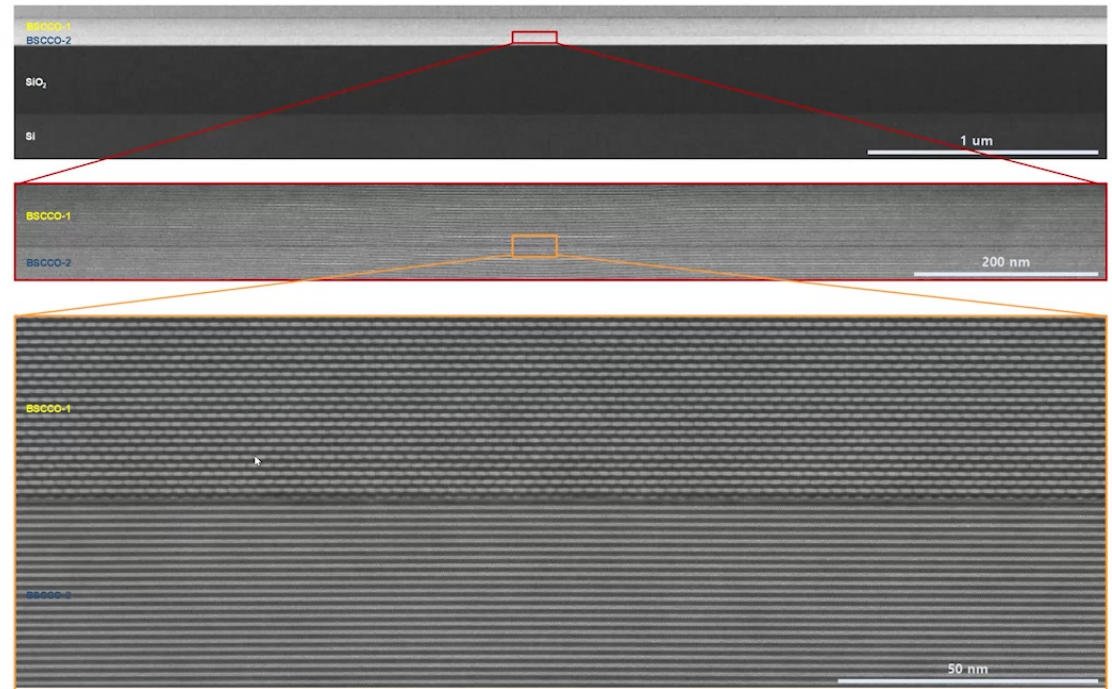
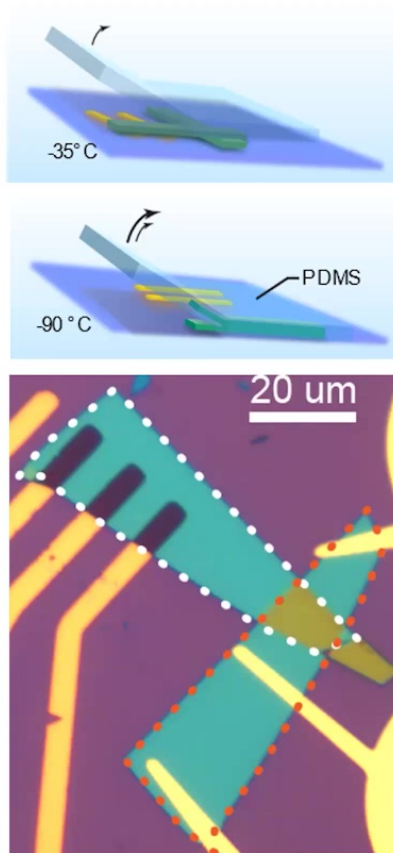
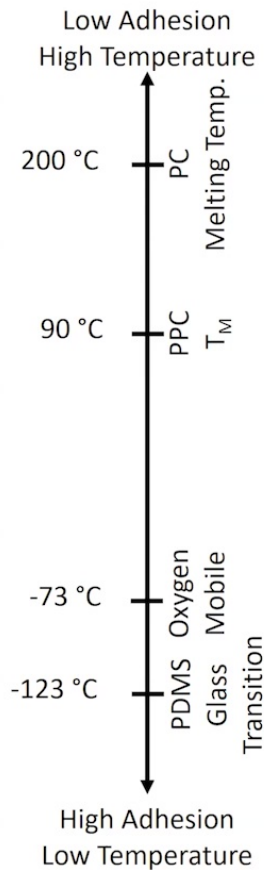
Resolution: < 200 nm

Alignment: < 1 μm

Ultra-Clean twisted BSCCO Interface



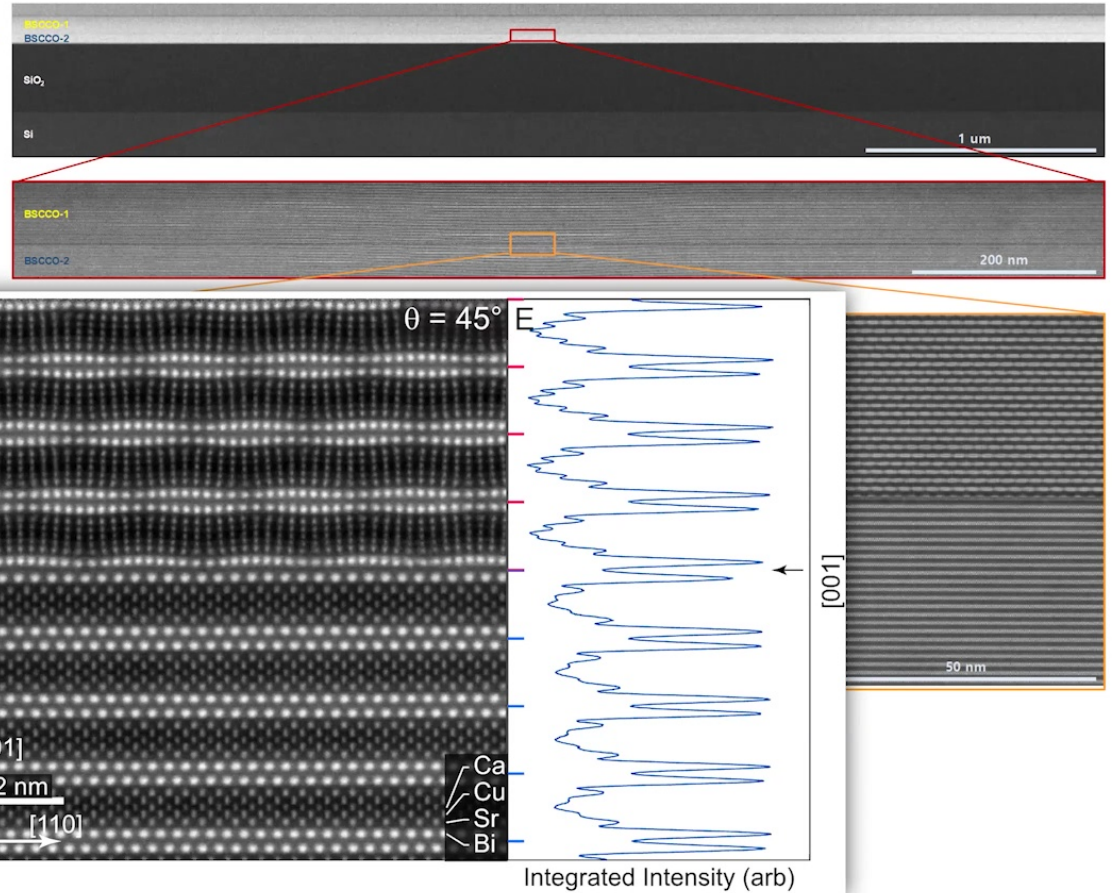
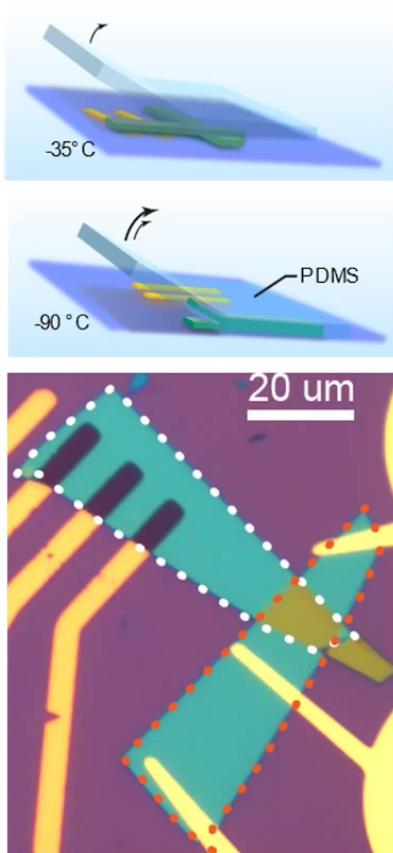
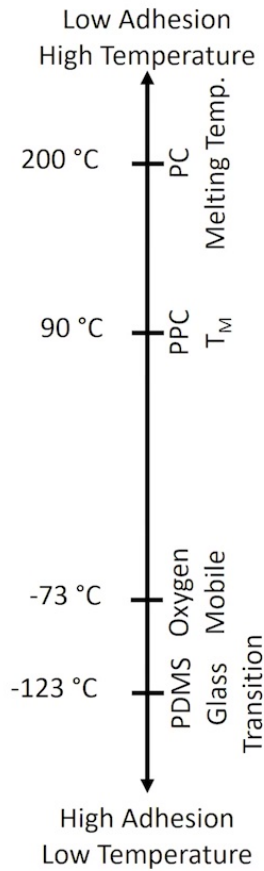
Stacking: Cryogenic PDMS Pickup



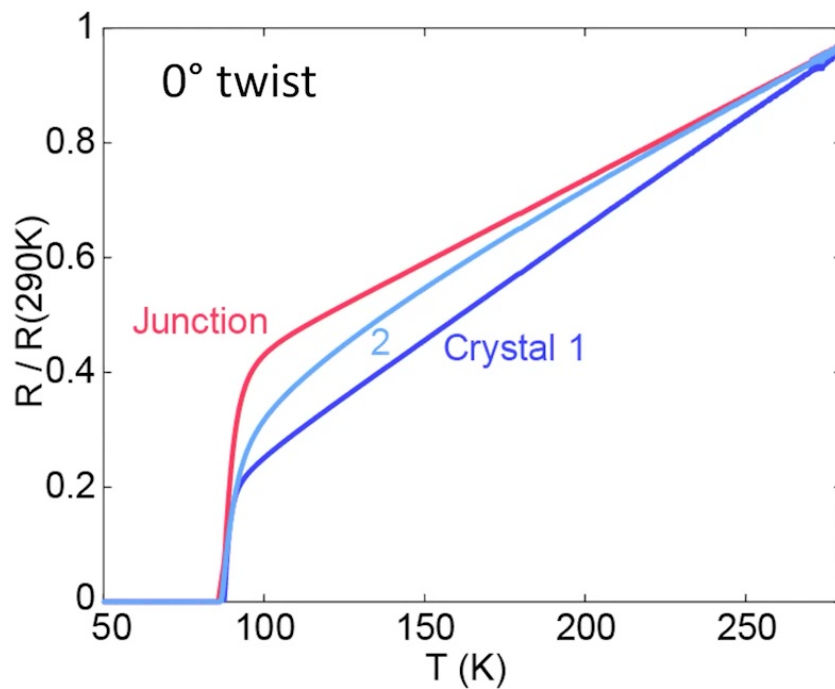
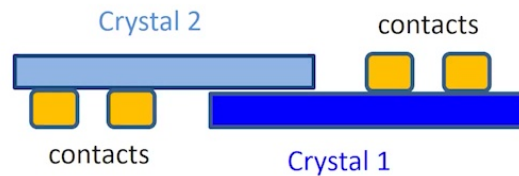
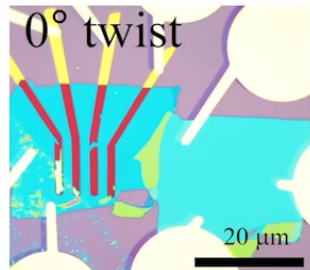
Ultra-Clean twisted BSCCO Interface



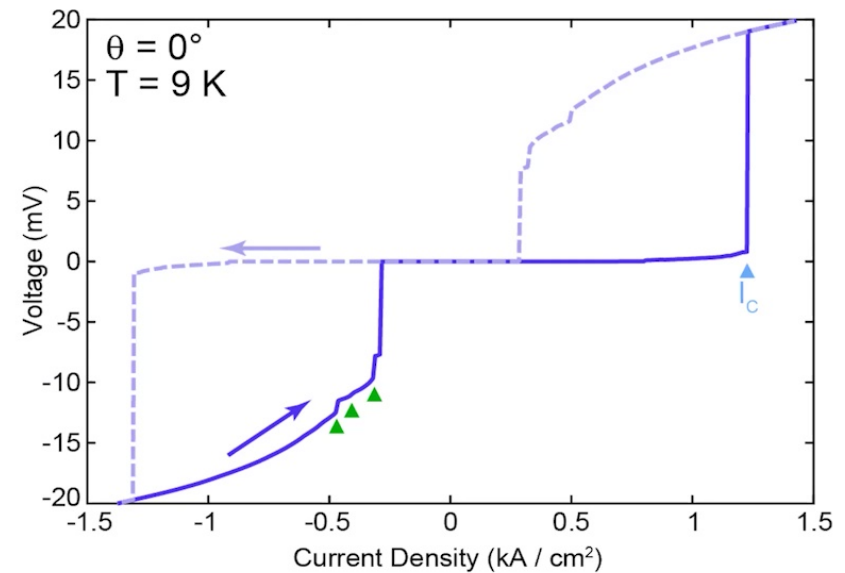
Stacking: Cryogenic PDMS Pickup



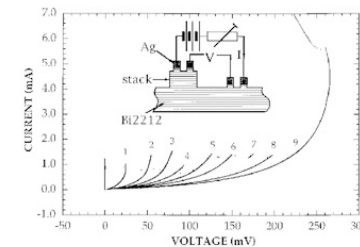
Artificial Vertical Josephson Junction Characteristics



Josephson Junction Characteristics: 0° twist



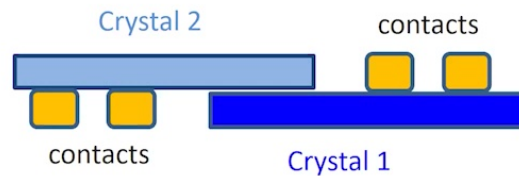
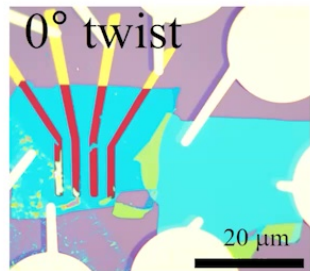
Critical switching current density is smaller than but comparable to intrinsic BSCCO JJ.



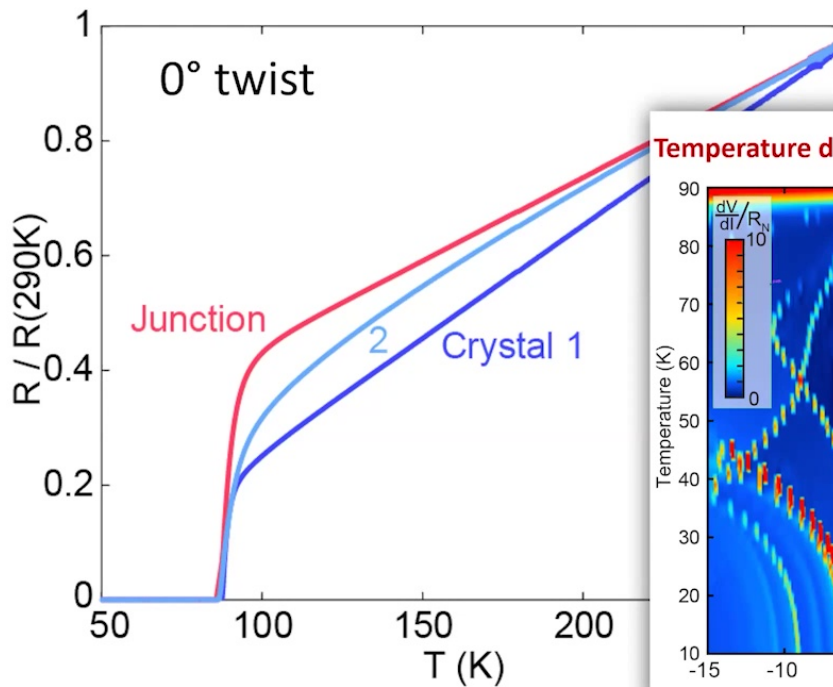
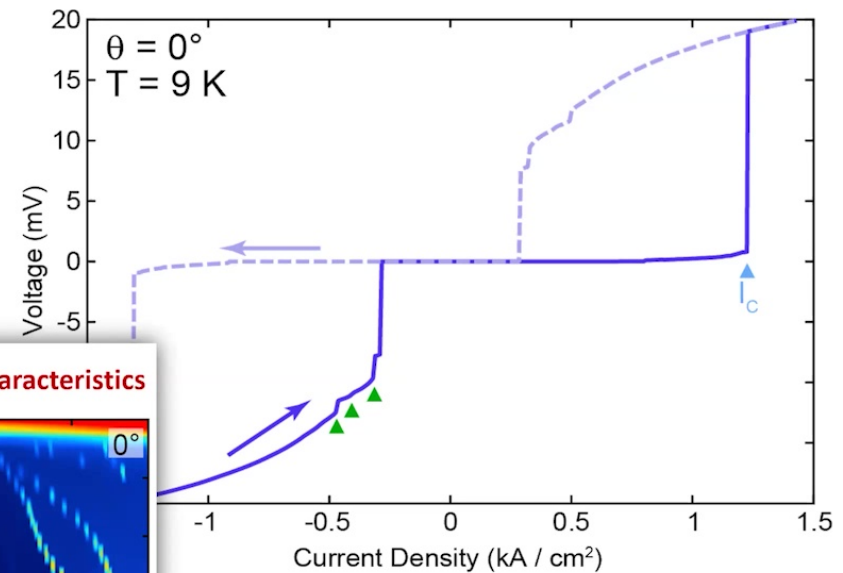
BSCCO intrinsic Junction
Junction Area $\sim 1000\text{ }\mu\text{m}^2$

Yurgens et. al., PRB 53 R8887 (1996)

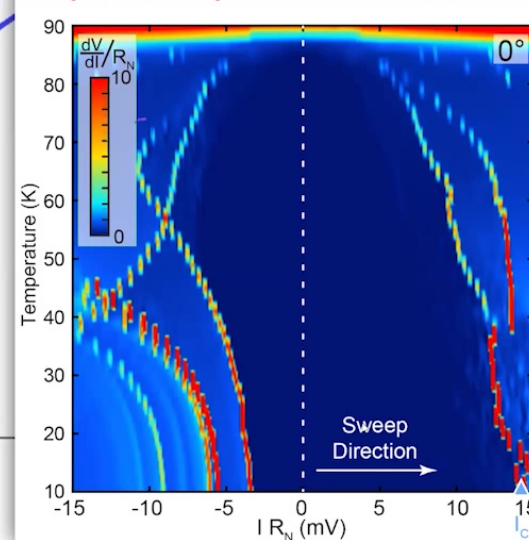
Artificial Vertical Josephson Junction Characteristics



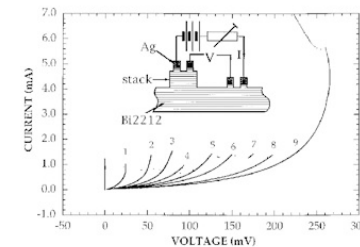
Josephson Junction Characteristics: 0° twist



Temperature dependent JJ Characteristics



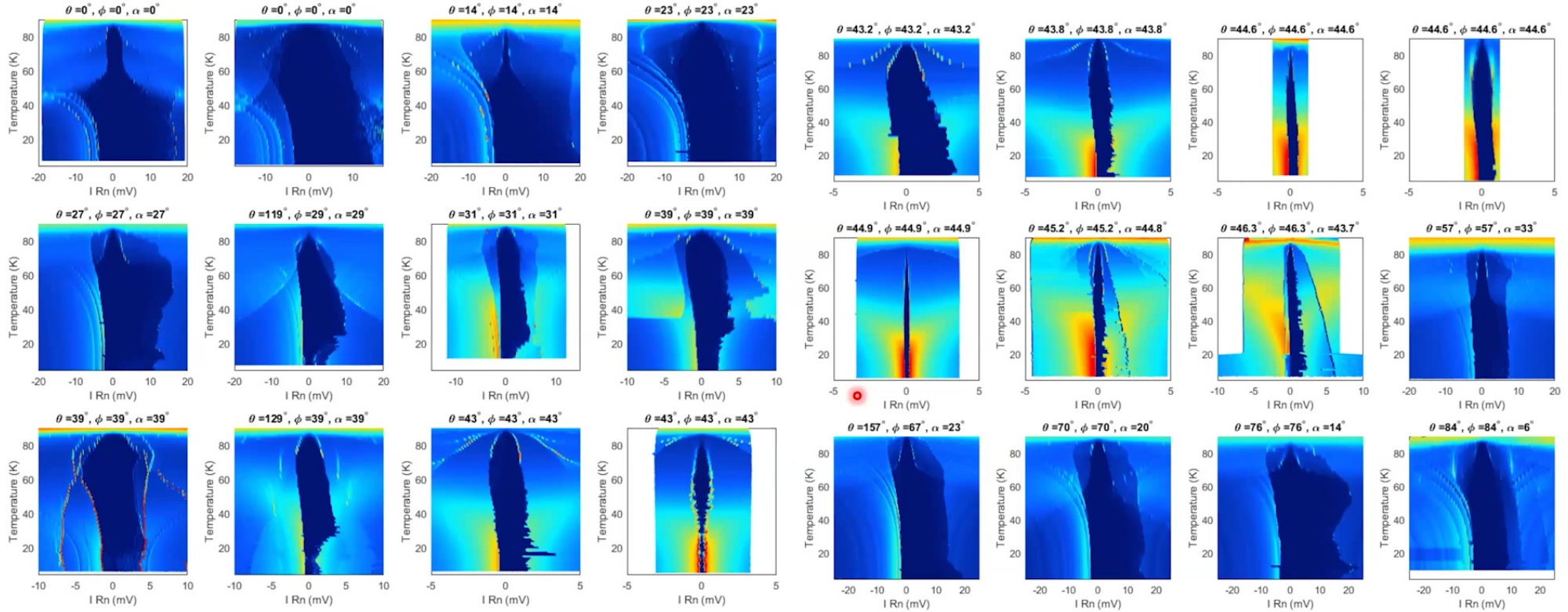
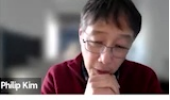
Critical switching current density is smaller than but comparable to intrinsic BSCCO JJ.



BSCCO intrinsic Junction
Junction Area $\sim 1000 \mu\text{m}^2$

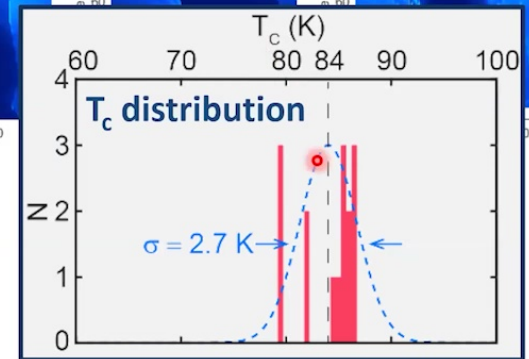
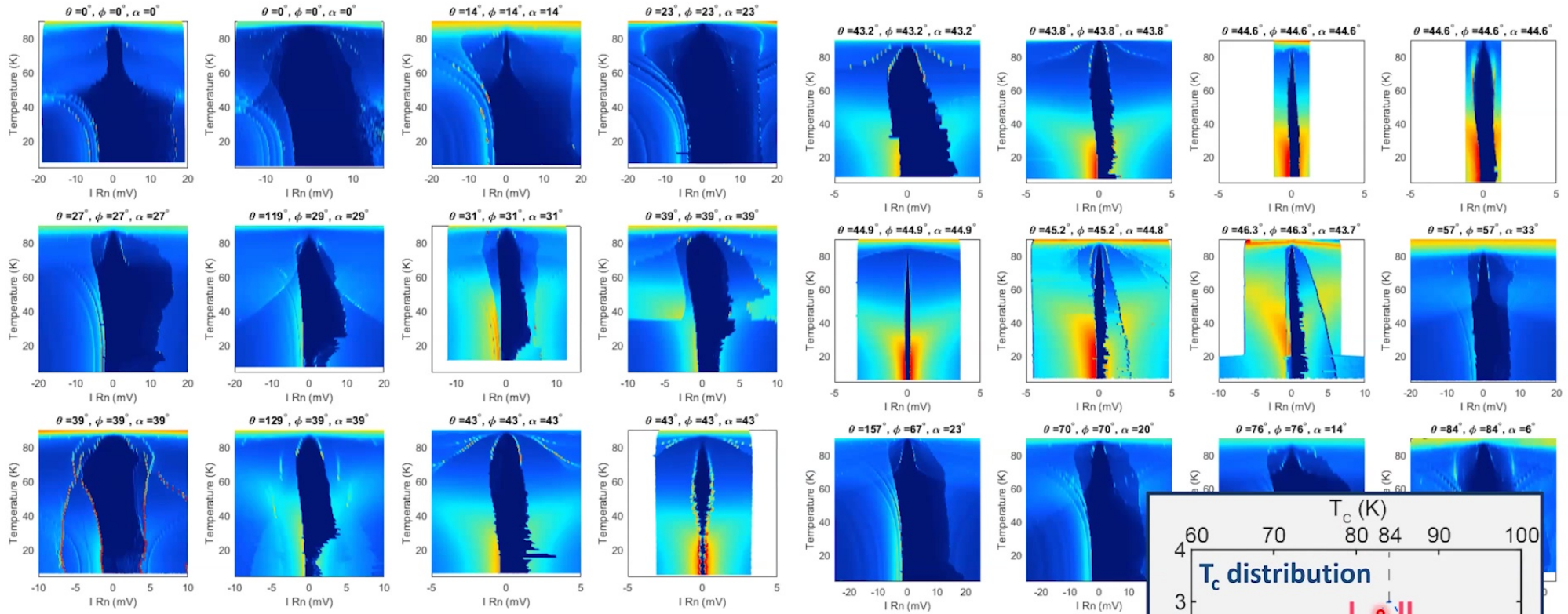
Yurgens et. al., PRB 53 R8887 (1996)

Twisted BSCCO vdW Josephson Junctions



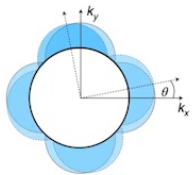
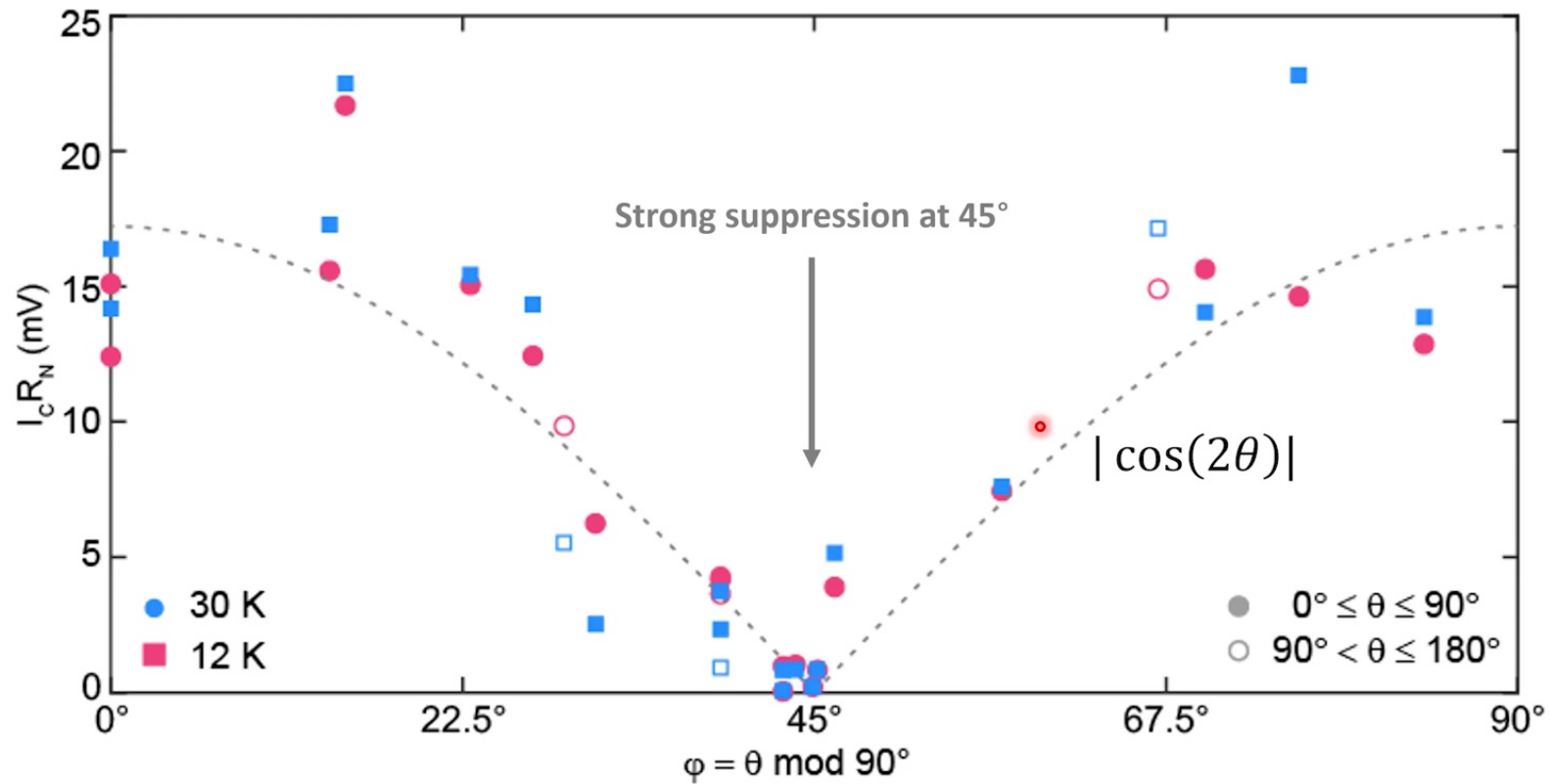
More than 24 junctions (so far): θ ranges 0° to 180°

Twisted BSCCO vdW Josephson Junctions



More than 24 junctions (so far): θ ranges 0° to 180°

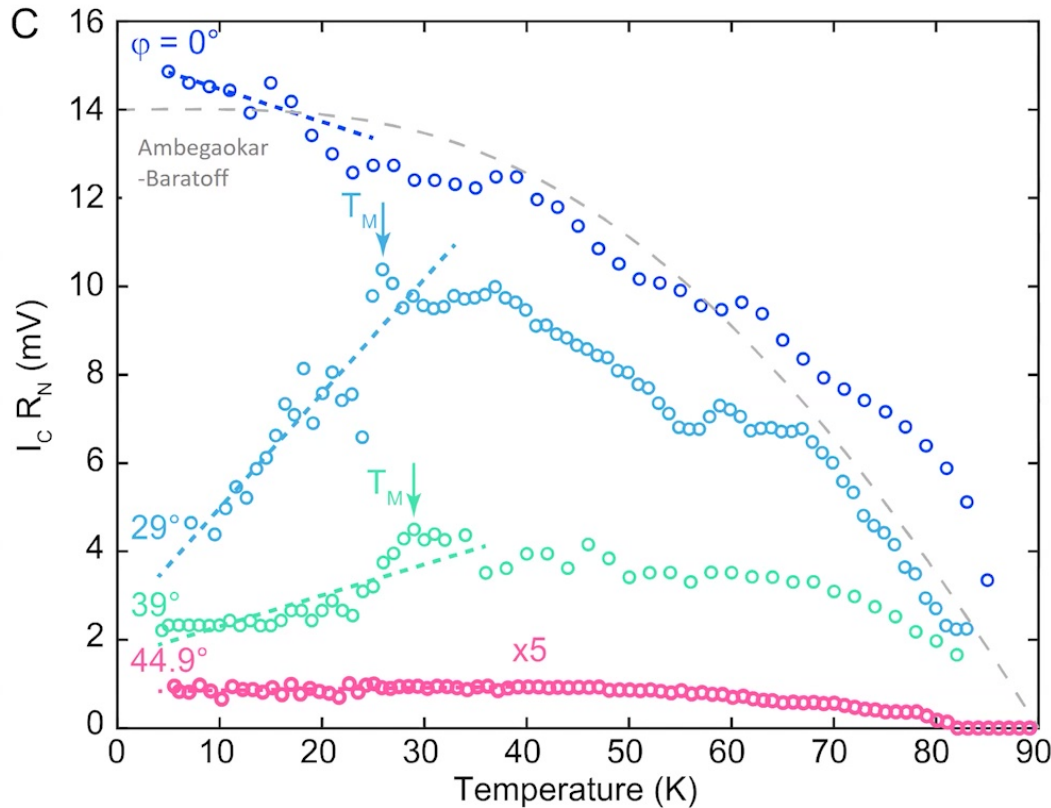
Twisting Angle Dependence $I_c R_N$



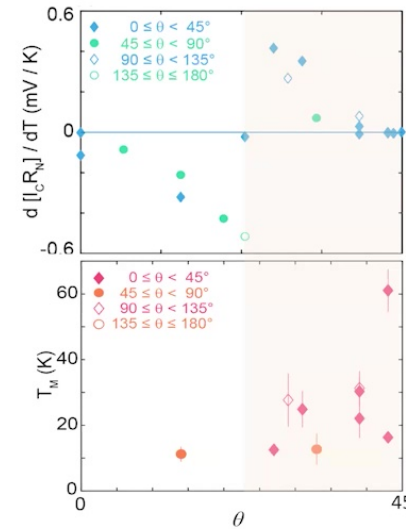
$$\Delta_{k1} = \Delta \cos(2\alpha_k + \theta), \quad \Delta_{k2} = \Delta \cos(2\alpha_k - \theta)$$

Twisted d-wave order parameters

Temperature Dependent Critical Current



Non-monotonic temperature dependent I_c



$0 \leq \theta \lesssim \pi/8$

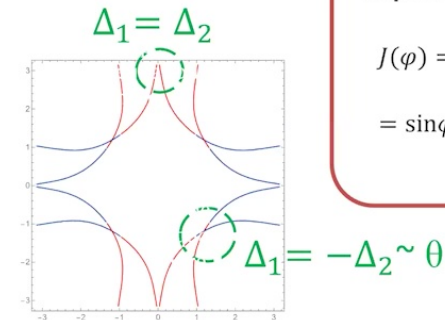
Decreasing I_c with increasing T

$\pi/8 \lesssim \theta \lesssim \pi/4$

Increasing I_c with increasing T
at low temperature (non-monotonic)

$\theta \approx \pi/4$

Suppressed (but finite) I_c



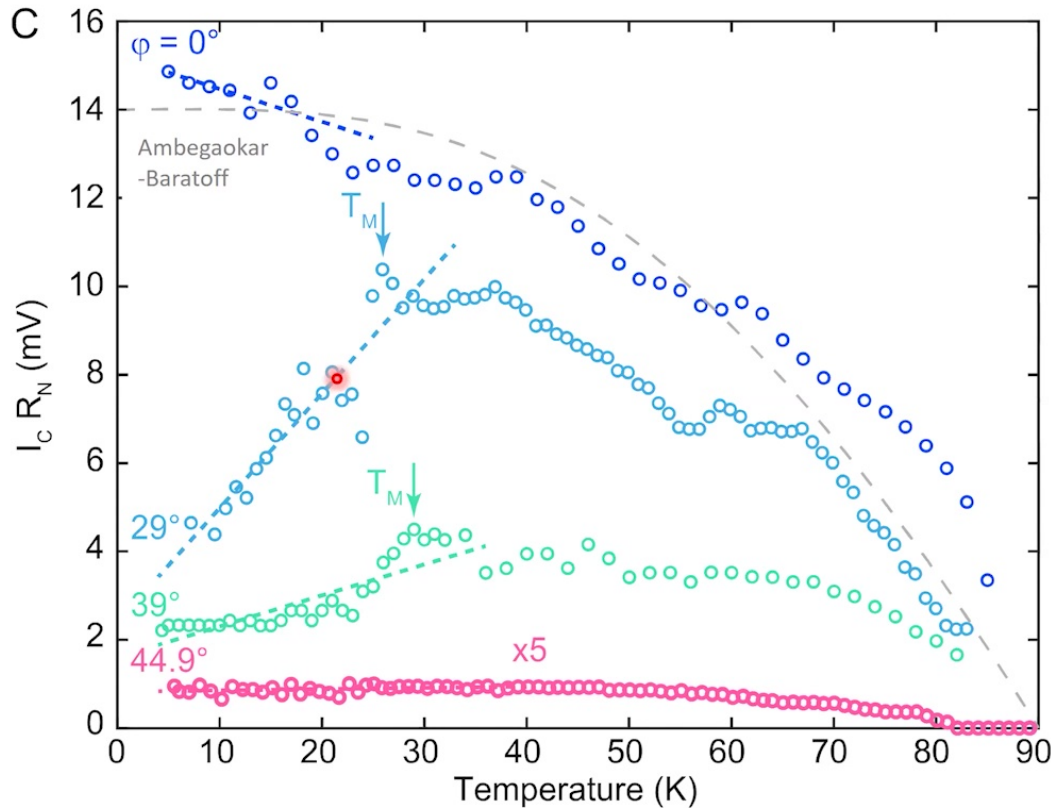
Supercurrent

$$J(\varphi) = \frac{2e}{\hbar} \frac{\partial F(\varphi)}{\partial \varphi}$$

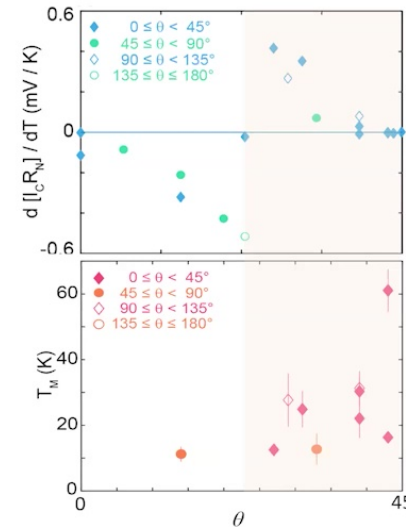
$$= \sin \varphi \frac{e g^2}{\hbar} \sum_k \Delta_{k1} \Delta_{k2} \sum_{a=\pm} \frac{a \tanh(E_{ka}/2k_B T)}{E_{ka}(E_{k-}^2 - E_{k+}^2)}$$

> 0

Temperature Dependent Critical Current



Non-monotonic temperature dependent I_c



$0 \leq \theta \lesssim \pi/8$

Decreasing I_c with increasing T

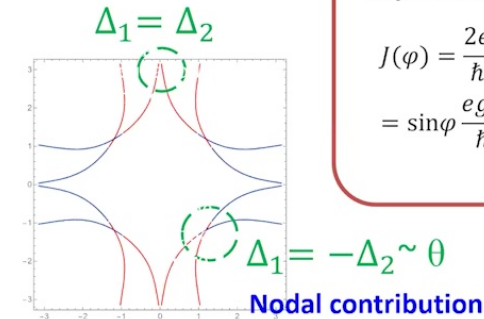
$\pi/8 \lesssim \theta \lesssim \pi/4$

Increasing I_c with increasing T
at low temperature (non-monotonic)

$\theta \approx \pi/4$

Suppressed (but finite) I_c

Anti-nodal contribution



Supercurrent

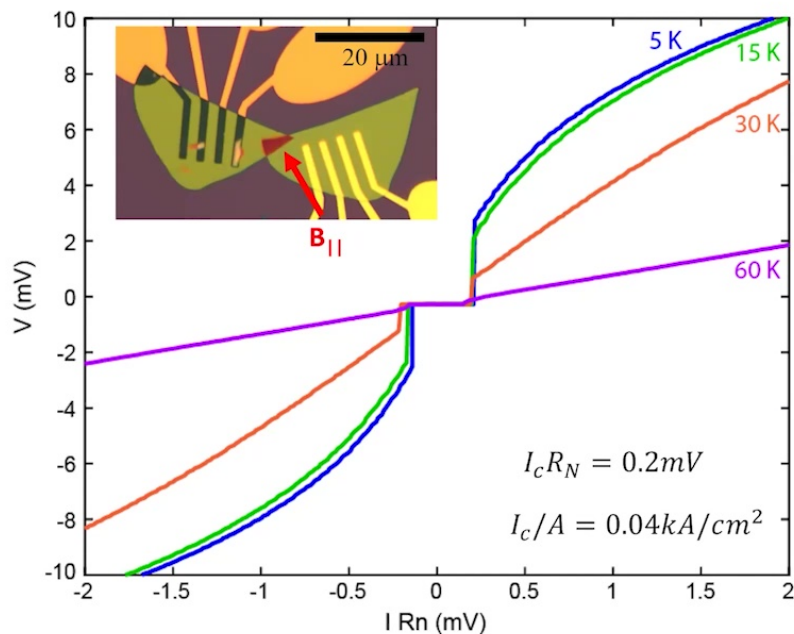
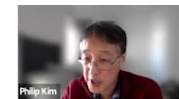
$$J(\varphi) = \frac{2e}{\hbar} \frac{\partial F(\varphi)}{\partial \varphi}$$

$$= \sin \varphi \frac{e g^2}{\hbar} \sum_k \Delta_{k1} \Delta_{k2} \sum_{a=\pm} \frac{a \tanh(E_{ka}/2k_B T)}{E_{ka}(E_{k-}^2 - E_{k+}^2)}$$

> 0

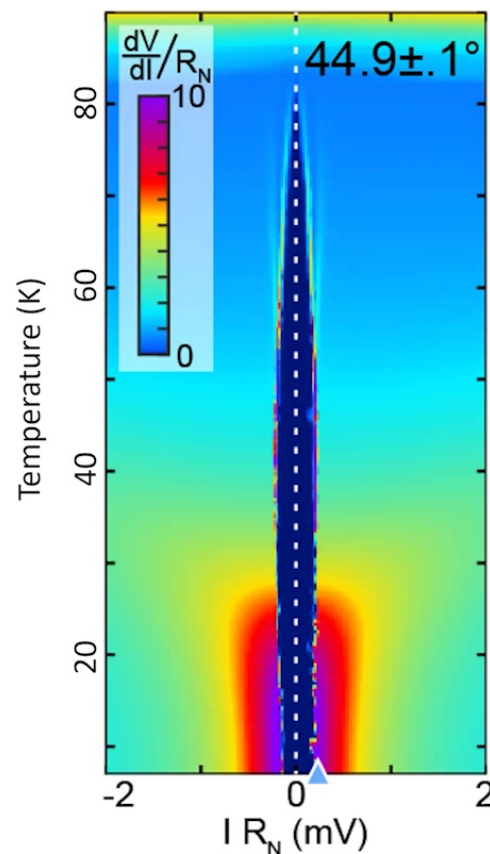
Competing nodal and anti-nodal supercurrent contributions!

Near 45° Twisted Josephson Junctions

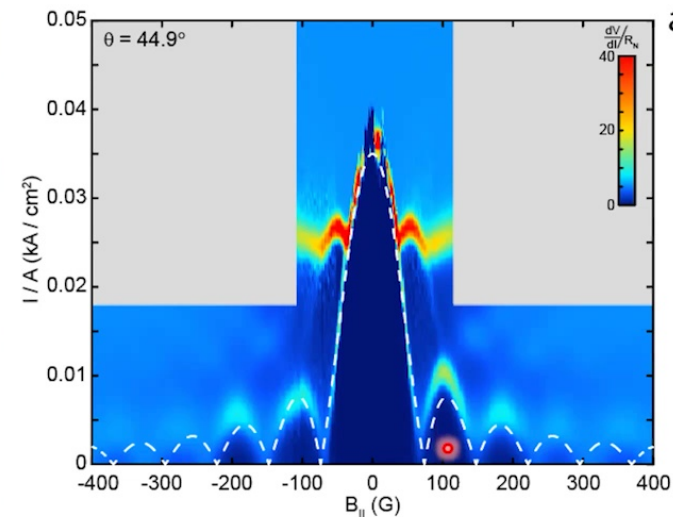


Reduced (~ 0.1 times of 0-twist), but finite $I_c R_N$

Temperature Dependent Critical Current



Fraunhofer Pattern with in-plane magnetic field



Well developed FP patterns suggest high quality Josephson Junction.

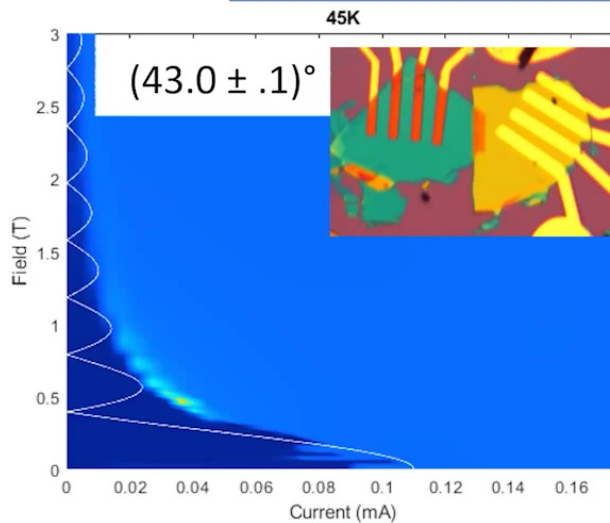
Fitting function: $I_c = I_{c0} \left| \frac{\sin(\frac{\pi\phi}{\phi_0})}{\frac{\pi\phi}{\phi_0}} \right|$

$\phi = Bwt$; $\phi_0 = 2.0678 \times 10^{-15} \text{ Tm}^2$

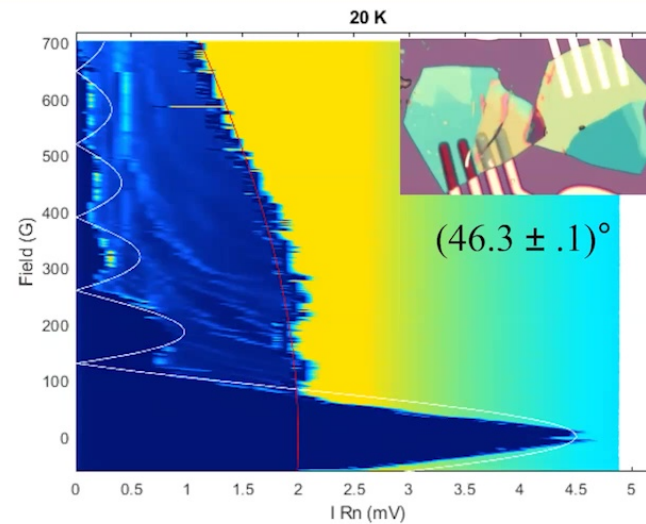
Width: $11 \mu\text{m}$ (measured)

Fitting Parameters: **Thickness: 27 nm**

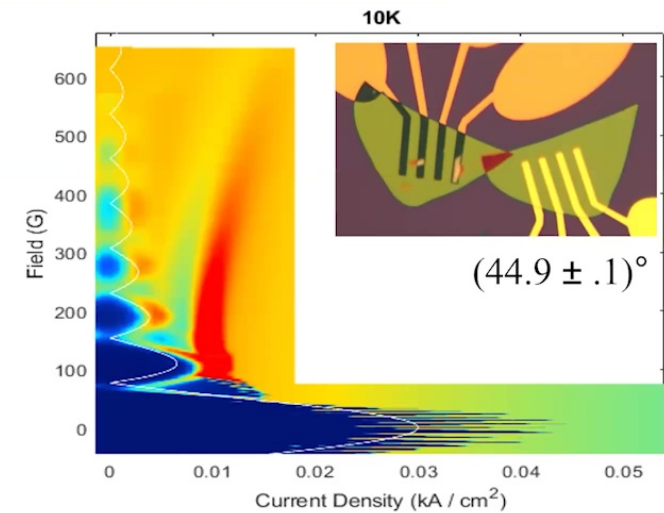
Fraunhofer Patterns in Twisted BSCCO Josephson Junction



Fitting Parameters: Width: $3.5 \mu\text{m}$
Thickness: 1.5 nm

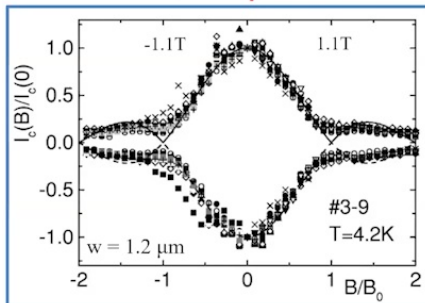


Fitting Parameters: Width: $11 \mu\text{m}$
Thickness: $14.5 / 1.5 \text{ nm}$

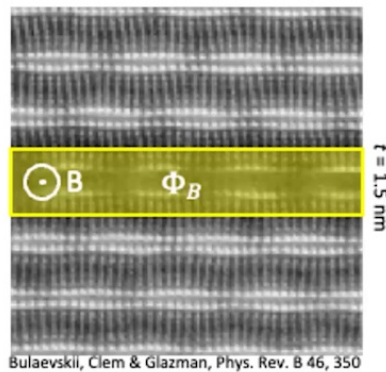


Fitting Parameters: Width: $10 \mu\text{m}$
Thickness: 27 nm

Intrinsic BSCCO Josephson Junction



Irie et al, Phys. Rev. B 62 6681

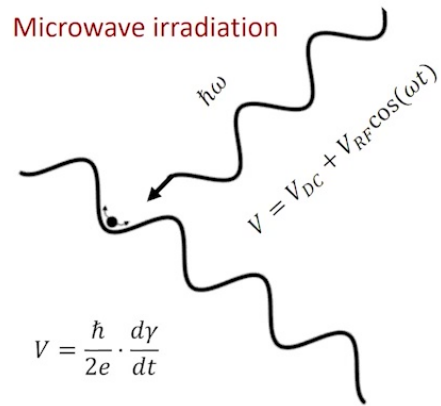


Bulaevskii, Clem & Glazman, Phys. Rev. B 46, 350

Two different Josephson current contributions

$$\mathcal{F}(\varphi) = \mathcal{F}_0 + 2B_0\psi^2[-\cos(2\theta)\cos\varphi + \mathcal{K}\cos(2\varphi)]$$

Shapiro Steps Measurement for JJ Current-Phase Relations



Integer Shapiro steps

$$n = 1, 2, 3, \dots$$

$$V_{DC} = \frac{n \hbar}{2e} \cdot \omega_{RF}$$

$$V = \frac{\hbar}{2e} \cdot \frac{d\gamma}{dt}$$

$$I(\gamma) = I_C \sin(\gamma)$$

$$I(\gamma) = I_C \sum_n (-1)^n J_n \left(\frac{2eV_{RF}}{\hbar\omega} \right) \sin \left(\gamma_0 + \frac{2eV_{DC}}{\hbar} t - n\omega_{RF} t \right)$$

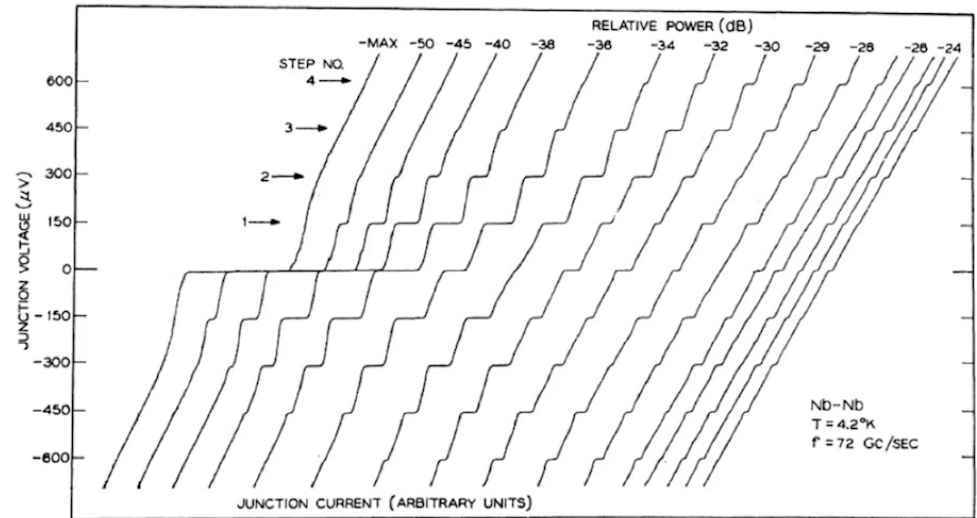
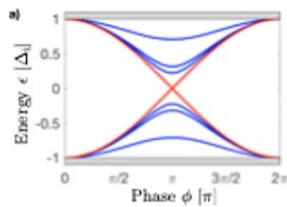
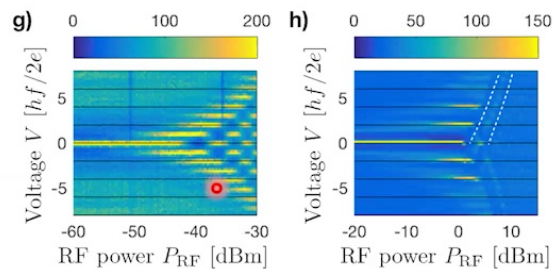


FIG. 1. Voltage-current curves for a Nb-Nb point-contact Josephson junction exposed to a 72-Gc/sec signal at various power levels. Grimes and Shapiro Phys. Rev. (1967)

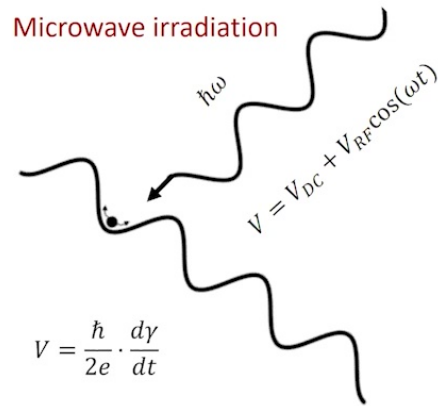
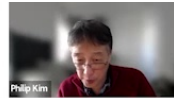
Fractional Josephson Effect: TI + Superconductors



Bocquillon et al., arXiv:1810.03862



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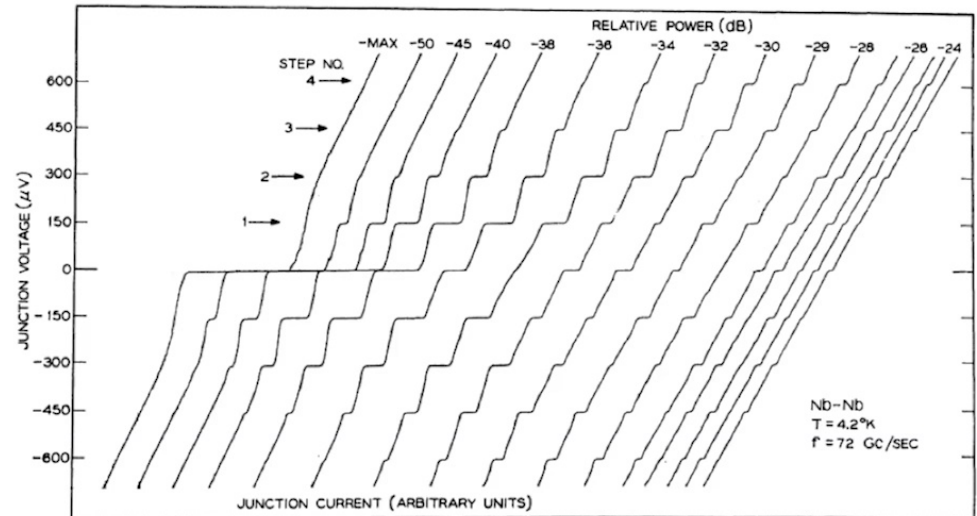
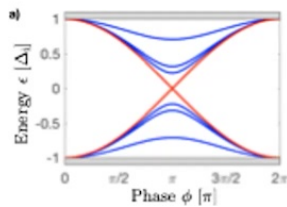
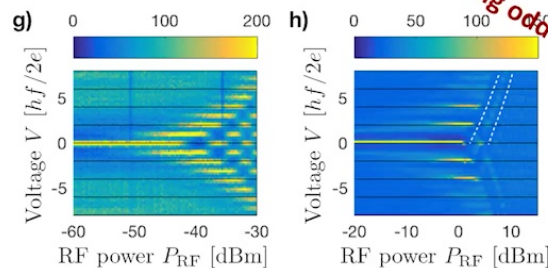


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Fractional Josephson Effect: TI + Superconductors

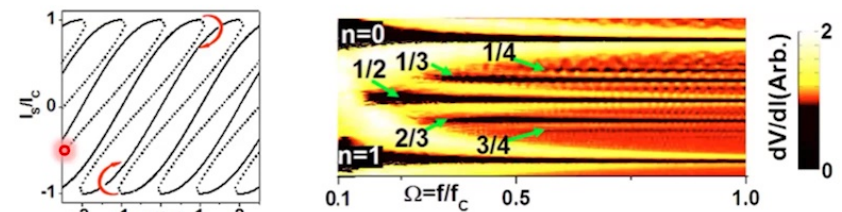


Bocquillon et al., arXiv:1810.03862



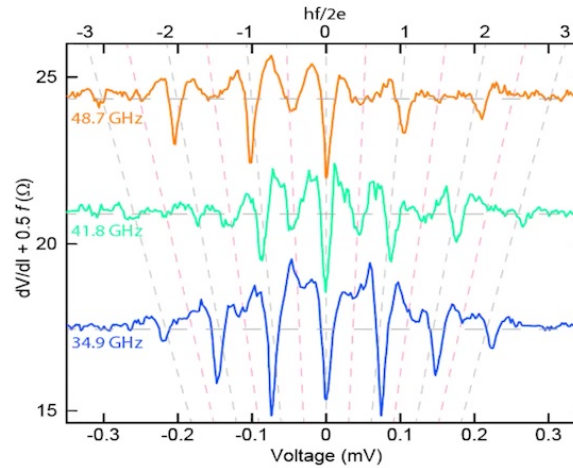
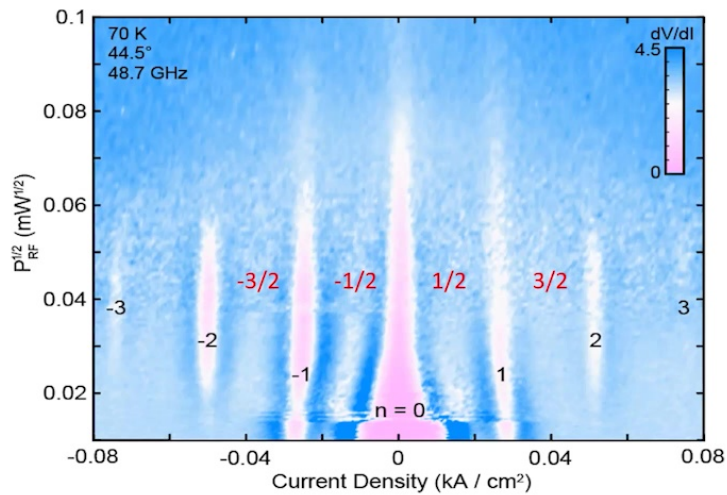
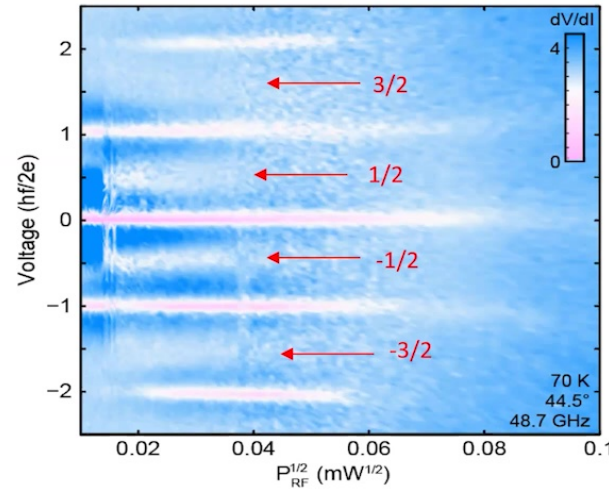
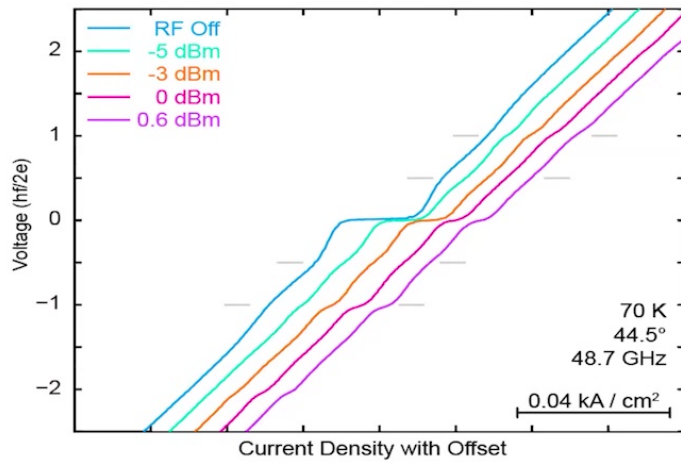
Missing odd integer n

Fractional Shapiro Steps: phase slip CPR in superconducting nanowire



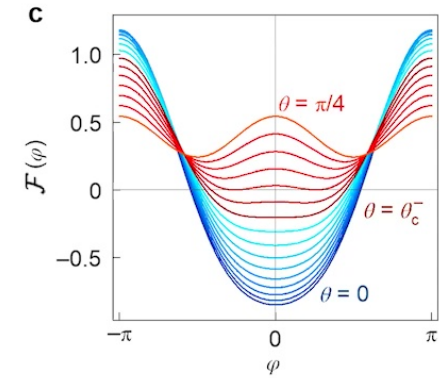
Dinsmore et al., APL 93, 192505 (2008)

Fractional Shapiro Steps: Presence of Second Harmonic Mode



Can et al. Nature Phys (2021)

$$\mathcal{F}(\varphi) = \mathcal{F}_0 + 2B_0\psi^2[-\cos(2\theta)\cos\varphi + \mathcal{K}\cos(2\varphi)]$$



JJ current phase relation

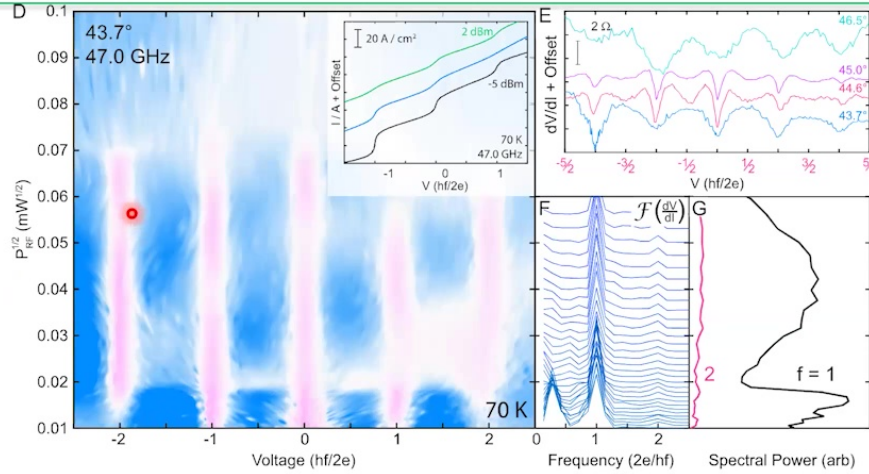
$$J(\varphi) = \frac{2e}{\hbar} \frac{\partial \mathcal{F}(\varphi)}{\partial \varphi}$$

Fractional Shapiro Steps: Presence of Second Harmonic Mode



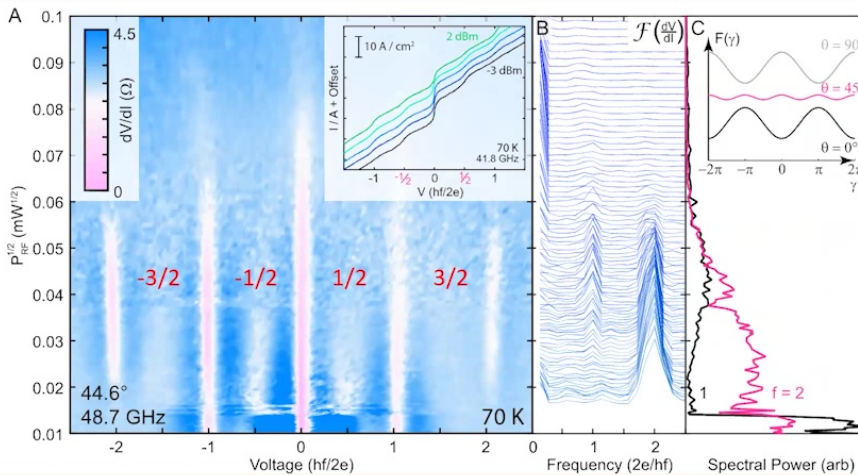
$(43.7 \pm .1)^\circ$

Integer Shapiro steps



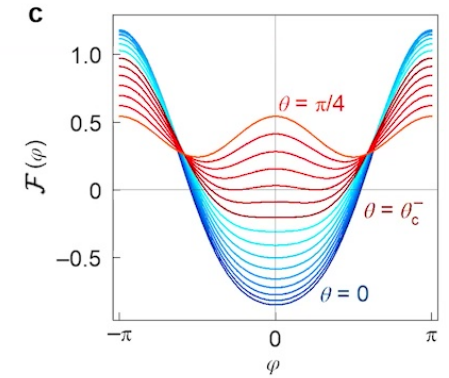
$(44.6 \pm .1)^\circ$

Dominant half Integer Shapiro steps

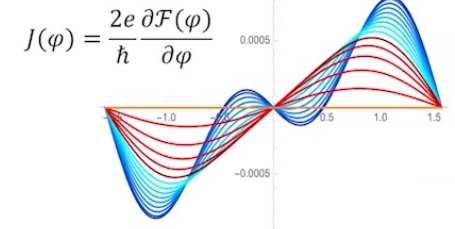


Can et al. Nature Phys (2021)

$$\mathcal{F}(\varphi) = \mathcal{F}_0 + 2B_0\psi^2[-\cos(2\theta)\cos\varphi + \mathcal{K}\cos(2\varphi)]$$



JJ current phase relation

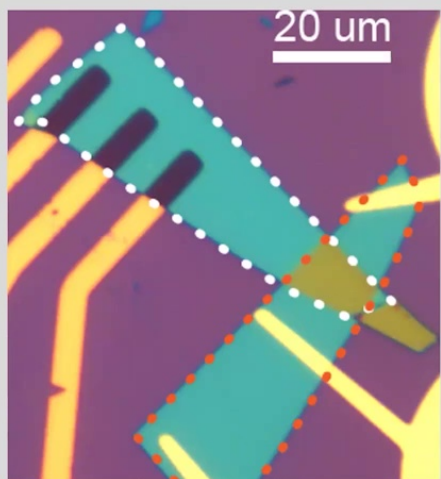
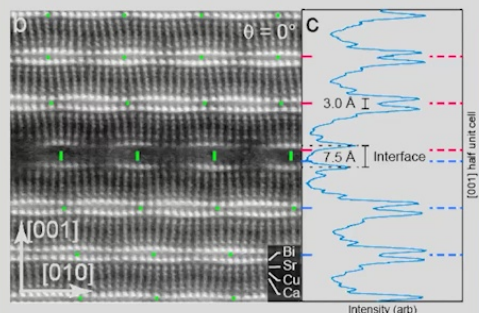


Observed fractional Shapiro steps suggest the presence of the second harmonic term!

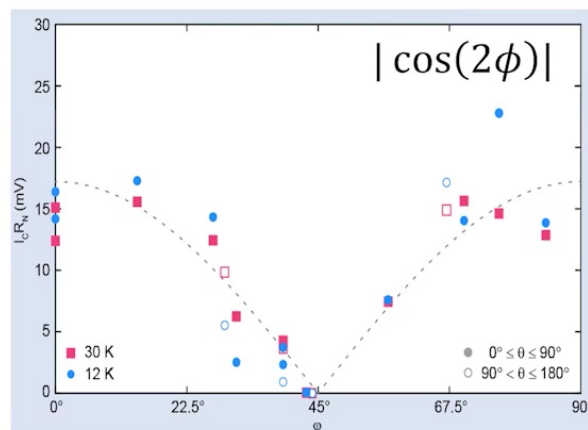
Summary



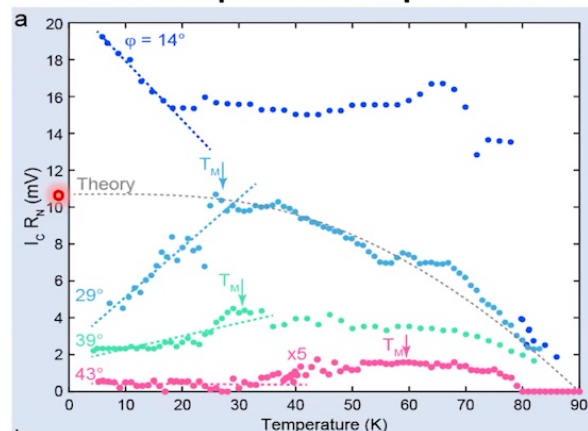
Twist Josephson Junctions Fabrication



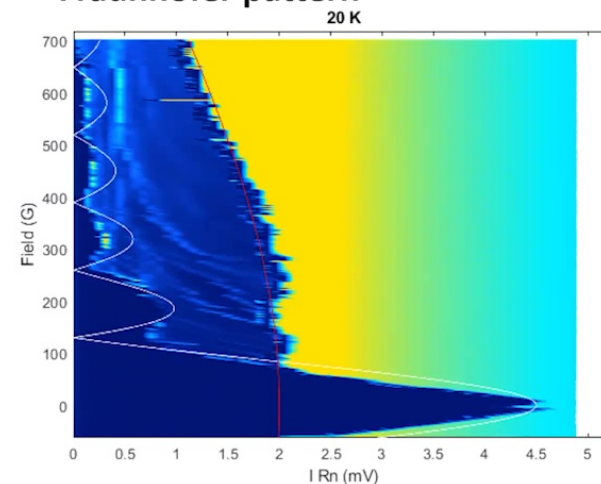
Twist Angle Dependent Critical Current



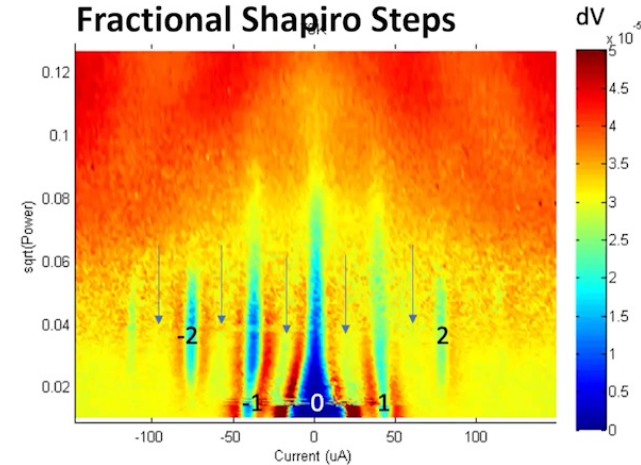
Unusual temperature dependent I_c



Fraunhofer pattern



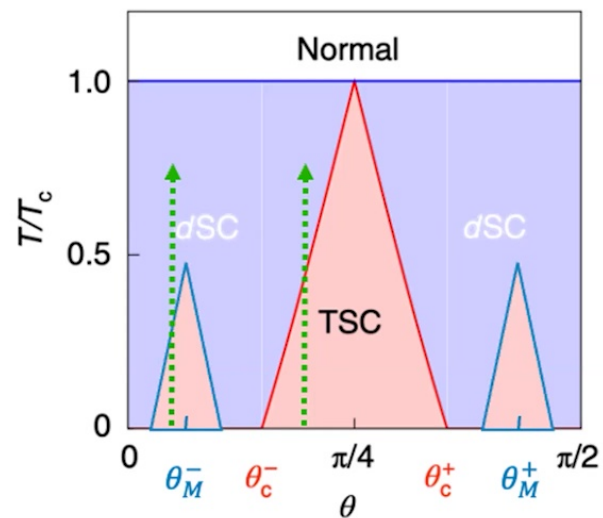
Fractional Shapiro Steps





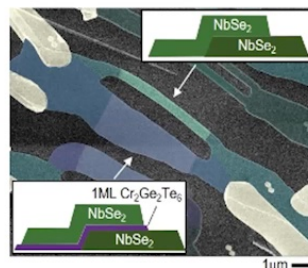
Outlook

Can we directly probe TSC with broken TRS?

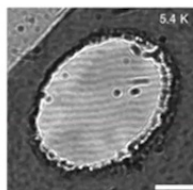


Can et al. Nature Phys (2021); Volkov et al. arXiv: 2012.07860

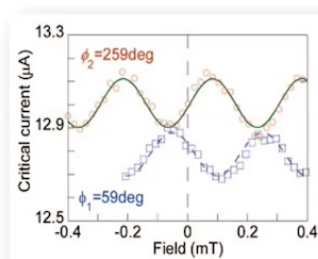
- Phase sensitive measurement: SQUID



NbSe₂: s-wave Ising SC
Cr₂Ge₂Te₆: magnetic insulator
NbSe₂: s-wave Ising SC



Magnetic domain structure of $\text{Cr}_2\text{Ge}_2\text{Te}_6$

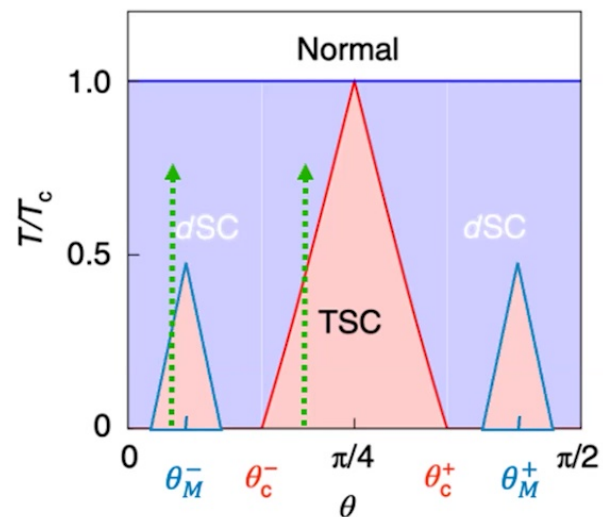


Idzuchi et al., Nature Comm. (2021)



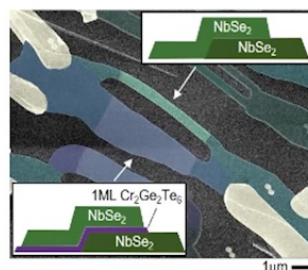
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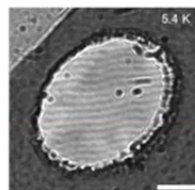


Can et al. Nature Phys (2021); Volkov et al. arXiv: 2012.07860

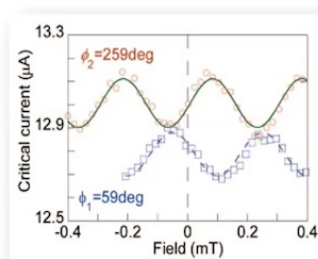
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NbSe₂: s-wave Ising SC
Cr₂Ge₂Te₆: magnetic insulator
NbSe₂: s-wave Ising SC

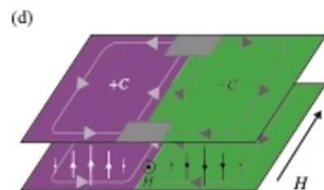


Magnetic domain structure of Cr₂Ge₂Te₆



Idzuchi et al., Nature Comm. (2021)

- Edge state measurement: thermal transport



- Other means of broken TRS measurement: Kerr effect and etc

Acknowledgment



Experiment

Frank Zhao, Nicolar Poccia, Xiaomeng Cui, Hyobin Yoo, Joon Young Park, Yuval Ronen, Rebecca Engelke



Samples

Rudan Zhong, Genda Gu (BNL); Kenji Watanabe, Takashi Taniguchi (NIMS)



Theory

Pavel Volkov, Jed Pixley (Rutgers); Tarun Tummuru, Stephan Plugge, Marcel Franz (UBC)

