Title: A New Perspective on Time Reversal Motivated by Quantum Gravity

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Abstract: Time Reversal T is usually discussed in the traditional framework of quantum mechanics in which T is represented by an anti-unitary operator. But quantum gravity may well need generalization of standard quantum mechanics which may not preserve even its linear structure, let alone the unitarity of dynamics and anti-unitarity of T. Then the currently used arguments to conclude that T violation is a fundamental aspect of Nature will break down. Fortunately, it turns out that one can analyze the T-violation experiments in a much more general setting, of which classical and quantum mechanics are special cases. The setting does not require a Hilbert space, or linearity of either dynamics or symmetry operations such as T. Nonetheless, somewhat surprisingly, one would still be to use the current experiments to conclude that there is T violation at a fundamental level under rather minimal assumptions on the structure of the final quantum gravity theory.

9:01 AM Tue Jun 15 Q 100% TODOGGO 5 ſĨ Q < new perspective on T-violation A motivated by quantum gravity Abbay Ashtekar Institute for Gravitation & the cosmos & Physics Dept The Pennsylvania State University Quantizing Time Workshop, PI, June 15th 2021 10  $\odot$ 



9:05 AM Tue Jun 15 TODOGBO 5 < ① + : 0 · Experimental progress over some 55 years tells us that a Fundamental force/interaction breaks T invariance. This is a genuinely quantum feature associated with time. But, observations \_\_\_\_\_ conclusion Heavy use of the detailed Minkowski QFT Framework These key elements will not survive when gravity is included. But we don't know the full QG framework ] what is the most general theoretical framework that will still allow us to draw a conclusion on breaking of T invaraince in Nature, using current experiments? Turns out that there is much excas baggage in the current analysis. A 'minimalistic framework suffices. Provides general pointers for QG as well. 





9:17 AM Tue Jun 15 ອ 🗢 🖓 99% 🗖 TODOGBO < ① 5 0 + : 0 conclusion: Assumption that interaction that binds quarks into 1k", and/or cause the decay is CP invariant is incorrect What does this have to do with T invariance? If we assume in addition that standard assumptions of OFT hold : / Locality, spin-statistics, I group is Unitarily implemented, vacuum I invariant) Then CPT is an exact symmetry of S : [CPT, S] = 0 1 [CP, S] ≠0 ⇒ [T, S] ≠0 (since T is Note: We do NOT need to know the Hamiltonian? Very general Thus: K2 -> T+T- => T violation in a fundamental law provided Nature is described by a local relativistic OFT. But this framework is too restrictive to include gravity. 

9:19 AM Tue Jun 15 🔹 穼 ନ 99% 🔳 TODOGBO 5 く① 0 + : 0 1B: Direct evidence for T violation 1998 CPLEAR Kº -10 Rº. 10 2012 BaBat Do not pass through CPT ( KO-KO ascillations ) CPLEAR obsorvations : K° -> Ko Transition occurs with a different probability than the -> x0 Transition Proporty TIXO>=1K> TIZO>=1K> key difference at the kinematical level tom, C, PaCP symmetrics: (2) T is represented by an anti-unitary operator  $(T(\alpha_1|\psi_1>+\alpha_2|\psi_2>=\alpha_1^*|\psi_1>+\alpha_2^*|\psi_2>, (T\phi_1T\psi_2)=(\phi_1\psi_2^*)$ 3 T is a dynamical symmetry : TS = ST (Recall  $T e^{-iHt} = e^{+iHt}$ , or  $T U_t T' = U_t$ ) (D) S is a unitary operator (D) d n't need this assumption for CP violation.) 

9:26 AM Tue Jun 15 97% TODOGED 5 < rî ) + : 0 The theoretical underpinning used to conclude that there is T violation in fundamental interactions is remarkably general. But it did assume: Linear Hilbert space structure of QM; Unitarity of dynamics (S) and the subtle anti-unitary representation of T All of these features may not survive in quantum gravity I will mention concrete ideas - Br generalizations later But simplest examples are Birula's a Weinberg's nonlinear quantum evolution. Already in part 2B, we saw that assumptions needed to go from observations to the startling conclusion that for dominity laws are not T invariant can be streamlined. Can we streamline & generalize further to make the assumptions weaker so that one can still draw conclusions from current observations, with quantum gravity perspective in mind? The answer is in the affirmative. 10 Minimal Bramphak

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Ninimal Framework  
1. There is a set X of states 6 of the system  
(Ex: 147 or P in Quantum Mechanics, both normalized)  
2. X is equipped with an overlap map 
$$\bigcirc$$
: X × X → Eo, 17  
that measures the overlap between 2 states st  
 $\bigcirc$ (Ex:  $\bigcirc$ (P, P') = Tr PP' or  $\bigcirc$ (4, Φ) =  $1 < \Psi | \Phi > 1^{2}$   
3. There is a 1.1 dynamical map S: X; → X<sub>P</sub> , (M:& X<sub>P</sub> are  
copies of X) st  $\bigcirc$ r(G; G;) =  $\bigoplus$ (SG; SG;)  
Prob G; evolveste G' = P(G; G;) =  $\bigoplus$ (SG; SG;)  
4. Symmetrifes : R :  $A_{i} \rightarrow A_{i}$  st  $\bigoplus$ (RG; RG;) =  $\bigoplus$ (G;S;)  
 $T : A_{i} \xrightarrow{\circ} A_{f}$   $\bigoplus$ (TG; TG;) =  $\bigoplus$ (G;S;)  
 $(Hpr) complex organism uncessed)$   
Def : T is a dynamical symmetry iff s'T = T'S Q

9:33 AM Tue Jun 15 5 TODOGKO Q < m considerations of 1A on CP violation go through in general mechanics straightforwardly. I will consider the more interesting cose 1.8 on T violation Result: If T is a dynamical symmetry. Then  $P(\sigma_r, \sigma_i) = P(T\sigma_i, T'\sigma_f)$ Notation:  $\lambda 5 = 5 \qquad \lambda 5' = 5'$  $\overline{\Theta}_{i}(\overline{\sigma}_{i},\overline{\sigma}_{i}') = \overline{\Theta}_{i}(\overline{S}\overline{\sigma}_{i}, \overline{S}\overline{\sigma}_{i}') = \overline{\Theta}_{i}(\overline{S}\overline{\sigma}_{i}, \overline{S}\overline{\sigma}_{i}') = \overline{\Theta}_{i}(\overline{\sigma}_{i}', \overline{S}\overline{\sigma}_{i}')$   $\overline{P}_{ropetty of S} \quad symmetry of \overline{\Theta}_{i} \quad betwore \overline{\sigma}_{i}'$ property N  $(T_{\sigma_i}, T_{\sigma_i}) = \bigoplus_{\substack{i \in \mathcal{F} \\ \forall \sigma_i \in \mathcal{F}}} (T_{\sigma_i}, T_{\sigma_i}) = \bigoplus_{\substack{i \in \mathcal{F} \\ \forall \sigma_i \in \mathcal{F}}} (T_{\sigma_i}, T_{\sigma_i}) = \bigoplus_{\substack{i \in \mathcal{F} \\ \forall \sigma_i \in \mathcal{F}}} (T_{\sigma_i}, S_{\tau_i}) = \bigoplus_{\substack{i \in \mathcal{F} \\ \forall \sigma_i \in \mathcal{F}}} (T_{\sigma_i}, S_{\tau_i}) = \bigoplus_{\substack{i \in \mathcal{F} \\ \forall \sigma_i \in \mathcal{F}}} (T_{\sigma_i}, S_{\tau_i}) = \bigoplus_{\substack{i \in \mathcal{F} \\ \forall \sigma_i \in \mathcal{F}}} (T_{\sigma_i}, S_{\tau_i}) = \bigoplus_{\substack{i \in \mathcal{F} \\ \forall \sigma_i \in \mathcal{F}}} (T_{\sigma_i}, S_{\tau_i}) = \bigoplus_{\substack{i \in \mathcal{F} \\ \forall \sigma_i \in \mathcal{F}}} (T_{\sigma_i}, S_{\tau_i}) = \bigoplus_{\substack{i \in \mathcal{F} \\ \forall \sigma_i \in \mathcal{F}}} (T_{\sigma_i}, S_{\tau_i}) = \bigoplus_{\substack{i \in \mathcal{F} \\ \forall \sigma_i \in \mathcal{F}}} (T_{\sigma_i}, S_{\tau_i}) = \bigoplus_{\substack{i \in \mathcal{F} \\ \forall \sigma_i \in \mathcal{F}}} (T_{\sigma_i}, S_{\tau_i}) = \bigoplus_{\substack{i \in \mathcal{F} \\ \forall \sigma_i \in \mathcal{F}}} (T_{\sigma_i}, S_{\tau_i}) = \bigoplus_{\substack{i \in \mathcal{F} \\ \forall \sigma_i \in \mathcal{F}}} (T_{\sigma_i}, S_{\tau_i}) = \bigoplus_{\substack{i \in \mathcal{F} \\ \forall \sigma_i \in \mathcal{F}}} (T_{\sigma_i}, S_{\tau_i}) = \bigoplus_{\substack{i \in \mathcal{F} \\ \forall \sigma_i \in \mathcal{F}}} (T_{\sigma_i}, S_{\tau_i}) = \bigoplus_{\substack{i \in \mathcal{F} \\ \forall \sigma_i \in \mathcal{F}}} (T_{\sigma_i}, S_{\tau_i}) = \bigoplus_{\substack{i \in \mathcal{F} \\ \forall \sigma_i \in \mathcal{F}}} (T_{\sigma_i}, S_{\tau_i}) = \bigoplus_{\substack{i \in \mathcal{F} \\ \forall \sigma_i \in \mathcal{F}}} (T_{\sigma_i}, S_{\tau_i}) = \bigoplus_{\substack{i \in \mathcal{F} \\ \forall \sigma_i \in \mathcal{F}}} (T_{\sigma_i}, S_{\tau_i}) = \bigoplus_{\substack{i \in \mathcal{F} \\ \forall \sigma_i \in \mathcal{F}}} (T_{\sigma_i}, S_{\tau_i}) = \bigoplus_{\substack{i \in \mathcal{F} \\ \forall \sigma_i \in \mathcal{F}}} (T_{\sigma_i}, S_{\tau_i}) = \bigoplus_{\substack{i \in \mathcal{F} \\ \forall \sigma_i \in \mathcal{F}}} (T_{\sigma_i}, S_{\tau_i}) = \bigoplus_{\substack{i \in \mathcal{F} \\ \forall \sigma_i \in \mathcal{F}}} (T_{\sigma_i}, S_{\tau_i}) = \bigoplus_{\substack{i \in \mathcal{F} \\ \forall \sigma_i \in \mathcal{F}}} (T_{\sigma_i}, S_{\tau_i}) = \bigoplus_{\substack{i \in \mathcal{F} \\ \forall \sigma_i \in \mathcal{F}}} (T_{\sigma_i}, S_{\tau_i}) = \bigoplus_{\substack{i \in \mathcal{F} \\ \forall \sigma_i \in \mathcal{F}}} (T_{\sigma_i}, S_{\tau_i}) = \bigoplus_{\substack{i \in \mathcal{F} \\ \forall \sigma_i \in \mathcal{F}}} (T_{\sigma_i}, S_{\tau_i}) = \bigoplus_{\substack{i \in \mathcal{F} \\ \forall \sigma_i \in \mathcal{F}}} (T_{\sigma_i}, S_{\tau_i}) = \bigoplus_{\substack{i \in \mathcal{F} \\ \forall \sigma_i \in \mathcal{F}}} (T_{\sigma_i}, S_{\tau_i}) = \bigoplus_{\substack{i \in \mathcal{F} \\ \forall \sigma_i \in \mathcal{F}}} (T_{\sigma_i}, S_{\tau_i}) = \bigoplus_{\substack{i \in \mathcal{F} \\ \forall \sigma_i \in \mathcal{F}}} (T_{\sigma_i}, S_{\tau_i}) = \bigoplus_{\substack{i \in \mathcal{F} \\ \forall \sigma_i \in \mathcal{F}}} (T_{\sigma_i}, S_{\tau_i}) = \bigoplus_{\substack{i \in \mathcal{F} \\ \forall \sigma_i \in \mathcal{F}}} (T_{\sigma_i}, S_{\tau_i}) = \bigoplus_{\substack{i \in \mathcal{F} \\ \forall \sigma_i \in \mathcal{F}}} (T_{\sigma_i}, S_{\tau_i}) = \bigoplus_{\substack{i \in \mathcal{F} \\ \forall \sigma_i \in \mathcal{F}}} (T_{\sigma_i}, S_{\tau_i}) = \bigoplus_{\substack{i \in \mathcal{F} \\ \forall \sigma_i \in \mathcal{F}}} (T_{\sigma_i}, S_{\tau_i}) = \bigoplus_{\substack{i \in \mathcal{F} \\ \forall \sigma_i \in \mathcal{F}}} (T_{\sigma_i}, S_{\tau_i}) = \bigoplus_{\substack{i \in \mathcal{F} \\ \forall \sigma_i \in \mathcal{F}}} (T_{\sigma_i}, S_{\tau_i}) = \bigoplus_{\substack{i \in \mathcal{F} \\ \forall \sigma_i \in \mathcal{F}}} (T_{\sigma_i}, S_{\tau_$ (of Soi) = Of (To; ST'of) Prob T'O' -> TG: 10 61 -> Or FOR K°, K° Prob K° -> K° = Prob T' K° -> T K° = Prob Hence observations => T is not a dynamical symmetry 

9:37 AM Tue Jun 15 😐 穼 🖸 94% 📩 5 TODOGED Q < ① + : 0 · The minimalistic framework suffices to conclude from observations that whatevar the fundamental laws are that govern the Ko- Ko (and B-Bo) system (laws that bind the quarks to form these mesons and that govern their decays they violate Time Reversal invariance provided the underlying Framework has our minimal ingredients X, @:ovalap, T, S satisfying elementary, physically notivated requirements. · standard Quantum mechanics provides them but with a lot of excress baggage, Hilbert space (not, just x). Scalar product (not just @), Anti unitary operator T (not just a 1-1, onto map on & presoning (D), and unitary S (not just 1-1, on to map preserving (). · Putative QG theory could realize the minimal ingredients via a generalization of the std QM framework, e.g. by recasting the Asymptotic Quantization scheme in a generalized QM framework I now describe, 10 

9:38 AM Tue Jun 15 TODOGBO Q < ſŢ Example of Generalized QM Relativity Motivation: special Relativity General Linear (or affine) structure Non-linear (Manifold) structure > Non-linear/Manifold (??) Hilbert space of GM special Diffees Linear/Antilinear operators (??) .----> concrete illustration to indicate the direction Geometric Reformulation of QM (~Mink Reformulation (AA + Troy schilling) of Special relativi ist step: of special relativity linearity @ forfort Non-linear geometric structures -> space of roys = Kähler manifold KINEMATICS Hilbert space KUIΦ/12 -> f(Geodesic distance between points) () Diffeos preserving kähler structure (C, P, T, as well as committing) Symmetries -Fris whose Hamiltonian NFs are killing VES 10 SA operators (PBs on co-dim kähler manifold X) (commutators) mechanics" are

9:42 AM Tue Jun 15 5 0000BD rT < lineanity a forfort Non-linear geometric structures  $(\mathbf{X})$ space of roys = kähler manifold KINEMATICS Hilbert space f(Geodesic distance between points) (D)  $|\langle \psi | \phi \rangle|^2$ Diffeos preserving kähler structure (C, P, T, as well as cosymmetries) Symmetries Fris whose Hamiltonian VFs are killing VES SA operators (PBs on co-dim kähler manifold X) (commutators) Note: kinematic structures of "general mechanics" are concretely realized using geometry. All necessary relations between X, @, C.P. CPT, T.... satisfied W/6 using linearity. Hamiltonian flow on V (generated  $i\hbar \partial_{+} |\Psi\rangle = H$ 7 by the function corresponding to H) Discrete (a cont.) diffeos on X symmetries eg C, P, T, .. (Rotations) preserving kähler studue and the Hamiltonian flow. din'n abaria