

Title: Aspects of Rotating Black Holes in Dynamical Chern-Simons Gravity

Speakers: Leah Jenks

Series: Cosmology & Gravitation

Date: December 16, 2021 - 1:00 PM

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Abstract: In this talk I will give an overview of recent and ongoing work regarding rotating black holes in dynamical Chern-Simons (dCS) gravity. dCS gravity is a well motivated modified theory of gravity which has been extensively studied in gravitational and cosmological contexts. I will first discuss unique geometric structures, 'the Chern-Simons caps,' which slowly rotating black holes in dCS gravity were recently found to possess. Motivated by the dCS caps, I will then discuss superradiance in the context of slowly rotating dCS black holes and show that there are corrections to the usual solution for a Kerr black hole. Lastly, I will comment on the observable implications for these corrections and point towards avenues for future work.

Zoom Link: <https://pitp.zoom.us/j/95228483630?pwd=dWk1c3p5dUU3RXJrNEhIT2M3Tk1Kdz09>



Leah Jenks

Aspects of Rotating Black Holes in Dynamical Chern-Simons Gravity

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arXiv:2104.00019 (PRD), arXiv:2201.XXXX
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BROWN Perimeter Institute Strong Gravity & Cosmology Seminar
December 16, 2021

Overview and Main Takeaways

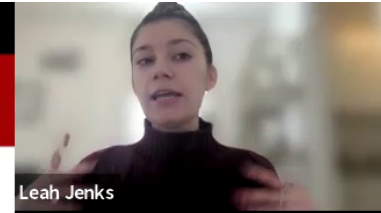
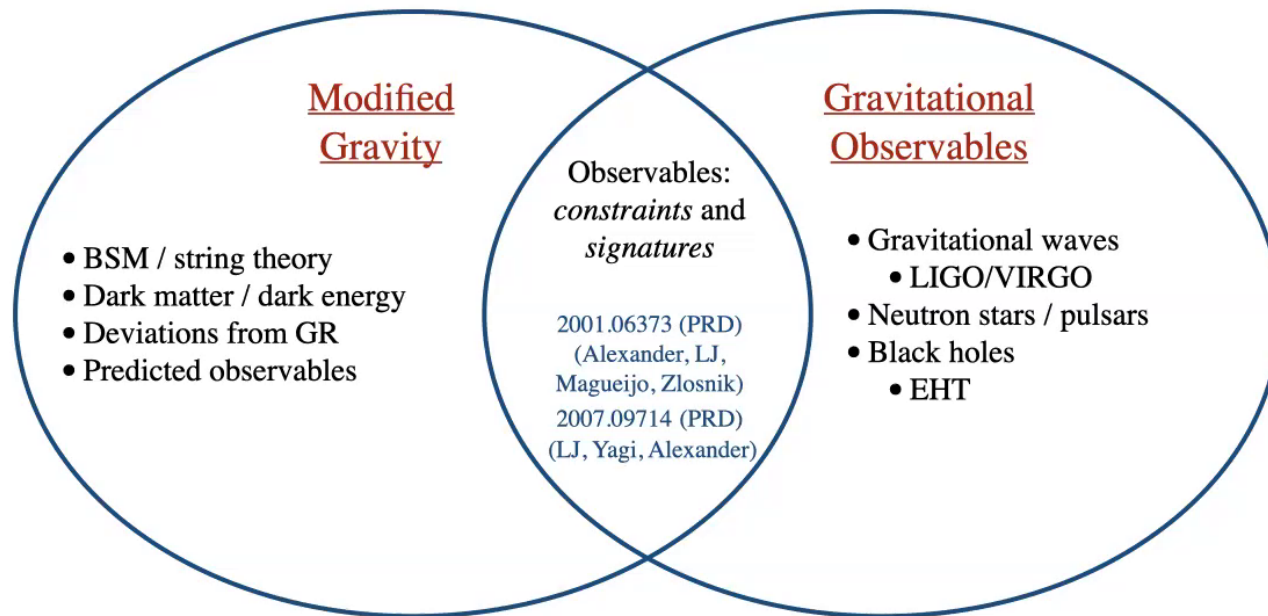


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Main Takeaways:

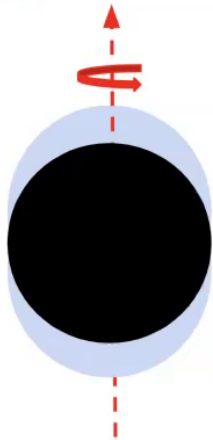
- 1. Dynamical Chern-Simons gravity (dCS) is a well-motivated theory with promising cosmological and gravitational implications**
- 2. dCS black holes possess unique geometric structures — the ‘dCS caps’ which may provide novel insights into deviations from GR**
- 3. Superradiance induced by dCS black holes deviates from GR and has potential observable consequences**

Big Picture

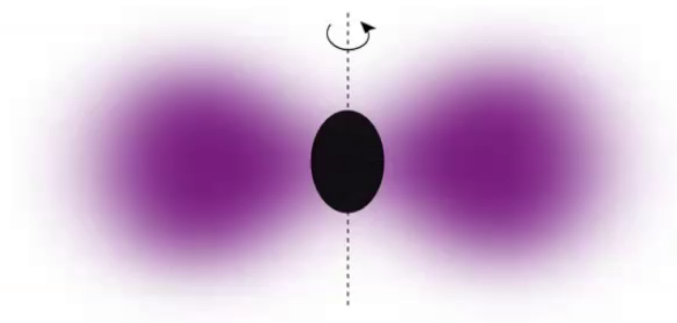


Essential Physics

dCS caps: Chern-Simons
pseudo scalar interaction
induces non-standard
behavior - analogy to
angular momentum



Superradiance: unstable
modes in a radiative system
- analogy to driven
resonance in harmonic
oscillator



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dCS Gravity: Motivations

- Chern-Simons term arises in string theory and effective field theory
 - Green-Schwarz mechanism (Green & Schwarz, 1984)
 - EFT of inflation (Weinberg, 2008)
- dCS gravity intrinsically includes an axion-like particle
- Cosmological, gravitational, and particle physics implications



dCS Gravity: Action

Jackiw & Pi 2003,
Alexander & Yunes 2007+

$$S_{\text{vac}} = \int d^4x \sqrt{-g} \left[\kappa R + \frac{\alpha}{4} \vartheta *R R - \frac{1}{2} (\nabla_a \vartheta) (\nabla^a \vartheta) \right]$$

Modified Einstein Eq.

$$G_{ab} + \frac{\alpha}{\kappa} C_{ab} = \frac{1}{2\kappa} T_{ab}$$

$$C^{ab} = (\nabla_c \vartheta) \epsilon^{cde(a} \nabla_e R^{b)}_d + (\nabla_c \nabla_d \vartheta) *R^{d(ab)c}$$

Pseudo-scalar EoM

$$\square \vartheta = -\frac{\alpha}{4\kappa} *R R$$



dCS Gravity: Effective Field Theory

We can write the dCS action as:

$$S = \int d^4x \sqrt{-g} \left[\kappa R + \frac{\sigma}{4\mu} {}^*R R - \frac{1}{2} (\nabla_a \sigma) (\nabla^a \sigma) \right]$$

In principle, we have an infinite sum of higher dimensional terms

- Truncate the series at the first term
- We never reach the scale at which dCS breaks down
- $\mu \ll M_{pl}$, highly unconstrained from observations

dCS Gravity: Slowly Rotating Solution

Yunes & Pretorius 2009,
Yagi, Yunes & Tanaka 2012,
Maselli, et al. 2017

$$ds^2 = ds_K^2 + \frac{5}{4} \zeta M \chi \frac{M^4}{r^4} \left(1 + \frac{12}{7} \frac{M}{r} + \frac{27}{10} \frac{M^2}{r^2} \right) \sin^2 \theta dt d\phi$$
$$\vartheta = \frac{5}{8} \alpha \chi \frac{\cos \theta}{r^2} \left(1 + 2 \frac{M}{r} + \frac{18}{5} \frac{M^2}{r^2} \right)$$

Deviations from Kerr

$$r_{\pm} = r_{\pm, K} \mp \frac{915}{28672} M \zeta \chi^2$$
$$r_{\text{ergo}} = r_{\text{ergo}, K} - \frac{915}{28672} M \zeta \chi^2 \left(1 + \frac{2836}{915} \sin^2 \theta \right)$$



dCS Gravity: Effective Field Theory

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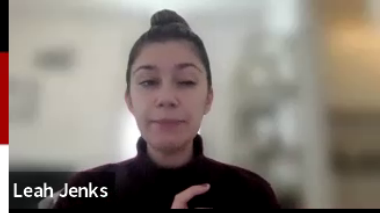
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The Chern-Simons Caps: Motivations



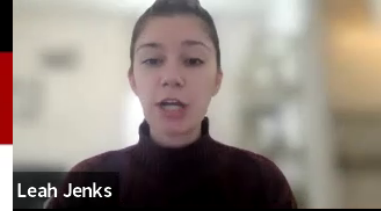
Recall:
$$R_{ab} = 8\pi \bar{T}_{ab} - 16\pi \alpha C_{ab}$$

↗

Geodesic focusing:
$$R_{ab} u^a u^b > 0$$

dCS gravity naturally has the structure to violate the focusing theorem!

The Chern-Simons Caps: Geodesic Focusing

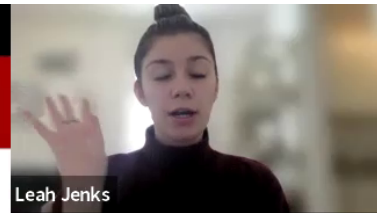


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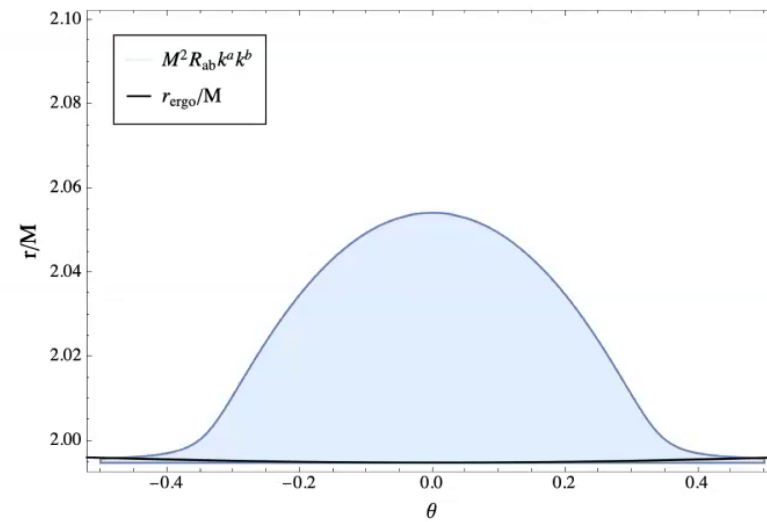
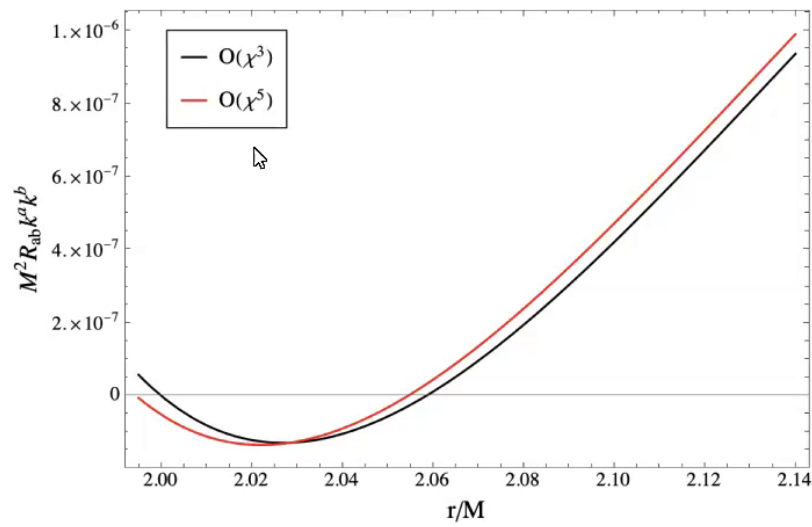
Consider a stationary timelike observer in the slowly rotating dCS spacetime:

$$R_{ab} k_{\text{st}}^a k_{\text{st}}^b = \frac{45}{4} \zeta \chi^2 f \frac{\gamma^8}{M^2} \left[1 + 2c_\theta^2 + \frac{40\gamma}{15} \left(1 + \frac{3}{4} c_\theta^2 \right) \right. \\ \left. + 6\gamma^2 \left(1 + \frac{1}{3} c_\theta^2 \right) - \frac{312}{5} \gamma^3 c_\theta^2 \right] + \mathcal{O}(\zeta \chi^4)$$

The Chern-Simons Caps: Geometry



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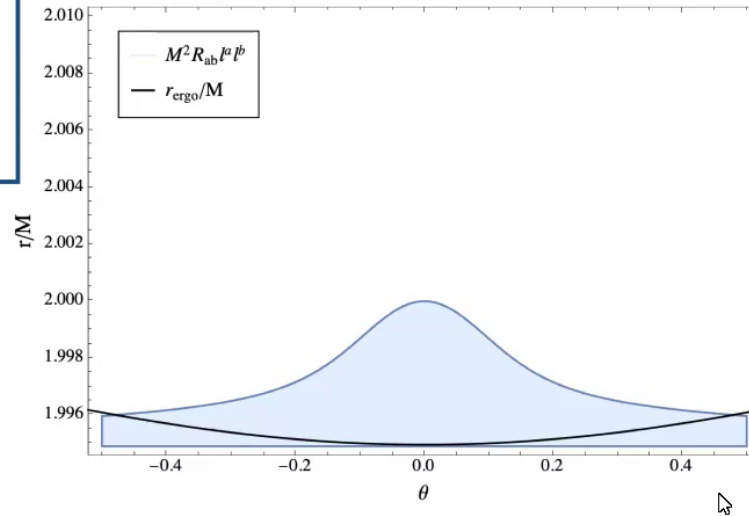
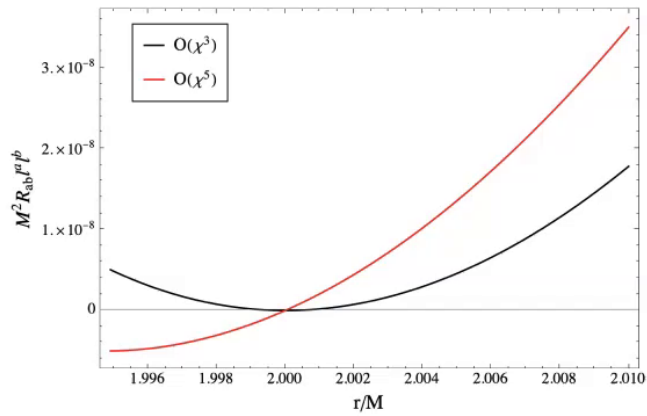


The Chern-Simons Caps: Radial Null Geodesics



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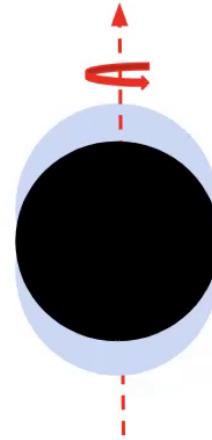
$$R_{ab} l^a l^b = \frac{25}{32} \zeta \chi^2 f \frac{\gamma^6}{M^2} \left[c_\theta^2 + 4\gamma c_\theta^2 + \frac{72}{5} \gamma^2 \left(1 + \frac{43}{24} c_\theta^2 \right) + \frac{192}{5} \gamma^3 \left(1 + \frac{13}{32} c_\theta^2 \right) + \frac{432}{5} \gamma^4 \left(1 - \frac{1}{15} c_\theta^2 \right) - \frac{19872}{25} \gamma^5 c_\theta^2 \right] + \mathcal{O}(\zeta \chi^4),$$



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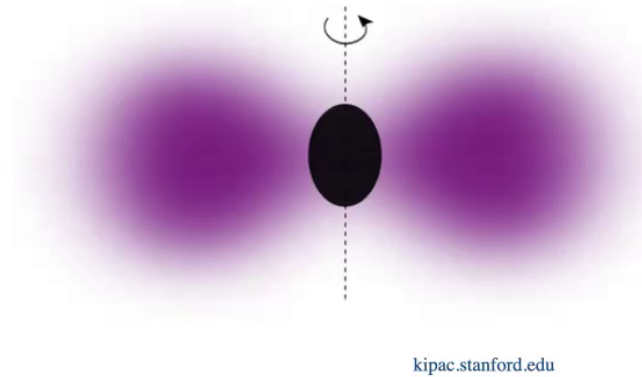
The Chern-Simons Caps: Implications

- Condition for Hawking-Penrose singularity theorem is violated
- Energy condition of the effective energy-momentum tensor
- Behavior of matter close to dCS black holes

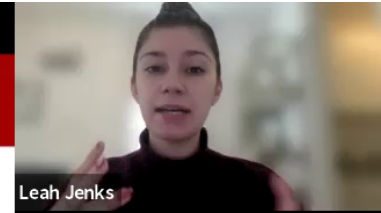


Superradiance in dCS Gravity: Motivations

- Scalar field amplified by rotation of black hole (Black hole bomb - Press & Teukolsky, 1972)
- “String Axiverse” - Arvanitaki et al., 2009 +
- Extended to vectors (eg East & Pretorius, 2017, East 2017 +, Baryakhtar et al., 2017) and tensors (eg Brito, 2020)



Superradiance in GR



Massive scalar field on a Kerr black hole background:

$$(\square_K - \mu^2)\varphi_K(t, r, \theta, \phi) = 0$$

Detweiler limit: $\omega M \ll 1, \mu M \ll 1$ (Detweiler, 1980):

$$\varphi_{\ell, m} = e^{-i\omega_K t} e^{im\phi} S_\ell(\theta) R_\ell(r)$$

Frequency spectrum: $\ell=1, m=1$:

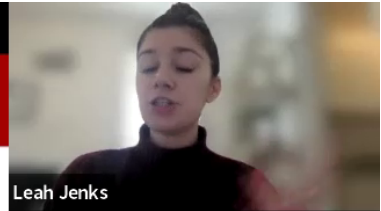
$$\text{Im}(\omega_K) = \mu \left(\frac{a}{M} \right) \frac{(\mu M)^8}{24}$$

Superradiance in dCS Gravity: Setup

Theory: dCS gravity + scalar field

$$S = S_{\text{EH}} + S_{\text{dCS}} + S_{\vartheta} + S_{\varphi}$$

$$S_{\varphi} = \int d^4x \sqrt{-g} \left[-\frac{1}{2} g^{ab} \nabla_a \varphi \nabla_b \varphi - \frac{1}{2} \mu^2 \varphi^2 \right]$$



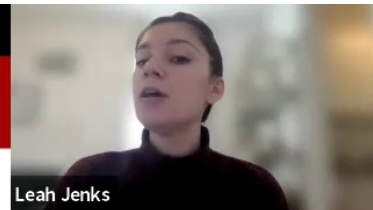
Superradiance in dCS Gravity

Klein-Gordon Equation up to $\mathcal{O}(\zeta\chi^2)$

$$(\square_K + \square_{\text{dCS}} - \mu^2)\varphi = 0$$

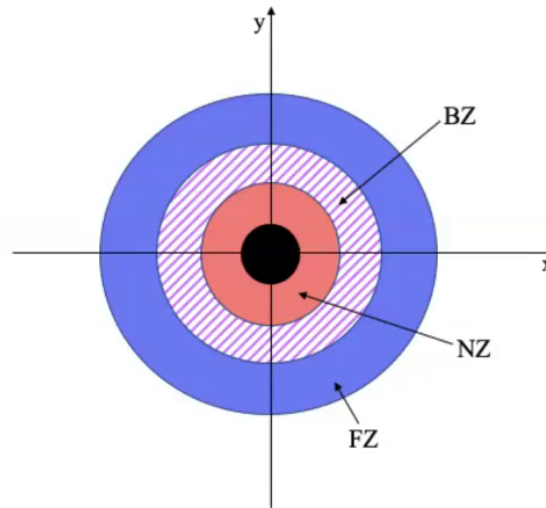
Ansatz with $\omega M \ll 1, \mu M \ll 1$

$$\begin{aligned} \varphi_{\ell,m} \propto & e^{-i\omega_\ell t} e^{im\phi} P_\ell^m(\cos\theta) f_{\ell,m}^0(r) + e^{im\phi} \zeta\chi^2 \left(e^{-i\omega_\ell t} P_\ell^m(\cos\theta) \tilde{f}_{\ell,m}^{\text{dCS}}(r) \right. \\ & \left. + e^{-i\omega_{\ell+2} t} P_{\ell+2}^m(\cos\theta) \tilde{f}_{\ell+2,m}^{\text{dCS}}(r) + e^{-i\omega_{\ell-2} t} P_{\ell-2}^m(\cos\theta) \tilde{f}_{\ell-2,m}^{\text{dCS}}(r) \right) \end{aligned}$$



Asymptotic Matching

- Far zone: $r \gg M$
- Near zone: $r - r_H \ll \max(\ell/\omega, \ell/\mu)$
- Buffer zone: $r \gg M$,
 $r - r_H \ll \min(\ell/\omega, \ell/\mu)$



Near zone metric resummation

dCS metric has spurious divergences at Schwarzschild and Kerr horizons, so we need to resum:

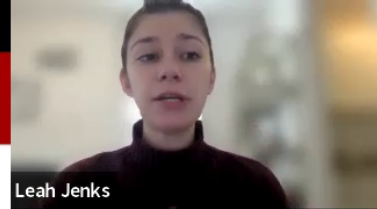
- 1) Resummed metric identically reduces to original metric when expanded in small $\chi \ll 1$
- 2) Each component remains finite everywhere outside the dCS horizon

$$\bar{\Delta} = \Delta + \frac{915}{14336} M \zeta \chi^2$$

$$\delta g_{rr} = \frac{915}{14336} \frac{M^2 r^2 \zeta \chi^2}{\bar{\Delta}^2}$$

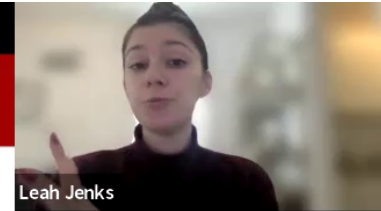
$$\delta g_{tt} = \frac{915}{14336} \frac{M^2 \chi^2 \zeta}{r^2}$$

$$\begin{aligned} g_{rr,\text{resum}} &= g_{rr,K}(\Delta \rightarrow \bar{\Delta}) + g_{rr,\text{dCS}}(f \rightarrow \bar{\Delta}/r^2) + \delta g_{rr} \\ g_{tt,\text{resum}} &= g_{tt,K}(\Delta \rightarrow \bar{\Delta}) + g_{tt,\text{dCS}}(f \rightarrow \bar{\Delta}/r^2) + \delta g_{tt} \end{aligned}$$



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Superradiance in dCS Gravity: Frequencies



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Frequency spectrum:

- $\ell, \ell+2$ maximized for $\ell=1, m=1$
- $\ell-2$ maximized for $\ell=3, m=1$

$\ell=1$

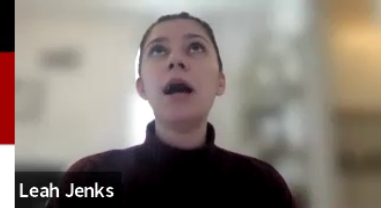
$$\text{Im}(\omega_\ell) \approx \mu \frac{(\mu M)^8}{24} (\chi - 2\mu r_+)$$

$\ell=3$

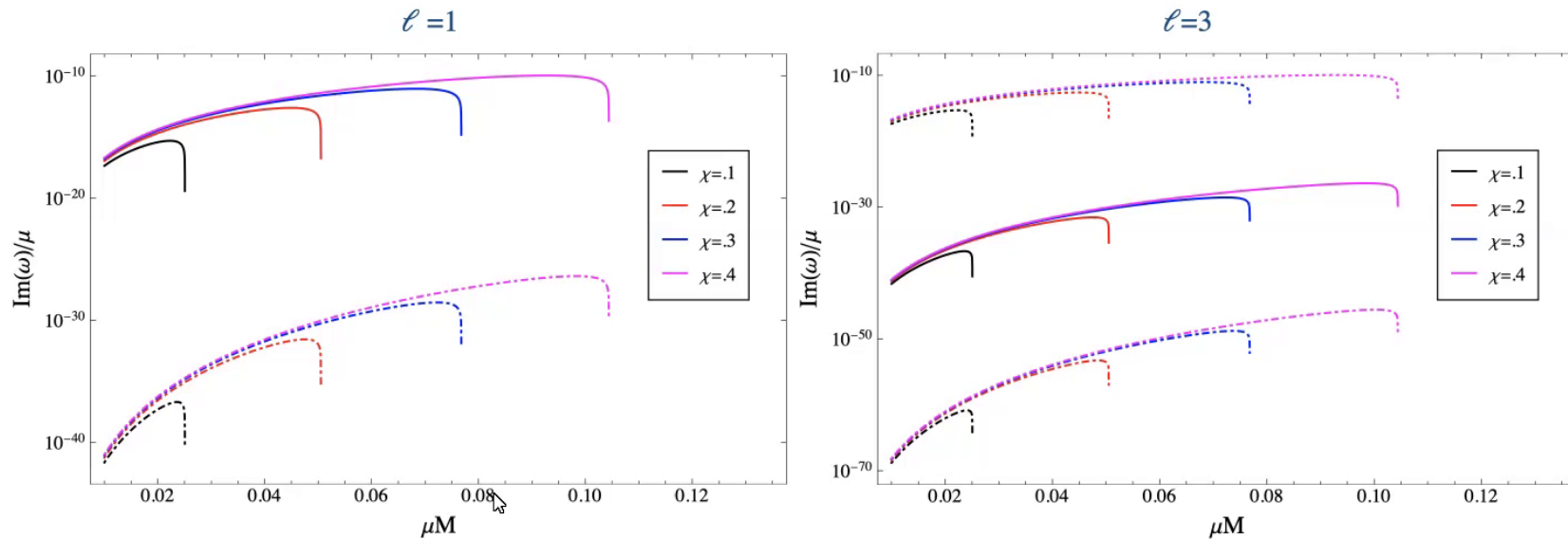
$$\text{Im}(\omega_{\ell-2}) \approx \mu \frac{(\mu M)^8}{24} (\chi - 2\mu r_+)$$

$$\text{Im}(\omega_{\ell+2}) \approx \mu \frac{(\mu M)^{16}}{129024000} (\chi - 2\mu r_+)$$

Superradiance in dCS Gravity: Frequency Spectrum



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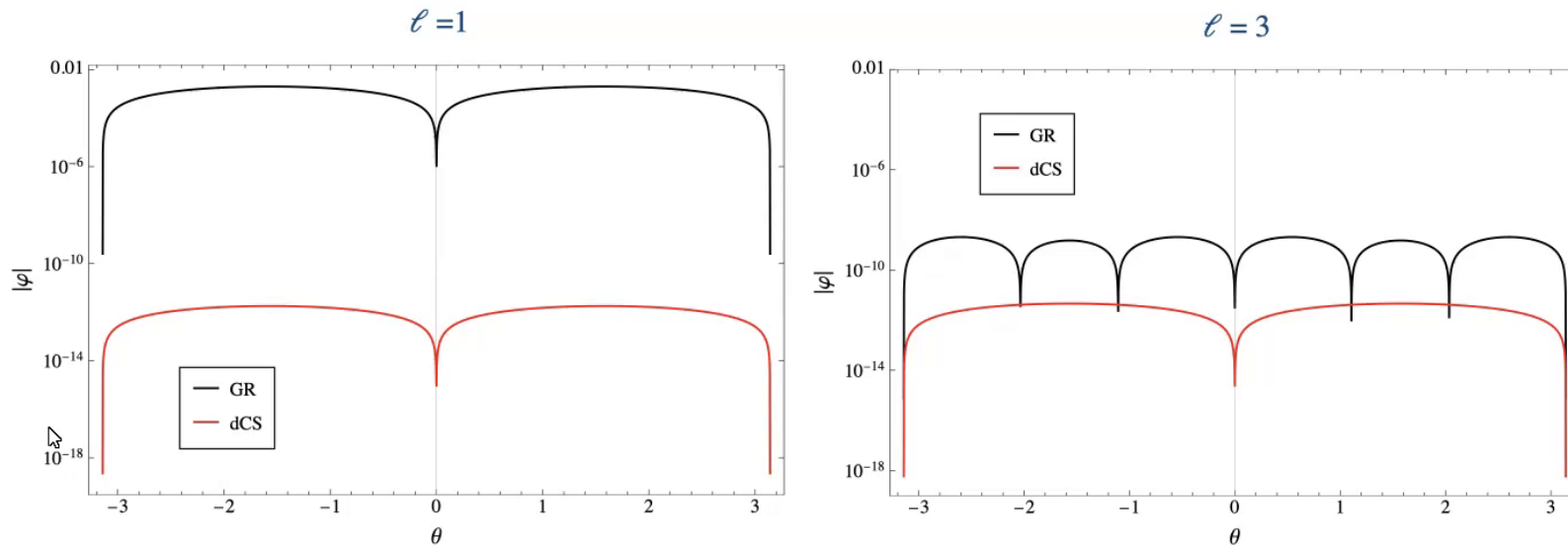


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Superradiance in dCS Gravity: Angular Dependence



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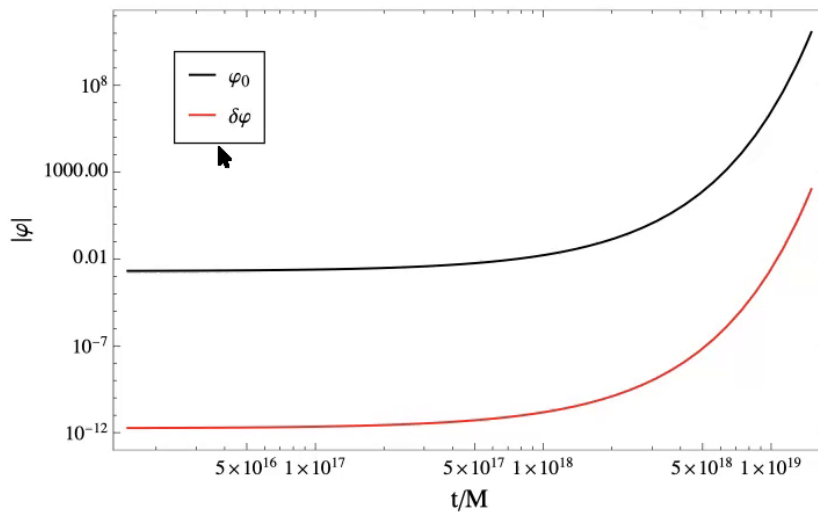


Superradiance in dCS Gravity: Time Evolution

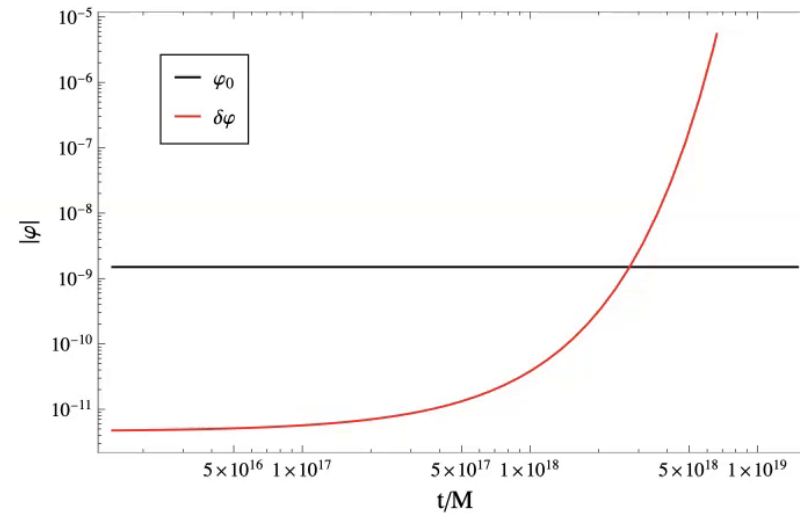


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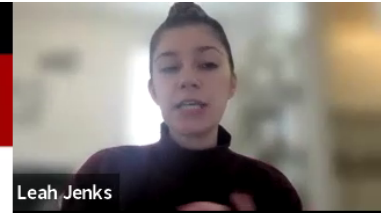
$\ell = 1$



$\ell = 3$



Superradiance in dCS Gravity: Observables & Future Work



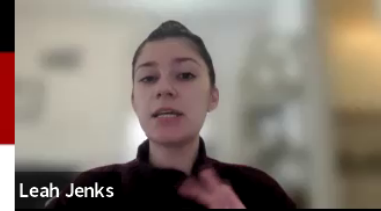
Observables:

- Black hole spin-down (Arvanitaki et al., 2009, Arvanitaki & Dubovsky, 2010 +)
 - Regge plane constraints
- Gravitational wave emission (Arvanitaki et al., 2009, Arvanitaki & Dubovsky, 2010 +)
- EHT observations (eg Davoudiasl & Denton, 2019)



Event Horizon Telescope

Superradiance in dCS Gravity: Observables & Future Work



Fermionic Superradiance in dCS?

- Fermions do NOT super radiate in GR - Unruh 1974
- Potential for presence of caps and dCS scalar to evade this - unlikely
- Potential for induced superradiance
 - Analogy to fermion preheating (Green & Kofman, 1998)
 - Analogy to induced EM superradiance (Boskovic et al., 2018)

↵

Summary & Conclusions

- dCS gravity provides a vast background to search for observations of modified gravity
- The dCS caps may provide a window into observable signatures
- Superradiance in dCS gravity is distinct from GR, leads to small corrections in the scalar profile
- More work is necessary!

Thank you!