

Title: QFT III 2021/2022

Speakers: Francois David

Collection: QFT III 2021/2022

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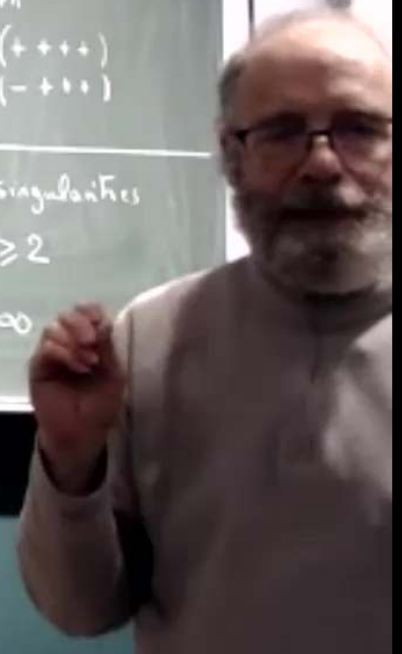
Free scalar field $\phi(x)$ Euclidean $X = (x^0, \vec{x})$ $t = -ix^0$ $K = (K_\mu) = (K_0, \vec{K})$ } $X \cdot K = X^\mu K_\mu = X^0 K_0 + \vec{x} \cdot \vec{K}$
 Minkowski $X = (t, \vec{x})$ $\omega = i K_0$ $K = (K_\mu) = (\omega, \vec{K})$ } both $= X^0 K_0 + \vec{x} \cdot \vec{K}$

Propagator vev of
 $\langle \phi(x) \phi(y) \rangle = G_0(x-y)$ Euclidean $\int \frac{d^d k}{(2\pi)^d} \frac{e^{ik(x-y)}}{k^2 + m^2}$ $K^2 = K_0^2 + \vec{K}^2$ Euclidean (+ + + +)
 $\langle 0|T[\phi(x)\phi(y)]|0\rangle = G_F(x-y)$ Minkowski $K^2 = -K_0^2 + \vec{K}^2$ Minkowski (- + + +)

Feynman propagator

what happens at short distances? UV singularities
 when $x=y$ divergent integral if $d \geq 2$
 $\int_0^\infty db \frac{k^{d-1}}{k^2}$ at $k \rightarrow \infty$

Euclidean Sym. Translation + Rotations mass of boson
 $G_0(|x-y|)$ $|x|^2 = (x^0)^2 + (\vec{x})^2$ \downarrow
 = Bessel function of $|x-y| \cdot m$



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Free scalar (real $\phi(x)$) Euclidean $X = (X^0, \vec{X})$ $t = -iX^0$ $K = (K_0) = (K_0, \vec{K})$ } $X \cdot K = X^\mu K_\mu$
 Minkowski $X = (t, \vec{x})$ $\omega = iK_0$ $K = (K_0) = (\omega, \vec{K})$ } $= X^\mu K_\mu = \vec{x} \cdot \vec{K}$
 Propagator $v.e.v.f$
 $\langle \phi(x) \phi(y) \rangle = G_0(x-y)$ Euclidean $\int \frac{d^d k}{(2\pi)^d} \frac{e^{iK(x-y)}}{K^2 + m^2}$ $K^2 = K_0^2 + \vec{K}^2$ Euclidean (+ + + +)
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what happens at short distances? UV singularities
 when $x=y$ divergent integral if $d \geq 2$
 $\int \frac{d^d k}{k^2}$ at $k \rightarrow \infty$

when $|x-y| \ll \frac{1}{m}$, then $|k| \sim \frac{1}{|x-y|}$ dominates
 forget about the mass some coefficient $2-d$
 $G_0(x) \approx \int \frac{d^d k}{(2\pi)^d} \frac{e^{iKx}}{K^2} \approx \frac{1}{2-d} |x|^{2-d}$
 Dimensional analysis $\approx \log|x|$ $d=2$
 $\approx |x|^{-1}$ $d=3$
 $\approx |x|^{-2}$ $d=4$
 etc...

Long distances $|x-y| \gg \frac{1}{m}$



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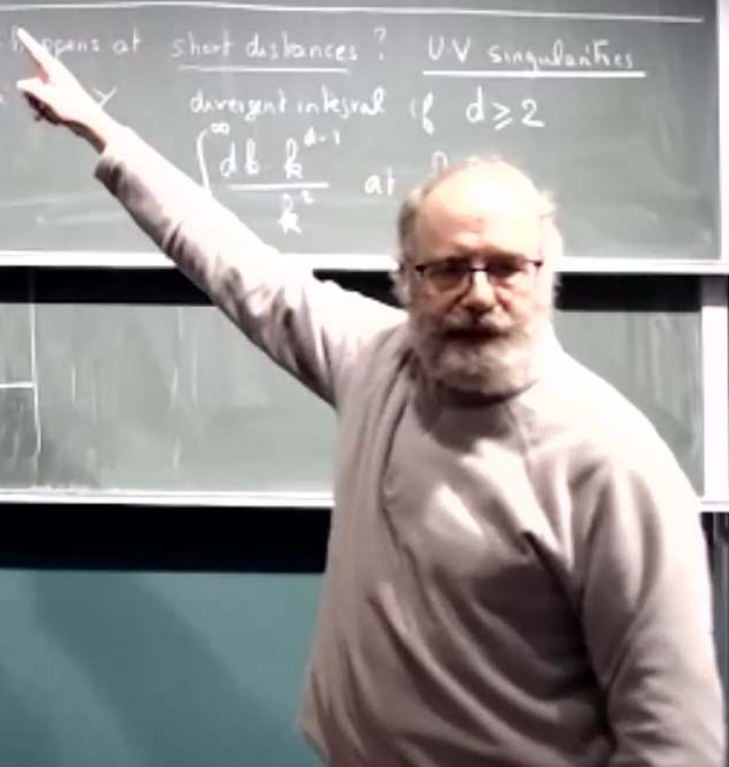
$\langle \phi(x) \phi(y) \rangle = G_0(x-y)$ Euclidean $\int \frac{d^d k}{(2\pi)^d} \frac{e^{ik(x-y)}}{k^2 + m^2}$ $K^2 = K_0^2 - \vec{K}^2$ Euclidean (+ + + +)
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Fourier propagator

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 = Bessel function of $|x-y|m$

what happens at short distances? UV singularities
 when $\int \frac{d^d k}{k^2}$ divergent integral if $d \geq 2$
 $\int \frac{d^d k}{k^2} \sim \int \frac{k^{d-1} dk}{k^2} \sim k^{d-2}$ at $k \rightarrow \infty$

Long distances $|x-y| \gg \frac{1}{m}$ $\sim |x|^{-d}$ etc...
 $G_0(x) \simeq \exp(-m|x|)$ exponential decay
 mass gap



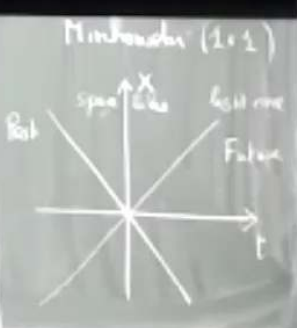
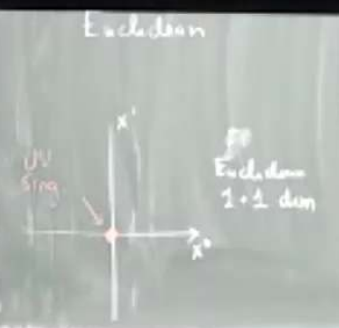
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Euclidean Sym. Translation + Rotations mass of boson
 $G_0(|x-y|)$ $|x|^2 = (x^0)^2 + (\vec{x})^2$
 = Bessel of $|x-y| \cdot m$

$G_0(x) \sim \log|x|$ $d=2$
 $\sim |x|^{-1}$ $d=3$
 $\sim |x|^{-2}$ $d=4$
 etc...
 Long distances $|x|$
 $G_0(x) \sim e^{-m|x|}$
 exponential decay
 mass gap $C=1$



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when $|x-y| \ll \frac{1}{m}$, then $|k| \sim \frac{1}{|x-y|}$ dominates

forget about the mass \sim some coefficient

$$G_0(x) = \int \frac{d^d k}{(2\pi)^d} \frac{e^{i k x}}{k^2} \approx \frac{1}{2-d} |x|^{2-d} \approx \log|x| \quad d=2$$

Dimensional analysis

$$\approx |x|^{-1} \quad d=3$$

$$\approx |x|^{-2} \quad d=4$$

etc...

Long distances $|x-y| \gg \frac{1}{m}$

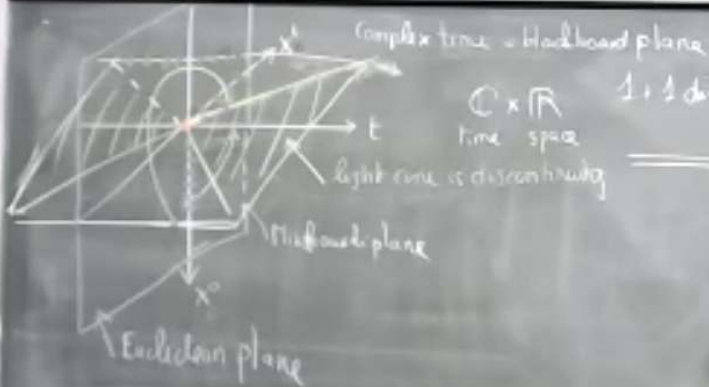
$$G_0(x) \approx \exp(-m|x|) \quad \text{exponential decay} \quad c=2$$

Euclidean



Euclidean
1+1 dim

Minkowski (1+1)



1+1 dim

$\mathbb{C} \times \mathbb{R}$
time space

Minkowski plane

Euclidean plane

causality
discontinuity

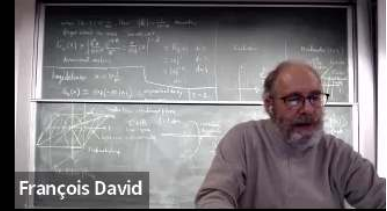
$$G_0(\text{black}) - G_0(\text{white}) \approx G_{\text{extended}}$$

$$\langle \phi(x) \phi(y) \rangle \xrightarrow{x \sim y} \langle \phi^2(x) \rangle \text{ infinite}$$

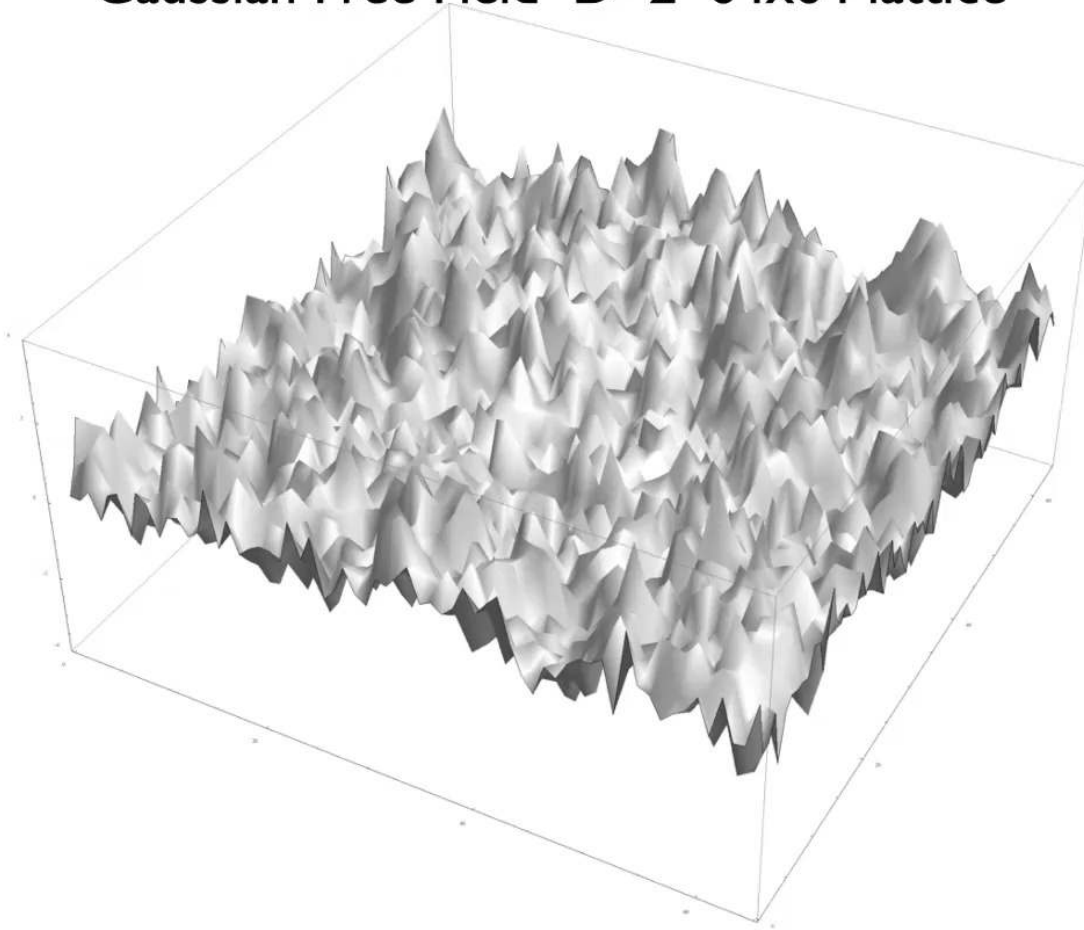
$$Q(t_1) Q(t_2) \xrightarrow{t_1 \rightarrow t_2} Q^2(t) \text{ finite}$$

Typical GFF configurations

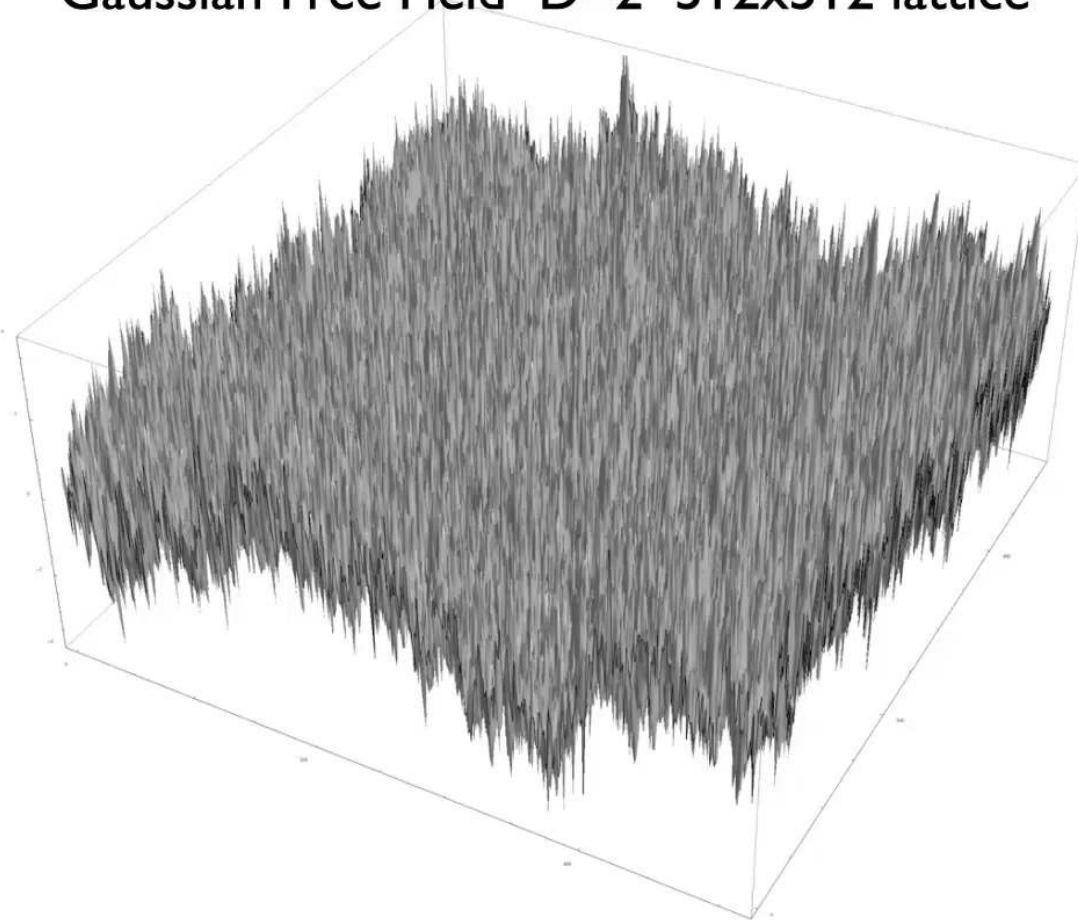
Gaussian Free Field = Massless Scalar Free Fields



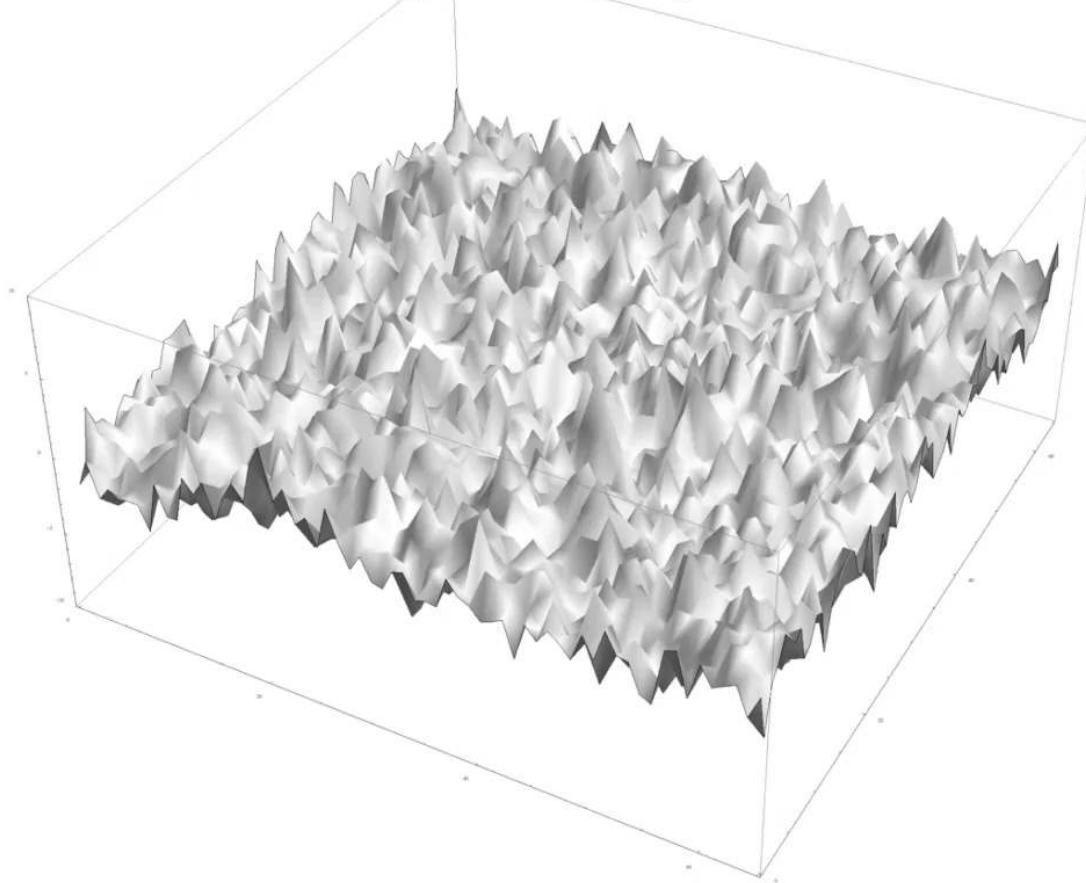
Gaussian Free Field $D=2$ 64x64 lattice



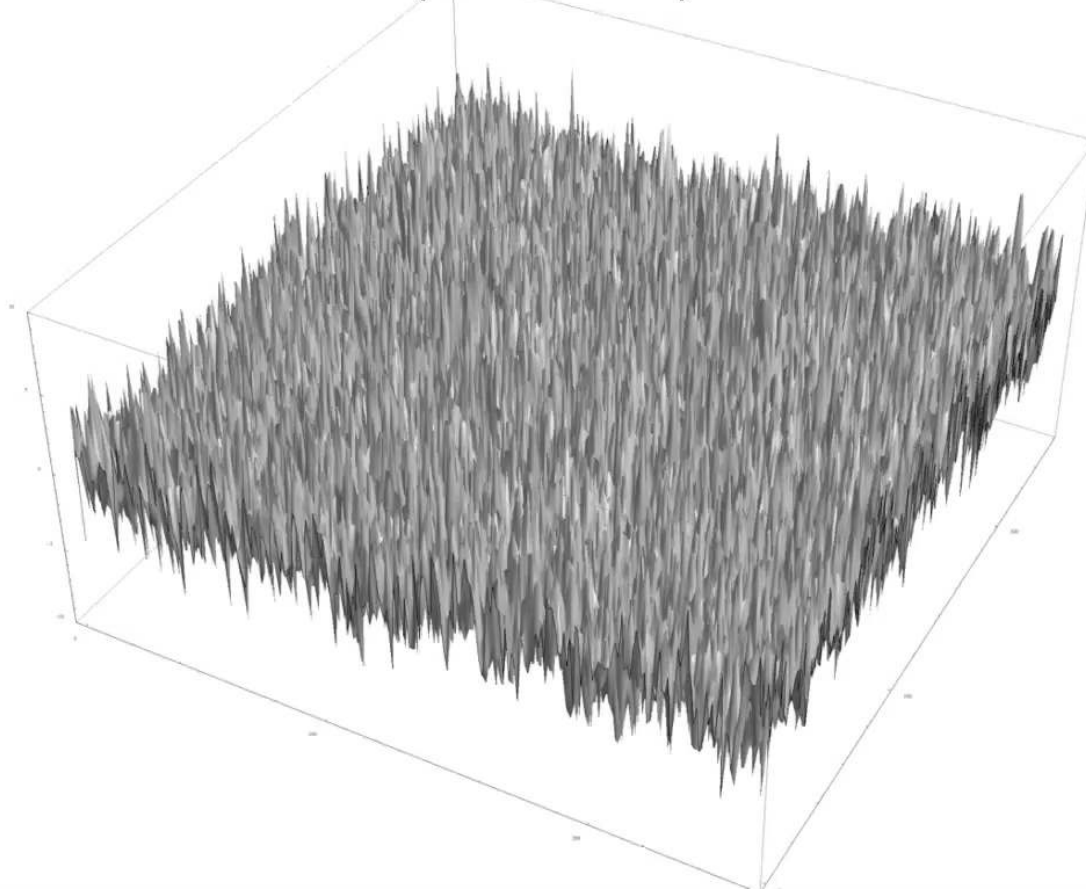
Gaussian Free Field $D=2$ 512x512 lattice



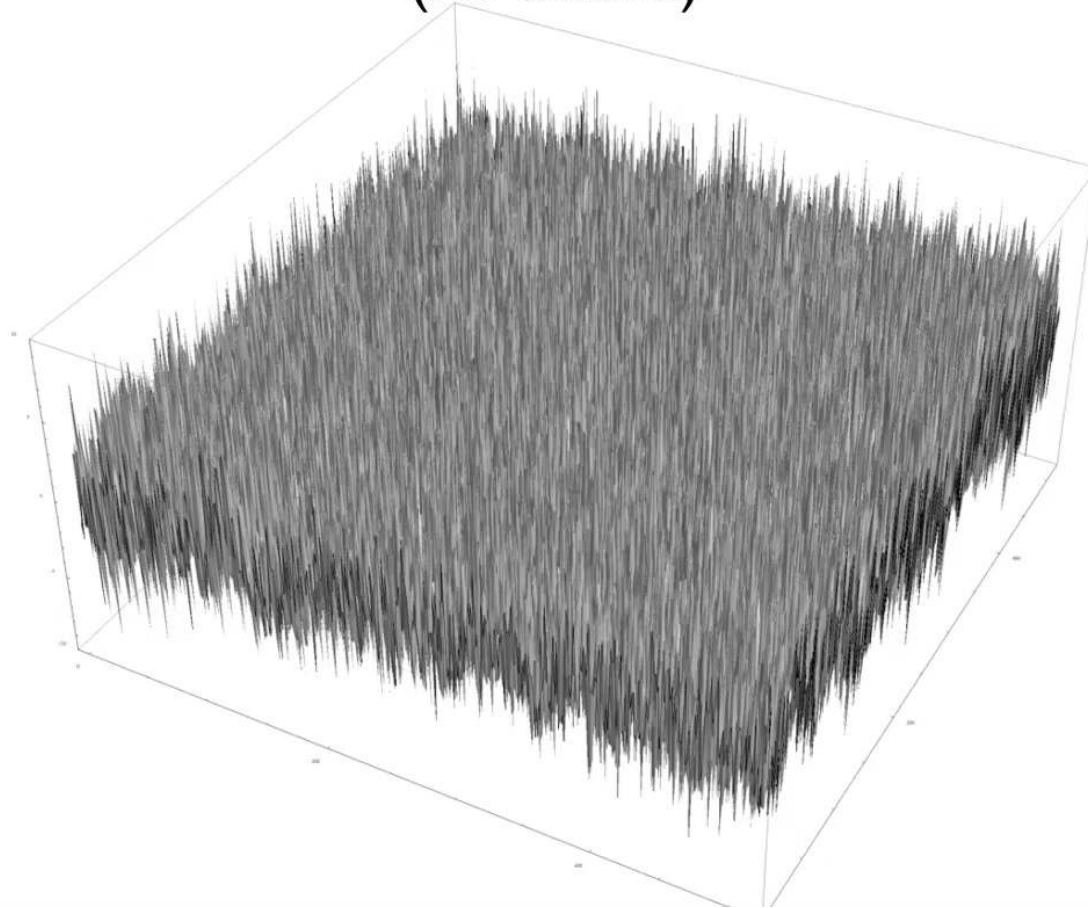
Gaussian Free Field $D=3$ $64 \times 64 \times 64$ lattice (2D section)



Gaussian Free Field $D=3$ 256x256x256 lattice (2D section)

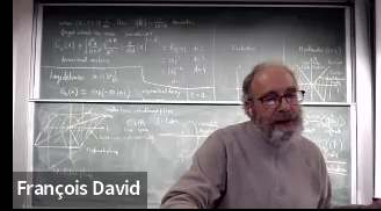
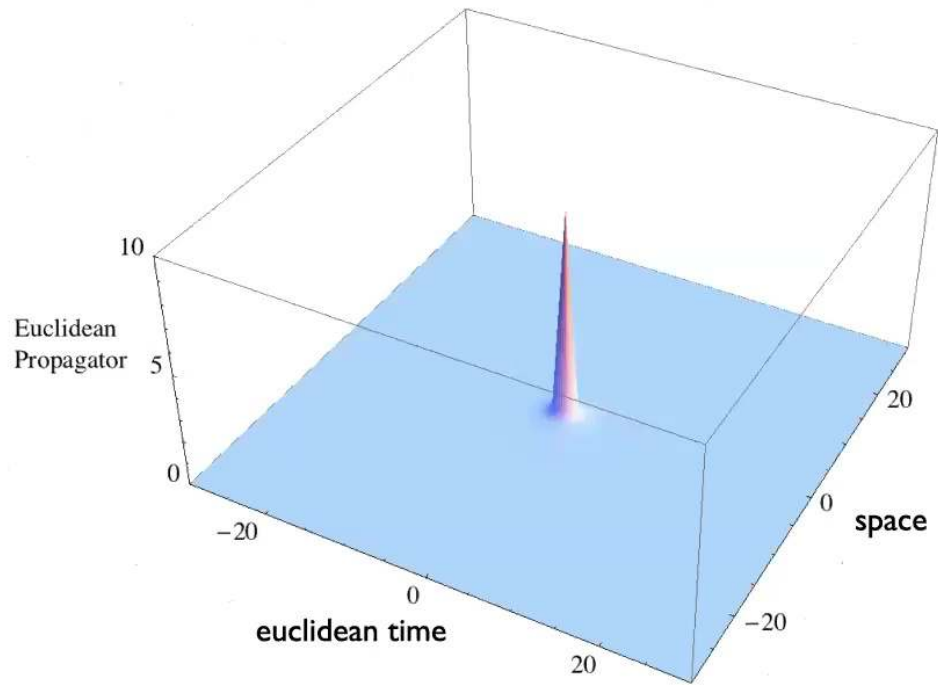


Gaussian Free Field $D=3$ $512 \times 512 \times 512$ lattice (2D section)

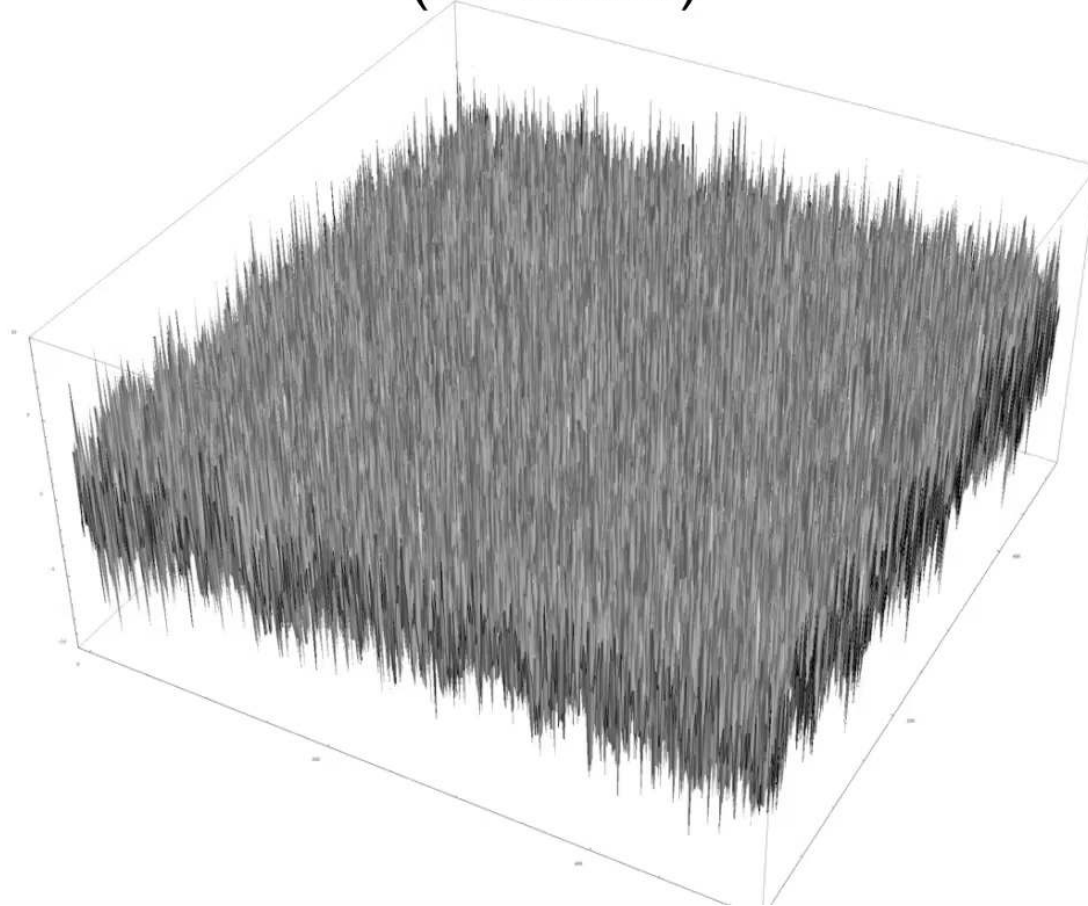


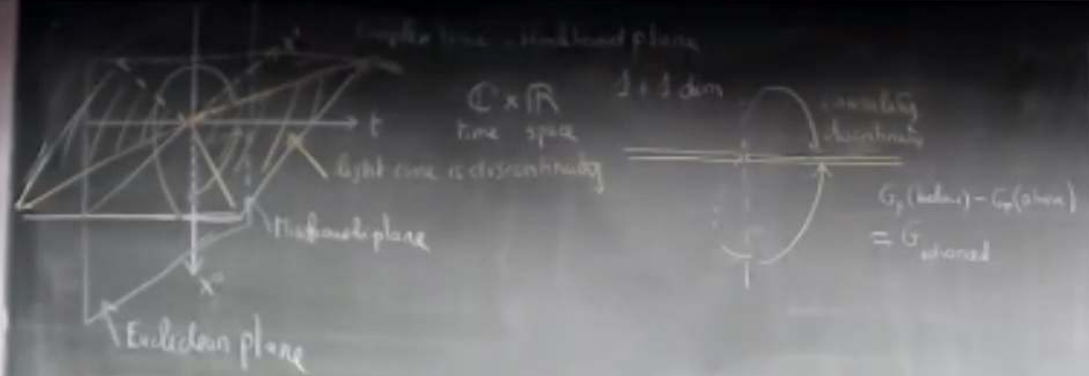
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Free field mass m , dimension $D=2$ propagator in Euclidean space



Gaussian Free Field $D=3$ $512 \times 512 \times 512$ lattice (2D section)





$$\langle \phi(x) \phi(y) \rangle \rightarrow \langle \phi^2(x) \rangle$$

well defined $x \neq y$ infinite

$$Q(t_1) Q(t_2) \rightarrow Q^2(t)$$

$\mathbb{Q}M$ $t_1 \neq t_2$ finite

Free Field is a random distribution

GFF (not continuous function)

Operator $\phi(x)$: naive product $\phi(x) = \phi(x) \cdot \phi(x)$ does not exist
 but the normal product operator

$$:\phi^2(x): = \lim_{x_1, x_2 \rightarrow x} \left[\phi(x_1) \phi(x_2) - \langle \phi(x_1) \phi(x_2) \rangle \cdot \mathbb{1}(x) \right]$$

exists

well defined operator located at X

identity operator

Wick Theorem: valid

$$\langle \phi(x_1) \phi(x_2) \phi(x_3) \rangle$$

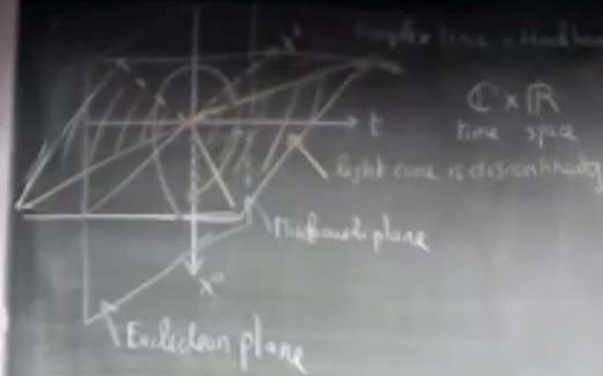
$$= \langle \phi(x_1) \cdot \phi(x_2) \phi(x_3) \rangle$$

$$\langle \phi(x_1) - \phi(x_2) \cdot \phi(x_3) \rangle$$

this is finite



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$\mathbb{C} \times \mathbb{R}$ time space
 $d = 3 \text{ dim}$
 light cone is discontinuity
 causality discontinuity
 $G_p(\text{below}) - G_p(\text{above}) = G_{\text{retarded}}$

$$\langle \phi(x) \phi(y) \rangle \xrightarrow{x \neq y} \langle \phi(x) \rangle$$

well defined

$$Q(t_1) Q(t_2) \xrightarrow{t_1 \rightarrow t_2} Q^2(t)$$

ΨM finite

Free Field is a random distribution
 GFF (not continuous function)

Operator $\phi(x)$? naive product
 but the normal product operator

$$:\phi^2(x): = \lim_{x_i \rightarrow x} \langle \phi(x_i) \phi(x_j) \rangle$$

well defined operator

does not exist
 identity operator

$$\langle \phi(x_i) \phi(x_j) \rangle \cdot \mathbb{1}(x)$$

$$\langle \phi(x_i) \phi(x_j) \phi(x_{i+j}) \rangle = \langle \phi(x_i) \phi(x_j) \rangle$$

$$\langle \phi(x_i) \phi(x_j) : \phi^2(x_{i+j}) : \rangle$$

this is finite

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Complex time = Minkowski plane

$\mathbb{C} \times \mathbb{R}$ time space
 1+1 dim
 light cone is discontinuity
 Minkowski plane
 Euclidean plane
 causality discontinuity
 $G_+(x, y) - G_+(y, x) = G_{\text{normal}}$
 Gaussian random variable
 $\Phi = (\phi_a)_{a \in \mathbb{N}}$
 $w(\Phi) = \exp\left(-\frac{1}{2} \sum_{a,b} \phi_a \phi_b K^{ab}\right)$
 variance of probability theory
 $\langle \phi_a \phi_b \rangle = (K^{-1})_{ab}$
 $E[\phi_a \phi_b]$

→ Generating functional Φ
 Large deviation theory M

$\langle \phi(x) \phi(y) \rangle \rightarrow \langle \phi^2(x) \rangle$
 well defined $x \neq y$ infinite

$Q(t_1) Q(t_2) \rightarrow Q^2(t)$
 $\mathbb{Q}M$ $t_1 \rightarrow t_2$ finite

Free Field is a random distribution
 GFF (not continuous function)

$E[\phi_a \phi_b \phi_c \phi_d] = \sum_{\text{pairs}} E[\phi \phi] E[\phi \phi]$

exists
 defined operator located at X

Wick Theorem: valid

$\langle \phi(x_1) \phi(x_2) \phi^2(x_{n+1}) \rangle$
 this is finite

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