

Title: Axion echos from supernovae remnants

Speakers: JiJi Fan

Series: Particle Physics

Date: February 22, 2022 - 1:00 PM

URL: <https://pirsa.org/22020069>

Abstract: Stimulated decays of axion dark matter, triggered by a source in the sky, could produce a photon flux along the continuation of the line of sight, pointing backward to the source. The strength of this so-called axion "echo" signal depends on the entire history of the source and could still be strong from sources that are dim today but had a large flux density in the past, such as supernova remnants (SNRs). This echo signal turns out to be most observable in the radio band. I will present the sensitivity of radio telescopes such as the Square Kilometer Array (SKA) to echo signals generated by SNRs that have already been observed. In addition, I will show projections of the detection reach for signals from newly born supernovae that could be detected in the future. Intriguingly, an observable echo signal could come from old "ghost" SNRs which were very bright in the past but are now so dim that they haven't been observed.

Zoom Link: <https://pitp.zoom.us/j/91076203387?pwd=UzNva3N4Zi9mV3BkMlJvUnhtRXRZdz09>

Axion Echos from Supernovae Remnants

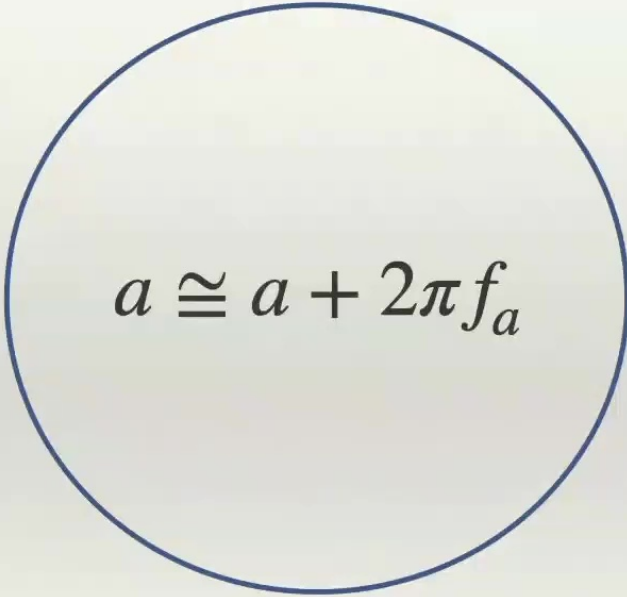
JiJi Fan

Brown University

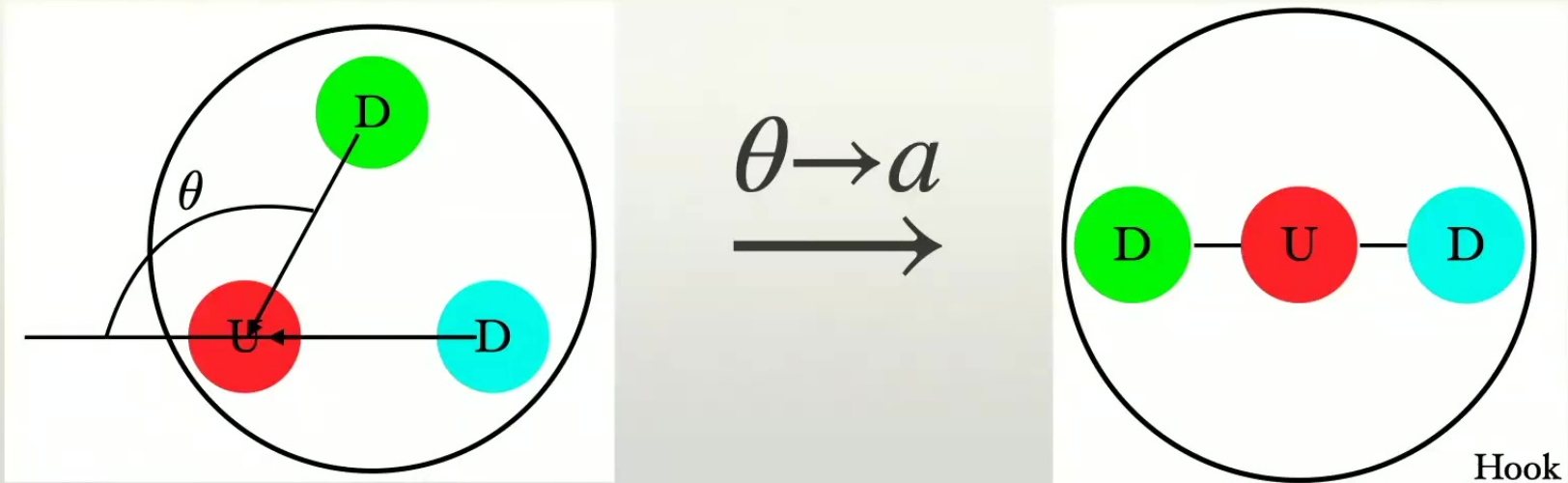
Perimeter Institute, 02/22/2022

What is an axion ?

a periodic pseudo-scalar field


$$a \cong a + 2\pi f_a$$

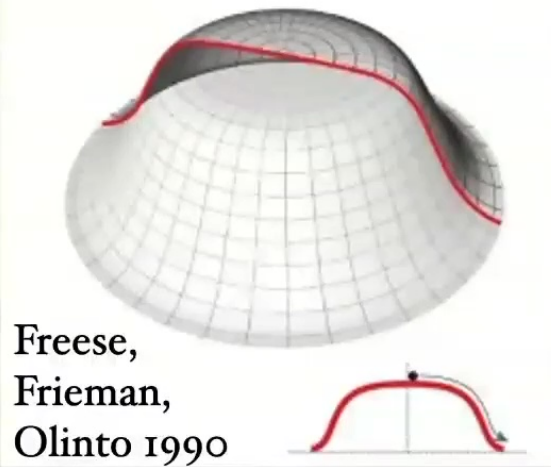
Strong CP problem: QCD axion



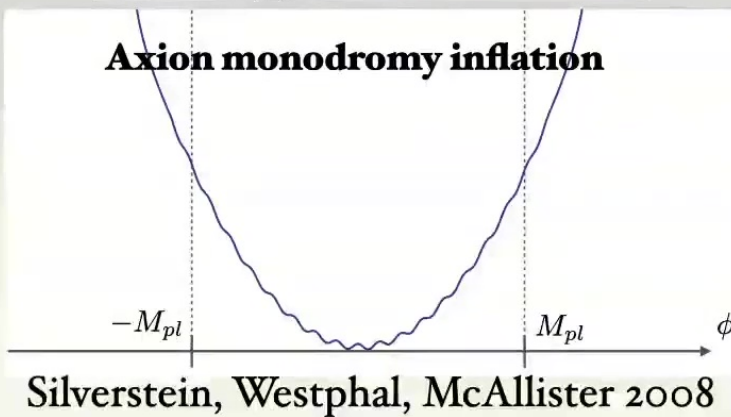
Peccei, Quinn; Weinberg; Wilczek; Kim; Shifman, Vainshtein, Zakharov; Zhitnitsky;
Dine, Fischler, Srednicki 1977 - 1981

Other applications

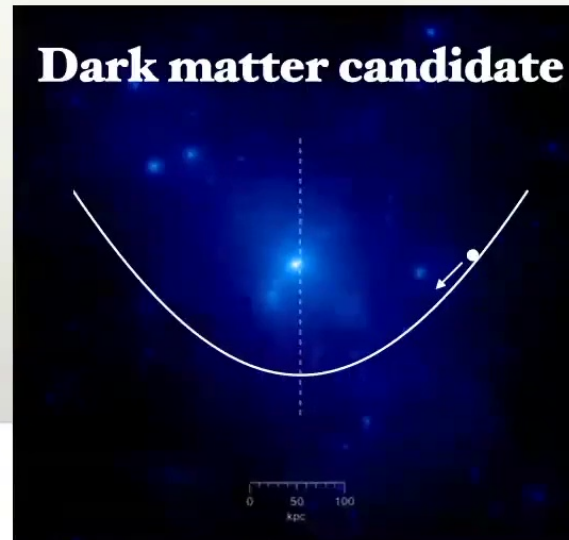
Natural inflation



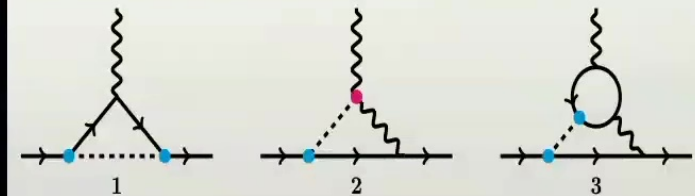
Axion monodromy inflation



Dark matter candidate



Muon $g-2$



Marciano, Masiero, Paradisi, Passera 2016; Bauer, Neubert, Thamm 2017; Buen-Abad, Fan, Reece, Sun 2021

**Terrestrial Experiments: DM radio,
Casper, ABRACADABRA, ADMX...**

**Model Building:
go beyond vanilla
models**

Axion (ALP and QCD axion)

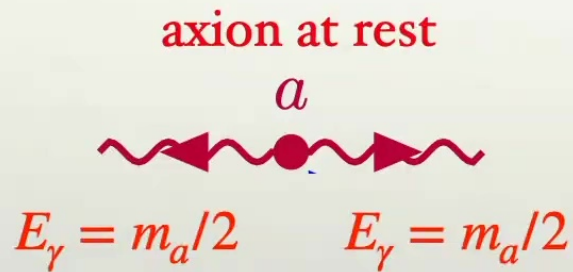
**Astrophysical/Cosmic Probes:
CMB, gravitational waves,
stars, galaxies....**

A new astrophysical probe of couplings of axion DM

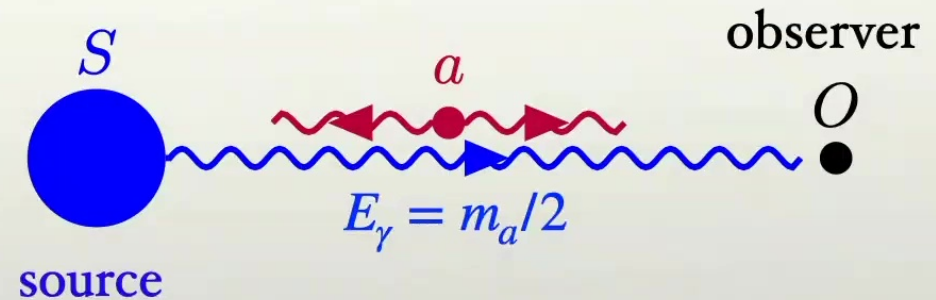
Buen-Abad, Fan, Sun; Sun, Schutz, Nambrath, Leung, Masui 2021

Axion-photon coupling: $-\frac{g_{a\gamma\gamma}}{4}aF_{\mu\nu}\tilde{F}^{\mu\nu}$

Spontaneous decays



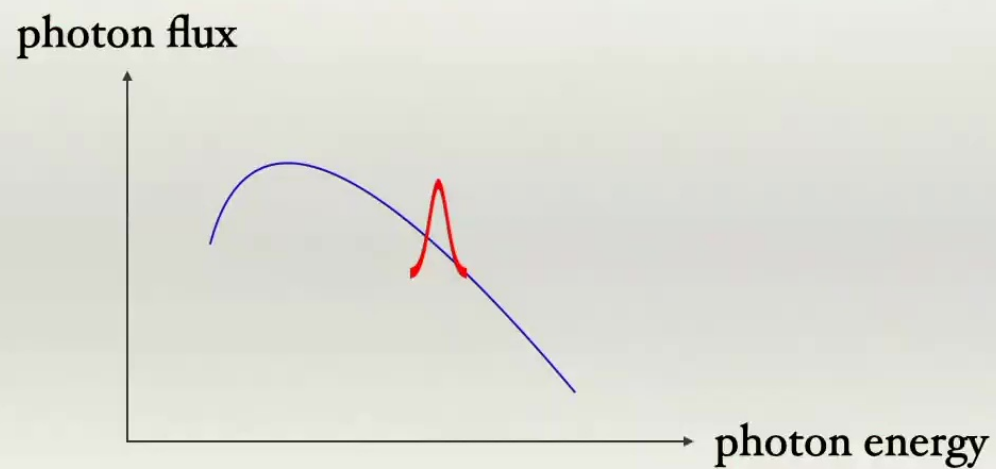
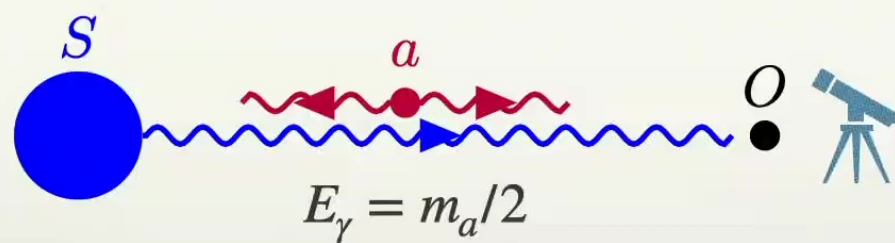
Stimulated decays



Analogous to laser, decay rate enhanced by the large occupation number of the external photon flux.

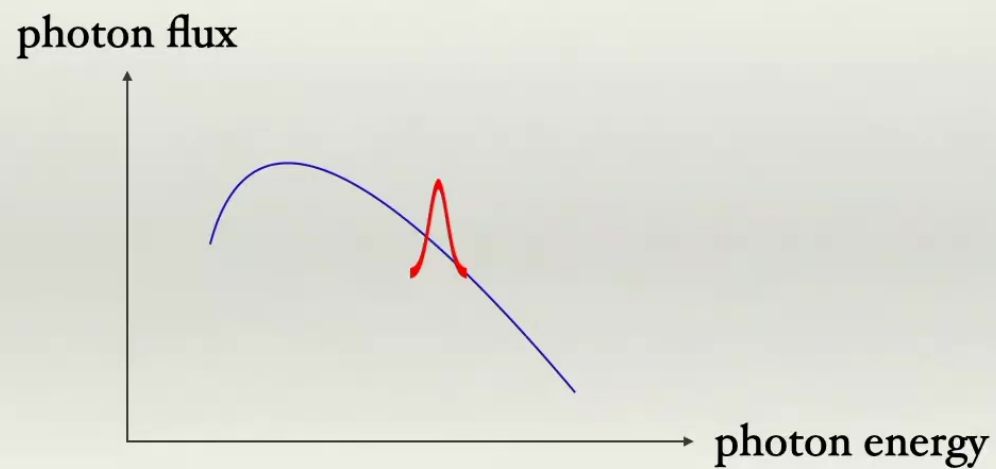
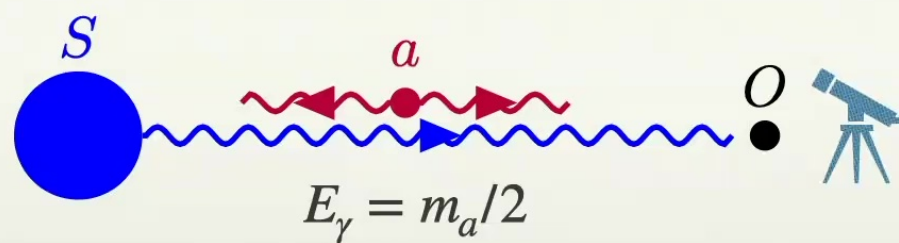
All wavy lines: photons; **red wavy lines: photons from axion decays**; **blue wavy lines: photons emitted by the source.**

Stimulated axion DM decays



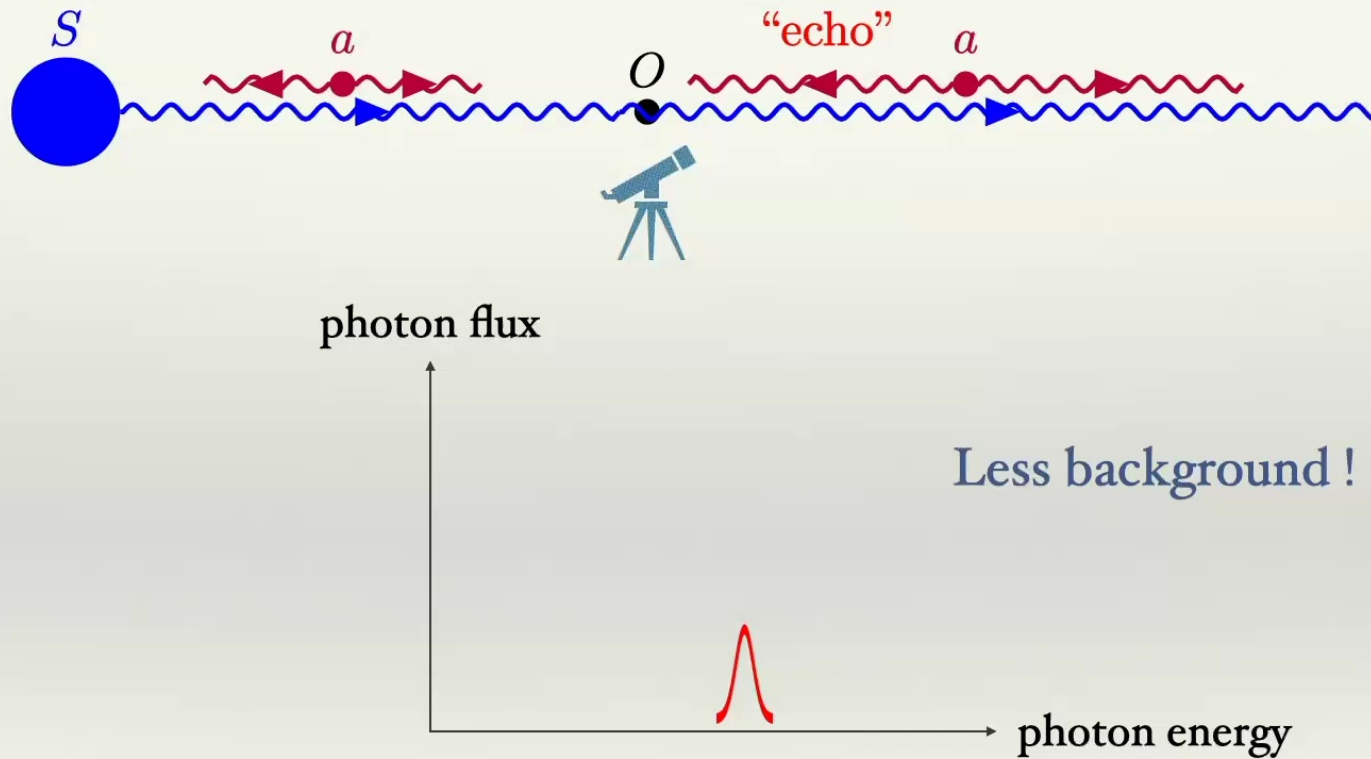
Caputo, Regis, Taoso and Witte 2018

Stimulated axion DM decays



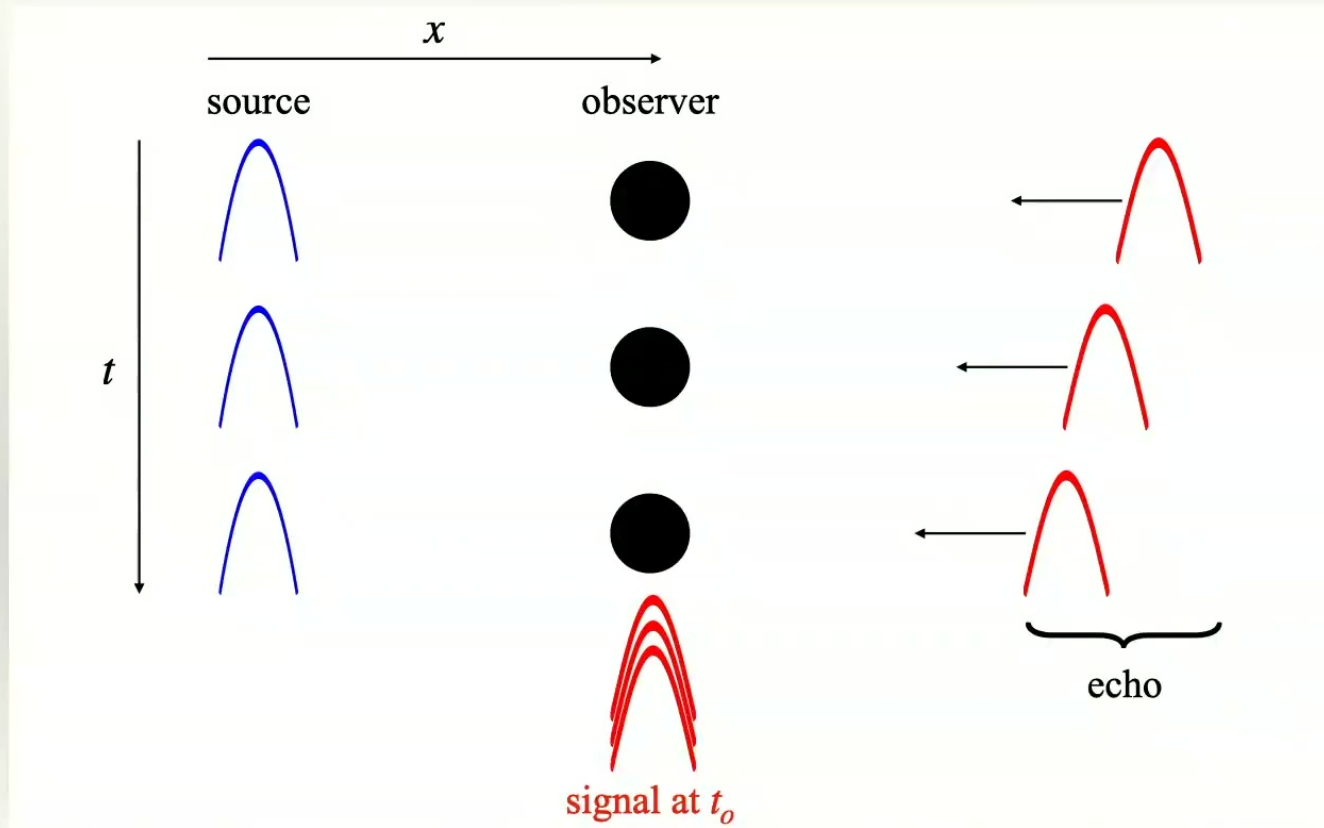
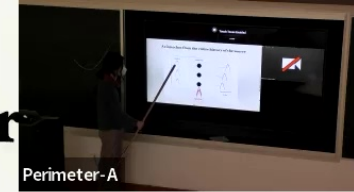
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Axion echos

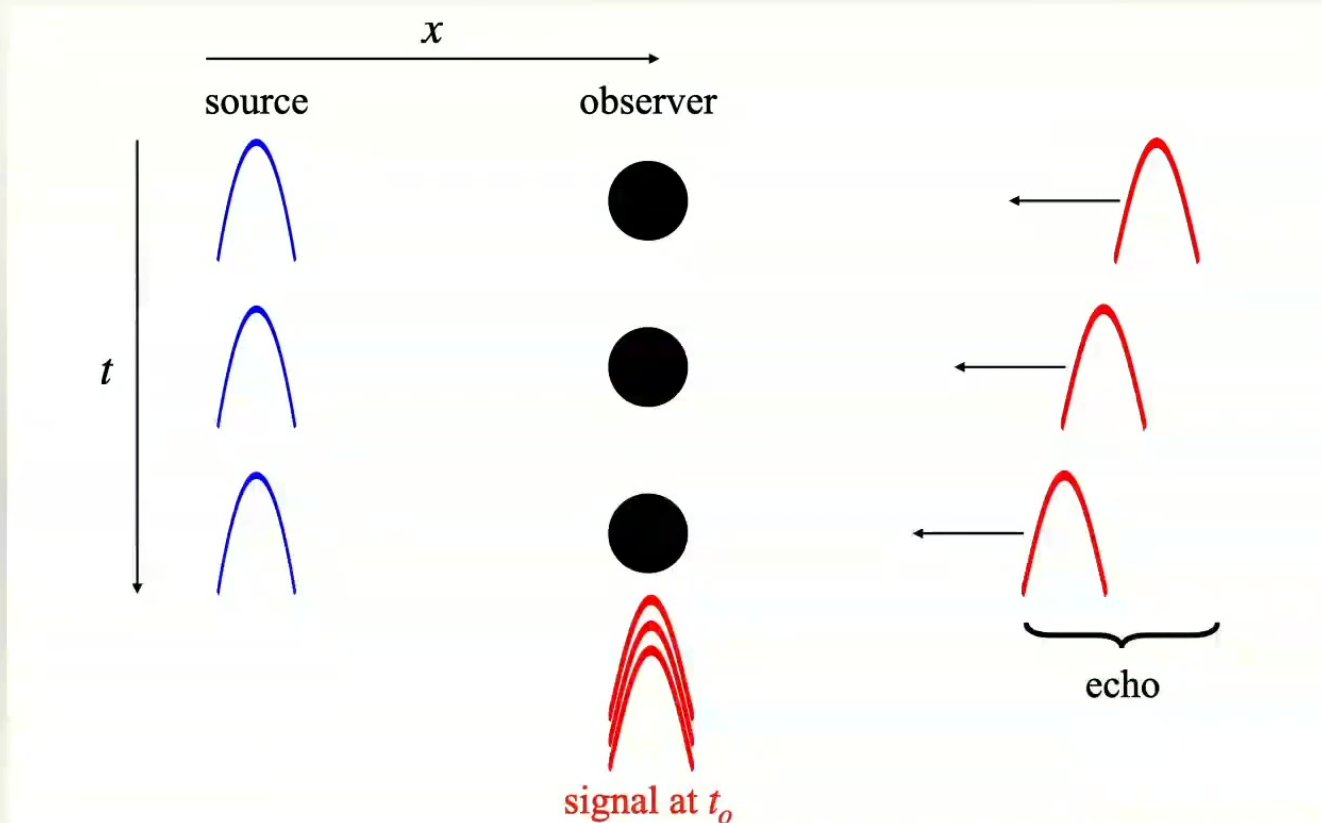
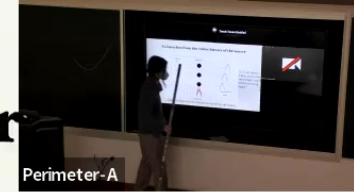


Arza, Sikivie 2019 (S : very powerful microwave beam); Ghosh, Salvado and Miralda-Escude 2020 (S : Cygnus A radio galaxy, one of the most powerful constant radio sources in the sky)

Axion echos from the entire history of the source



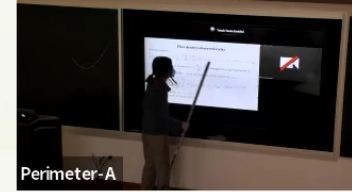
Axion echos from the entire history of the source



Dim old source
which used to be
very bright could
lead to a strong
signal !

Signal is integrated over the entire history of the source

Flux density observed today



Boltzmann equation:

$$\frac{d}{dt} f_\gamma = C[f_\gamma] \quad C[f_1] = \frac{1}{2E_1} \int \frac{DP_a}{2E_a} \int \frac{DP_2}{2E_2} |\overline{\mathcal{M}}|^2 \left(f_a(1 + f_1 + f_2) - f_1 f_2 \right) (2\pi)^4 \delta^4(p_a - p_1 - p_2)$$

\Rightarrow Echo flux density at frequency $\nu_a = \frac{E_a}{2\pi} = \frac{m_a}{4\pi}$, averaged over the bandwidth Δ_ν

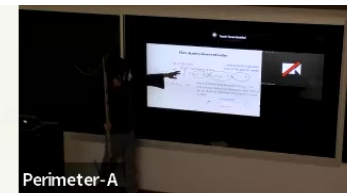
fraction of signal within
bandwidth

axion spontaneous decay width ($m_a, g_{a\gamma\gamma}$)

$$\overline{S}_{\nu_a,e} = \frac{f\Delta}{2E_a^3 \Delta_\nu} \int_0^{t_{\text{age}}/2} dx \rho_a(x, -\hat{n}_*) S_{\nu_a,s}(t_{\text{age}} - 2x) e^{-\tau(\nu,x,-\hat{n}_*)}$$

bandwidth $\Delta_\nu \sim \nu_a \sigma_\nu$ (σ_ν :
DM velocity dispersion)

Flux density observed today



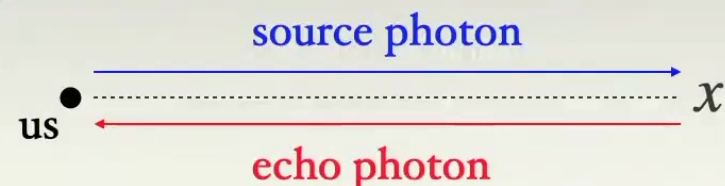
age of the source

τ : optical depth (important close to the galactic center)

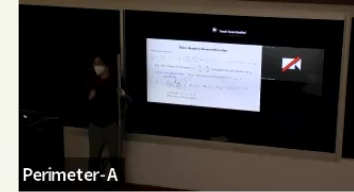
$$\bar{S}_{\nu_{a,e}} = f_{\Delta} \frac{\pi \Gamma_a}{2 E_a^3 \Delta \nu} \int_0^{t_{\text{age}}/2} dx \rho_a(x, -\hat{n}_*) S_{\nu_{a,s}}(t_{\text{age}} - 2x) e^{-\tau(\nu, x, -\hat{n}_*)}$$

$$t_{\text{age}} \sim (10^4 - 10^5) \text{ years} \Rightarrow x_{\text{max}} \sim 10 \text{ kpc}$$

source flux density at an earlier time $t_{\text{age}} - 2x$:
echo photon emitted at a distance x away from us
was produced by light that first passed us a time
 $2x$ away.



Flux density observed today



Boltzmann equation:

$$\frac{d}{dt} f_\gamma = C[f_\gamma] \quad C[f_1] = \frac{1}{2E_1} \int \frac{DP_a}{2E_a} \int \frac{DP_2}{2E_2} |\overline{\mathcal{M}}|^2 \left(f_a(1 + f_1 + f_2) - f_1 f_2 \right) (2\pi)^4 \delta^4(p_a - p_1 - p_2)$$

⇒ Echo flux density at frequency $\nu_a = \frac{E_a}{2\pi} = \frac{m_a}{4\pi}$, averaged over the bandwidth Δ_ν

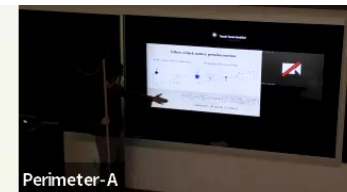
fraction of signal within
bandwidth

axion spontaneous decay width ($m_a, g_{a\gamma\gamma}$)

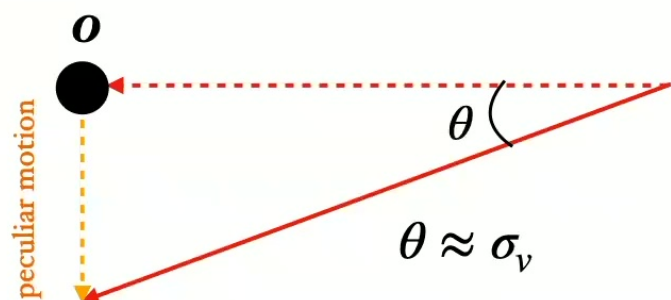
$$\overline{S}_{\nu_a,e} = \frac{f\Delta}{2E_a^3 \Delta_\nu} \pi \Gamma_a \int_0^{t_{\text{age}}/2} dx \rho_a(x, -\hat{n}_*) S_{\nu_a,s}(t_{\text{age}} - 2x) e^{-\tau(\nu,x,-\hat{n}_*)}$$

bandwidth $\Delta_\nu \sim \nu_a \sigma_\nu$ (σ_ν :
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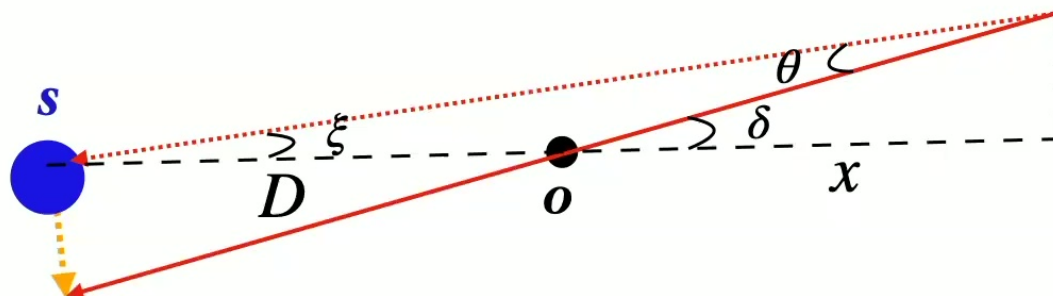
Effect of dark matter peculiar motion



a) Effect of dark matter peculiar motion



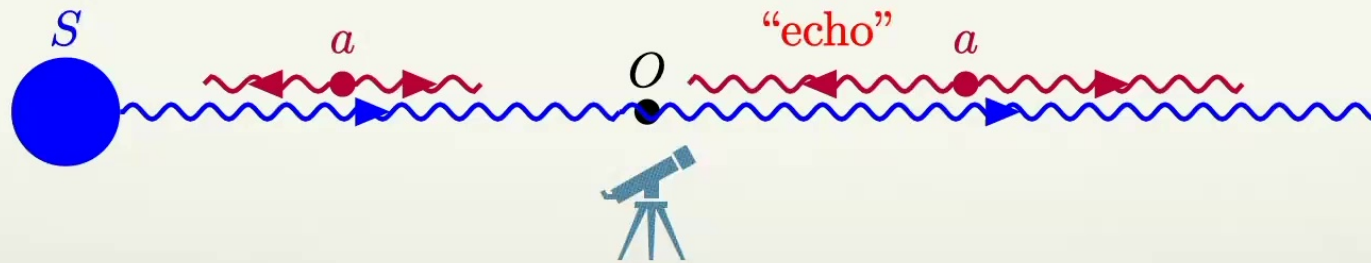
b) Enlarging collecting solid angle



$$2\delta \approx 2\sigma_v \frac{x}{D} \approx 2\sigma_v \frac{t_{\text{age}}/2}{D} \approx 10 \text{ arcmin} \left(\frac{\sigma_v}{10^{-3}} \right) \left(\frac{t_{\text{age}}}{10^4 \text{ years}} \right) \left(\frac{1 \text{ kpc}}{D} \right)$$

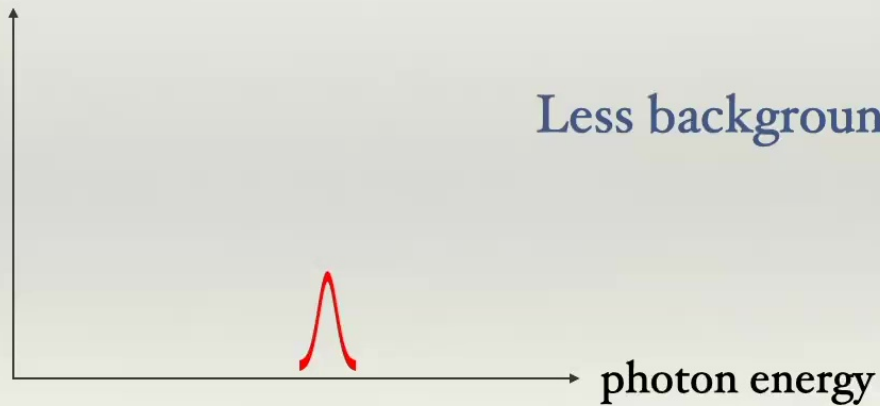
source size: ~ 85 arcmin (1.7 degree)

Axion echos



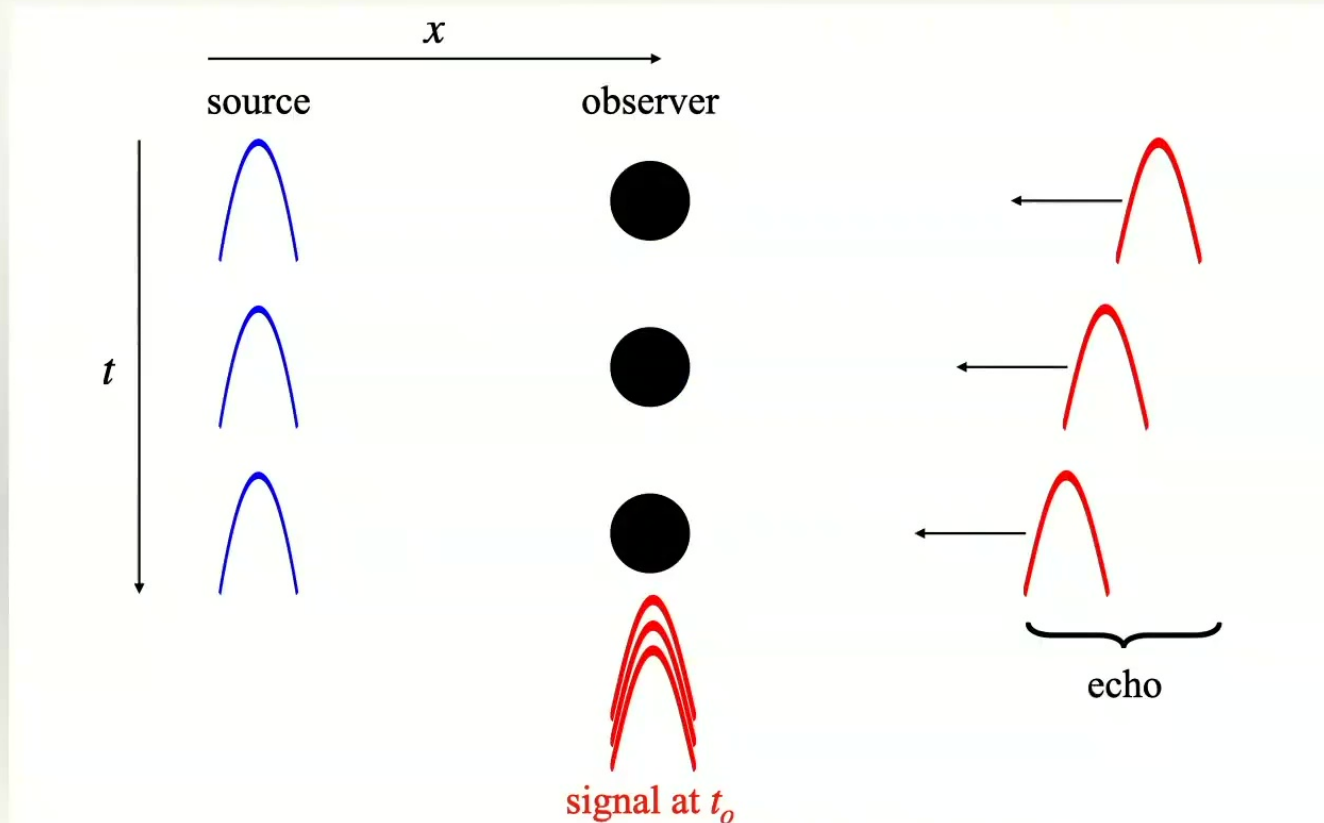
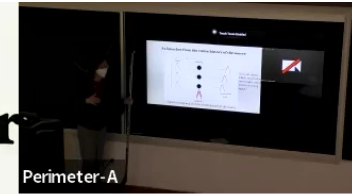
photon flux

Less background !



Arza, Sikivie 2019 (S: very powerful microwave beam); Ghosh, Salvado and Miralda-Escude 2020 (S: Cygnus A radio galaxy, one of the most powerful constant radio sources in the sky)

Axion echos from the entire history of the source



Dim old source
which used to be
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lead to a strong
signal !

Signal is integrated over the entire history of the source

Supernova: powerful stellar explosion

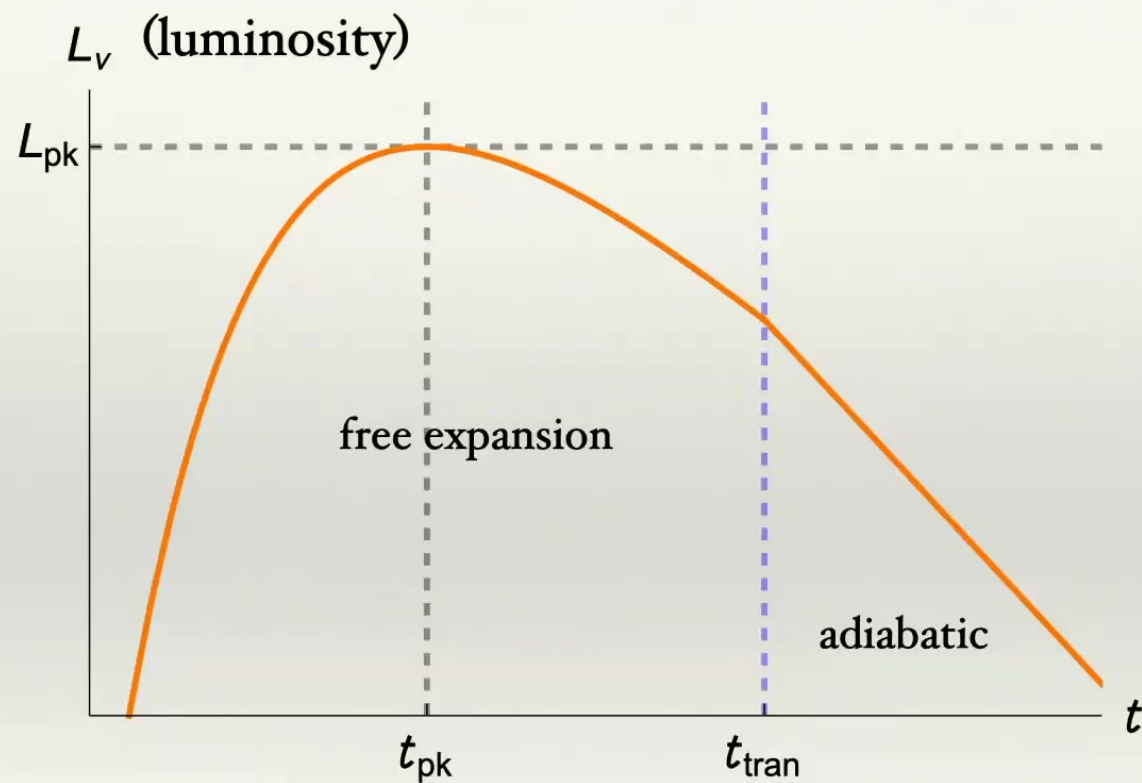
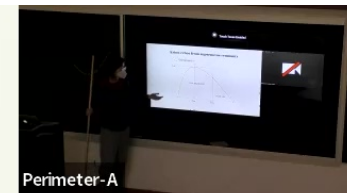


SN2016aps

M. Weiss



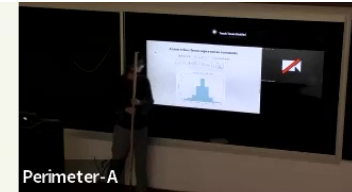
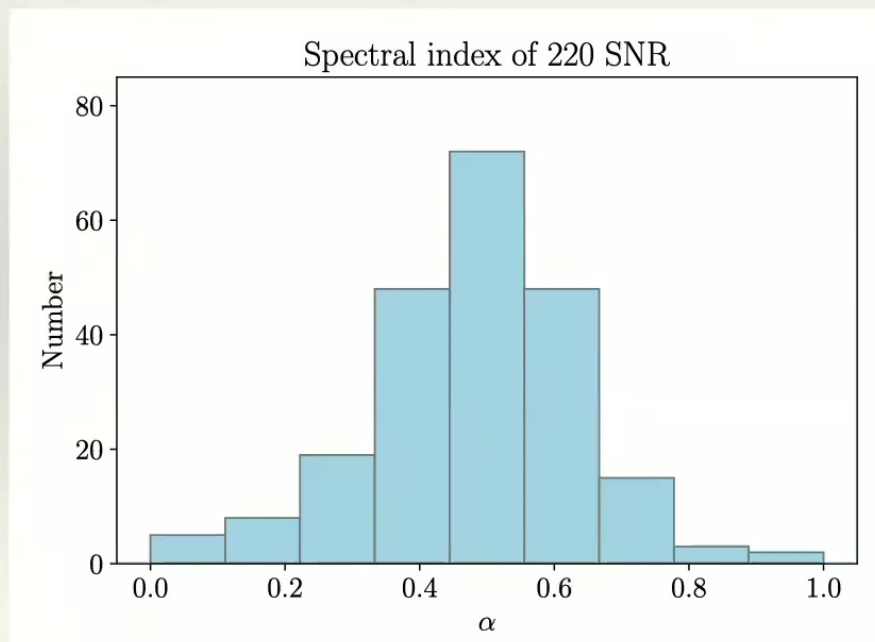
Axion echos from supernovae remnants



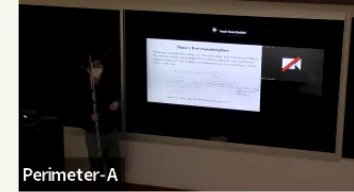
Axion echos from supernovae remnants

Spectrum: $S_\nu \propto E_\gamma^{-\alpha}$ α : spectral index

$$L_\nu = S_\nu (4\pi D^2) = L_{1\text{GHz}} \left(\frac{\nu}{1 \text{ GHz}} \right)^{-\alpha}$$



Phase I: Free expansion phase



Shock wave expands with a radius $r \propto t$, the time elapsed after the explosion. Before the power law falling, optical depth due to free-free absorption and synchrotron radiation drops over time, leading to an exponential rise in the luminosity. Can last $\mathcal{O}(\text{few} \times 100)$ years.

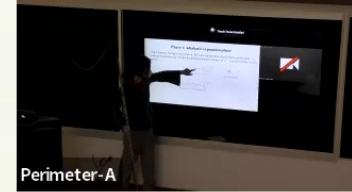
$$L_{\nu, \text{free}}(t) \equiv L_{\nu, \text{pk}} e^{\frac{3}{2}(1-t_{\text{pk}}/t)} \left(\frac{t}{t_{\text{pk}}} \right)^{-1.5}$$

Parameter	mean (μ)	standard deviation (σ)
$\log_{10} \left(\frac{L_{\nu, \text{pk}}}{\text{erg s}^{-1} \text{Hz}^{-1}} \right)$	25.5	1.5
$\log_{10} \left(\frac{t_{\text{pk}}}{\text{days}} \right)$	1.7	0.9

from $\sim \mathcal{O}(100)$ young supernovae in the radio band

Bietenholz, Bartel, Argo, Dua, Ryder and Soderberg 2011.11737

Phase 2: Adiabatic expansion phase



Approximate energy conservation. Pressure behind the shock waves push away surrounding material. Hydrodynamical estimates show $r \propto t^{2/5}$. Can last $\mathcal{O}(\text{few} \times 10^4)$ years.

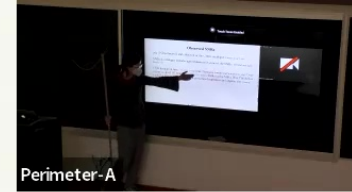
$$L_{\nu,\text{ad}}(t) \equiv L_{\nu,\text{tran}} \left(\frac{t}{t_{\text{tran}}} \right)^{-\gamma} \quad \gamma = \frac{4}{5}(2\alpha + 1) > 0$$
$$L_{\nu,\text{tran}} \equiv L_{\nu,\text{free}}(t_{\text{tran}}) ,$$

spectral index

Or

$$L_{\nu,\text{ad}}(t) \equiv L_{\nu,0} \left(\frac{t}{t_{\text{age}}} \right)^{-\gamma} ,$$
$$L_{\nu,\text{ad}}(t_{\text{tran}}) = L_{\nu,\text{free}}(t_{\text{tran}}) .$$

Observed SNRs

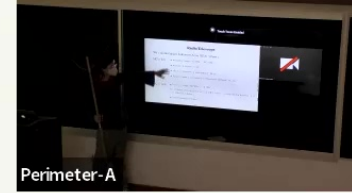


294 SNRs observed and collected in the Green catalogue. Green 2014-2019

SNRcat catalogue includes age estimates for some of the SNRs. Ferrand and Safi-Harb 2012

SNR formation rate: $\sim 0.02 - 0.03$ / year (Tammann, Loeffler and Schroeder 1994). Over a time scale of 10^5 years, about 2000 - 3000 SNRs in the Milky Way. The deficit could be due to selection criteria on surface brightness and angular size (Green, 2005 - 2019).

Radio Telescope



We consider Square Kilometer Array (SKA) - phase 1.

- SKA1-low:
- frequency range: 50 MHz – 350 MHz;
 - number of stations: 512;
 - effective diameter of each station: 38 m;
 - longest baseline (separation between two stations): 80 km,

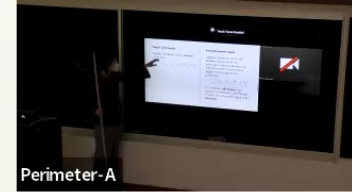
- SKA1-mid:
- frequency range: 350 MHz – 15.4 GHz;
 - 133 dishes with a diameter of 15 m (SKA dishes), 64 dishes with a diameter of 13.5 m (MeerKAT dishes);
 - longest baseline: 150 km.

Single-dish mode

Angular resolution set by *diameter* of the dish:

$$\theta_{\text{res}} = 1.22 \left(\frac{\lambda}{d} \right) \text{ rad} \approx 1.4^\circ \left(\frac{\text{GHz}}{\nu} \right) \left(\frac{15 \text{ m}}{d} \right)$$

Interferometer mode



Angular resolution set by the *baseline* (separation of two telescopes): finer resolution; but will fail to detect an extended signal with size above the resolution.

$$\theta_b = \left(\frac{\lambda}{B} \right) \text{ rad} = 0.17^\circ \left(\frac{\text{GHz}}{\nu} \right) \left(\frac{100 \text{ m}}{B} \right)$$

As frequency ***increases***, the number of baselines contributing to detection of a given signal could ***decrease!***

Single-dish mode

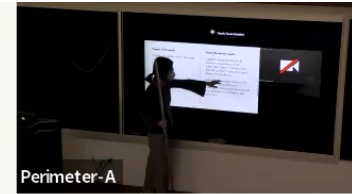
Angular resolution set by diameter of the dish;

Signal to noise ratio enhanced by \sqrt{N} if there are N single dishes.

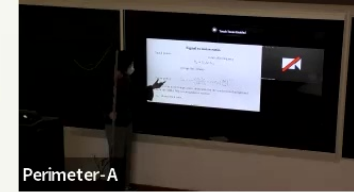
Interferometer mode

Angular resolution set by the baseline (separation of two telescopes): finer resolution; but will fail to detect extended signal above the resolution.

Signal to noise ratio enhanced by $\sqrt{N(N-1)/2}$ if there are N telescopes in the array, provided all the baselines contribute to the detection.



Signal to noise ratio



Signal power:

$$P_{\text{sig}} \propto \bar{S}_{\nu,e} \Delta\nu A_{\text{tot}}$$

average flux density total collecting area

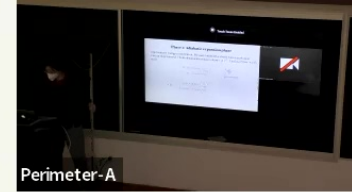
Noise power:

$$P_{\text{noi, unit}} = \frac{2 T_{\text{sys}} \Delta\nu}{\sqrt{2 \Delta\nu t_{\text{obs}}}} = \sqrt{2} T_{\text{sys}} \left(\frac{\Delta\nu}{t_{\text{obs}}} \right)^{1/2}.$$

T_{sys} : system noise temperature, dominated by the synchrotron background from the Milky Way or extragalactic sources

t_{obs} : observation time.

Radio Telescope

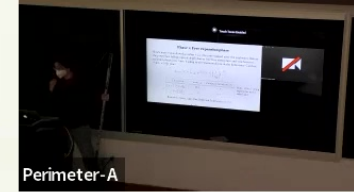


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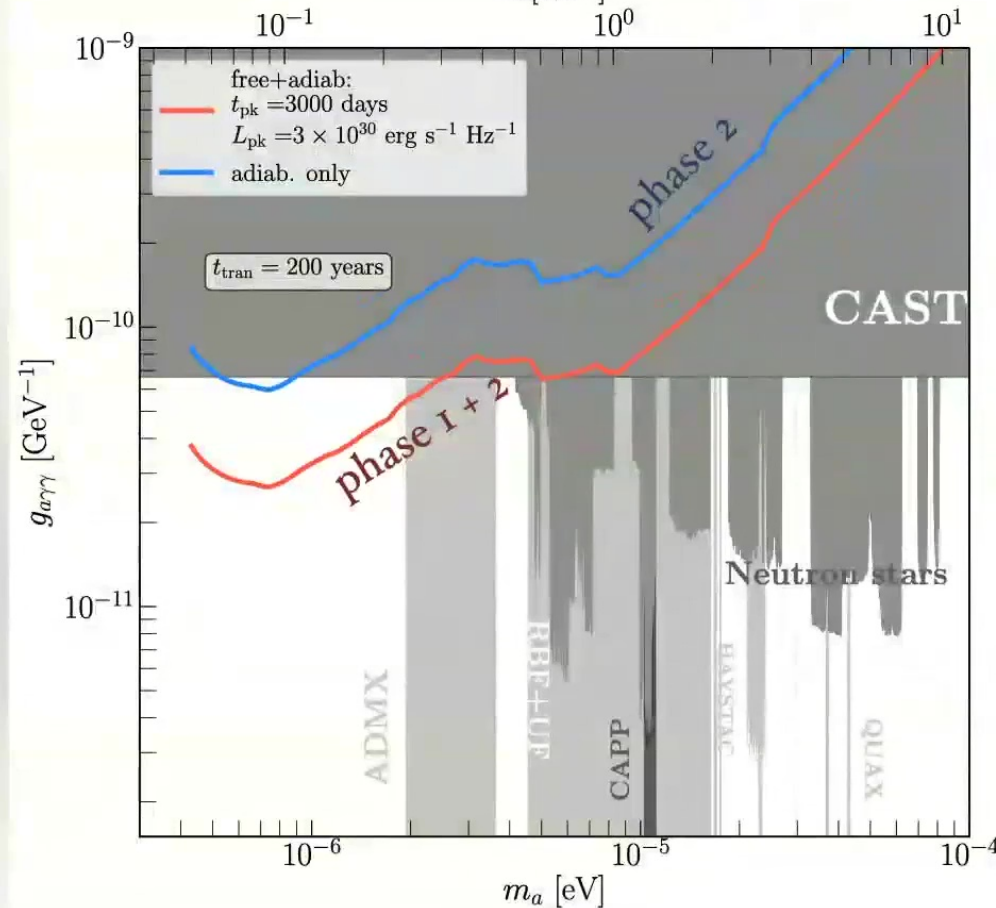
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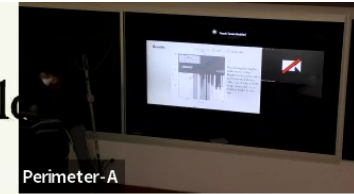
Bietenholz, Bartel, Argo, Dua, Ryder and Soderberg 2011.11737

Results

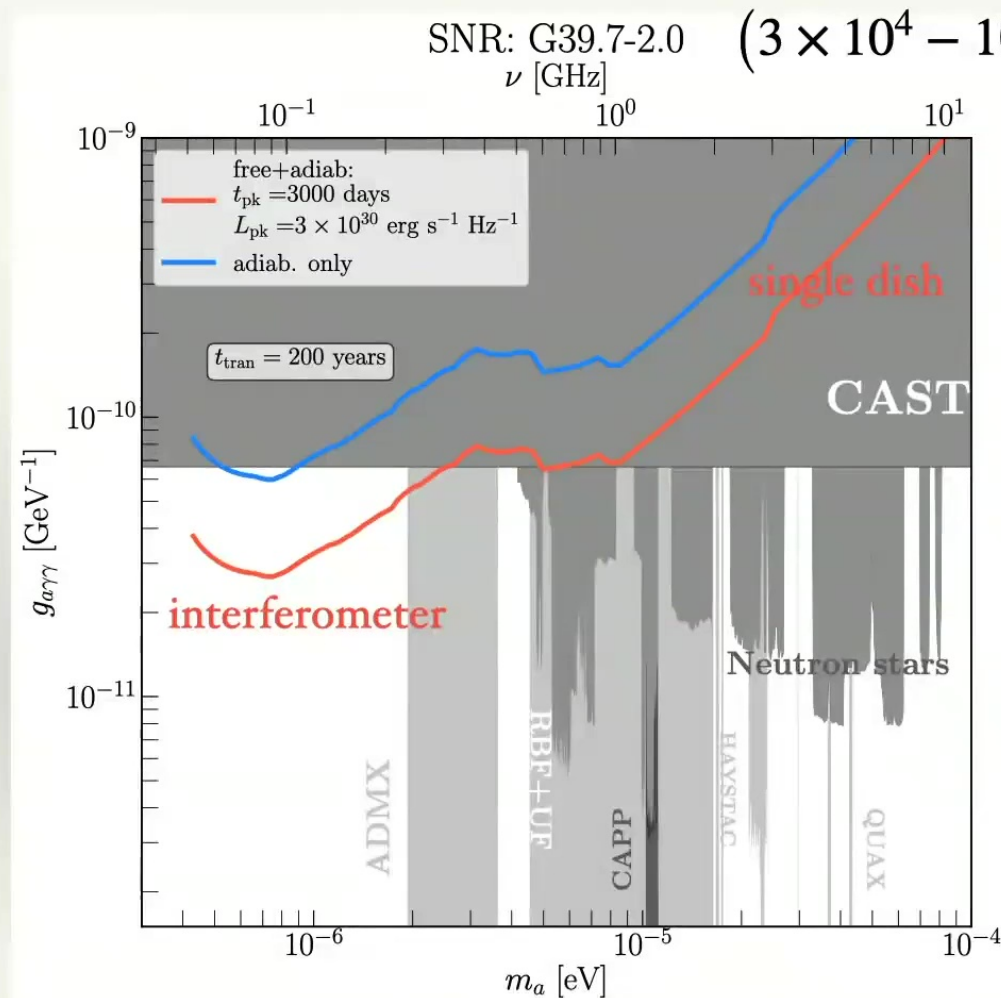
SNR: G39.7-2.0 $(3 \times 10^4 - 10^5)$ years old



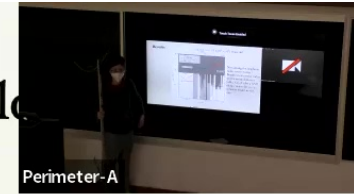
Not among the brightest radio source today !
 Brightest extrasolar radio source around/above 1 GHz, Cas A (also a SNR of 340 years old), is ~ 30 times as bright as this one.



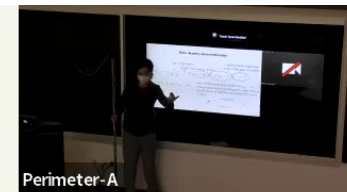
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Flux density observed today



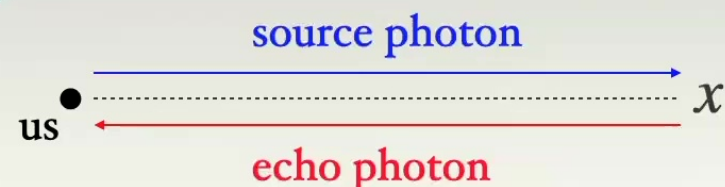
age of the source

τ : optical depth (important close to the galactic center)

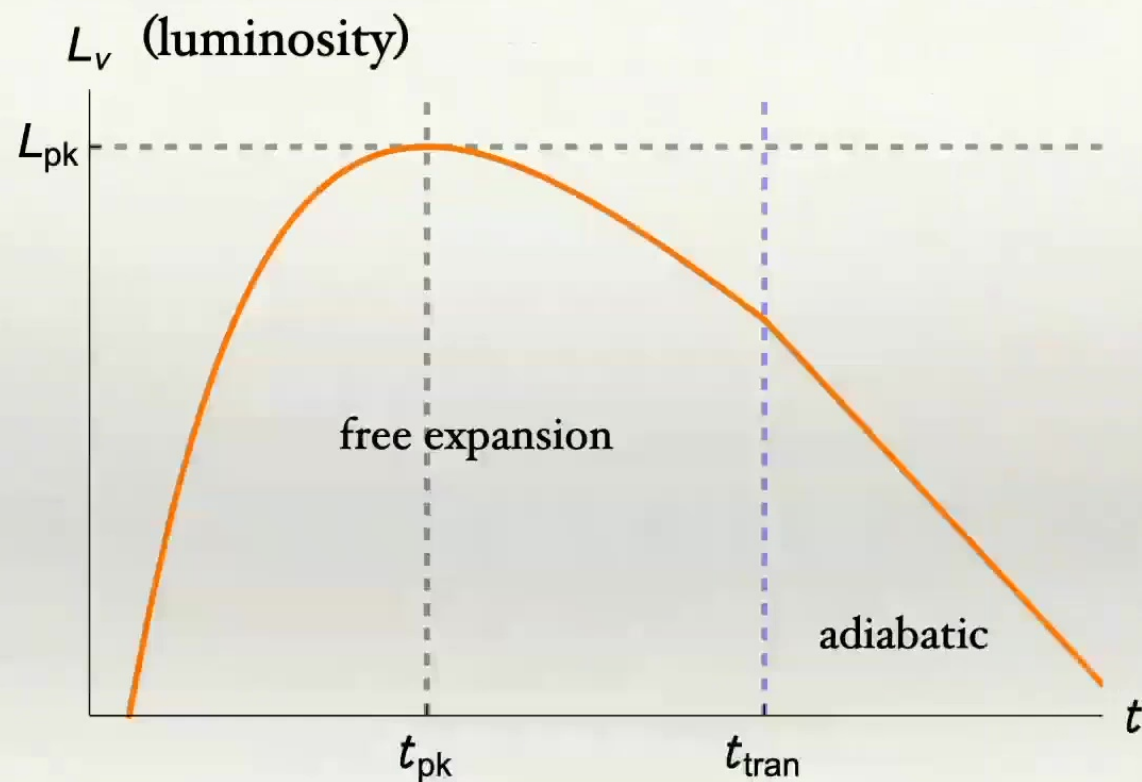
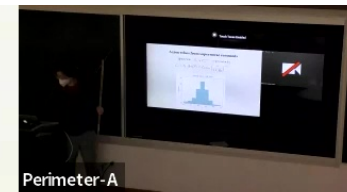
$$\bar{S}_{\nu_{a,e}} = f_{\Delta} \frac{\pi \Gamma_a}{2 E_a^3 \Delta \nu} \int_0^{t_{\text{age}}/2} dx \rho_a(x, -\hat{n}_*) S_{\nu_{a,s}}(t_{\text{age}} - 2x) e^{-\tau(\nu, x, -\hat{n}_*)}$$

$$t_{\text{age}} \sim (10^4 - 10^5) \text{ years} \Rightarrow x_{\text{max}} \sim 10 \text{ kpc}$$

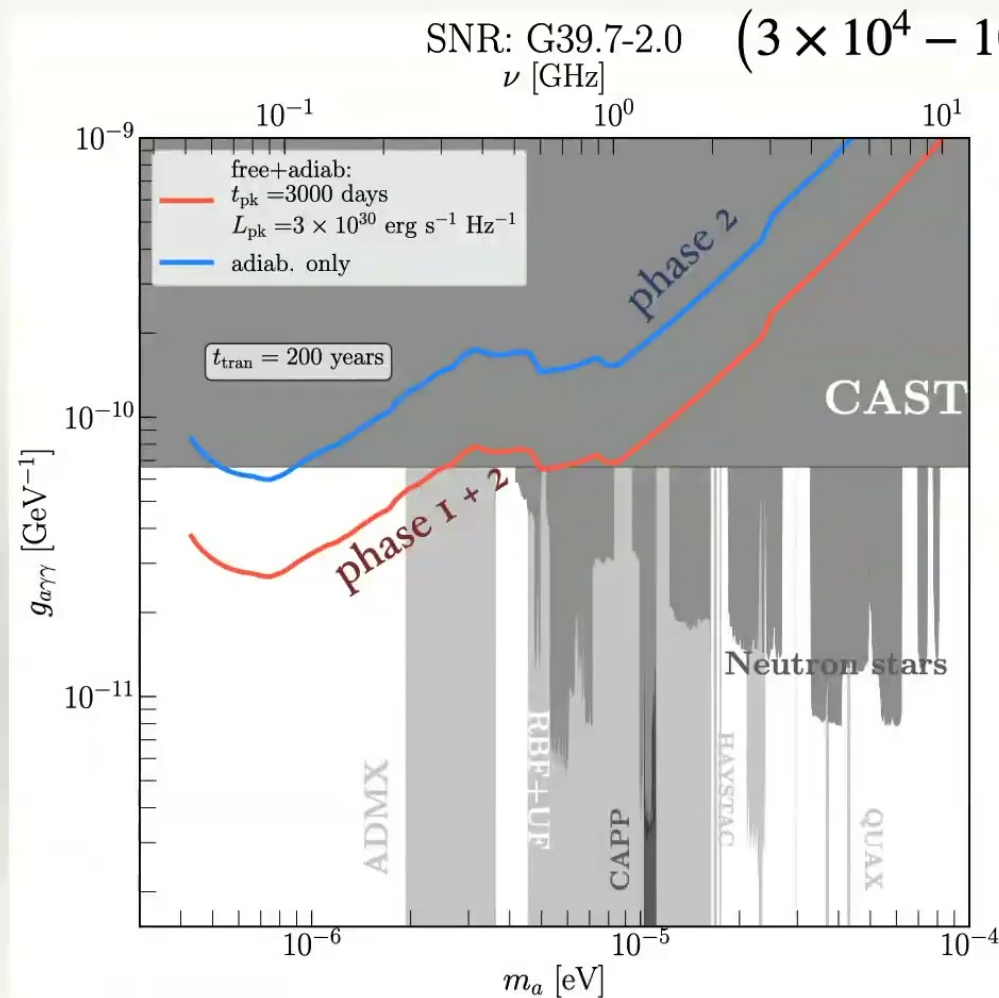
source flux density at an earlier time $t_{\text{age}} - 2x$:
echo photon emitted at a distance x away from us
was produced by light that first passed us a time
 $2x$ away.



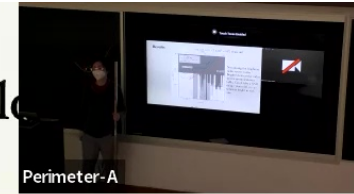
Axion echos from supernovae remnants



Results

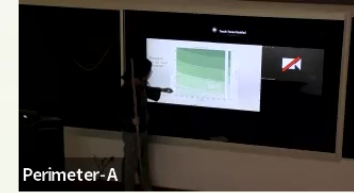
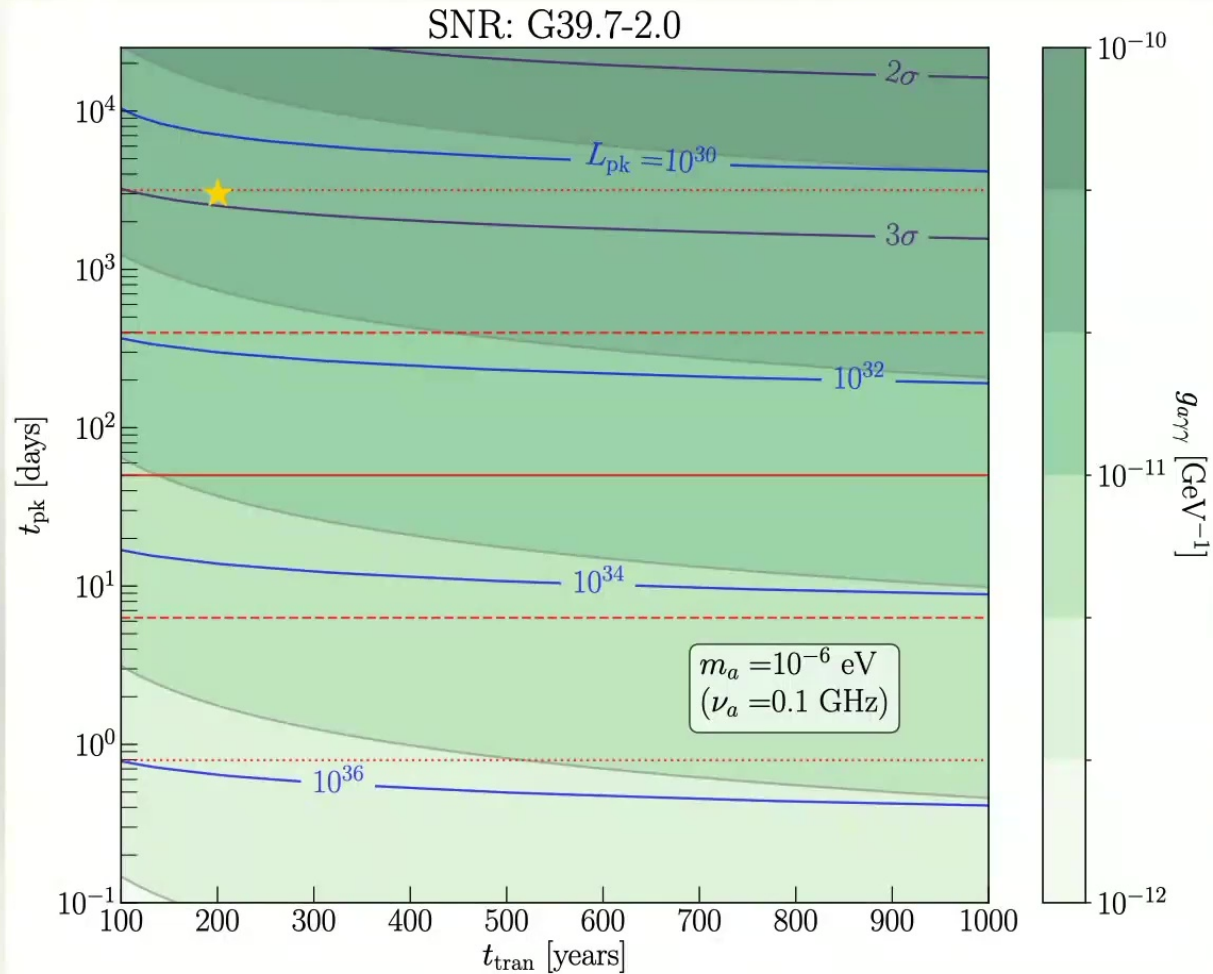


Not among the brightest radio source today !
 Brightest extrasolar radio source around/above 1 GHz, Cas A (also a SNR of 340 years old), is ~ 30 times as bright as this one.



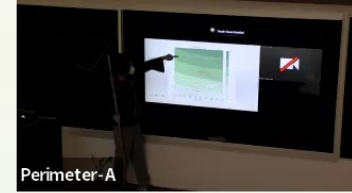
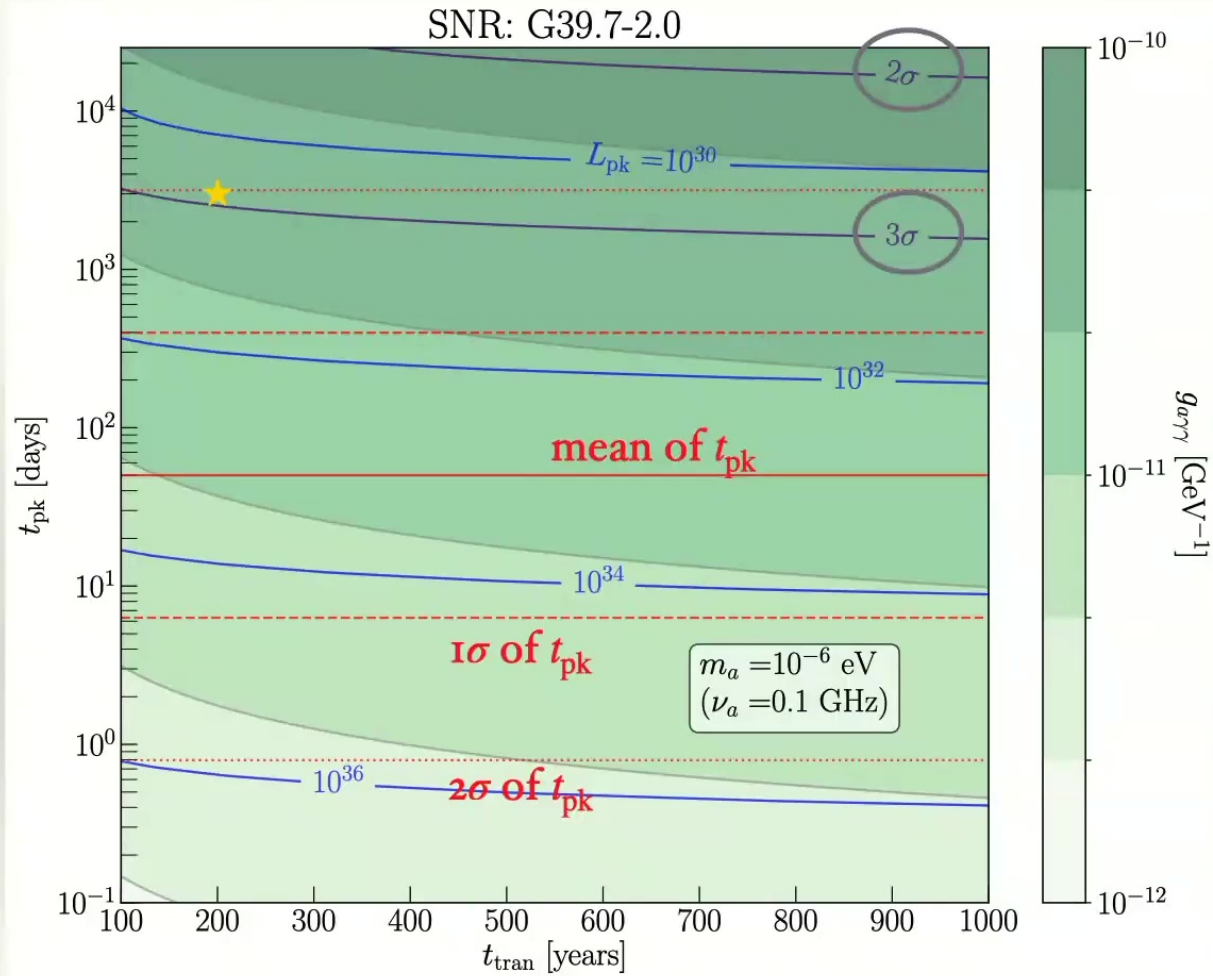
For a given
 $(t_{\text{tran}}, t_{\text{pk}})$, L_{pk} is
determined.

$$L_{\text{pk}} : \text{erg s}^{-1} \text{ Hz}^{-1}$$



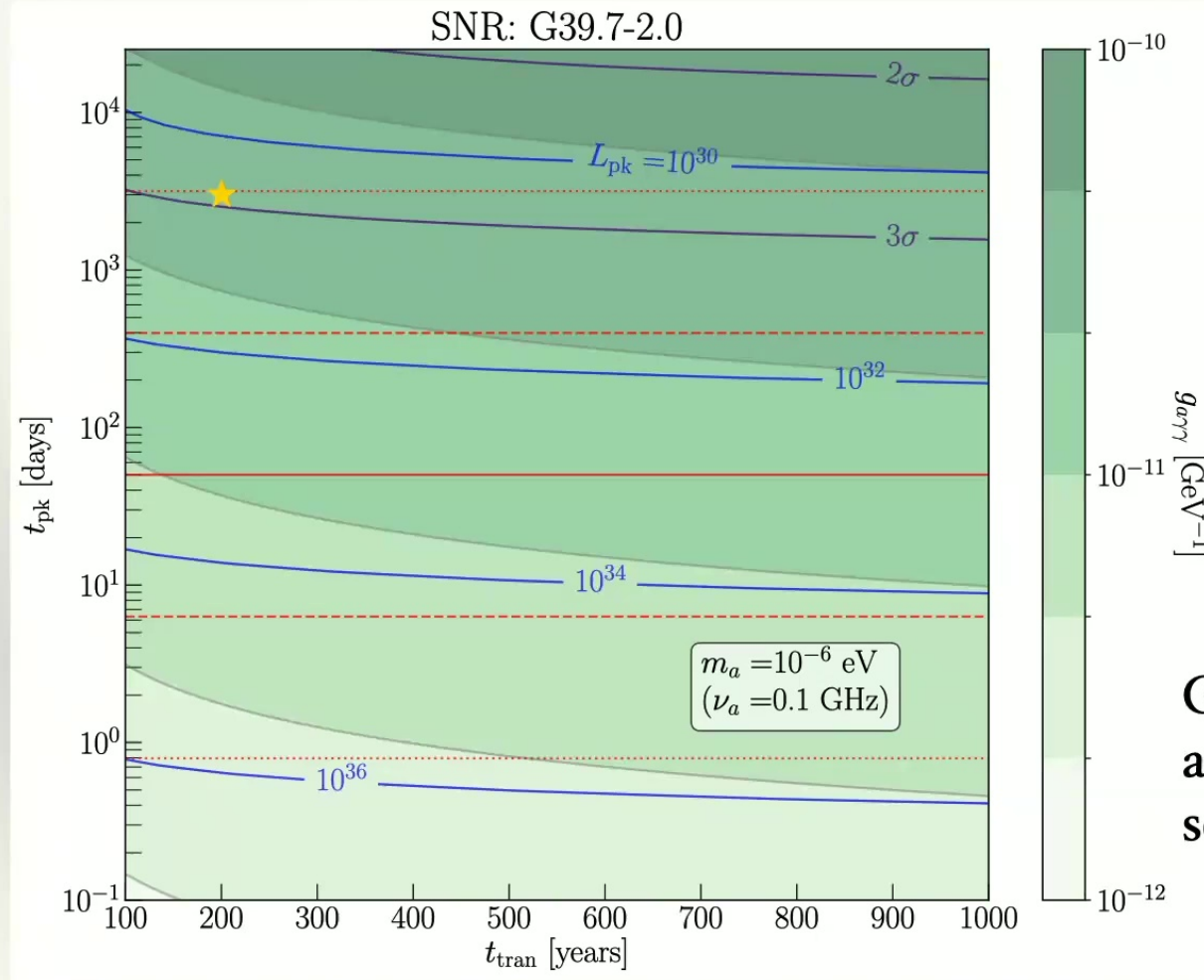
For a given
 $(t_{\text{tran}}, t_{\text{pk}})$, L_{pk} is
determined.

$$L_{\text{pk}} : \text{erg s}^{-1} \text{ Hz}^{-1}$$

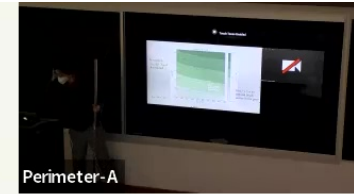


For a given
 $(t_{\text{tran}}, t_{\text{pk}})$, L_{pk} is
 determined.

$$L_{\text{pk}} : \text{erg s}^{-1} \text{ Hz}^{-1}$$

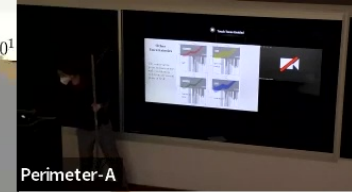
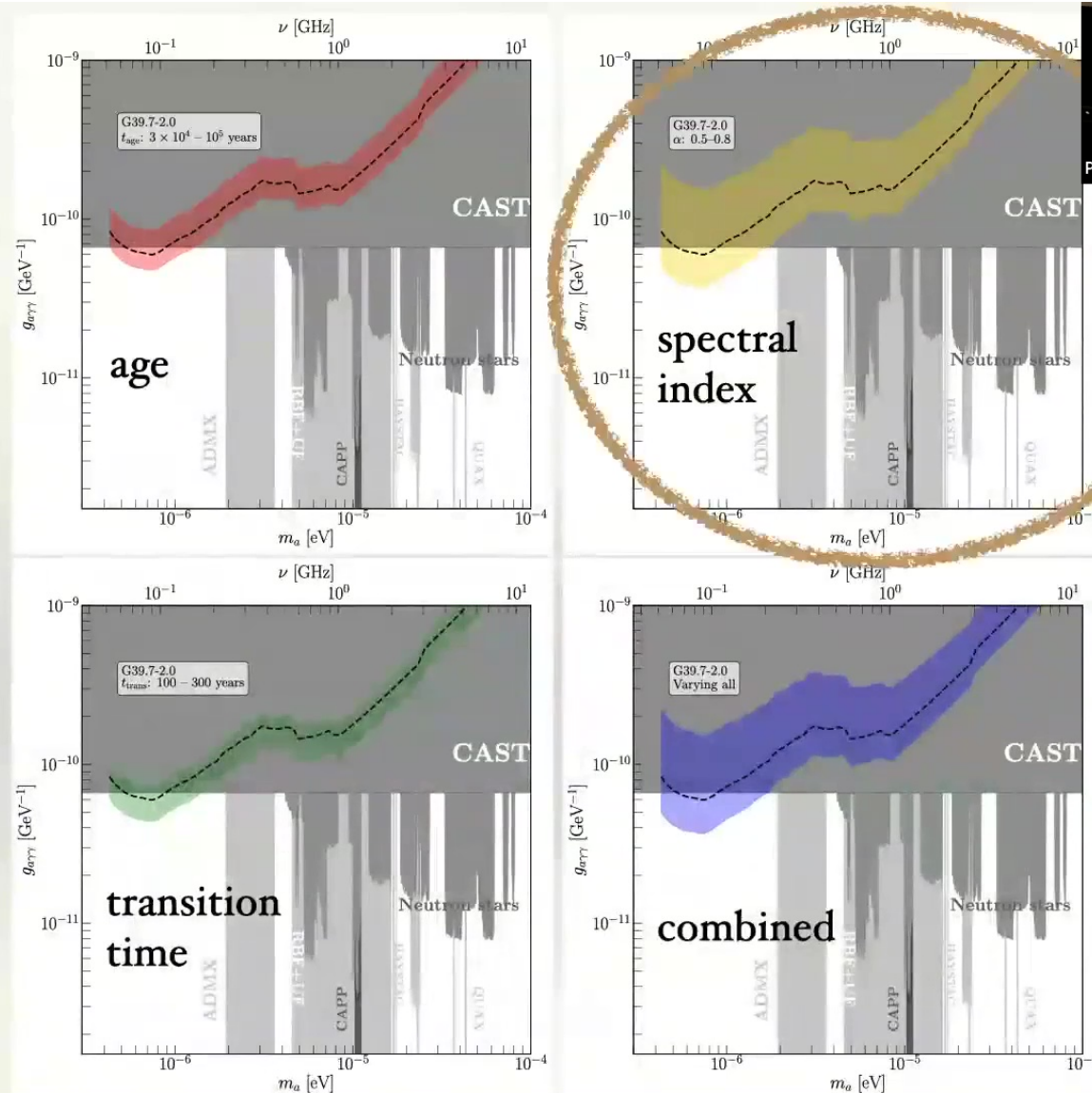


G39.7-2.0 is an
 atypical bright
 source in the past!

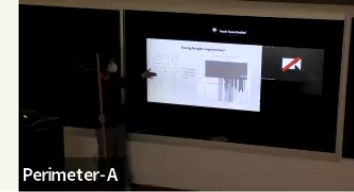


Other Uncertainties

The conservative
projected sensitivities
only from the second
phase of SNR

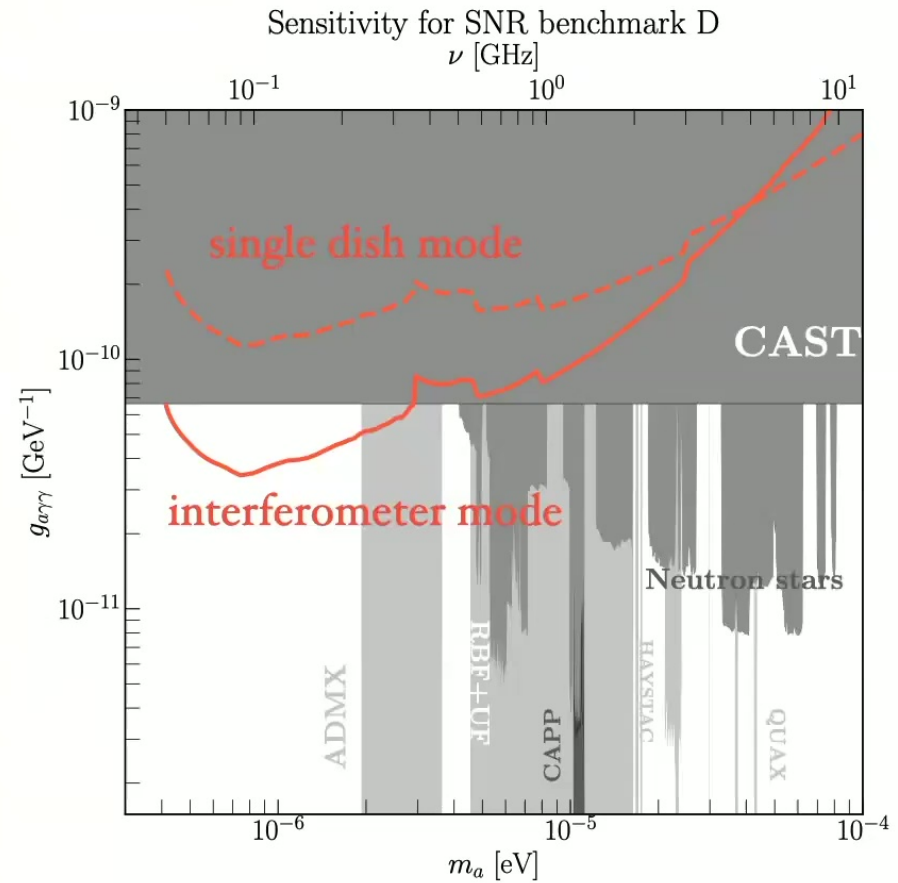


Young bright supernovae?

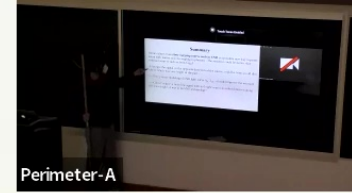


Benchmark	D (new)
(l, b)	$(40^\circ, 0^\circ)$
θ_s [arcmin]	1.0(*)
$S_{1\text{GHz},s}^{(0)}$ [Jy]	$2.1 \times 10^6(*)$
$L_{1\text{GHz},\text{pk}}$ [cgs]	1.2×10^{29}
t_{pk} [day]	50
t_{tran} [year]	4.1
D [kpc]	0.5
t_{age} [year]	10
α	0.65
γ	1.84(*)

Flux density
comparable
to the solar
burst



Summary



Axion echoes from *time-varying source such as SNR* is a possible new way to probe axion dark matter and its coupling to photons. The sensitive could be better than constant sources such as from Cyg A.

A strong echo signal, in the opposite direction of the source, could be from an old dim source which was very bright in the past.

- More precise modeling of SNR light curve (t_{pk}, L_{pk}, α) could improve the estimate.
- “Ghost” source: a radio line signal with no bright source in either direction along the line of sight. A way to do SNR archaeology ?