Title: TBA

Speakers: Daniel Ranard

Series: Perimeter Institute Quantum Discussions

Date: February 02, 2022 - 3:30 PM

URL: https://pirsa.org/22020044

Abstract: Abstract: TBD

Zoom Link:

Pirsa: 22020044 Page 1/35



# Coarse-grained entropy, microstates, and the quantum marginal problem

Daniel Ranard (MIT)

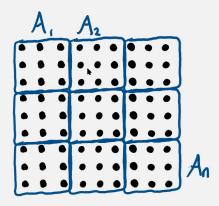
Perimeter Institute, February 2022

Pirsa: 22020044 Page 2/35



#### Outline

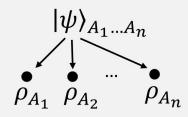
- Review:
  - Coarse-grained entropy
  - Quantum marginal problem
- Main results
  - $\log M < S_{CG}(\rho)$  (Osborne 2008)
  - $S_{CG}(\rho) \approx \log M$  (Today)
- Proof sketches
- Open questions



$$S_{CG}(\rho) = \sum S(\rho_{A_i})$$

$$S_{CG}(\rho) \approx \log M$$

M = # of solutions  $|\psi\rangle$  to marginal problem



Pirsa: 22020044 Page 3/35



Classically, entropy satisfies the 2<sup>nd</sup> Law of Thermo. What's the analog in quantum many–body systems?

Take closed quantum system with state  $\rho(t)$ 

$$\rightarrow S(\rho(t)) = \text{constant}$$

Global von Neumann entropy not a great candidate for the 2<sup>nd</sup> Law!

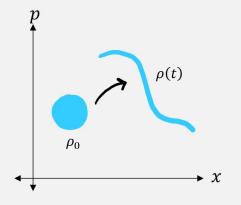
Pirsa: 22020044 Page 4/35



Analogy from classical physics:

Classical "state" 
$$\rho(x,p)$$
 = probability distribution over phase space  $S(\rho(t)) = -\int dx \ dp \ \rho(x,p,t) \log \rho \ (x,p,t)$ 

That entropy doesn't change either!



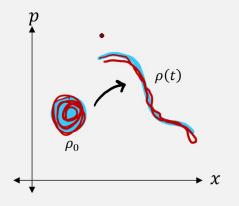
Pirsa: 22020044 Page 5/35



Analogy from classical physics:

Classical "state" 
$$\rho(x,p)$$
 = probability distribution over phase space  $S(\rho(t)) = -\int dx \ dp \ \rho(x,p,t) \log \rho \ (x,p,t)$ 

That entropy doesn't change either!



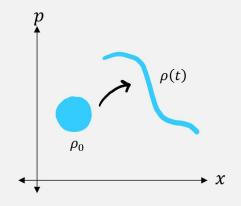
Pirsa: 22020044 Page 6/35



Analogy from classical physics:

Classical "state"  $\rho(x,p)$  = probability distribution over phase space  $S(\rho(t)) = -\int dx \ dp \ \rho(x,p,t) \log \rho \ (x,p,t)$ 

That entropy doesn't change either! Instead....





(4) (b) (2) (b) (c) (c) (c) (c)

Pirsa: 22020044 Page 7/35



Analogy from classical physics:

Classical "state" 
$$\rho(x,p)$$
 = probability distribution over phase space  $S(\rho(t)) = -\int dx \ dp \ \rho(x,p,t) \log \rho \ (x,p,t)$ 

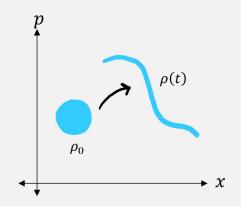
That entropy doesn't change either! Instead....



For microstate  $\alpha$ ,

$$S(\alpha) = \log \Omega$$

 $\Omega$  = # microstates with same coarse-grained properties as  $\alpha$  e.g. total energy





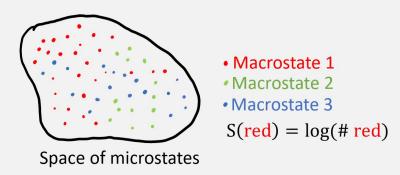
Pirsa: 22020044 Page 8/35



#### Coarse-grained entropy in QM?

#### Lessons from classical case:

- Useful notion of entropy involves coarse-graining
- Entropy associated to microstate counts how many *other* microstates have the same coarse-grained properties



Define analogous "coarse-grained entropy"  $S_{CG}$  in quantum? Lots of options. We'll study a very simple one!

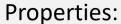
Pirsa: 22020044 Page 9/35



# Quantum coarse-grained entropy $S_{CG}$

Partition the lattice into regions  $A_i$ . Define **coarse-grained entropy**:

$$S_{CG}(\rho) = \sum_{i} S(\rho_{A_i})$$

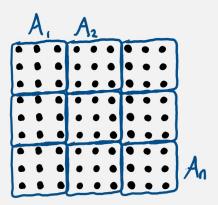


- Depends on partition, but often insensitive to exact choice.
- If  $\rho$  is pure product state,  $S_{CG}(\rho) = 0$ .

• If 
$$\rho$$
 is thermal state  $\rho = e^{-\beta H}/Z$  for local Hamiltonian, 
$$S_{CG}(\rho) \approx \sum_i S\left(\frac{e^{-\beta H}_{A_i}}{Z_i}\right) \approx S\left(\frac{e^{-\beta H}}{Z}\right) + boundary\ terms \approx \ \text{thermal entropy.}$$

- If  $\rho$  is a pure state that looks thermal on regions  $A_i$ , same as above.
- Under natural time-evolution,  $S_{CG}$  tends to increase.

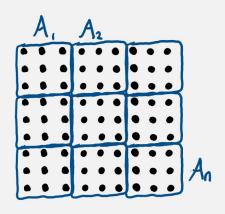
But what does this entropy *count*?





#### Coarse-grained entropy as microstate counting

Given  $\rho$ , ask about states  $|\phi\rangle$  that match  $\rho$  on all regions  $A_i$ , i.e.  ${\rm Tr}_{\overline{A_i}}|\phi\rangle\langle\phi|\approx \rho_{A_i}$ 



Such states  $|\phi\rangle$  like "microstates," with the same local properties as  $\rho$ . If you can only observe  $\rho$  locally, then for all you know the system could really be in state  $|\phi\rangle$ .

We show: The quantity  $S_{CG}(\rho) = \sum_i S(\rho_{A_i})$  counts microstates.

Pirsa: 22020044 Page 11/35



#### Quantum marginal problem

"Marginal" = "Reduced density matrix" = "RDM".

#### **Quantum marginal problem:**

Given list of marginals  $\rho_{A_i}$ , does there exist a global state  $\sigma$  consistent with them?

Any  $\sigma$  with  $\sigma_{A_i} = \rho_{A_i}$  is a "solution" to that marginal problem.

#### **Example:**

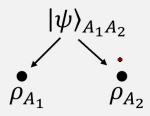
Given  $\rho_{A_1}$  and  $\rho_{A_2}$  on two qubits  $A_1$  and  $A_2$ , does there exist a pure state  $|\psi\rangle$  with those marginals?

(Answer: Yes iff  $\rho_{A_1}$  and  $\rho_{A_2}$  have same spectrum.)

General marginal problem: hard! (QMA-complete.)

But we're interested in approximate solutions for large systems: easier.

We want to **count how many solutions**: what's the largest number of pure, orthogonal states we can find, such that each state has marginals approximately given by the  $\rho_{A_i}$ ?



Pirsa: 22020044 Page 12/35



#### Prior work

Coarse-grained entropy  $S_{CG}(\rho) = \sum_i S(\rho_{A_i})$  and related quantities discussed by many.

- Quantum thermodynamics: Gell-Mann + Hartle
- Holography: Susskind, Kelly + Wall, Englehardt + Wall

Tobias Osborne noted the upper bound

$$\log M \leq S_{CG}(\rho)$$

where M is the max number of pure, orthogonal solutions to marginal problem. We'll be showing (approximate) equality!

Brandao + Dalzell show the existence of *one* approximate solution to the marginal problem, given arbitrary marginals on overlapping local regions. (With an MPS!) We'll be interested in showing *many* solutions.

Pirsa: 22020044 Page 13/35



#### Main results and intuition



Pirsa: 22020044 Page 14/35



### Main result (for case of disjoint regions)

Partition the system into disjoint regions  $A_i$ . Define *coarse-grained entropy:* 

$$S_{CG}(\rho) = \sum_{i=1}^{n} S(\rho_{A_i})$$



Let M = "number of orthogonal, pure, approximate solutions to marginal problem"

= size of the largest set of orthogonal states 
$$\{|\phi_{\alpha}\rangle\}_{\alpha=1}^{M}$$
 such that  $\|\mathrm{Tr}_{\overline{A_i}}|\phi_{\alpha}\rangle\langle\phi_{\alpha}|-\rho_{A_i}\|_1\leq \epsilon$  for each  $A_i$ ,  $|\phi_{\alpha}\rangle$ .

Then

$$S_{CG} \approx \log M$$

In particular:

$$S_{CG} - \sqrt{n \log \epsilon^{-1}} \log d \le \log M \le S_{CG} + \epsilon n \log d$$

$$d = \max_{i} \dim(A_i)$$

Note generally  $S_{CG} \propto n$  is extensive.  $S_{CG} \approx \log M$  up to "sub-extensive corrections" if e.g.  $\epsilon \sim \frac{1}{n^2}$ .

Pirsa: 22020044 Page 15/35



## Main result (for case of disjoint regions)

Partition the system into disjoint regions  $A_i$ . Define *coarse-grained entropy:* 

$$S_{CG}(\rho) = \sum_{i=1}^{n} S(\rho_{A_i})$$



= size of the largest set of orthogonal states  $\{|\phi_{\alpha}\rangle\}_{\alpha=1}^{M}$  such that  $\|\mathrm{Tr}_{\overline{A_i}}|\phi_{\alpha}\rangle\langle\phi_{\alpha}|-\rho_{A_i}\|_1\leq \epsilon$  for each  $A_i$ ,  $|\phi_{\alpha}\rangle$ .



Then

In particular:

$$S_{CG} \approx \log M$$

$$S_{CG} - \sqrt{n \log \epsilon^{-1}} \log d \le \log M \le S_{CG} + \epsilon n \log d$$

$$d = \max_{i} \dim(A_i)$$

Note generally  $S_{CG} \propto n$  is extensive.  $S_{CG} \approx \log M$  up to "sub-extensive corrections" if e.g.  $\epsilon \sim \frac{1}{n^2}$ .

Pirsa: 22020044



# Intuition for $S_{CG} \approx \log M$

If each  $\rho_{A_i} = |\psi_{A_i}\rangle\langle\psi_{A_i}|$  is pure, the only solution is  $|\psi\rangle = |\psi_{A_1}\rangle \dots |\psi_{A_n}\rangle$ .  $S_{CG} = 0, \quad M = 1.$ 

If each  $\rho_{A_i}$  maximally mixed, Haar-random  $|\psi\rangle$  gives approximate solution. So "most" pure states work.  $S_{CG} = n \log d$ ,  $M \approx \dim H = d^n$ .

Pirsa: 22020044 Page 17/35



### Generalization to overlapping marginals

With caveats:

 $S_{CG}(\rho)$  also counts the number of orthogonal pure states that match the marginals  $\rho_A$  on all regions of some length scale, i.e. including overlapping regions.

Pirsa: 22020044 Page 18/35

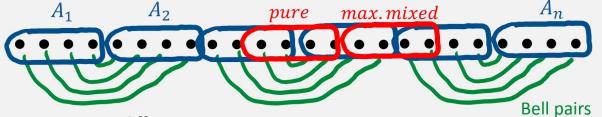


#### Generalization to overlapping marginals

#### With caveats:

 $S_{CG}(\rho)$  also counts the number of orthogonal pure states that match the marginals  $\rho_A$  on *all* regions of some length scale, i.e. including overlapping regions.

#### Example of issues:



 $A_i$  maximally mixed

 $S_{CG}$  for blue and red partitions very different.

We'll define a version of  $S_{CG}$  that's minimized over all partitions of fixed length scale.

Then we'll count pure states that match the marginals for all regions (independent of partition) below some **smaller** scale.

Pirsa: 22020044 Page 19/35

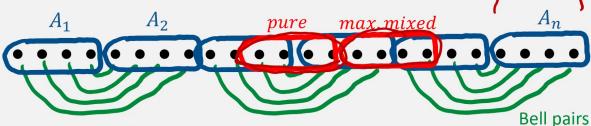


### Generalization to overlapping marginals

#### With caveats:

 $S_{CG}(\rho)$  also counts the number of orthogonal pure states that match the marginals  $\rho_A$  on *all* regions of some length scale, i.e. including overlapping regions.

#### Example of issues:



 $A_i$  maximally mixed



 $S_{CG}$  for blue and red partitions very different.

We'll define a version of  $S_{CG}$  that's minimized over all partitions of fixed length scale.

Then we'll count pure states that match the marginals for all regions (independent of partition) below some smaller scale.

Pirsa: 22020044 Page 20/35



**Statement:** [specialized to case  $\epsilon = 0$ ]

Partition the system into disjoint regions  $A_i$ . Assume there exist orthogonal states  $\{|\phi_{\alpha}\rangle\}_{\alpha=1}^{M}$  such that

$$\operatorname{Tr}_{\overline{A_i}} |\phi_{\alpha}\rangle\langle\phi_{\alpha}| = \rho_{A_i}$$

for each  $|\phi_{\alpha}\rangle$  and  $A_i$ . Then  $\log M \leq S_{CG}(\rho) \equiv \sum_i S(\rho_{A_i})$ .

Pirsa: 22020044 Page 21/35



**Statement:** [specialized to case  $\epsilon = 0$ ]

Partition the system into disjoint regions  $A_i$ . Assume there exist orthogonal states  $\{|\phi_{\alpha}\rangle\}_{\alpha=1}^{M}$  such that

$$\operatorname{Tr}_{\overline{A_i}} |\phi_{\alpha}\rangle\langle\phi_{\alpha}| = \rho_{A_i}$$

for each  $|\phi_{\alpha}\rangle$  and  $A_i$ . Then  $\log M \leq S_{CG}(\rho) \equiv \sum_i S(\rho_{A_i})$ .

Proof (Osborne):

Take the average:  $\sigma = \frac{1}{M} \sum_{\alpha} |\phi_{\alpha}\rangle \langle \phi_{\alpha}|$ . Note  $\sigma_{A_i} = \rho_{A_i}$ .

Then

$$\log M = S(\sigma) \leq \sum_{subadditivity} S(\sigma_{A_i}) = \sum_{i} S(\rho_{A_i}).$$

Easy!





**Statement:** [specialized to case  $\epsilon = 0$ ]

Partition the system into disjoint regions  $A_i$ . Assume there exist orthogonal states  $\{|\phi_{\alpha}\rangle\}_{\alpha=1}^{M}$  such that

$$\operatorname{Tr}_{\overline{A_i}} |\phi_{\alpha}\rangle\langle\phi_{\alpha}| = \rho_{A_i}$$

for each  $|\phi_{\alpha}\rangle$  and  $A_i$ . Then  $\log M \leq S_{CG}(\rho) \equiv \sum_i S(\rho_{A_i})$ .

Proof (Osborne):

Take the average:  $\sigma = \frac{1}{M} \sum_{\alpha} |\phi_{\alpha}\rangle\langle\phi_{\alpha}|$ . Note  $\sigma_{A_i} = \rho_{A_i}$ .

Then

$$\log M = S(\sigma) \leq \sum_{subadditivity} S(\sigma_{A_i}) = \sum_{i} S(\rho_{A_i}) > \int_{CG} (\varsigma) \epsilon_{asy!}$$

Pirsa: 22020044 Page 23/35



#### **Statement:**

Partition the system into disjoint regions  $A_i$ . Let  $S_{CG}(\rho) \equiv \sum_i S(\rho_{A_i})$ .

Then there exist M orthogonal pure states  $\{|\phi_{\alpha}\rangle\}_{\alpha=1}^{M}$  such that

$$\operatorname{Tr}_{\overline{A_i}} |\phi_{\alpha}\rangle\langle\phi_{\alpha}| \approx \rho_{A_i}$$

for some M with

$$\log M > S_{CG}(\rho) - o\left(\sqrt{n\log \epsilon^{-1}}\log d\right). \qquad d = \max_{i} \dim(A_i)$$

#### Extra desideratum:

Would be nice for the  $|\phi_{\alpha}\rangle$  to be low complexity, as possible "microstates."

Pirsa: 22020044 Page 24/35



**Statement:**  $\log M > S_{CG}(\rho) - o(\sqrt{n \log \epsilon^{-1}} \log d)$ 

#### **Proof sketch:**

Specialize to case  $A_1, ..., A_n$  are each single qubits, with same  $\rho_{A_i} = p_1 |0\rangle\langle 0| + p_2 |1\rangle\langle 1|$ .

Want to construct  $M \approx e^{S_{CG}}$  states  $|\phi\rangle$  with marginals  $\approx \rho_{A_i}$ .  $S_{CG} = nS(\rho_A) = nS(\{p_1, p_2\})$ .

Let  $T = \{ \text{product states } | 011011 \dots \} \text{ with } \approx p_1 \text{ zeros and } \approx p_2 \text{ ones} \}$ 

Number of such states is  $|T| \approx 2^{nS(\{p_1,p_2\})}$ .

Randomly choose k states from T, denoted  $\{|\psi_i\rangle\}_{i=1}^k$ . Take  $k \propto \log n$ . The state  $\frac{1}{k}\sum_{i=1}^k |\psi_i\rangle \langle \psi_i|$  has marginals close to  $\rho_{A_i}$ , with "sampling error" like  $1/\sqrt{k}$ . Note  $|\phi_1\rangle \equiv \frac{1}{\sqrt{k}}\sum_{i=1}^k |\psi_i\rangle$  has same marginals as the  $\frac{1}{k}\sum_{i=1}^k |\psi_i\rangle \langle \psi_i|$ , because the terms are unlikely to interfere.

 $\rightarrow$  We already found one approximate solution  $|\phi_1\rangle$ 

Now repeat: Pick k more random states from states from T (avoiding the ones we already picked). Repeat to obtain another solution  $|\phi_2\rangle$ .

Can repeat process for  $M \approx \exp(S_{CG}(\rho) - o(n))$  iterations without significantly depleting T.

Pirsa: 22020044 Page 25/35



**Statement:**  $\log M > S_{CG}(\rho) - o(\sqrt{n \log \epsilon^{-1}} \log d)$ 

#### **Proof sketch:**

P(error) on SA. Specialize to case  $A_1, \dots, A_n$  are each single qubits, with same  $\rho_{A_i} = p_1 |0\rangle\langle 0| + p_2 |1\rangle\langle 1|$ 

Want to construct  $M \approx e^{S_{CG}}$  states  $|\phi\rangle$  with marginals  $\approx \rho_{A_i}$ .  $S_{CG} = nS(\rho_A) = nS(\{p_1, p_2\})$ .

Let  $T = \{\text{product states } | 011011 \dots \}$  with  $\approx p_1$  zeros and  $\approx p_2$  ones $\}$ 

Number of such states is  $|T| \approx 2^{nS(\{p_1,p_2\})}$ .

Randomly choose k states from T, denoted  $\{|\psi_i\rangle\}_{i=1}^k$ . Take  $k \propto \log n$ . The state  $\frac{1}{k}\sum_{i=1}^k |\psi_i\rangle\langle\psi_i|$  has marginals close to  $\rho_{A_i}$ , with "sampling error" like  $1/\sqrt{k}$ . Note  $|\phi_1\rangle\equiv\frac{1}{\sqrt{k}}\sum_{i=1}^k |\psi_i\rangle$  has same marginals as the  $\frac{1}{k}\sum_{i=1}^k |\psi_i\rangle\langle\psi_i|$ , because the terms are unlikely to interfere.

 $\rightarrow$  We already found one approximate solution  $|\phi_1\rangle$ .

Now repeat: Pick k more random states from states from T (avoiding the ones we already picked). Repeat to obtain another solution  $|\phi_2\rangle$ .

Can repeat process for  $M \approx \exp(S_{CG}(\rho) - o(n))$  iterations without significantly depleting T.

Pirsa: 22020044 Page 26/35

**Statement:**  $\log M > S_{CG}(\rho) - o(\sqrt{n \log \epsilon^{-1}} \log d)$ 

#### **Proof sketch:**

Specialize to case  $A_1, ..., A_n$  are each single qubits, with same  $\rho_{A_i} = p_1|0\rangle\langle 0| + p_2|1\rangle\langle 1|$ .

Want to construct  $M \approx e^{S_{CG}}$  states  $|\phi\rangle$  with marginals  $\approx \rho_{A_i}$ .  $S_{CG} = nS(\rho_A) = nS(\{p_1, p_2\})$ .

Let  $T = \{\text{product states } | 011011 \dots \}$  with  $\approx p_1$  zeros and  $\approx p_2$  ones $\}$ 

Number of such states is  $|T| \approx 2^{nS(\{p_1,p_2\})}$ .

Randomly choose k states from T, denoted  $\{|\psi_i\rangle\}_{i=1}^k$ . Take  $k \propto \log n$ . The state  $\frac{1}{k}\sum_{i=1}^k |\psi_i\rangle\langle\psi_i|$  has marginals close to  $\rho_{A_i}$ , with "sampling error" like  $1/\sqrt{k}$ . Note  $|\phi_1\rangle\equiv \frac{1}{\sqrt{k}}\sum_{i=1}^k |\psi_i\rangle$  has same marginals as the  $\frac{1}{k}\sum_{i=1}^k |\psi_i\rangle\langle\psi_i|$ , because the terms are unlikely to interfere.

 $\rightarrow$  We already found one approximate solution  $|\phi_1\rangle$ .

Now repeat: Pick k more random states from states from T (avoiding the ones we already picked). Repeat to obtain another solution  $|\phi_2\rangle$ .

Can repeat process for  $M \approx \exp(S_{CG}(\rho) - o(n))$  iterations without significantly depleting T.



**Statement:**  $\log M > S_{CG}(\rho) - o(\sqrt{n \log \epsilon^{-1}} \log d)$ 

#### **Proof sketch:**

Specialize to case  $A_1, ..., A_n$  are each single qubits, with same  $\rho_{A_i} = p_1 |0\rangle\langle 0| + p_2 |1\rangle\langle 1|$ .

Want to construct  $M \approx e^{S_{CG}}$  states  $|\phi\rangle$  with marginals  $\approx \rho_{A_i}$   $\rho_{A_i}$   $\rho_{A_i}$ 

Number of such states is  $|T| \approx 2^{nS(\{p_1,p_2\})}$ .

Randomly choose k states from T, denoted  $\{|\psi_i\rangle\}_{i=1}^k$ . Take  $k \propto \log n$ . The state  $\frac{1}{k}\sum_{i=1}^k |\psi_i\rangle\langle\psi_i|$  has marginals close to  $\rho_{A_i}$ , with "sampling error" like  $1/\sqrt{k}$ . Note  $|\phi_1\rangle\equiv\frac{1}{\sqrt{k}}\sum_{i=1}^k |\psi_i\rangle$  has same marginals as the  $\frac{1}{k}\sum_{i=1}^k |\psi_i\rangle\langle\psi_i|$ , because the terms are unlikely to interfere.

 $\rightarrow$  We already found one approximate solution  $|\phi_1\rangle$ .

Now repeat: Pick k more random states from states from T (avoiding the ones we already picked). Repeat to obtain another solution  $|\phi_2\rangle$ .

Can repeat process for  $M \approx \exp(S_{CG}(\rho) - o(n))$  iterations without significantly depleting T.

Pirsa: 22020044 Page 28/35



#### Statement for "overlapping" case:

Choose two length scales,  $L'\gg L$ . Define  $S_{CG}^{(L')}(\rho)=\min_{partitions}\sum_{i=1}^n S(\rho_{A_i})$  with minimum over partitions into **disjoint** regions  $A_i$  of size  $|A_i|=L'$ 

Then  $\exists$  M<sub>L</sub> orthogonal pure states that match  $\rho$  (to error L/L') on all regions of size  $\leq L$  (including overlapping regions), with

$$\log M_L > S_{CG}^{(L')} - o\left(n\frac{L'}{L}\right)$$

#### **Proof sketch:**

Fix a partition into regions  $A_i$  of size L'. Already showed how to find  $\approx \exp(S_{CG}^{(L')})$  pure states with marginals  $\rho_i$ . Call this set of pure states  $W_1$ . Then shift regions  $A_i$  by one site and repeat: call this set of states  $W_2$ . Repeat to  $W_{L'}$ .

Choose one state from each  $W_i$  and take uniform superposition. New state has correct marginals on every region of size  $\leq L$ , to error L'/L. (How do we ensure the states from different  $W_i$  don't interfere? Trick adapted from Brandao + Dalzell.)

Then repeat, choosing a different state from each  $W_i$ .

Pirsa: 22020044 Page 29/35



#### Summary

We took a common proxy for thermodynamic entropy in quantum many-body systems, the coarse-grained entropy:  $A_i$   $A_i$ 

$$S_{CG}(\rho) = \sum_{i} S(\rho_{A_i})$$

What "microstates" does this quantity count?

We showed

$$S_{CG}(\rho) \approx \log M$$

counts the number M of orthogonal pure states with the same marginals as  $\rho$ .

Pirsa: 22020044 Page 30/35



#### Concrete open questions

- Can we count **exact** solutions to the marginal problem? For non-overlapping marginals, I suspect there are  $\exp(S_{CG}(\rho))$  exact solutions. Sketchy proof, with numerical support.
- Can we construct O(1)-complexity solutions to the marginal problem? Brandao + Dalzell already showed how to construct a *single* approximate solution of O(1)-complexity. Can you construct  $\exp(S_{CG}(\rho))$  such solutions?
- Can we extend our bounds to the alternative "max-entropy" definition of  $S_{CG}$  (e.g. as used in holography)?

Pirsa: 22020044 Page 31/35



#### Open open questions

- Connection to quantum PCP conjecture?
- What are dynamics of  $S_{CG}(\rho(t))$ ? Microstate interpretation suggests it should increase?

Point of tension: recent work (Cotler, Jones, DR) suggests  $S_{CG}$  fluctuates in time much *less* in quantum case than classical.

• What's structure of space of consistent marginals? To specify all k-body RDMs on an n-body system only takes  $\approx k \, n \log n$  bits, rather than naïve  $k^n$ . What's the right "data structure" to encode the set of k-body RDMs?

Pirsa: 22020044 Page 32/35



#### Open open questions

- Connection to quantum PCP conjecture?
- What are dynamics of  $S_{CG}(\rho(t))$ ? Microstate interpretation suggests it should increase?

Point of tension: recent work (Cotler, Jones, DR) suggests  $S_{CG}$  fluctuates in time much *less* in quantum case than classical.

• What's structure of space of consistent marginals? To specify all k-body RDMs on an n-body system only takes  $\approx k \, n \log n$  bits, rather than naïve  $k^n$ . What's the right "data structure" to encode the set of k-body RDMs?

Pirsa: 22020044 Page 33/35



### Acknowledgments

Thanks to Michael Walter for discussions on the quantum marginal problem, and Geoff Penington for discussions on holographic coarsegrained entropies.

•

Pirsa: 22020044 Page 34/35



# Thank you!

•

Pirsa: 22020044 Page 35/35