

Title: Cosmological and Astrophysical Probes of Sterile Neutrinos - Nashwan Sabti, King's College London

Speakers:

Series: Particle Physics

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Abstract: Sterile neutrinos have been proposed to tackle a number of outstanding questions in physics, including the phenomena of dark matter, neutrino masses and the baryon asymmetry of the Universe. I will review current limits on GeV-, MeV- and keV-scale sterile neutrinos from cosmology and astrophysics. In particular, I will focus on how primordial abundance determinations, Cosmic Microwave Background observations and stellar kinematic inferences from dwarf spheroidal galaxies allow us to set robust constraints across this mass scale. I will then end with a short discussion on how sterile neutrinos could relax cosmological bounds on neutrino masses and what implications this would have for experiments like KATRIN.

# Cosmological and Astrophysical Bounds on Sterile Neutrinos

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Nashwan Sabti

Perimeter Institute – 19/11/2021



# I-SLIDE SUMMARY OF STERILE NEUTRINOS

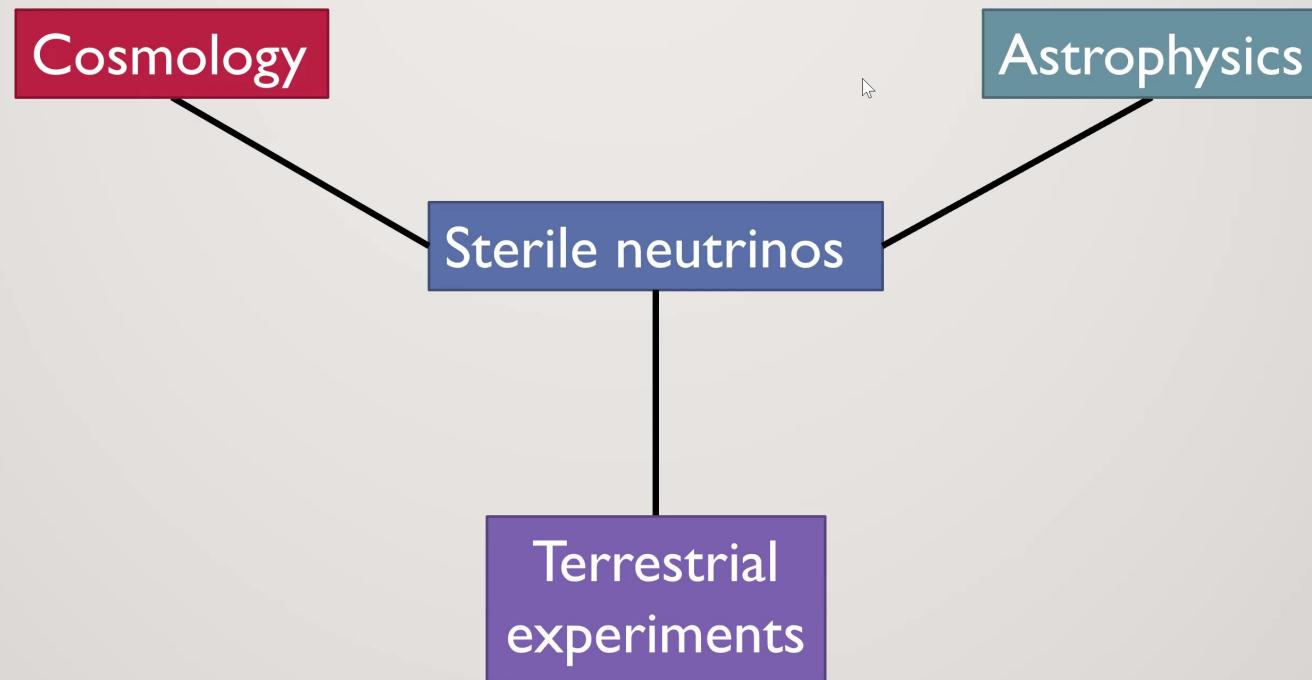
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- Dark matter, neutrino masses, baryon asymmetry all require new physics
- Sterile neutrinos can accommodate for all three phenomena
- Neutrino portal:

$$L_{\text{neutrino-portal}} = F_\alpha (\overline{L}_\alpha \tilde{H}) N + h.c.$$

- Causes mixing between active neutrinos and sterile neutrinos
- Induces neutrino-like interactions for sterile neutrinos, suppressed by small mixing angles

$$\Gamma_{\text{sterile neutrino}} \sim \theta^2 \Gamma_{\text{active neutrino}}$$



# Constraints on the Mass of Fermionic Dark Matter from Dwarf Spheroidal Galaxies



In collaboration with:

James Alvey

Victoria Tiki

Diego Blas

Kyrylo Bondarenko



Alexey Boyarsky



Miguel Escudero



Malcolm Fairbairn



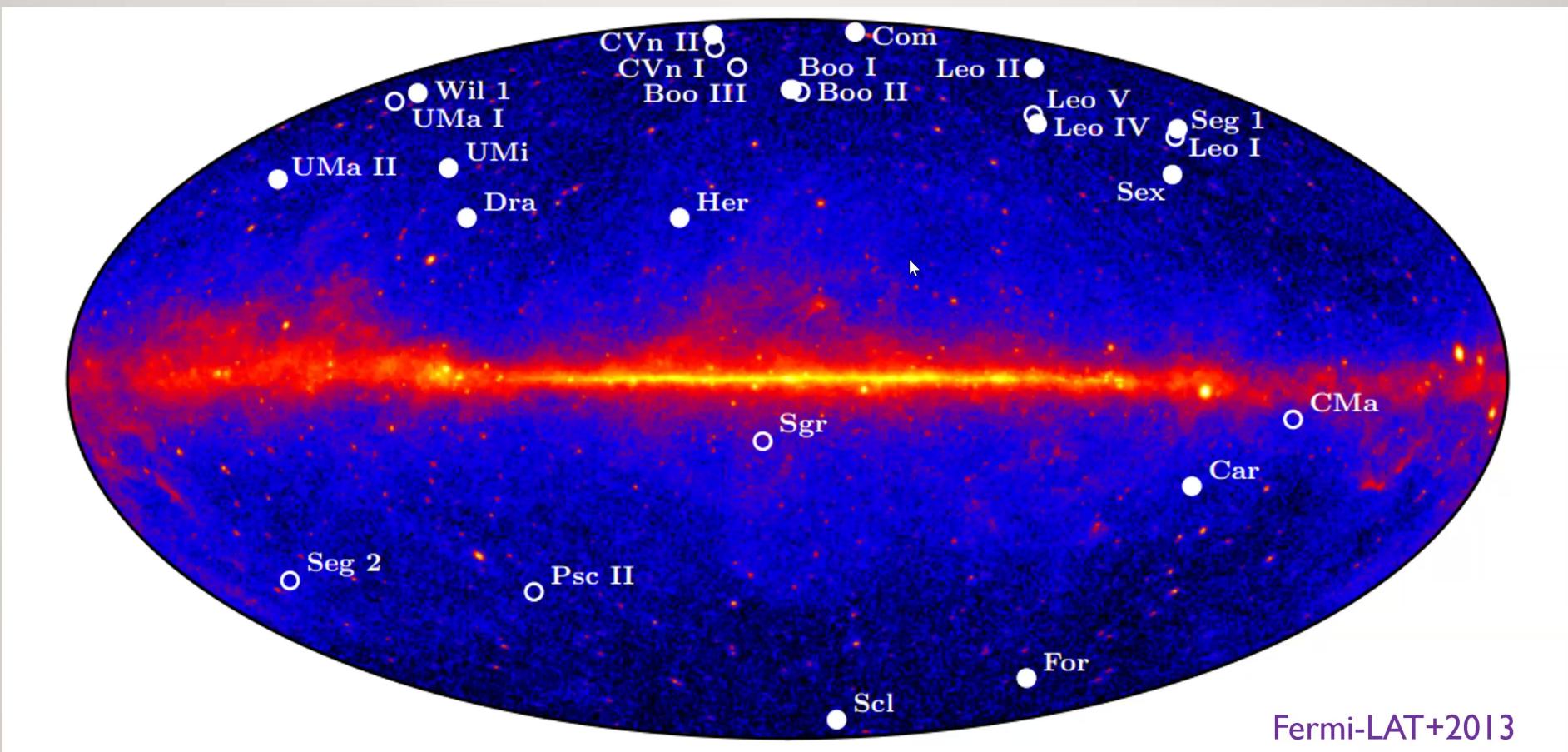
Matthew Orkney

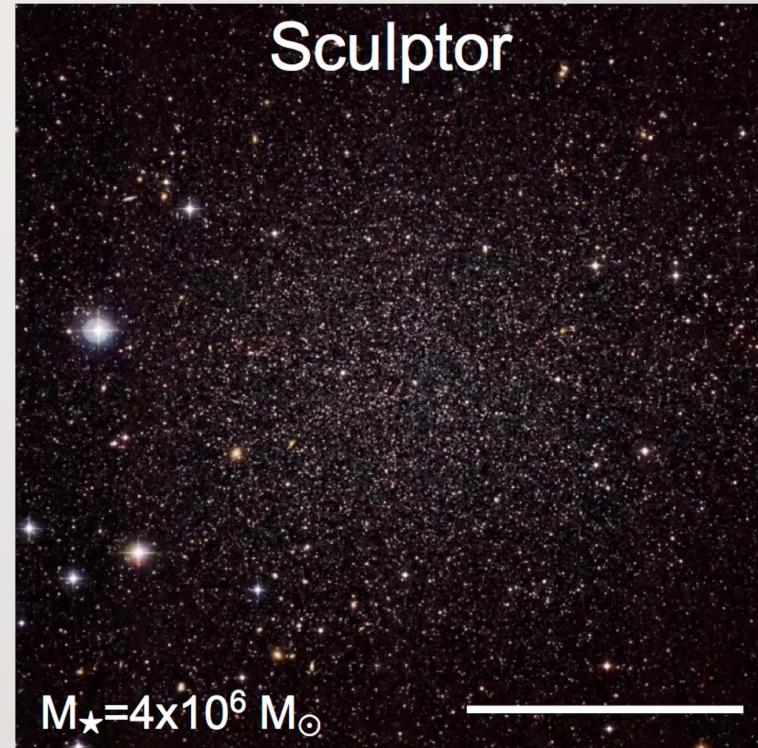
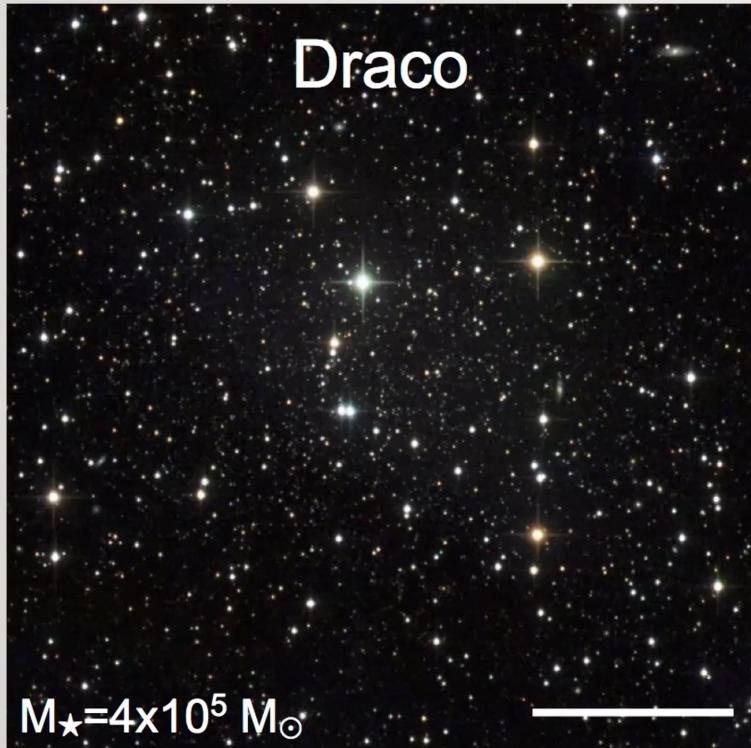


Justin Read

Based on [2010.03572](#)

# Milky Way Dwarf Spheroidal Galaxies





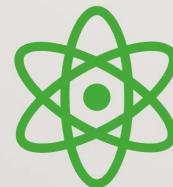
Bullock+2017



Dwarf spheroidal galaxies are  
DM dominated

Measure DM density and  
velocity dispersion

Compute coarse-grained  
phase-space density



Assume a (class of) particle  
physics model(s)

Compute fine-grained phase-  
space density

## OBSERVING THE DM DENSITY AND VELOCITY DISPERSION

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- To compute PS density, we need: *i*) DM density profile and *ii*) velocity dispersion profile
- Simplest approach: assume e.g. isothermal profile

$$\rho_{\text{DM}}(r \ll r_c) \approx \rho_0(r_c, \sigma_{\text{DM}})$$

and  $\sigma_{\text{DM}} \approx \sigma_{\text{star}}$

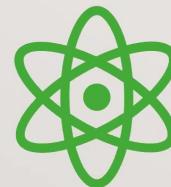
- $m_\nu \gtrsim 0.1 \text{ keV}$



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## OBSERVING THE DM DENSITY AND VELOCITY DISPERSION

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- Jeans equation:

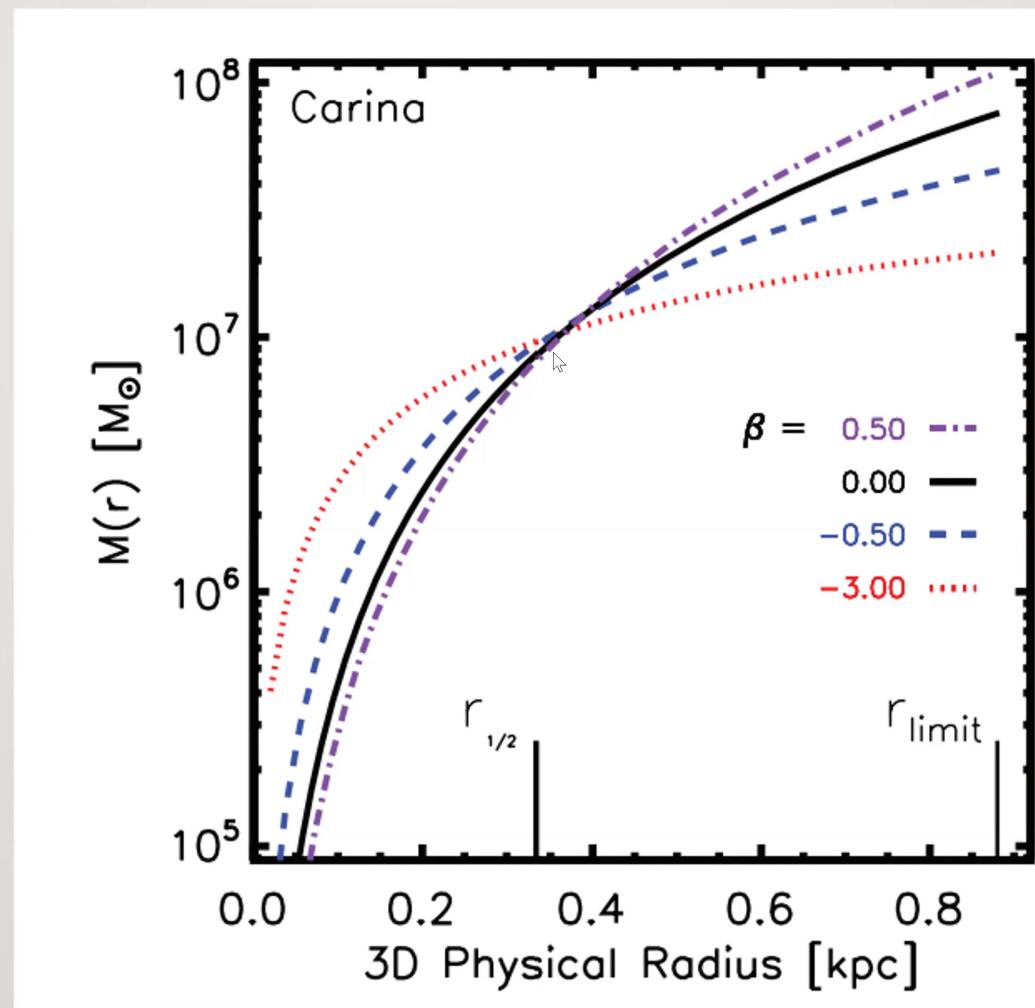
$$\frac{1}{\nu} \frac{\partial}{\partial r} (\nu \sigma_r^2) + 2 \frac{\beta(r) \sigma_r^2}{r} = - \frac{GM(r)}{r^2} , \quad \beta = 1 - \frac{\sigma_t^2}{\sigma_r^2}$$



- Observe line-of-sight velocity and projected tracer density:

$$\sigma_{\text{LOS}}^2(R) = \frac{2}{\Sigma(R)} \int_R^\infty \left( 1 - \beta \frac{R^2}{r^2} \right) \nu \sigma_r^2 \frac{r dr}{\sqrt{r^2 - R^2}}$$

10



Wolf+2009

## OBSERVING THE DM DENSITY AND VELOCITY DISPERSION

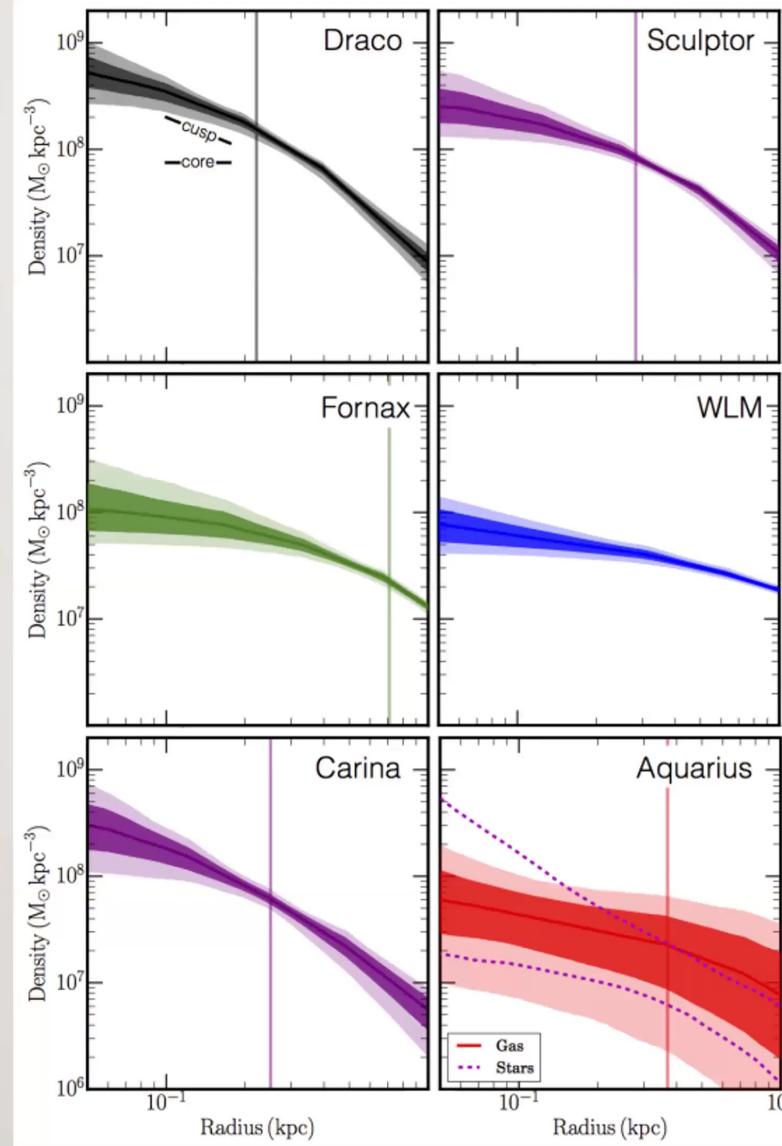
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- Introduce virial shape parameters:

$$\text{VSP1} = \int_0^\infty \Sigma \langle v_{\text{LOS}}^4 \rangle R dR = \frac{2}{5} \int_0^\infty \nu(5 - 2\beta) \sigma_r^2 GMR dR$$

$$\text{VSP2} = \int_0^\infty \Sigma \langle v_{\text{LOS}}^4 \rangle R^3 dR = \frac{4}{35} \int_0^\infty \nu(7 - 6\beta) \sigma_r^2 GM R^3 dR$$

Merrifield & Kent 1990



Read+2018

## OBSERVING THE DM DENSITY AND VELOCITY DISPERSION

- Introduce virial shape parameters:

$$\text{VSP1} = \int_0^\infty \Sigma \langle v_{\text{LOS}}^4 \rangle R dR = \frac{2}{5} \int_0^\infty \nu(5 - 2\beta) \sigma_r^2 GMR dR$$

$$\text{VSP2} = \int_0^\infty \Sigma \langle v_{\text{LOS}}^4 \rangle R^3 dR = \frac{4}{35} \int_0^\infty \nu(7 - 6\beta) \sigma_r^2 GM R^3 dR$$

Merrifield & Kent 1990

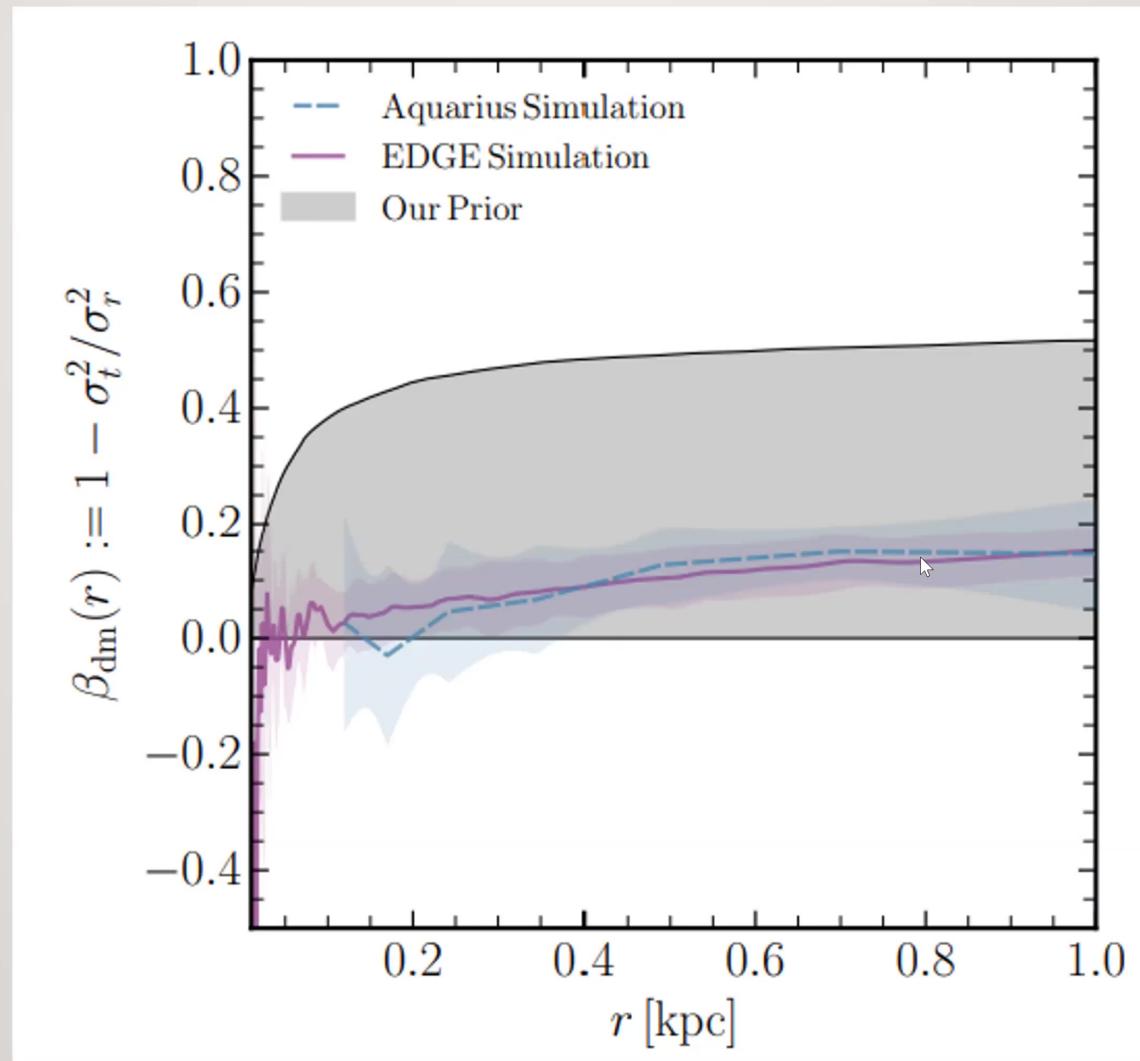
- GravSphere Read+2017

## OBSERVING THE DM DENSITY AND VELOCITY DISPERSION

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- Solve 2<sup>nd</sup> Jeans equation for DM:

$$\frac{1}{\nu_{\text{DM}}} \frac{\partial}{\partial r} (\nu_{\text{DM}} (\sigma_r^{\text{DM}})^2) + 2 \frac{\beta_{\text{DM}}(r) (\sigma_r^{\text{DM}})^2}{r} = - \frac{GM(r)}{r^2}$$



## OBTAINING PHASE-SPACE CONSTRAINTS

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- Bounds from Pauli's principle – model independent

Fermi velocity:  $v_F = \left( \frac{6\pi^2 \rho(r)}{gm^4} \right)^{1/3}$

$$v_F < v_{\text{esc}} \rightarrow m_{\text{deg}} > \left( \frac{6\pi^2 \rho(r)}{g v_{\text{esc}}^3(r)} \right)^{\frac{1}{4}} \Big|_{r_{\min}} = 0.27^{+0.30}_{-0.14} \text{ keV (2}\sigma\text{)}$$

## OBTAINING PHASE-SPACE CONSTRAINTS

- Bounds from Liouville's theorem – model dependent

$$F_{\text{coarse}}^{\text{late}} \leq F_{\text{fine}}^{\text{late}} \leq F_{\text{fine}}^{\text{primordial}} = m^4 f_{\max}$$

Averaging

Liouville's th.

## OBTAINING PHASE-SPACE CONSTRAINTS

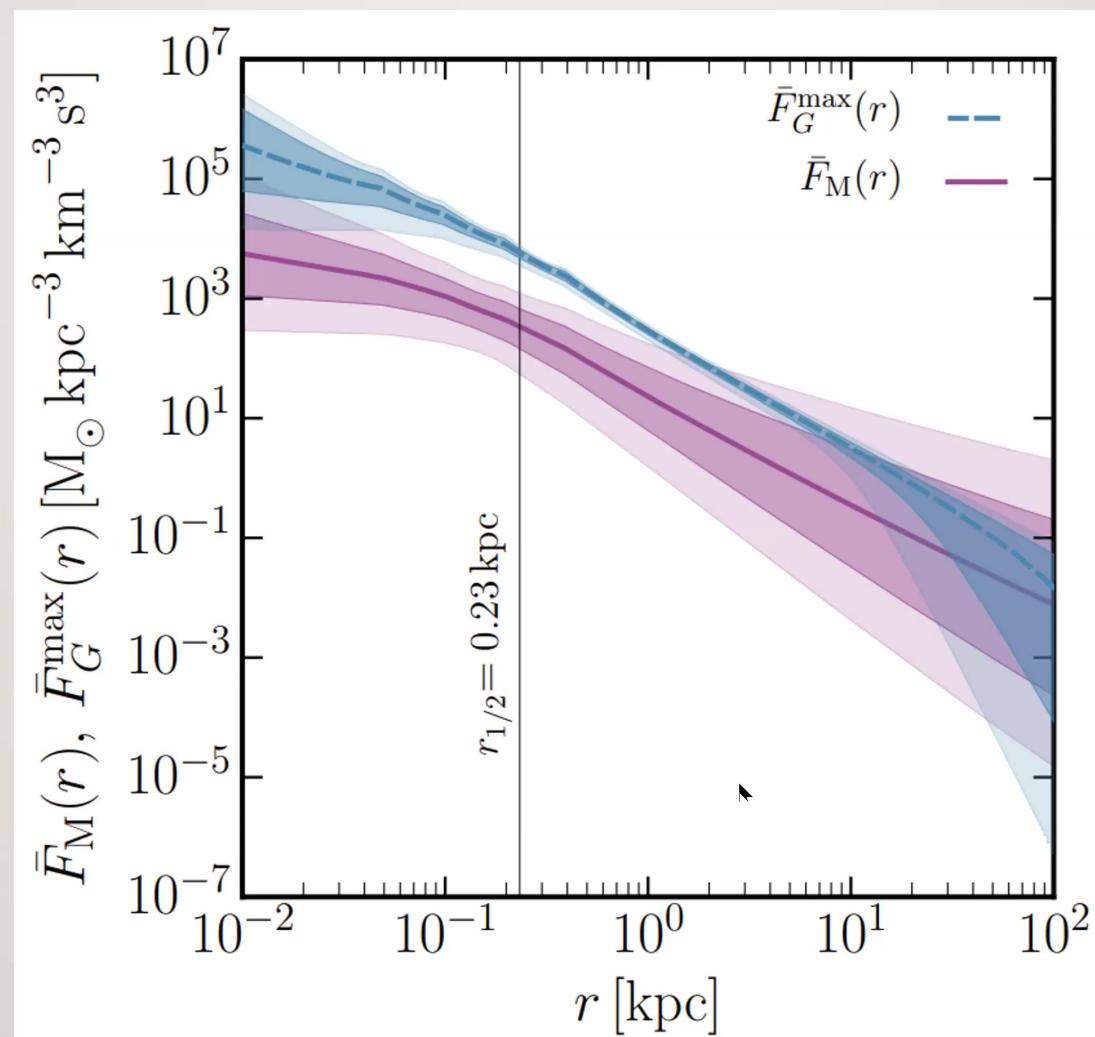
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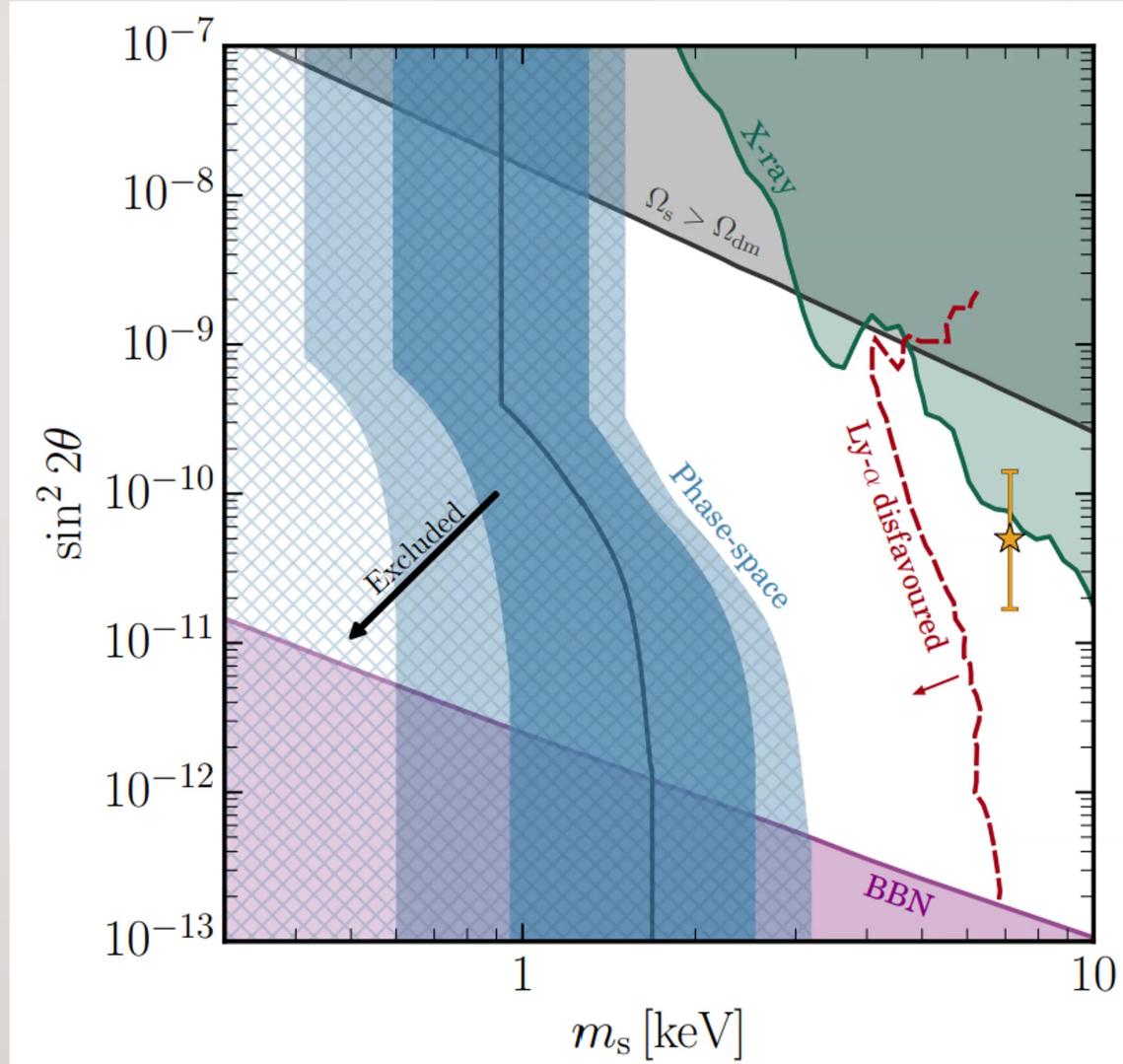
- Maximal coarse-graining:

$$\overline{F_M}(r) = \frac{\rho(r)}{\frac{4\pi}{3} v_{\text{esc}}^3(r)} \leq m^4 f_{\max}$$

- Gaussian coarse-graining:

$$\overline{F_G}(r) = \frac{\rho(r)}{(2\pi)^{3/2} \sigma_r \sigma_t^2} \leq m^4 f_{\max}$$





## SUMMARY

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- Improved Jeans analysis allows to probe the DM density profile in dwarf spheroidal galaxies at small radii
- Phase-space density is higher in the inner regions and allows for stronger limits
- This method provides robust lower bounds on fermionic dark matter

# Bounds on the Lifetime of Sterile Neutrinos from primordial nucleosynthesis

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In collaboration with



Alexey Boyarsky



Maksym Ovchynnikov



Vsevolod Syvolap



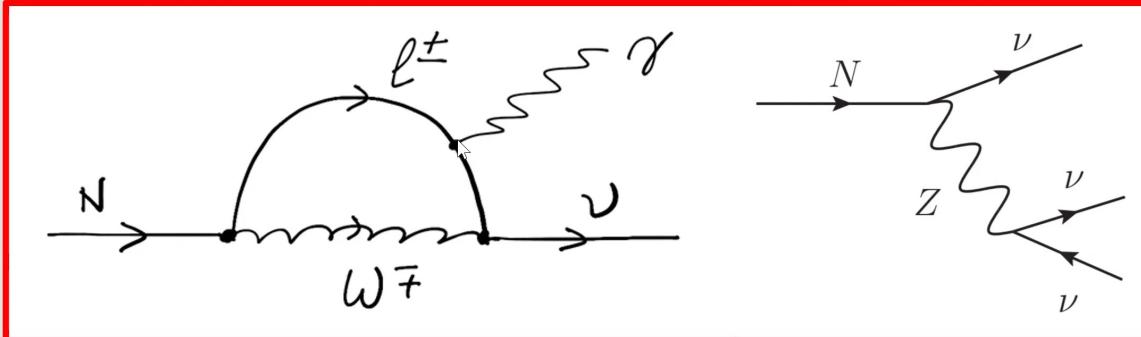
Anastasiia Filimonova



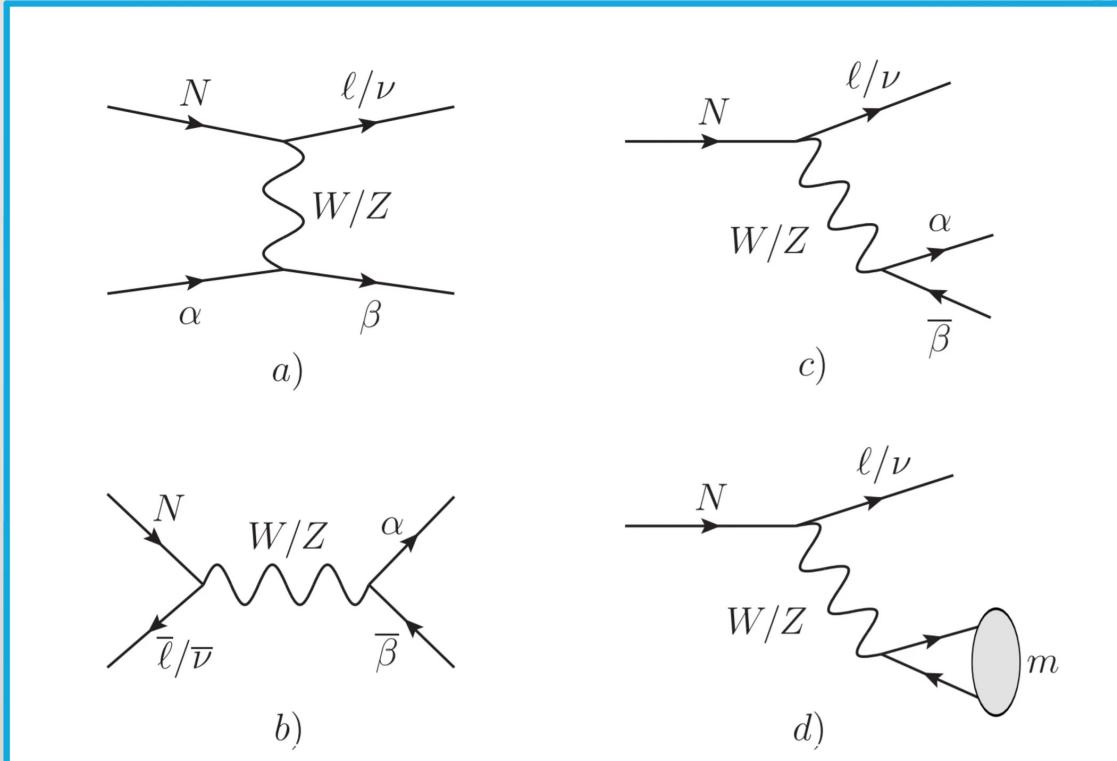
Andrii Magalich

Based on [2006.07387](#) and [2103.09831](#)

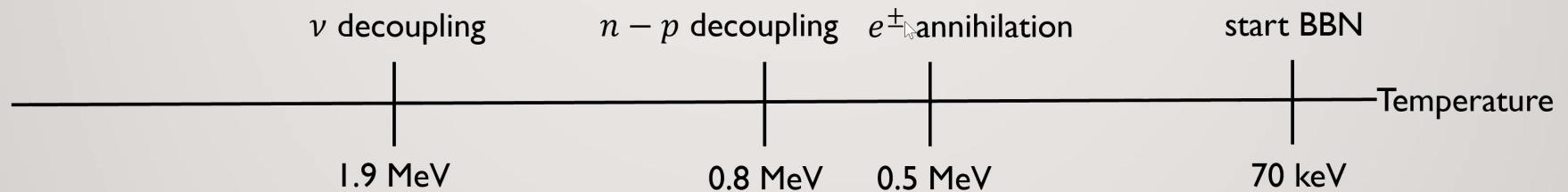
keV scale:



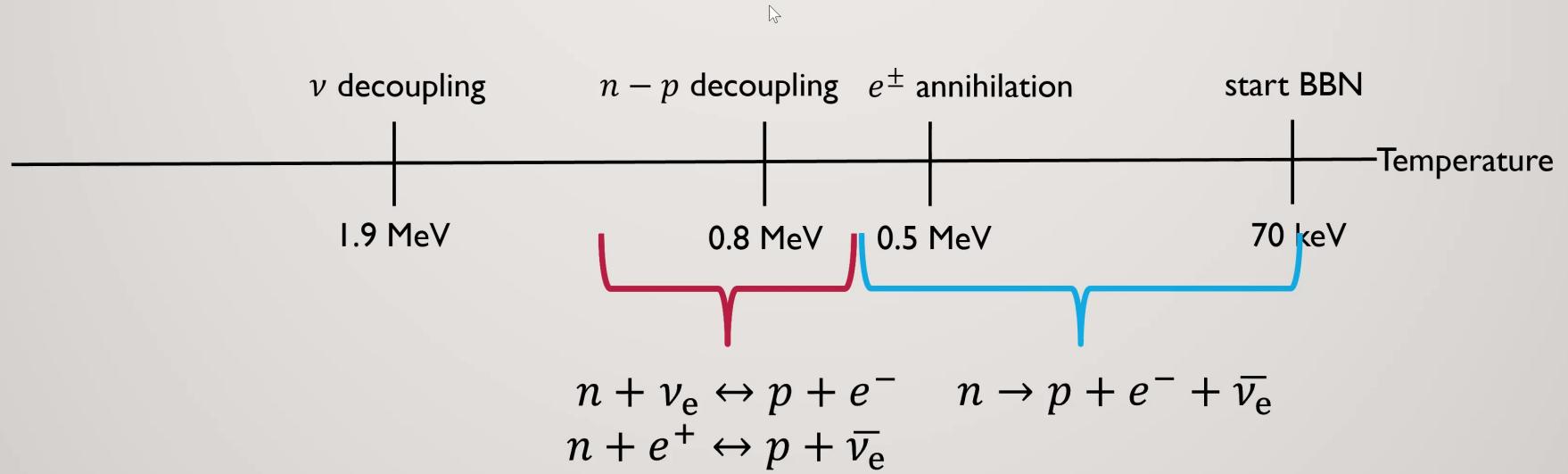
MeV - GeV scale:



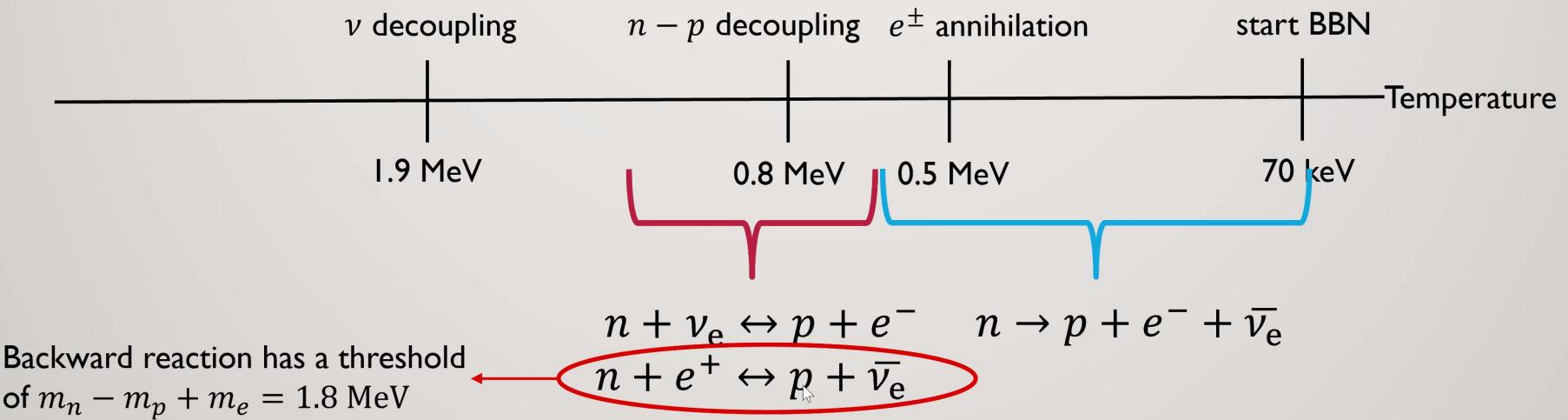
## Timeline of relevant events in the early Universe



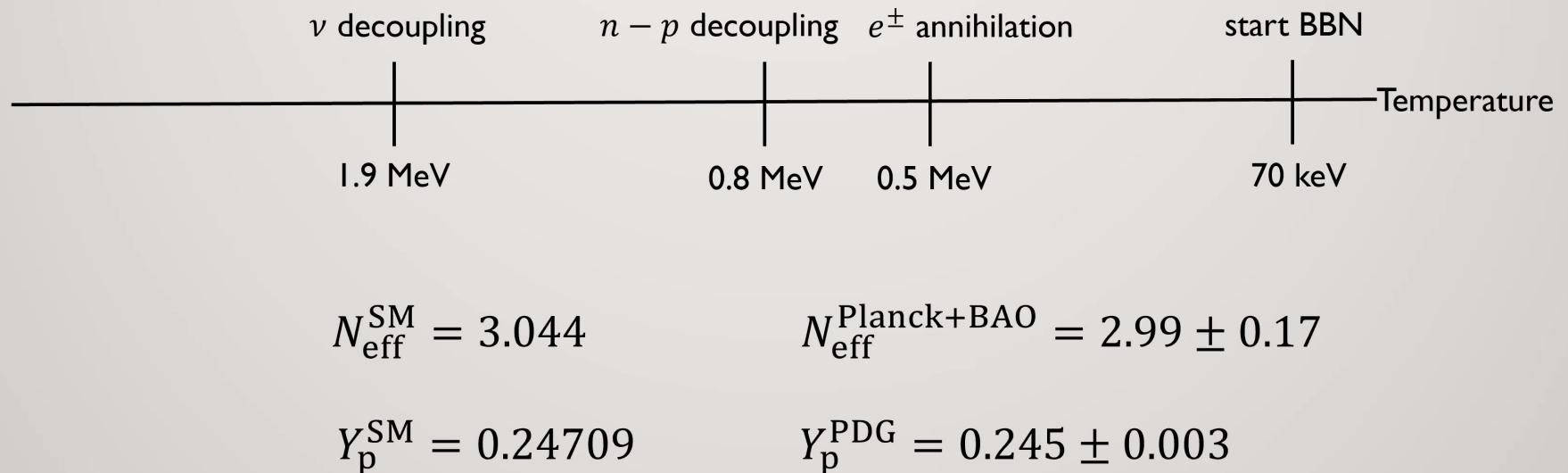
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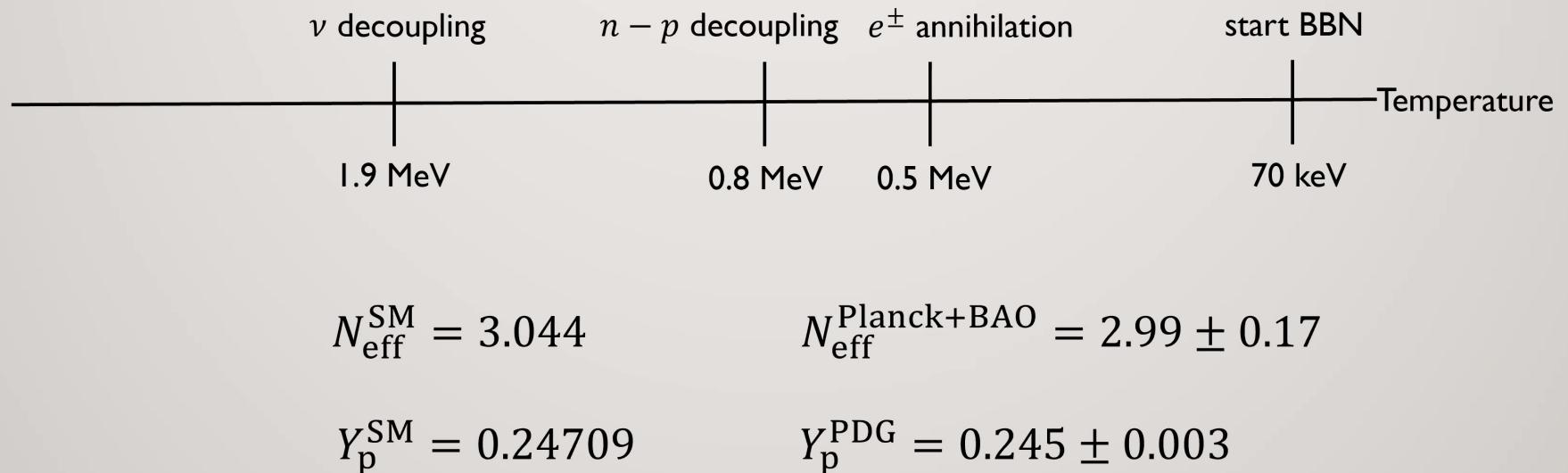
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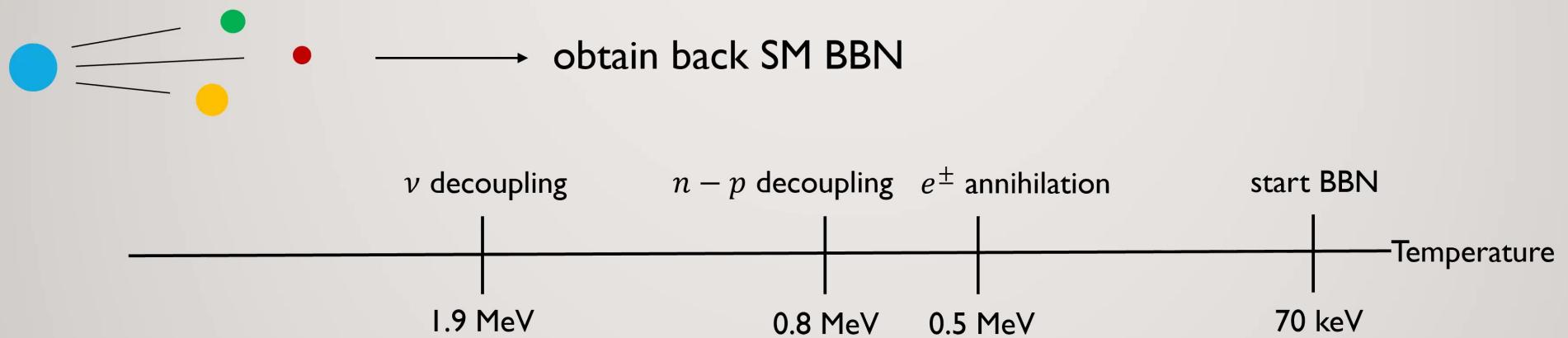
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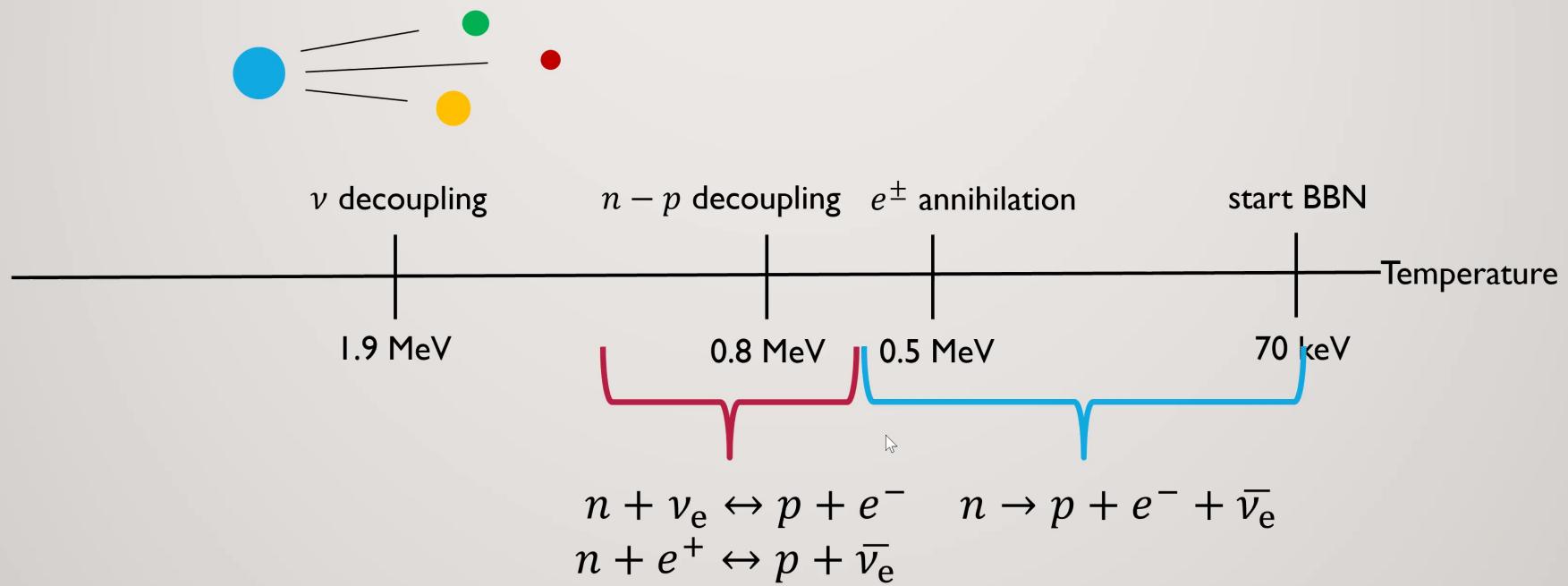
## Timeline of relevant events in the early Universe



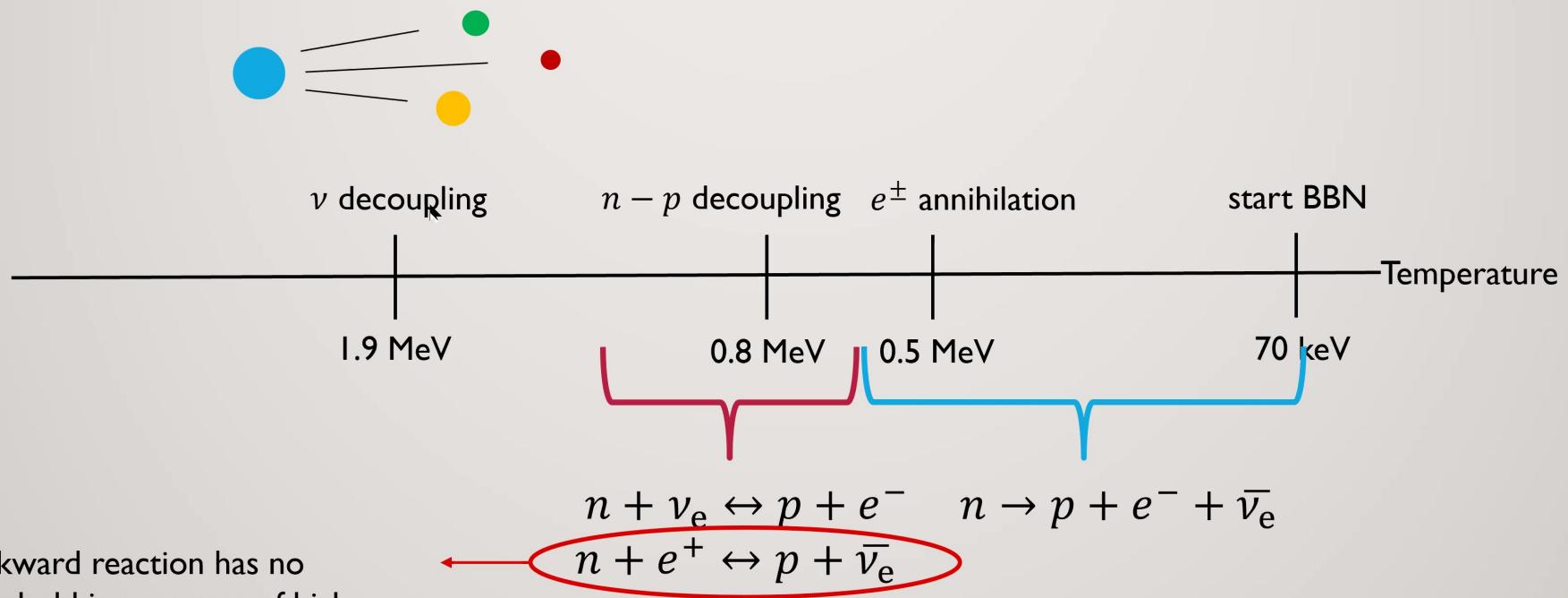
## What happens if a SN decays before $\nu$ decoupling?



## What happens if a SN decays during $\nu$ decoupling?

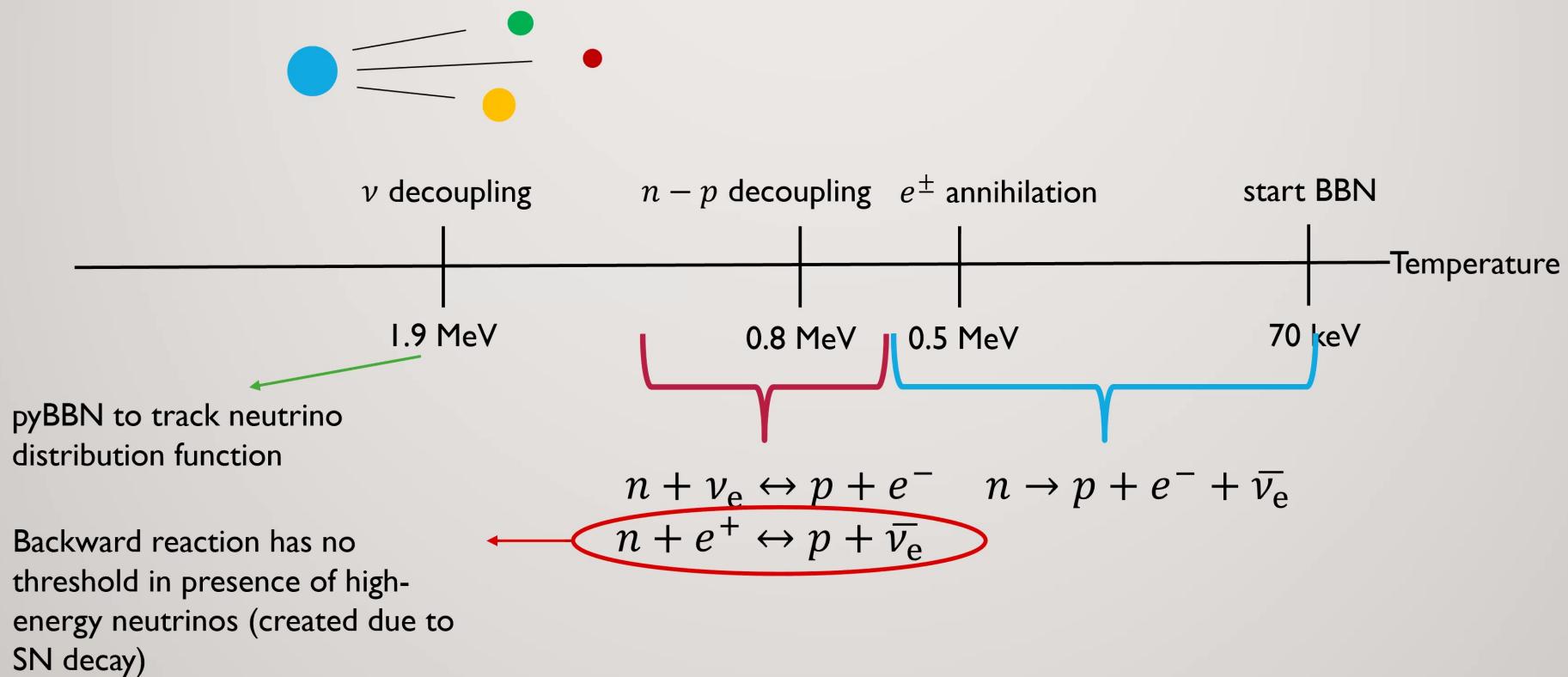


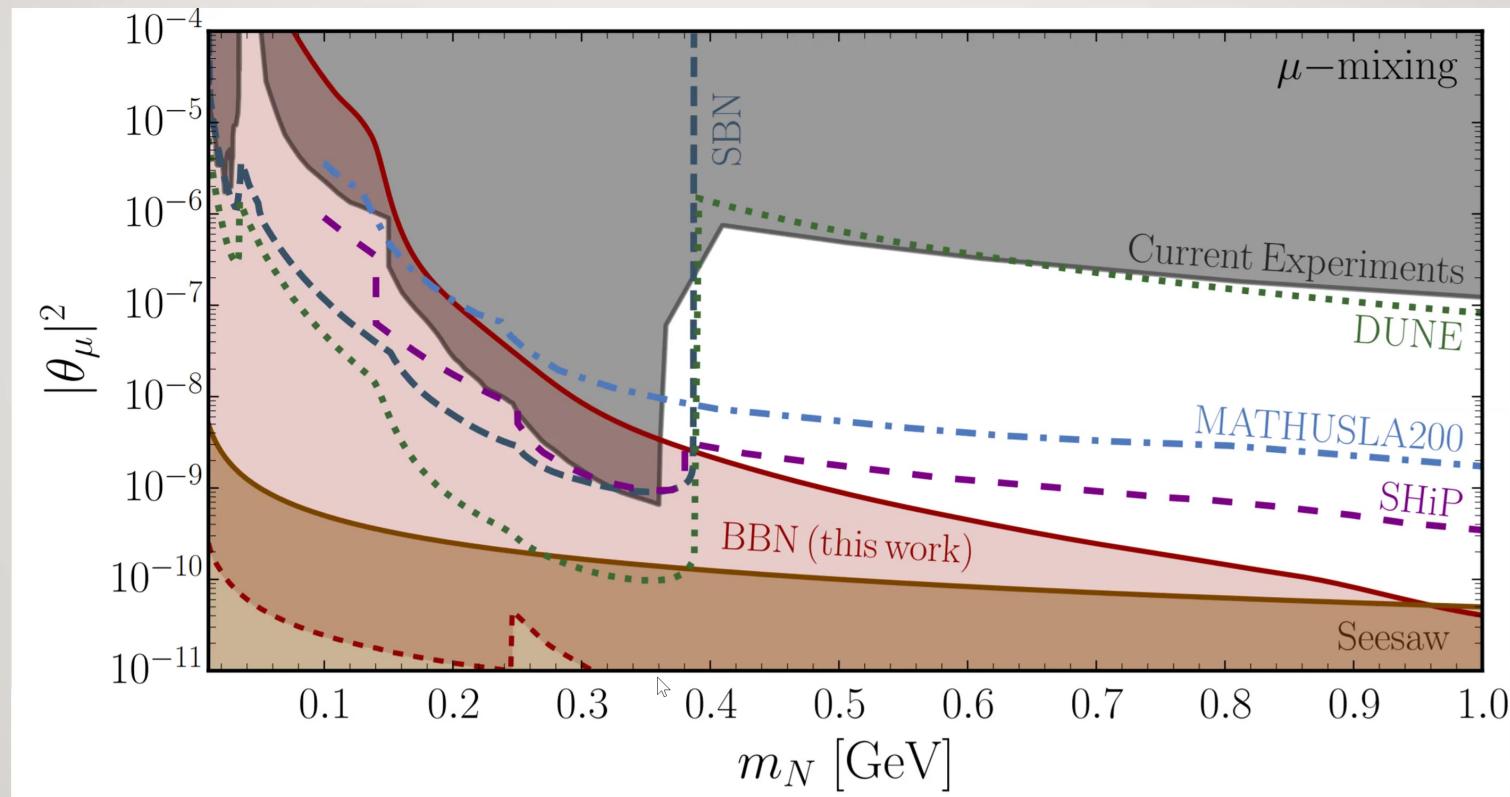
## What happens if a SN decays during $\nu$ decoupling?



Backward reaction has no threshold in presence of high-energy neutrinos (created due to SN decay)

## What happens if a SN decays during $\nu$ decoupling?

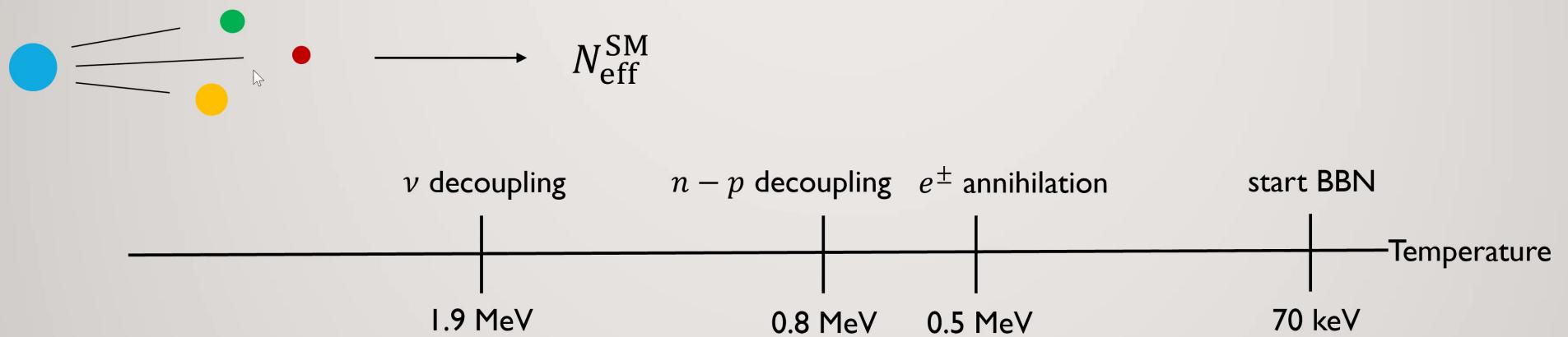




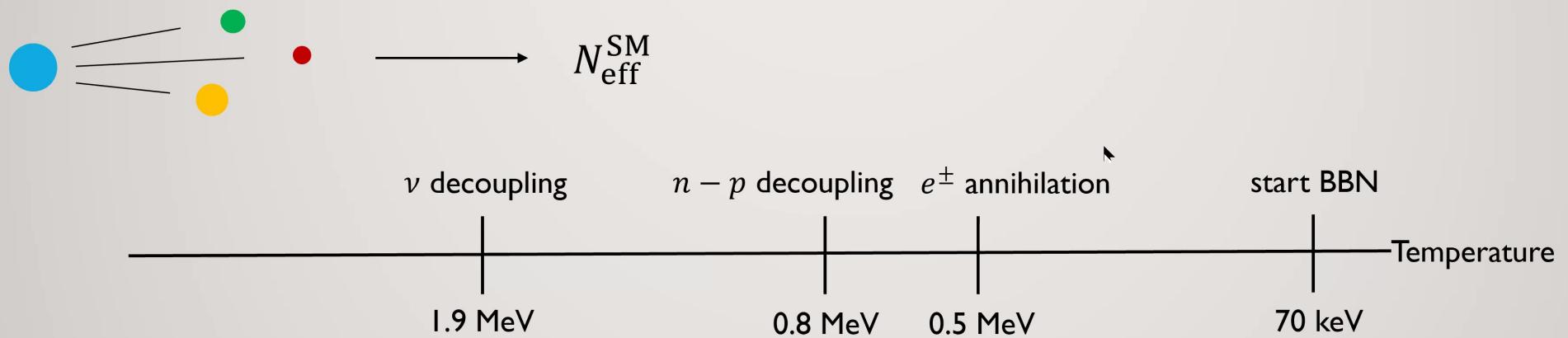
2006.07387

# Detour: How do SNe that decay mostly into active neutrinos affect $N_{\text{eff}}$ ?

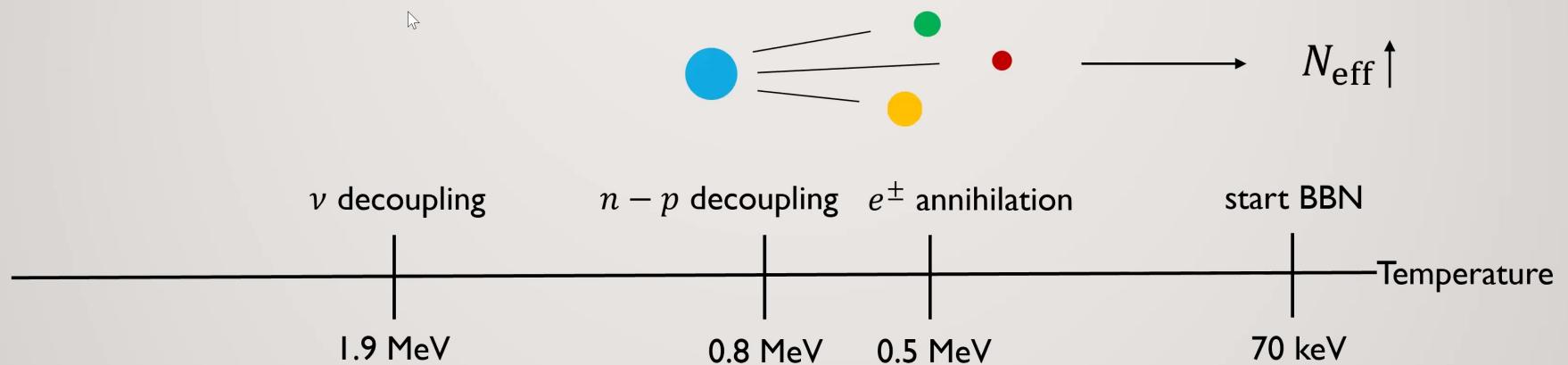
What happens if a FIMP decays into  $\nu$  before  $\nu$  decoupling?



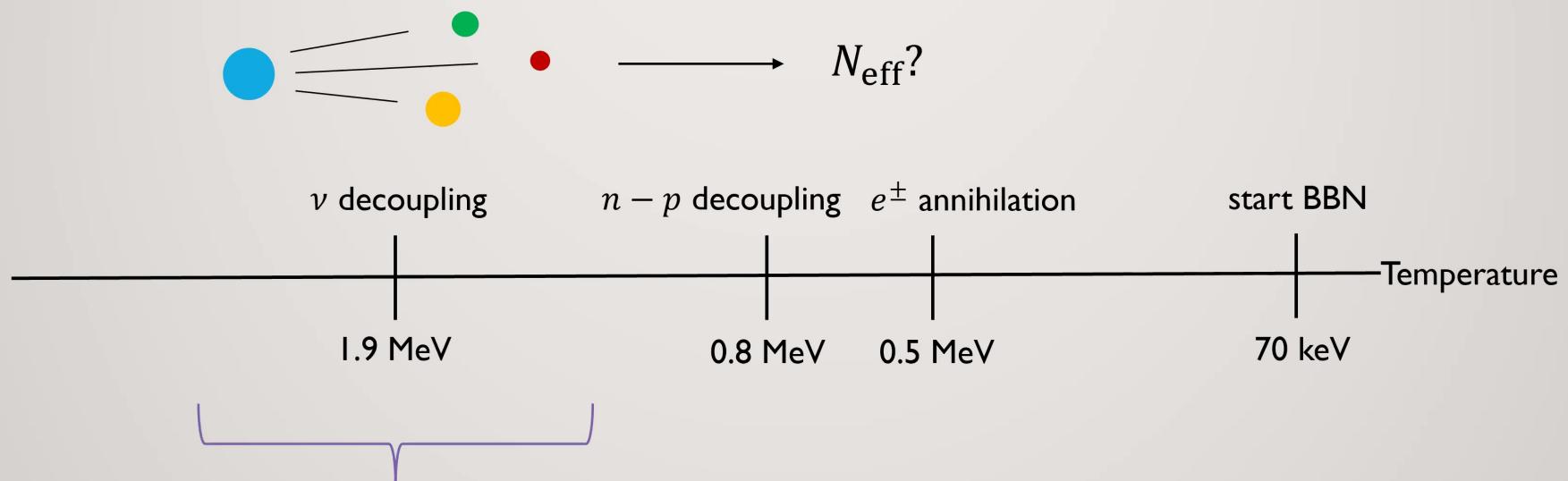
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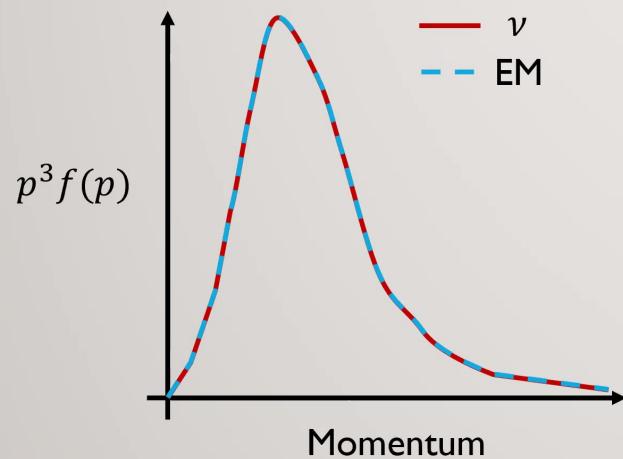
## What happens if a FIMP decays into $\nu$ during $\nu$ decoupling?



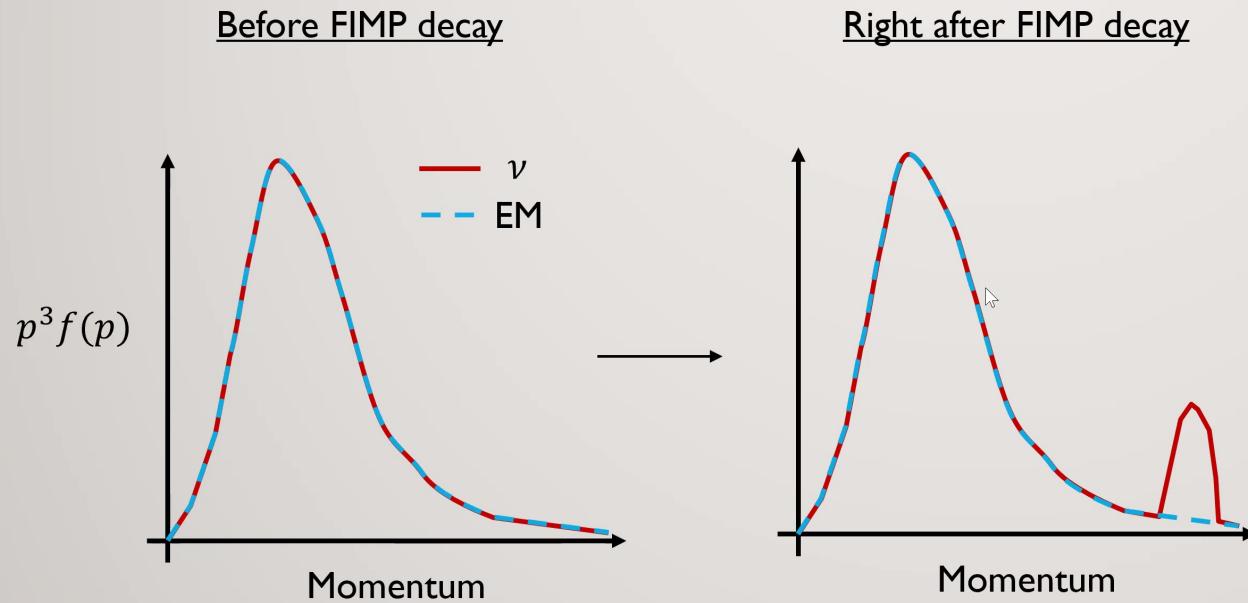
- $\nu$  decoupling not instantaneous
- cross-section  $\sigma \propto E_\nu^2$
- $\nu$ s with higher energies stay longer in equilibrium

What is the actual mechanism behind this effect?

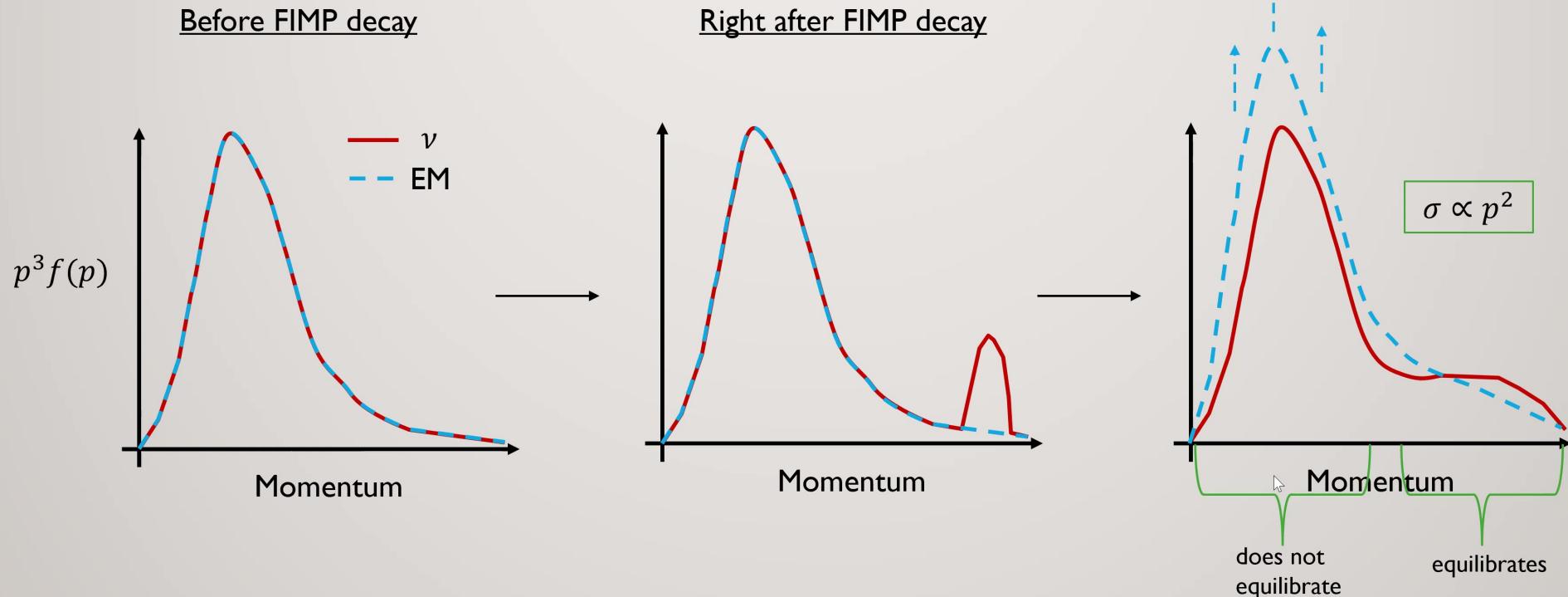
Before FIMP decay

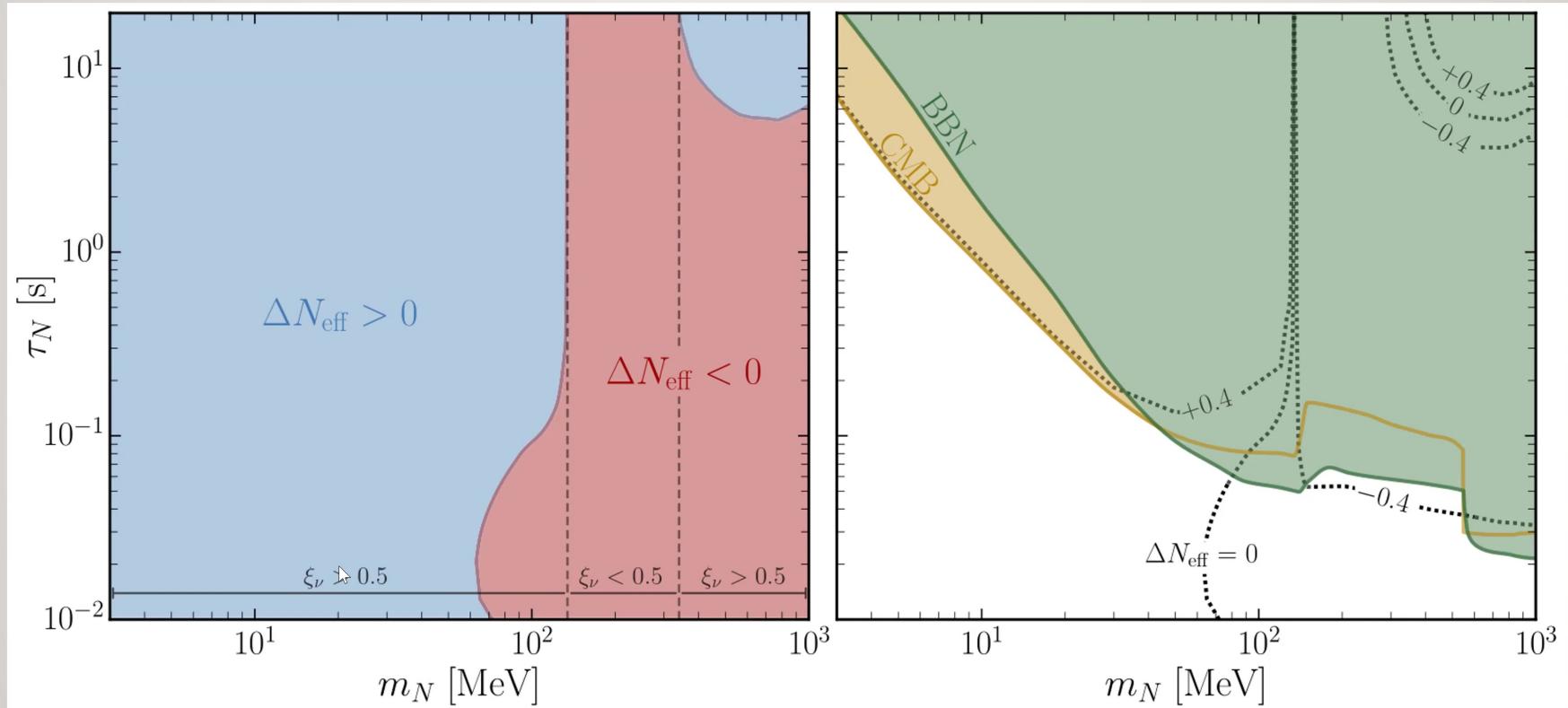


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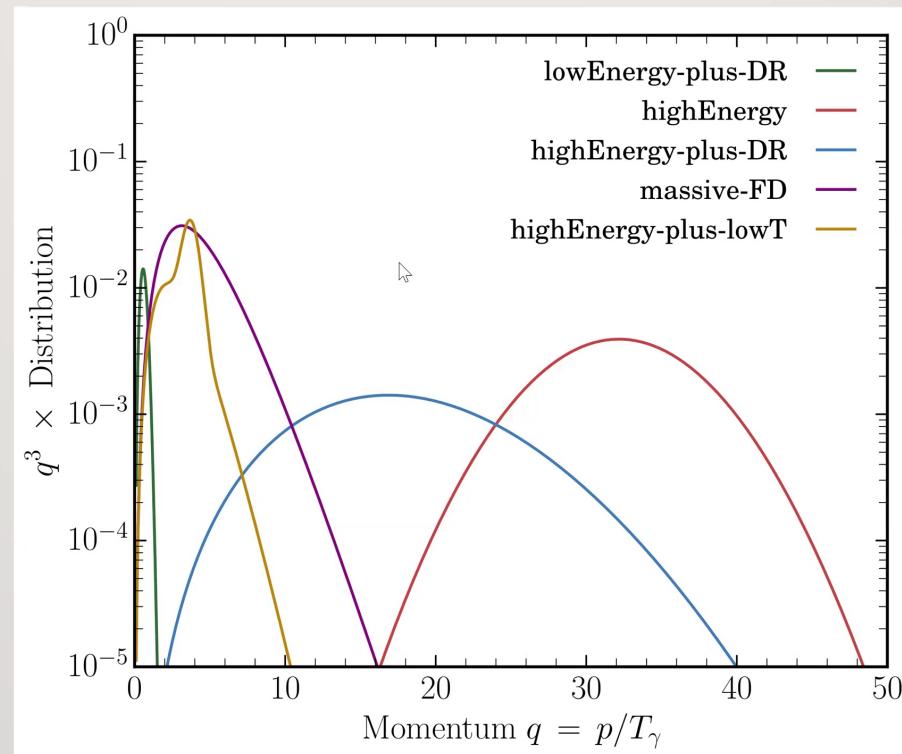
2103.09831

# Signatures at neutrino experiments?

- With current CMB data, we are not able to measure sum of neutrino masses

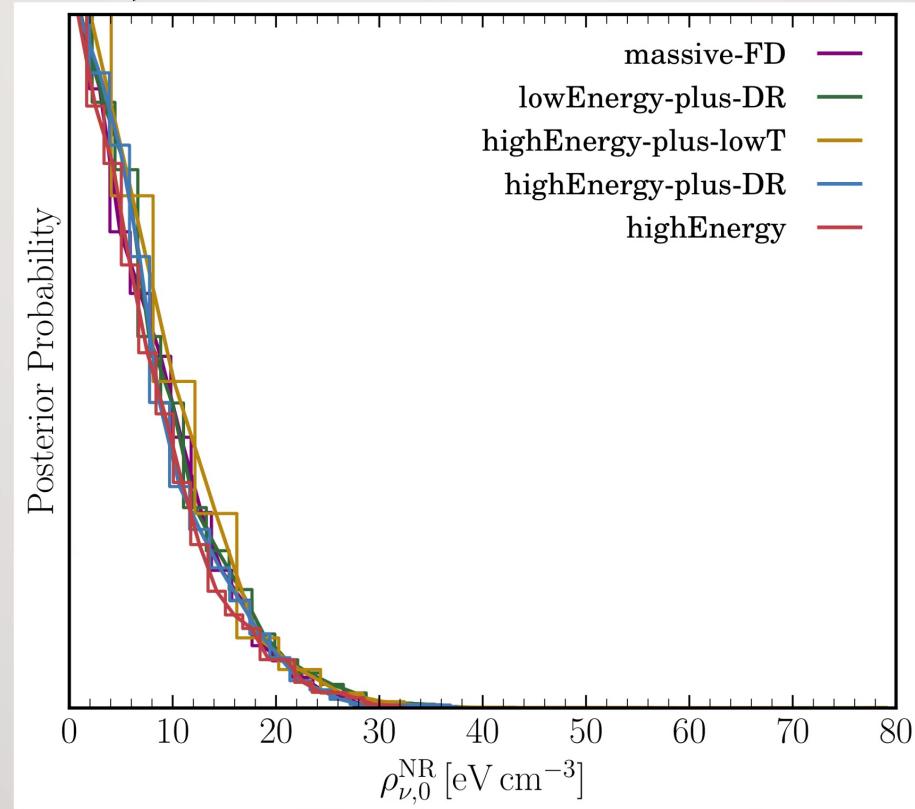
# Signatures at neutrino experiments?

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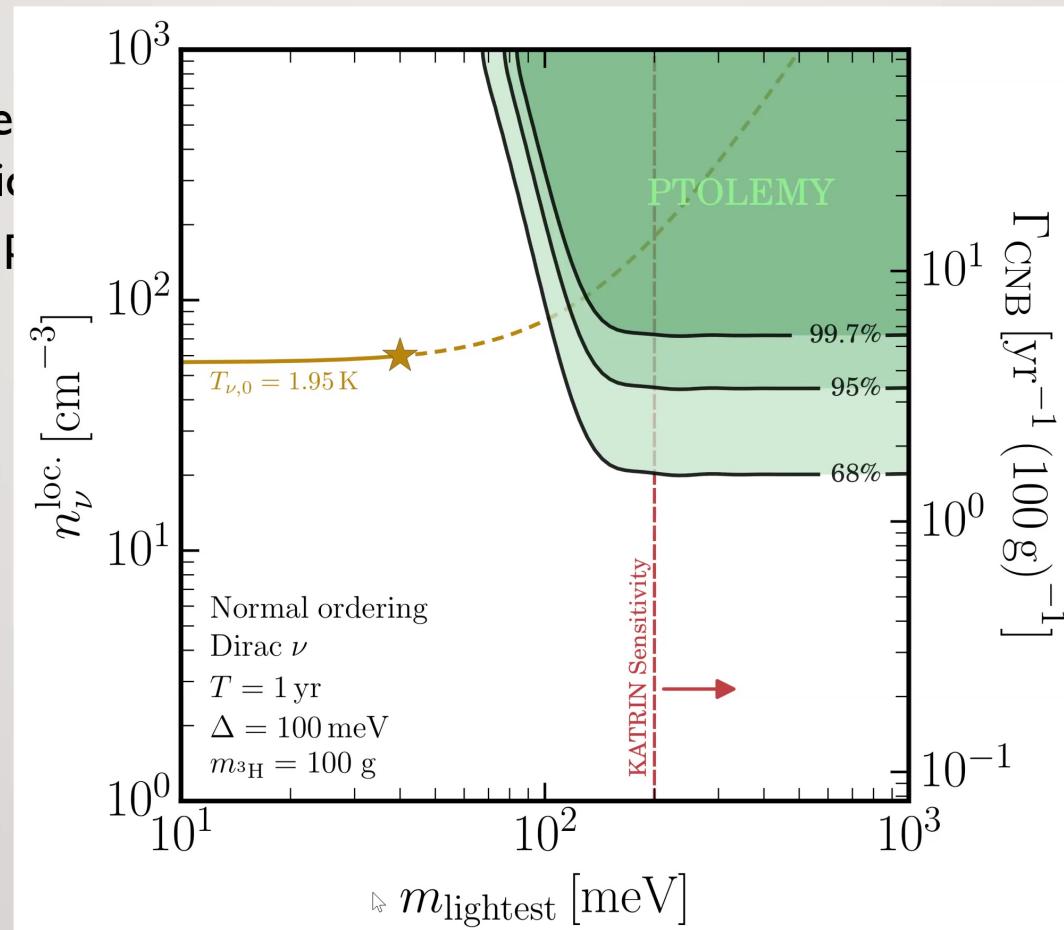


# Signatures at neutrino experiments?

- With current CMB data, we are not able to measure sum of neutrino masses
- Cosmological mass bound can thus be relaxed with SN (NS, Boyarsky, Lesgourgues et al., to appear)

# Signatures at neutrino experiments?

- With current
- Cosmologic
- Detection p



neutrino masses  
(S. Lesgourges et al., to appear)

, PTOLEMY,...

(NS, Alvey, Escudero,  
Schwetz, to appear)

## SUMMARY

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- BBN and the CMB provide robust lower bounds on the lifetime of sterile neutrinos
- SN (and FIMPs in general) can decrease  $N_{\text{eff}}$  even if they inject most of their energy into neutrinos
- Possible signatures at KATRIN/PTOLEMY in the near future?

THANK YOU 😊

