

Title: PSI Lecture - Condensed Matter - Lecture 15

Speakers: Aaron Szasz

Collection: PSI Lecture - Condensed Matter

Date: February 10, 2022 - 9:00 AM

URL: <https://pirsa.org/22020012>

Last class! - Finish Ising model

Announcements

- All remaining lecture notes online by Mon due Feb 28

Last time: what happened to degeneracy?

Periodic or anti-periodic BCs on our chain of fermions:

- Periodic if # of spin $\uparrow$ is odd $\rightarrow k = \frac{2\pi}{N}n$
AP if even $\rightarrow k = \frac{2\pi}{N}(n + \frac{1}{2})$

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AP

$$H = \begin{pmatrix} PBC & \\ & APBC \end{pmatrix} \leftarrow \text{fermionic } H$$

$$H = \sum_k E_k (a_k^\dagger a_k - \frac{1}{2})$$

key: not all eigenstates of $\rightarrow$ are allowed!

$$H = \begin{pmatrix} ABC \\ \hline APBC \end{pmatrix} \quad \leftarrow \text{APBC}$$

$$H = \sum_k \epsilon_k \left(a_k^\dagger + a_k - \frac{1}{2} \right)$$

$$\omega_k \rightarrow \epsilon_k = \frac{2\pi}{N} (n+k)$$

key: not all eigenstates of are allowed!

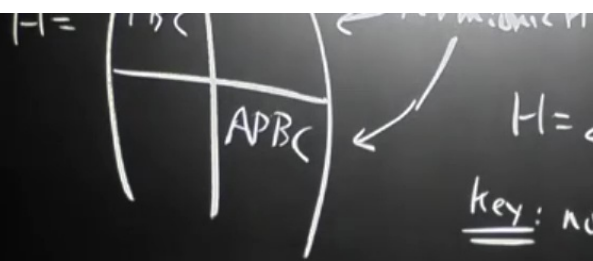
$|4\rangle \rightarrow a_k^\dagger |4\rangle \rightarrow \text{change parity}$

I can only act with $a_k^\dagger, a_{k_2}^\dagger$

$$H = \sum_k \left(\frac{1}{2} \hbar \omega_k (a_k^\dagger + a_k) + \frac{1}{2} \hbar \omega_k \right)$$

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$|4\rangle \rightarrow a_k^\dagger |4\rangle \rightarrow \text{change parity}$
 I can only act with $a_k, a_k^\dagger$

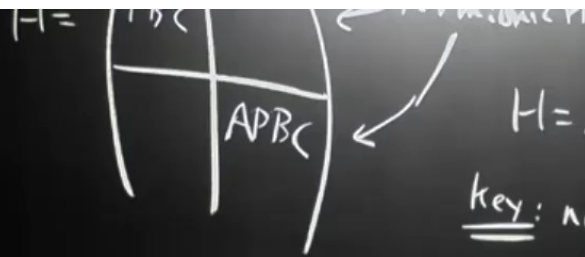


$H = \sum_r E_r (a_r^\dagger + a_r - \frac{1}{2})$
 key: not all eigenstates of H are allowed!

can only act with $a_n, a_n^\dagger$
 Our actual energy gap will be between the ground states in the two sectors.

One sector large h : $|\uparrow\uparrow\ldots\uparrow\rangle$ (N be even)
 APBC
 large J : $|\rightarrow\rightarrow\ldots\rightarrow\rangle + |\leftarrow\leftarrow\ldots\leftarrow\rangle$
 $\sqrt{2}$

PBC large h : 1 spin flip, $E = E_{GS} + 2h$
 large J : $|\rightarrow\rightarrow\ldots\rightarrow\rangle - |\leftarrow\leftarrow\ldots\leftarrow\rangle$
 $\sqrt{2}$
 $E \sim \left(\frac{h}{J}\right)^N + E_{GS}$



$$H = \sum_k E_k \left(\frac{1}{2} a_k - \frac{1}{2} \right)$$

key: not all eigenstates of are allowed!

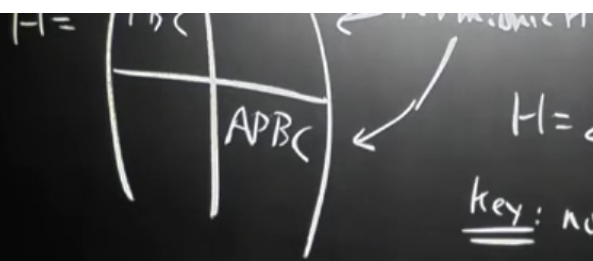
can only act with $a_k, a_k^\dagger$

Our actual energy gap will be between the ground states in the two sectors.

One sector large h : $|\uparrow\uparrow\cdots\uparrow\rangle$ (N be even)
APBC

large J : $|\rightarrow\rightarrow\cdots\rightarrow\rangle + |\leftarrow\leftarrow\cdots\leftarrow\rangle$
 H^N $\frac{\sqrt{2}}$

PBC large h : 1 spin flip, $E = E_{GS} + 2h$
large J : $|\rightarrow\rightarrow\cdots\rightarrow\rangle - |\leftarrow\leftarrow\cdots\leftarrow\rangle$
 $E \sim \left(\frac{h}{J}\right)^N \sqrt{2} + E_{GS}$



$$H = \sum_k \epsilon_k (a_k^\dagger + a_k - \frac{1}{2})$$

key: not all eigenstates of are allowed!

$$\hbar \omega \rightarrow \hbar = \frac{2\pi}{N} (n + \frac{1}{2})$$

Our actual energy gap will be between the ground states in the two sectors

One sector large h : $|\uparrow\uparrow\cdots\uparrow\rangle$ (N be even)
 APBC
 $|\psi_1\rangle \xrightarrow{H^N} |\psi_2\rangle$
 $|\psi_1\rangle \xrightarrow{H^N} |\psi_2\rangle$
 H^N
 $H_{\text{eff}} = \begin{pmatrix} 1 & 1 \\ 1 & 1 \end{pmatrix} \begin{pmatrix} |\psi_1\rangle \\ |\psi_2\rangle \end{pmatrix}$

$$|\rightarrow\rangle^{\otimes N} = |\uparrow\rangle^{\otimes N} + (1 \text{ flip}) + (2 \text{ flips})$$

$$|\leftarrow\rangle^{\otimes N} = |\uparrow\rangle^{\otimes N} - (1 \text{ flip})$$

PBC large h : 1 spin flip, $E = E_{GS} + 2h$
 large J : $|\rightarrow\rightarrow\cdots\rightarrow\rangle + |\leftarrow\leftarrow\cdots\leftarrow\rangle$
 $E \sim \left(\frac{h}{J}\right)^N \sqrt{2} + E_{GS}$

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Translate to fermions:

$$|\uparrow\rangle^{\otimes N} \rightarrow |0\rangle_c \text{ for } c^\dagger$$

$$|GS\rangle \text{ is } \left(\prod_k a_k^\dagger \right) |0\rangle_a, \quad a_k |0\rangle_a = 0 \quad \forall k$$

Parity sector of $|0\rangle_a$:

$$\underbrace{\left(\prod_k a_k \right)}_{\text{even \# of factors}} |0\rangle_c$$

$\Rightarrow$ Parity of $|0\rangle_a = \text{parity of } |0\rangle_c = (-1)^{E_N}$

• $|0\rangle_a$ has anti-periodic BCs

• APBC sector, GS energy is $\sum_k E_{k/2}$

• PBC sector, GS energy is $\left(\sum_k E_{k/2} \right) + \min_k (E_{k/2})$

Large h : $|\uparrow\uparrow\uparrow \dots \uparrow\rangle$ (PBC)

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$$|\uparrow\rangle^{\otimes N} \rightarrow |0\rangle_c \text{ for } c^\dagger$$

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Parity sector of $|0\rangle_a$:

$$\left| \left(\prod_k a_k \right) |0\rangle_c \right|$$

even # of factors

$\Rightarrow$ Parity of $|0\rangle_a = \text{parity of } |0\rangle_c = |\uparrow\rangle^{\otimes N}$

$|0\rangle_a$ has anti-periodic BCs

• APBC sector, GS energy is $\sum_k E_{k/2}$

• PBC sector, GS energy is $\left(\sum_k E_{k/2} \right) + \min_{0 < k < \pi} (E_k)$

Energy gap is

$$\min_k (E_k) + \sum_{k=\frac{2\pi}{N}} E_{k/2} - \sum_{k=\frac{2\pi}{N}(n+1)} E_{k/2}$$

$L \rightarrow \infty$: $\min_k (E_k)$

$$= 2J|1-g|$$

ThR is correct for $g > 1$

$g < 1$ At $k=\pi$

$$g < 1 \quad a_\pi = c_\pi$$

$$g > 1 \quad a_\pi = c_\pi$$

$$\forall g, a_0 = c_0$$

CAUTION

even # of factors
 $\Rightarrow$ Parity of $|0\rangle_a = \text{parity of } |0\rangle_c = (-1)^N$
 • PBC sector, GS energy is $\sum_k E_{k/2}$
 • GS energy is $\left(\sum_k E_{k/2}\right) + \min_k (E_{1/2})$
 (g>1) $a_{\pi} = c_{\pi}$
 $\forall g, a_0 = c_0$

Goal $a_k |0\rangle_a \neq 0$, and $\langle 0|_a |0\rangle_a \neq 0$

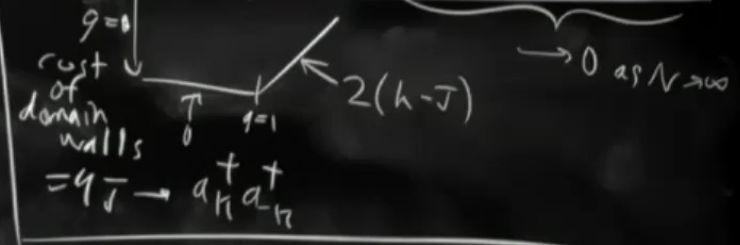
$$\left(\prod_k a_k\right) |0\rangle_c = |0\rangle_a$$

but if we include a_0 or a_{π}

Issue only occurs w/ PBC, b/c $k=0, \pi$ are not in APBC case
 Conclusion: For APBC, $E_0 = \sum_k E_{k/2}$

For PBC, when $g < 1$, $|0\rangle_a$ is in the correct sector, $\left(\prod_{k \neq 0} a_k\right) |0\rangle_c$
 $g > 1$, $|0\rangle_a$ is in the wrong sector $\leftarrow E_{\pi} + \sum E_{k/2}$

$$\Delta E = E_{\pi} \theta(g-1) + \sum_{k=\frac{2\pi}{N}} E_{k/2} - \sum_{k=\frac{2\pi}{N}(h+1)} E_{k/2}$$



$\sum E_{k/2}$
 odd #



For PBC, when $g < 1$, $|0\rangle_a$ is in the correct sector, $\left(\prod_{k \neq 0} a_k\right) |0\rangle_c$
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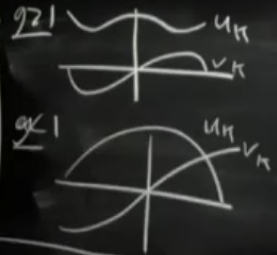
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$$d_k = u_k \sigma_k^x - i v_k \sigma_k^y$$

$$u_k = \cos(\theta_k/2)$$

$$v_k = \sin(\theta_k/2)$$

$$\tan(\theta_k) = \frac{\sin(k)}{g + \cos(k)}$$

Energy gap is

$$\min_k (E_k) + \sum_{k=\frac{2\pi}{N}} E_{k/2} - \sum_{k=\frac{2\pi}{N}(n+1)} E_{k/2}$$

$$L \rightarrow \infty: \min_k (E_k)$$

$$= 2J|1-g|$$

ThR is correct for $g > 1$

$$g \leq 1 \text{ At } k=\pi$$

$$g \leq 1 \rightarrow a_{\pi} = c_{\pi}$$

$$g > 1 \rightarrow a_{\pi} = c_{\pi}$$

$$\forall g, a_0 = c_0$$

sector of $|0\rangle_a$:

$|0\rangle_a$ has anti-periodic BCs

• APBC sector, GS energy is $\sum_k E_{k/2}$

• PBC sector, GS energy is $\left(\sum_k E_{k/2}\right) + \min_k (E_k)$

if $|0\rangle_a$ = parity of $|0\rangle_c = \prod_k \sigma_k^z$

For PBC, when $g < 1$, $|0\rangle_a$ is in the correct sector $\left(\prod_{k \neq 0} a_k |1\rangle \right)$

$E_0 = \sum_k E_{k/2}$

$\sum E_{k/2}$

odd $\#$

Claim $g=1$ is a quantum phase transition.

Last thing to do: compute order parameter / correlation function

$$M = \sqrt{\langle \sigma_0^x \sigma_n^x \rangle}, \quad M^2 = \langle \underbrace{(c_0^+ - c_0)}_{B_0} \underbrace{(c_1^+ + c_1)}_{A_1} \underbrace{(c_1^+ - c_1)}_{B_1} \dots \underbrace{(c_{n-1}^+ - c_{n-1})}_{B_{n-1}} \underbrace{(c_n^+ + c_n)}_{A_n} \rangle$$

and H is quadratic $\rightarrow$ Wick's theorem

$$\Rightarrow M^2 = \det \begin{pmatrix} G_{-1} & G_{-2} & \dots & G_{-n} \\ G_0 & & & \\ & & & \\ & & & G_{n-2} \\ & & G_{n-1} & \end{pmatrix}, \quad G_n = \frac{1}{\pi} \int_0^\pi \frac{g \cos(kn) + \cos(k(n+1))}{\sqrt{1+g^2+2g\cos(k)}} dk$$

Tueplitz determinant

$\rightarrow$ some hypergeometric fct.

Correlations - Table of Not r