Title: Melonic large N limit of 5-index irreducible random tensors

Speakers: Sabine Harribey

Series: Quantum Gravity

Date: November 25, 2021 - 2:30 PM

URL: https://pirsa.org/21110040

Abstract: The main feature of tensor models is their melonic large N limit, leading to applications ranging from random geometry and quantum gravity to many-body quantum mechanics and conformal field theories. However, this melonic limit is lacking for tensor models with ordinary representations of O(N) or Sp(N). We demonstrate that random tensors with sextic interaction transforming under rank-5 irreducible representations of O(N) have a melonic large N limit. This extends the recent proof obtained for rank-3 models with quartic interaction. After giving an introduction to random tensors, I will present the main ideas of our proof relying on recursive bounds derived from a detailed combinatorial analysis of the Feynman graphs.

Zoom Link: https://pitp.zoom.us/j/94691275506?pwd=RGFaN0NZR0FScFdOTXFzeFVXaXUvUT09



Sabine Harribey

Joint work with Sylvain Carrozza - arXiv:2104.03665

Perimeter Institute - November 25th 2021









ഥ

DEPARTMENT OF PHYSICS AND





History of random tensors

- First introduced in zero dimension: random geometry and quantum gravity [Ambjorn Durhuus Jonsson '90, Boulatov '92, Ooguri '92, ...]
- Generalisation of matrix models in dimension higher than two
- Provide one possible definition of discretized Euclidean quantum gravity
- Gluing of tetrahedrons for rank 3.



History of random tensors

- First introduced in zero dimension: random geometry and quantum gravity [Ambjorn Durhuus Jonsson '90, Boulatov '92, Ooguri '92, ...]
- Generalisation of matrix models in dimension higher than two
- Provide one possible definition of discretized Euclidean quantum gravity
- Gluing of tetrahedrons for rank 3.
- No large N limit
- Very degenerate ways of gluing tetrahedra are allowed

 \Rightarrow Revived when a large *N* expansion was found [Gurau, Bonzom, Rivasseau, '10, ...]



Melon diagrams



Different types of melonic limit

• **Colored** tensor models: 4 tensor fields, $O(N)^6$ symmetry

• Uncolored tensor models: 1 tensor field, $O(N)^3$ symmetry

 $\varphi_{\textit{abc}}$

 φ^{i}_{abc}

SYK model

 $J_{i_1 i_2 i_3 i_4} \psi_{i_1} \psi_{i_2} \psi_{i_3} \psi_{i_4}$

• Multi matrix models with a large number of matrices [Ferrari, Schaposnik Massolo, Valette ...], models in higher rank [Bonzom,Lionni, ...]



Melonic limit and quantum gravity

- Colored tensor models: colored triangulations
- Limitation: melons \sim branched polymers [Gurau, Ryan, 2013]
- Open challenge [Bonzom, Oriti]



Melonic limit and quantum gravity

- Colored tensor models: colored triangulations
- Limitation: melons \sim branched polymers [Gurau, Ryan, 2013]
- Open challenge [Bonzom, Oriti]

Can be used in another setting: large N field theory



General motivations: renormalization

Physics change with the energy scale: Renormalization group



Flow in the space of theories with respect to the energy scale [Wilson 1972, Polchinski 1984, \dots]:

 \rightarrow fixed points and trajectories



Melonic limit and quantum gravity

- Colored tensor models: colored triangulations
- Limitation: melons \sim branched polymers [Gurau, Ryan, 2013]
- Open challenge [Bonzom, Oriti]

Can be used in another setting: large N field theory



4/24

Ð.

General motivations: renormalization

Physics change with the energy scale: Renormalization group



Flow in the space of theories with respect to the energy scale [Wilson 1972, Polchinski 1984, \dots]:

ightarrow fixed points and trajectories



Weak coupling	Strong coupling
Perturbation theory	?
Non perturbative ?	?

- Study non-trivial fixed points of the renormalization group [Wilson]
- Gain control over non-perturbative phenomena
- Spontaneous symmetry breaking [Coleman, Jackiw, Politzer]
- Dynamical mass generation [Gross, Neveu]



Weak coupling	Strong coupling
Perturbation theory	?
Non perturbative ?	?

- Study non-trivial fixed points of the renormalization group [Wilson]
- Gain control over non-perturbative phenomena
- Spontaneous symmetry breaking [Coleman, Jackiw, Politzer]
- Dynamical mass generation [Gross, Neveu]
- Large N limit: helpful approximation scheme
- Vast array of applications: statistical mechanics, QCD, quantum gravity



From vector to tensor models





Melonic quantum mechanics

- Strongly coupled QFTs and holography (d = 1): SYK model without disorder [Witten, Klebanov, Tarnopolsky, ...]
- Tensor models in higher dimension: new class of conformal field theories
- Problem: divergences
- Renormalization group: computation of beta functions, stable IR fixed point? Unitarity ?





Main features of melonic CFTs

- Melonic dominance
 - Closed Schwinger-Dyson e_{μ}^{a} uation
 - Simplification of the four-point kernel



Main features of melonic CFTs

- Melonic dominance
 - Closed Schwinger-Dyson equation
 - Simplification of the four-point kernel
- Two fixed points mechanisms:
 - Usual $d = 4 \epsilon$ regularization for short-range models
 - Exact marginality for long-range models
- Unitarity range: real scaling dimensions above unitarity bounds and real OPE coefficients





Irreducible tensor models

- Random geometry: colored tensor models
- Field theory: uncolored tensor models
- What about other tensor representations ?
- Completely symmetric tensors: no melonic large N limit
- Conjecture: melonic large N limit for irreducible representations of O(N) or Sp(N) [Klebanov, Tarnopolsky]
- Proof in rank 3 [Benedetti, Carrozza, Gurau, Kolanowski]
- Generalization in higher rank ?

 \Rightarrow Here for rank 5



Outline

Sabine Harribey

The model

2 Perturbative expansion

3 Sketch of the proof

⊕,

The model

Free energy:

$$F_{\mathbf{P}}(\lambda) = \frac{6}{N^5} \lambda \partial_{\lambda} \ln \left\{ \left[e^{\frac{1}{2} \partial_{\tau} \mathbf{P} \partial_{\tau}} e^{\frac{\lambda}{6N^5} \delta^{h}_{abcdef} T_{a} T_{b} T_{c} T_{d} T_{e} T_{f}} \right]_{T=0} \right\}$$

- **P** one of seven orthogonal projectors on irreducible representations
- Interaction vertex

 $T_{a_1 a_2 a_3 a_4 a_5} T_{a_5 b_2 b_3 b_4 b_5} T_{b_5 a_4 c_3 c_4 c_5} T_{c_5 b_4 a_3 d_4 d_5} T_{d_5 c_4 b_3 a_2 e_5} T_{e_5 d_4 c_3 b_2 a_1}$





Perturbative expansion

• $F_{P}(\lambda)$: sum over rooted connected combinatorial maps \mathcal{G}



- Half-edge: represented with five strands
- 945 ways to connect two half-edges
- Projector: combination of those terms with different weights and signs
- Stranded map G: combinatorial map with a choice of one term per edge

$$F_{P}(\lambda) = \sum_{G \text{ connected, rooted}} \lambda^{V(G)} \mathcal{A}(G)$$





Types of edges

• Unbroken: all strands traverse



• Broken: a pair of corners is connected by a strand at each end of the edge. Rescaled by 1/N



• Doubly broken: two pairs of corners are connected by a strand at each end of the edge. Rescaled by $1/N^2$



1/N expansion

Amplitude of a stranded map:

$$\mathcal{A}(G) = \mathcal{K}(G) \mathcal{N}^{-\omega(G)}(1 + \mathcal{O}(1/\mathcal{N}))$$

- K(G) non-vanishing rational number independent of N
- One free sum = One factor of N per face
- Degree of a stranded map:

$$\omega(G) = 5 + 5V(G) + B(G) + 2B_2(G) - F(G)$$



The model

Free energy:

$$F_{\mathbf{P}}(\lambda) = \frac{6}{N^5} \lambda \partial_{\lambda} \ln \left\{ \left[e^{\frac{1}{2} \partial_{\tau} \mathbf{P} \partial_{\tau}} e^{\frac{\lambda}{6N^5} \delta^{h}_{abcdef} T_{a} T_{b} T_{c} T_{d} T_{e} T_{f}} \right]_{T=0} \right\}$$

- P one of seven orthogonal projectors on irreducible representations
- Interaction vertex

 $T_{a_1 a_2 a_3 a_4 a_5} T_{a_5 b_2 b_3 b_4 b_5} T_{b_5 a_4 c_3 c_4 c_5} T_{c_5 b_4 a_3 d_4 d_5} T_{d_5 c_4 b_3 a_2 e_5} T_{e_5 d_4 c_3 b_2 a_1}$





1/N expansion

Amplitude of a stranded map:

 $\mathcal{A}(G) = \mathcal{K}(G) \mathcal{N}^{-\omega(G)}(1 + \mathcal{O}(1/\mathcal{N}))$

- K(G) non-vanishing rational number independent of N
- One free sum = One factor of N per face
- Degree of a stranded map:

$$\omega(G) = 5 + 5V(G) + B(G) + 2B_2(G) - F(G)$$

Goals:

- \rightarrow Non-negative degree
- \rightarrow Maps of zero degree are melonic



Problematic cases

- Graphs with only long faces: positive degree
- Simplify the degree by considering number of faces of length p

$$\omega(G) = 5 + B(G) + 2B_2(G) + \sum_p F_p\left(\frac{p}{3} - 1\right)$$

- \rightarrow Can be negative iff faces of length p = 1 or p = 2 (short-faces):
 - p = 1: Tadpoles, double-tadpoles
 - p = 2: Melons, dipoles



Bad double tadpoles



Chain of p double tadpoles:

- 4 faces per vertex
- Factor N^{-5} per vertex
- 2 faces when we glue two double-tadpoles

$$\left(\frac{1}{N}\right)^p N^{2p-1} = N^{p-1}$$

Unbounded from above

 \rightarrow non-trivial cancellations



Bounds on combinatorial maps

- Stranded maps with negative degree
- Use irreducibility of the representation to bound the amplitude of the full combinatorial maps
- Double-tadpoles combinatorial maps well-behaved
- Melons: contribute to leading-order



Bounds on combinatorial maps

- Stranded maps with negative degree
- Use irreducibility of the representation to bound the amplitude of the full combinatorial maps
- Double-tadpoles combinatorial maps well-behaved
- Melons: contribute to leading-order
- Problem: generalized double-tadpoles \rightarrow arbitrarily negative degree
- Need to subtract both melons and double-tadpoles





Main theorem

Theorem

We have (in the sense of perturbation series):

$$F_{\boldsymbol{P}}(\lambda) = \sum_{\omega \in \mathbb{N}} N^{-\omega} F_{\boldsymbol{P}}^{(\omega)}(\lambda) \,. \tag{1}$$

- Subtract double-tadpoles and melons
- Restrict to unbroken edges
- Induction: remaining graphs have positive degree



Step 1: Subtraction of double-tadpoles and melons

- Consider a theory with modified covariance KP and subtracted interaction
- Show that we can choose K s.t. this is our original theory
- Partial resummation of the perturbative series: closed and algebraic SDE
- New perturbative expansion in terms of graphs with no double-tadpoles or melons



Step 2: Restriction to unbroken edges

- $\bullet\,$ Cut and glue: changes the number of faces by $-1,\,0$ or +1
- From doubly broken to unbroken propagator:





Step 2: Restriction to unbroken edges

- $\bullet\,$ Cut and glue: changes the number of faces by $-1,\,0$ or +1
- From doubly broken to unbroken propagator:



Ð.

Remove broken edge:

- decrease the number of faces by at most one
- decrease the number of broken edges by one
- the degree can only decrease



Let G be a stranded graph. If G has no double-tadpole and no melon, then $\omega(G) \ge 0$.

 \rightarrow proof by induction

- Look for a strict subgraph that can be deleted without increasing the degree and preserving the constraints
- Exhaustive graph-theoretic distinction of cases
- High number of particular two-point subgraphs to consider



Stranded graphs with no melon can have vanishing degree

- \rightarrow Non-trivial cancellations
- Bounds on combinatorial maps
- Maps with no melons are subleading
- Conclusion: A Feynman map is leading order iff it is melonic





Conclusion and outlook

- Irreducible tensor models with 5-simplex interactions: melonic large-*N* expansion
- Recursive bounds from a detailed combinatorial analysis of the Feynman graphs.
- Estimated scaling of four and eight-point functions: could include other effective interactions
- Application to large N QFT
- Fermionic fields with *Sp*(*N*) symmetry
- Generalization in arbitrary rank $r \ge 6$?



Conclusion and outlook

- Irreducible tensor models with 5-simplex interactions: melonic large-*N* expansion
- Recursive bounds from a detailed combinatorial analysis of the Feynman graphs.
- Estimated scaling of four and eight-point functions: could include other effective interactions
- Application to large N QFT
- Fermionic fields with *Sp*(*N*) symmetry
- Generalization in arbitrary rank $r \ge 6$?



