

Title: PSI Lecture - Condensed Matter - Lecture 9

Speakers: Aaron Szasz

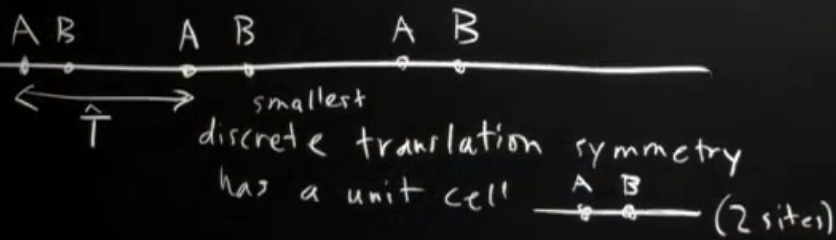
Collection: PSI Lecture - Condensed Matter

Date: January 27, 2022 - 9:00 AM

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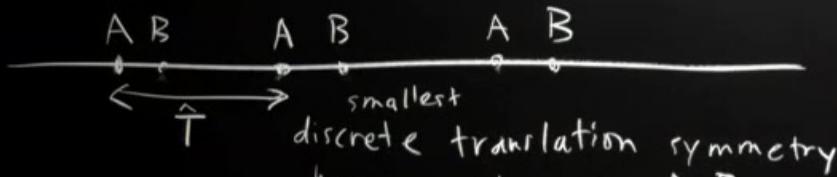
Today! Multiple sites per unit cell

Start with simplest example: 1D



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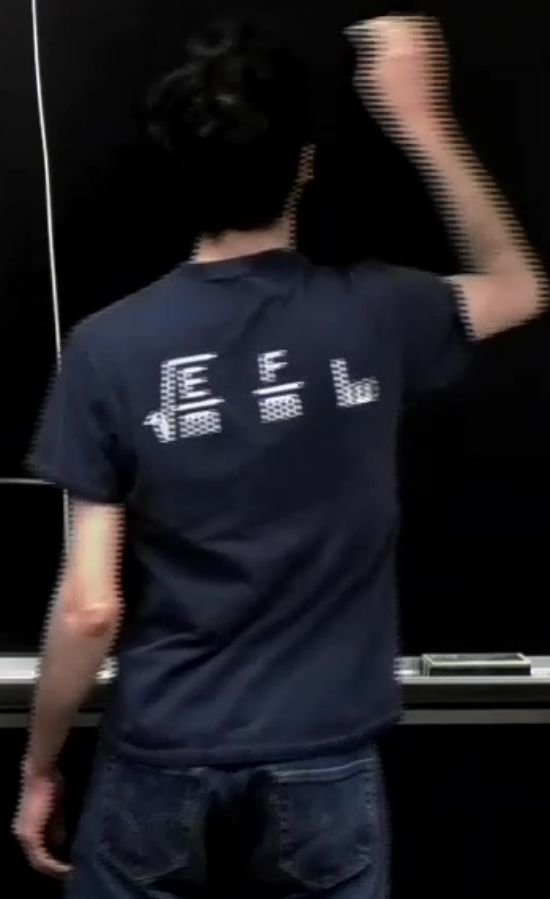


Probably "t" has a unit cell (2 sites)

is larger for $\frac{A B}{t_1}$ than for $\frac{B A}{t_2}$

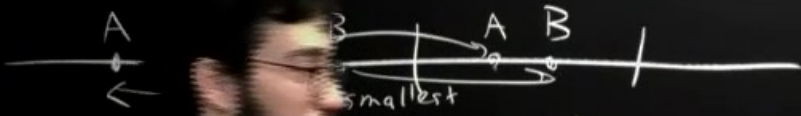
$$H = - \sum_j^{\text{unit cell}} t_1 (c_{jA}^\dagger c_{jB} + c_{jB}^\dagger c_{jA}) + t_2 (c_{j-1,B}^\dagger c_{jA} + c_{j+1,A}^\dagger c_{jB})$$

To solve (find single-



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Start with simplest example: 1D

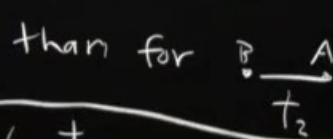


crete translation symmetry

has a unit cell $\overline{A \ B}$ (2 sites)

Probab

is larger for



than for

$$H = t_1 (c_{jB}^\dagger + c_{jB}^\dagger c_{jA}) + t_2 (c_{j-1,B}^\dagger c_{jA} + c_{j+1,A}^\dagger c_{jB})$$

To solve (find single-particle eigenstates)

Use \hat{T} , since $[H, T] = 0$
find eig. states of \hat{T}

$$\hat{T} c_{jA} = c_{j+1,A}$$

$$\hat{T} c_{jB} = c_{j+1,B}$$

$$H = - \sum_j^{\text{unit cell}} t_1 (c_{jA}^\dagger c_{jB} + c_{jB}^\dagger c_{jA}) + t_2 (c_{j-1,B}^\dagger c_{jA} + c_{j+1,A}^\dagger c_{jB})$$

Write an eigenstate as the linear combination Conclusion

$$a^\dagger = \sum_j a_{jA} c_{jA}^\dagger + a_{jB} c_{jB}^\dagger, \quad \hat{T} a^\dagger |0\rangle = \lambda a^\dagger |0\rangle$$

$$\hat{T} a^\dagger = \sum_j a_{jA} c_{j+1,A}^\dagger + a_{jB} c_{j+1,B}^\dagger$$

$$= \sum_j a_{j+1,A} c_{jA}^\dagger + a_{j-1,B} c_{jB}^\dagger = \sum_j \lambda a_{jA} c_{jA}^\dagger + \lambda a_{jB} c_{jB}^\dagger$$

equate coeffs: $a_{j-1,A} = \lambda a_{jA}$

$a_{j-1,B} = \lambda a_{jB}$

$\hat{T}^N = 1$ ← (# of unit cells)

$$\Rightarrow \lambda^N = 1 \Rightarrow \lambda = e^{i \frac{2\pi}{N} n}, \quad n = 0, \dots, N-1$$

$$H = - \sum_j^{\text{unit cell}} (c_{jA}^\dagger c_{jB} + c_{jB}^\dagger c_{jA}) + t_2 (c_{j-1,B}^\dagger c_{jA} + c_{j+1,A}^\dagger c_{jB})$$

write an eigenstate as the linear combination

$$a^\dagger = \sum_j a_{jA} c_{jA}^\dagger + a_{jB} c_{jB}^\dagger, \quad T a^\dagger = \lambda a^\dagger$$

$$\begin{aligned} T a^\dagger &= \sum_j a_{jA} c_{j+1,A}^\dagger + a_{jB} c_{j+1,B}^\dagger \\ &= \sum_j a_{j+1,A} c_{jA}^\dagger + a_{j-1,B} c_{jB}^\dagger = \sum_j \lambda a_{jA} c_{jA}^\dagger + \lambda a_{jB} c_{jB}^\dagger \end{aligned}$$

equate coeffs: $a_{j-1,A} = \lambda a_{jA}$
 $a_{j-1,B} = \lambda a_{jB}$

(# of unit cells)

$$\lambda^N = 1 \Rightarrow \lambda = e^{i \frac{2\pi}{N} n}, \quad n=0, \dots, N-1$$

Conclusion: T eig states look like

$$a_{\hbar\alpha}^\dagger = a_{0A} \sum_j e^{ikaj} c_{jA}^\dagger + a_{0B} \sum_j e^{ikaj} c_{jB}^\dagger$$

$$c_{\hbar\alpha}^\dagger = \frac{1}{\sqrt{N}} \sum_j e^{ikaj} c_{j\alpha}^\dagger, \quad \alpha=A,B$$

$$a_{\hbar\alpha}^\dagger = a_{\hbar A} c_{\hbar A}^\dagger + a_{\hbar B} c_{\hbar B}^\dagger$$

$$|a_{\hbar A}|^2 + |a_{\hbar B}|^2 = 1$$

$$H = - \sum_j^{\text{unit cell}} (c_{jA}^+ c_{jB}^+ + c_{jB}^+ c_{jA}^+) + \frac{t_2}{2} (c_{j-1,B}^+ c_{jA}^+ + c_{j+1,A}^+ c_{jB}^+)$$

$$a^+ = \sum_j a_{jA}^+ c_{jA}^+ + a_{jB}^+ c_{jB}^+, \quad T a^+ = \lambda a^+ |0\rangle$$

$$\hat{T} a^+ = \sum_j a_{jA}^+ c_{j+1,A}^+ + a_{jB}^+ c_{j+1,B}^+$$

$$= \sum_j a_{j+1,A}^+ c_{jA}^+ + a_{j-1,B}^+ c_{jB}^+ = \sum_j \lambda a_{jA}^+ c_{jA}^+ + \lambda a_{jB}^+ c_{jB}^+$$

equating coeffs: $a_{j-1,A} = \lambda a_{jA}$
 $a_{j-1,B} = \lambda a_{jB}$

$$\hat{T}^N = 1 \quad (\text{\# of unit cells}) \Rightarrow \lambda^N = 1 \Rightarrow \lambda = e^{i \frac{2\pi}{N} n}, \quad n=0, \dots, N-1$$

like

$$a_n^+ = a_{0A} \sum_j e^{ik_{nj}} c_{jA}^+ + a_{0B} \sum_j e^{ik_{nj}} c_{jB}^+$$

$$c_{k\alpha} = \frac{1}{\sqrt{N}} \sum_j e^{ik_{nj}} c_{j\alpha}^+, \quad \alpha=A,B$$

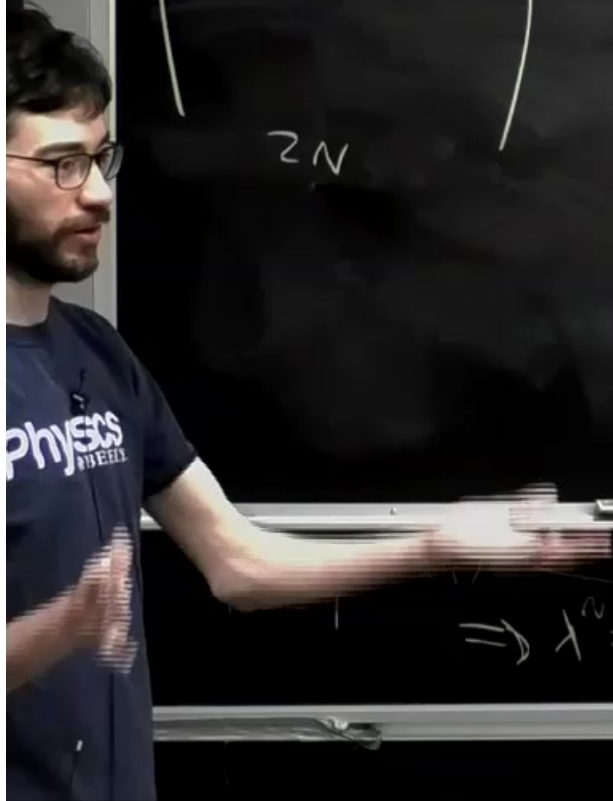
$$|a_n^+ = a_{nA} c_{nA}^+ + a_{nB} c_{nB}^+|$$

$$|a_{nA}|^2 + |a_{nB}|^2 = 1$$

$$\langle 0 | a_{k\alpha} a_n^+ | 0 \rangle = \| a_n^+ | 0 \rangle \|^2$$

Perimeter-A

CAUTION
 DO NOT TOUCH THE SURFACE OF THE BOARD
 IN A LABORATORY OR OTHER
 HIGH-VOLTAGE AREA
 YOUR PROPERTY RISK



$2N$

$\begin{pmatrix} \dagger |0\rangle \\ \dagger |0\rangle \\ \vdots \\ \dagger |N-1A\rangle \\ \dagger |N-1B\rangle \end{pmatrix}$

$2N$

\hat{T} has N eig. values, $\lambda = e^{i\frac{2\pi}{N}n}$
 \Rightarrow each eig. val has a 2-dimensional eigenspace with basis

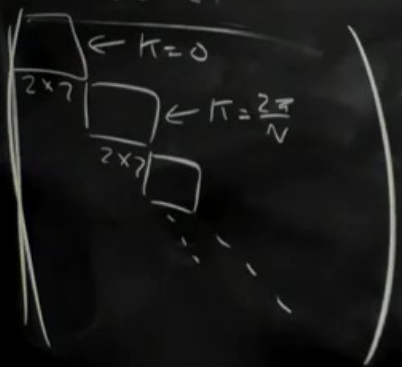
$\begin{pmatrix} \dagger |nA\rangle \\ \dagger |nB\rangle \end{pmatrix}, \begin{pmatrix} \dagger |nA\rangle \\ \dagger |nB\rangle \end{pmatrix}$

For each n , any linear comb is a trans. eig. state, and I could choose to use the same 2 vectors for each n or 2 different ones

$\Rightarrow \lambda^N = 1 \Rightarrow \lambda = e^{i\frac{2\pi}{N}n}, n=0, \dots, N-1$



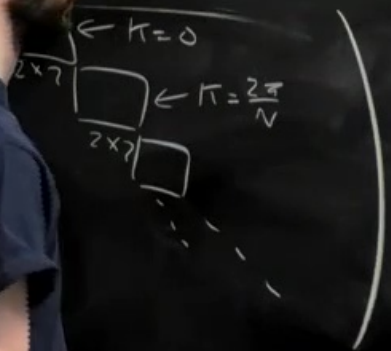
We found eig vectors of T ,
 since $[H, T] = 0$, this divides
 H into 2-dimensional subspaces
 i.e. Block diagonal, with blocks of
 size 2:



What we still need to do:
 diagonalize each block, this
 will fix " a_{kA} " and " a_{kB} " in order for
 a_k^\dagger to be an eig vector

$$\sqrt{\frac{E}{m}} \quad \frac{F}{m} \quad \omega$$

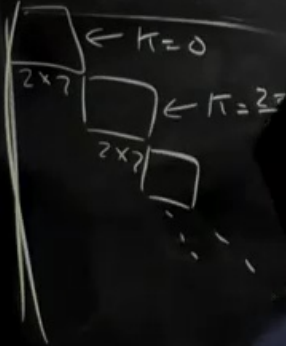
We found eig vectors of T ,
 since $[H, \hat{T}] = 0$, this divides
 to 2-dimensional subspaces
 block diagonal, with blocks of
 size 2:



What we still need to do:
 diagonalize each 2×2 block, this
 will fix " a_{kA} " and " a_{kB} " in order for
 a_k^\dagger to be an eig state not just of \hat{T} but
 also H



Since $[H, T] = 0$, this divides H into 2-dimensional subspaces
 Block diagonal, with blocks of size 2:



diagonalize each 2×2 block, this will fix a_k^+ and a_k^- in order for a_k^+ to be an eig state not just of T but also H

Our approach: substitute for c_j^+ in terms of c_k^+ :

$$c_{j\alpha}^+ = \frac{1}{\sqrt{N}} \sum_j e^{-ikaj} c_{j\alpha}^+, \alpha = A, B$$

$\hat{T}^N = 1$

$R = \frac{1}{2}(a_{j+1} + a_{j-1})$

$\lambda = e^{i \frac{2\pi}{N} n}, n = 0, \dots, N-1$

$|a_{kA}|^2 + |a_{kB}|^2 = 1$

$\langle 0 | a_k a_k^+ | 0 \rangle = \langle 0 | a_k^+ | 0 \rangle^2$

of c_k^+ : $c_j^\alpha = \frac{1}{\sqrt{N}} \sum_j e^{ikaj} c_j^\alpha, \alpha=A,B$

$$H = - \sum_{k, k'} \left[t_1 (c_{kA}^+ c_{kB} + c_{kB}^+ c_{kA}) + t_2 (e^{-ika} c_{kA}^+ c_{k'B} + e^{ik'a} c_{k'B}^+ c_{kA}) \right] \left[\frac{1}{N} \sum_j e^{i(k-k')aj} \right]$$

$$= - \sum_k \underbrace{\begin{pmatrix} c_{kA}^+ & c_{kB}^+ \end{pmatrix}}_{\vec{c}_k^+} \begin{pmatrix} 0 & t_1 + t_2 e^{-ika} \\ t_1 + t_2 e^{ika} & 0 \end{pmatrix} \begin{pmatrix} c_{kA} \\ c_{kB} \end{pmatrix}$$

2x2 blocks from H
 $\equiv H_k$

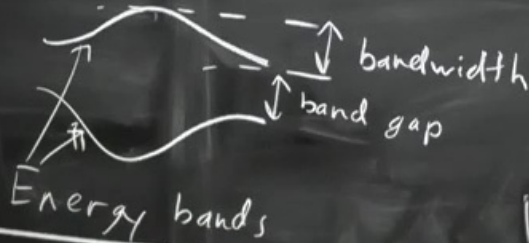
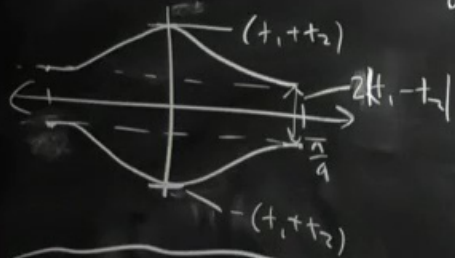
$E(k)$ as eig. valr of H_k

$$E(k) = \pm \sqrt{(t_1 + t_2 e^{ika})(t_1 + t_2 e^{-ika})}$$

$$= \pm \sqrt{t_1^2 + t_2^2 + 2t_1 t_2 \cos(ka)}$$

Plot $E(k)$

(a is the size of one unit cell)

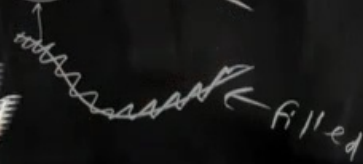


orbitals
~~2 sites~~ per cell
 \Downarrow
 2×2 blocks for each k
 \Downarrow
 2 energy bands

 #bands = # orbitals per unit cell,
 b/c every orbital has its own c_j
 \uparrow orbital index

Metals vs Insulators

← empty



insulator b/c
 E must supply
 energy = band gap

← partially filled band



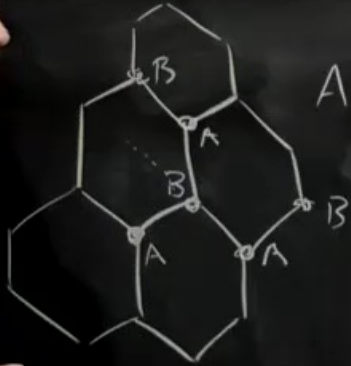
partially filled band
 → metal

insulator b/c filled \vec{E} must supply energy = band gap

partially filled band \rightarrow metal

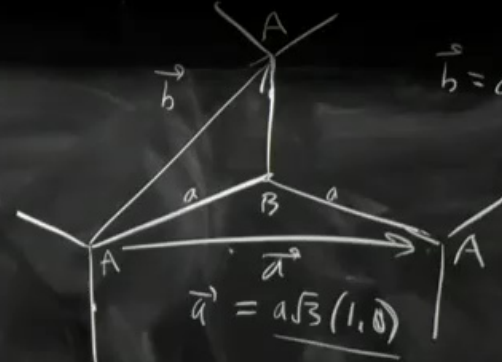
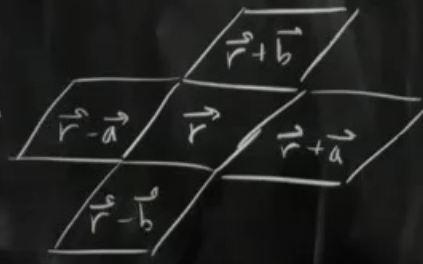
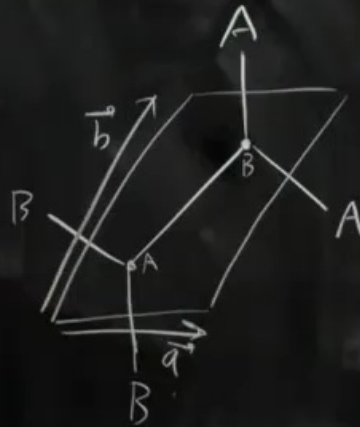
orbital index

In 2D: Honeycomb lattice



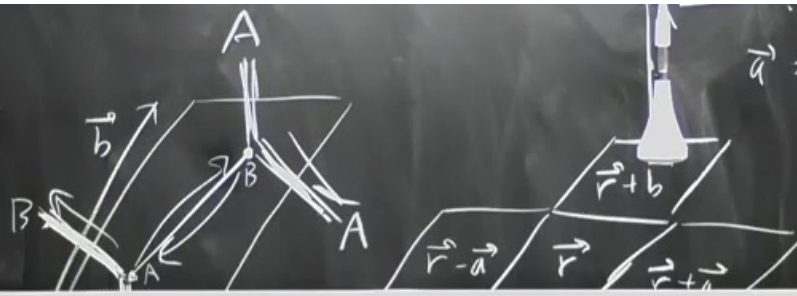
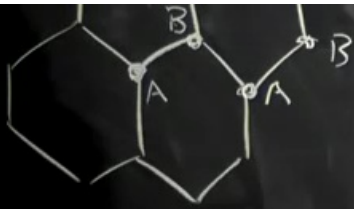
A, B distinct

Translation vectors.



$$\vec{b} = a\sqrt{3} \left(\frac{1}{2}, \frac{\sqrt{3}}{2} \right)$$

$$\vec{a} = a\sqrt{3} (1, 0)$$



$$\vec{a} = a\sqrt{3}(1, 0)$$

$$H = -t \sum_{\vec{r}} \left[c_{\vec{r}-\vec{a}, B}^\dagger c_{\vec{r}, A} + c_{\vec{r}-\vec{b}, B}^\dagger c_{\vec{r}, A} + c_{\vec{r}, A}^\dagger c_{\vec{r}, B} + c_{\vec{r}, B}^\dagger c_{\vec{r}, A} \right. \\ \left. + c_{\vec{r}+\vec{a}, A}^\dagger c_{\vec{r}, B} + c_{\vec{r}+\vec{b}, A}^\dagger c_{\vec{r}, B} \right]$$

Substitute: $c_{\vec{r}, \alpha} = \frac{1}{\sqrt{N}} \sum_{\vec{k}} e^{-i\vec{k} \cdot \vec{r}} c_{\vec{k}, \alpha}$

$$H = -t \sum_{\vec{k}} e^{-i\vec{k} \cdot \vec{a}} \left[c_{\vec{k}, B}^\dagger c_{\vec{k}, A} + e^{-i\vec{k} \cdot \vec{b}} \left[c_{\vec{k}, B}^\dagger c_{\vec{k}, A} + c_{\vec{k}, B}^\dagger c_{\vec{k}, A} + c_{\vec{k}, A}^\dagger c_{\vec{k}, B} \right] \right. \\ \left. + e^{i\vec{k} \cdot \vec{a}} \left[c_{\vec{k}, A}^\dagger c_{\vec{k}, B} + e^{i\vec{k} \cdot \vec{b}} \left[c_{\vec{k}, A}^\dagger c_{\vec{k}, B} + c_{\vec{k}, A}^\dagger c_{\vec{k}, B} \right] \right] \right]$$



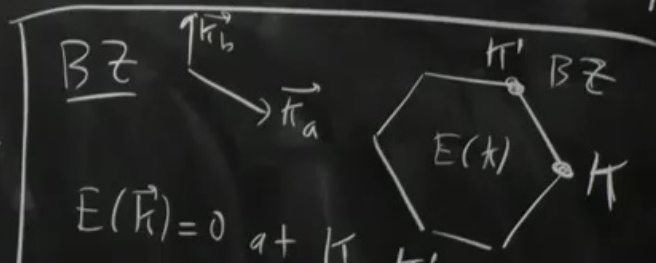
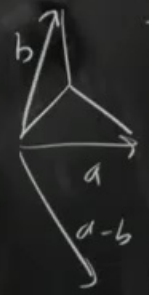
$$H = -t \sum_{\vec{r}} e^{-i\vec{k} \cdot \vec{a}} c_{\vec{r}B} c_{\vec{r}A} + e^{-i\vec{k} \cdot \vec{b}} c_{\vec{r}B} c_{\vec{r}A} + e^{i\vec{k} \cdot \vec{a}} c_{\vec{r}A} c_{\vec{r}B} + e^{i\vec{k} \cdot \vec{b}} c_{\vec{r}A} c_{\vec{r}B}$$

$$H = -t \sum_{\vec{k}} \vec{c}_{\vec{k}}^\dagger \cdot H_{\vec{k}} \cdot \vec{c}_{\vec{k}}$$

$$H_{\vec{k}} = \begin{pmatrix} 0 & 1 + e^{i\vec{k} \cdot \vec{a}} + e^{-i\vec{k} \cdot \vec{b}} \\ 1 + e^{i\vec{k} \cdot \vec{a}} + e^{-i\vec{k} \cdot \vec{b}} & 0 \end{pmatrix}$$

Diagonalize & simplify

$$E(\vec{k}) = \pm \sqrt{3 + 2(\cos(\vec{k} \cdot \vec{a}) + \cos(\vec{k} \cdot \vec{b}) + \cos(\vec{k} \cdot (\vec{a} - \vec{b})))}$$



$E(\vec{k}) = 0$ at K, K' points
 @ $K, \vec{k} = \frac{2}{3}\vec{k}_a + \frac{1}{3}\vec{k}_b$
 @ $K', \vec{k} = \frac{1}{3}\vec{k}_a + \frac{2}{3}\vec{k}_b$

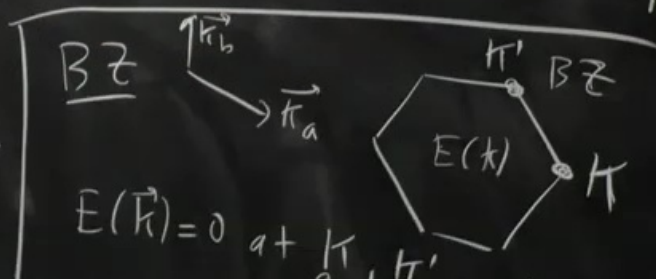
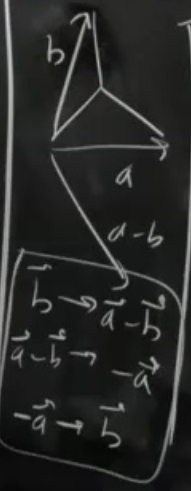
$$\vec{k} \cdot \vec{a} + e^{-i\vec{k} \cdot \vec{b}} + \vec{k} \cdot \vec{b} + e^{i\vec{k} \cdot \vec{a}} + \vec{k} \cdot (\vec{a} - \vec{b}) + e^{i\vec{k} \cdot \vec{a}} + e^{-i\vec{k} \cdot \vec{b}}$$

$$H = -t \sum_{\vec{k}} \vec{c}_{\vec{k}}^\dagger H_{\vec{k}} \vec{c}_{\vec{k}}$$

$$H_{\vec{k}} = \begin{pmatrix} 0 & e^{i\vec{k} \cdot \vec{a}} + e^{-i\vec{k} \cdot \vec{b}} \\ t & 0 \end{pmatrix}$$

Diagonalize & simplify $\rightarrow \pm \sqrt{3 + 2 \left(\sum_{\vec{v}} e^{i\vec{k} \cdot \vec{v}} \right)}$

$$E(\vec{k}) = \pm \sqrt{3 + 2 \left(\cos(\vec{k} \cdot \vec{a}) + \cos(\vec{k} \cdot \vec{b}) + \cos(\vec{k} \cdot (\vec{a} - \vec{b})) \right)}$$



$E(\vec{k}) = 0$ at K, K' points

@ $K, \vec{k} = \frac{2}{3} \vec{k}_a + \frac{1}{3} \vec{k}_b$

@ $K', \vec{k} = \frac{1}{3} \vec{k}_a + \frac{2}{3} \vec{k}_b$