Title: Pseudospectrum and black hole quasi-normal mode (in)stability

Speakers: Rodrigo Panosso Macedo

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Abstract: We study the stability of quasi-normal modes (QNM) in asymptotically flat black hole spacetimes by means of a pseudospectrum analysis. The construction of the Schwarzschild QNM pseudospectrum reveals: i) the stability of the slowest decaying QNM under perturbations respecting the asymptotic structure, reassessing the instability of the fundamental QNM discussed by Nollert (1996); ii) the instability of all overtones under small scale perturbations of sufficiently high frequency, that migrate to a universal class of QNM branches along pseudospectra boundaries, shedding light on Nollert & Price's analysis (1996).

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PSEUDOSPECTRUM AND BLACK-HOLE QUASI-NORMAL MODE (IN)STABILITY

Rodrigo Panosso Macedo J. L. Jaramillo, L. Al Sheik







2004.06434 (to appear in PRX) - 2105.03451





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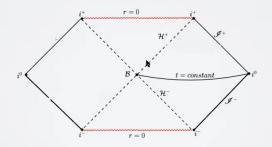
 Wave equation on (spherically sym.) black-hole background

$$-\Psi_{,\bar{t}\bar{t}} + \Psi_{,xx} - \mathcal{P}\Psi = 0$$

Phenomenological approach:

Ansatz for time dependence (Fourier modes)

$$\Psi = e^{-i\omega \bar{t}} \psi(x) \qquad \Longrightarrow \quad \left[\partial_{xx}^2 - \mathcal{P}\right] \psi = -\omega^2 \psi$$



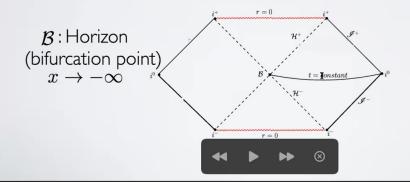
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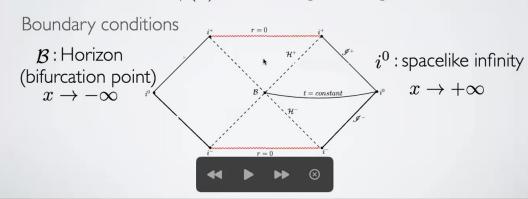
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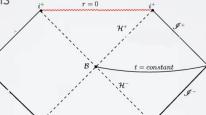
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Boundary conditions

B: Horizon (bifurcation point)

 $\psi \sim e^{-i\omega x}$



 i^0 : spacelike infinity

$$x \to +\infty$$

$$\psi \sim e^{i\omega x}$$

•Comparison to normal Dege value problem modes:

Def. Given matrix L and its adjoint L^{\dagger} . Then L is <u>normal</u> iff $[L,L^{\dagger}]=0$

Ex. Symmetric, hermitian, orthogonal, unitary...

Spectral Theo. (matrices): L is normal iff is unitarily diagonalisable

Obs. Theorem extends to normal operators in Hilbert (or Banach) space



Hermitian Physics (self adjoint operators):

Eigenvectors are orthogonal and form complete basis Eigenvalues are stable $L \to L + \epsilon \delta L \Rightarrow \lambda \to \lambda + \epsilon \delta \lambda$

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Completeness more difficult to study If $[L,L^\dagger]
eq 0$ Eigenvectors not necessarily orthogonal Spectral Instabilities

Non-Hermitian Physics (non-self adjoint operator):

$$\epsilon < 1$$

$$L(\epsilon) = L + \epsilon \delta L, \quad |\delta L| = 1$$
$$|\lambda_i(\epsilon) - \lambda_i| \le \kappa_i \epsilon$$

Eigenvalue Condition Number:
$$\kappa_i = \frac{||r_i|| \, ||l_i||}{|r_i \cdot l_i|}$$

$$L\,r_i = \lambda_i\,r_i$$
 (Right eigenvector) self adjoint operator:

$$L^\dagger \, l_i = ar{\lambda}_i \, l_i$$
 (Left eigenvector) $\underline{\qquad L = L^\dagger \Rightarrow r_i \parallel l_i \Rightarrow \kappa_i = 1}$

Small perturbation in operator; small dis lacement in eigenvalue

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Completeness more difficult to study If $[L,L^{\dagger}] \neq 0$ Eigenvectors not necessarily orthogonal Spectral Instabilities

Non-Hermitian Physics (non-self adjoint operator):

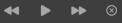
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 (Left eigenvector) $L
eq L^\dagger \stackrel{?}{\Rightarrow} r_i
mid l_i \Rightarrow \kappa_i > 1$



Completeness more difficult to study If $[L,L^{\dagger}]
eq 0$ Eigenvectors not necessarily orthogonal Spectral Instabilities

Non-Hermitian Physics (non-self adjoint operator): $\operatorname*{Im}(\lambda)$

$$\epsilon < 1$$

$$L(\epsilon) = L + \epsilon \delta L, \quad |\delta L| = 1$$

Spectra:

$$\sigma(L): ||L-\lambda_i \mathbb{I}|| = 0$$
 (Eigenvalue)

 $Re(\lambda)$

Completeness more difficult to study If $[L,L^{\dagger}]
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Non-Hermitian Physics (non-self adjoint operator):

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 $\operatorname{Im}(\lambda)$

Spectra:

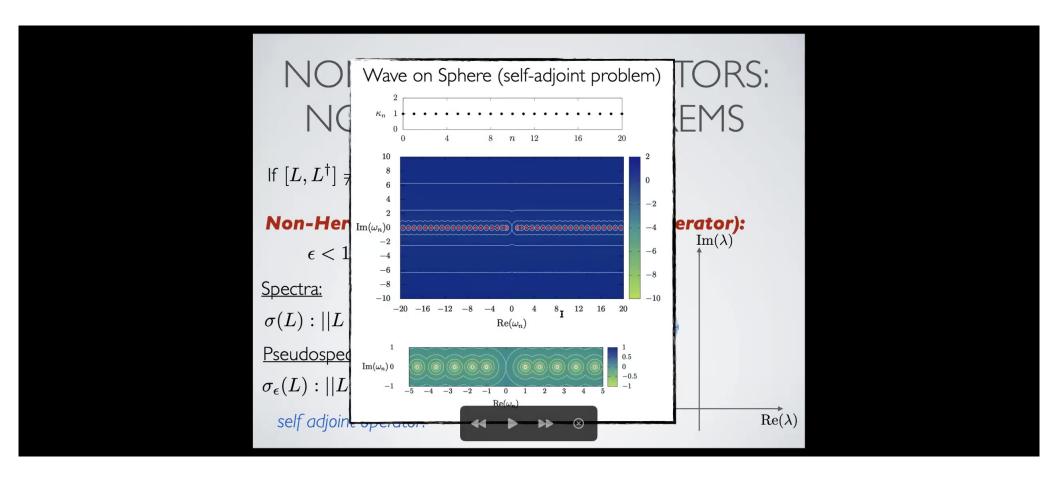
$$\sigma(L): ||L - \lambda_i \mathbb{I}|| = 0$$
 (Eigenvalue)

 λ_i

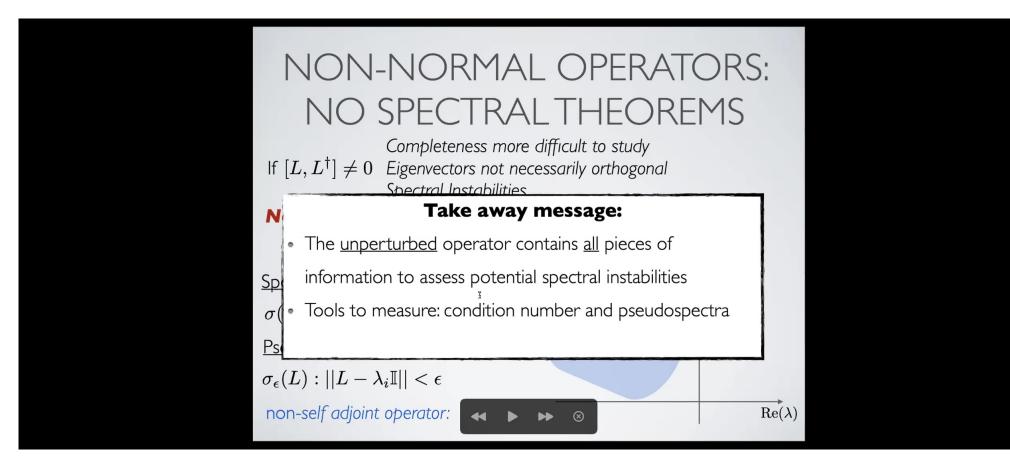
Pseudospectra:

$$\sigma_{\epsilon}(L): ||L - \lambda_i \mathbb{I}|| < \epsilon$$





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Completeness more difficult to study If $[L,L^{\dagger}] \neq 0$ Eigenvectors not necessarily orthogonal Spectral Instabilities

Non-Hermitian Physics (non-self adjoint operator):

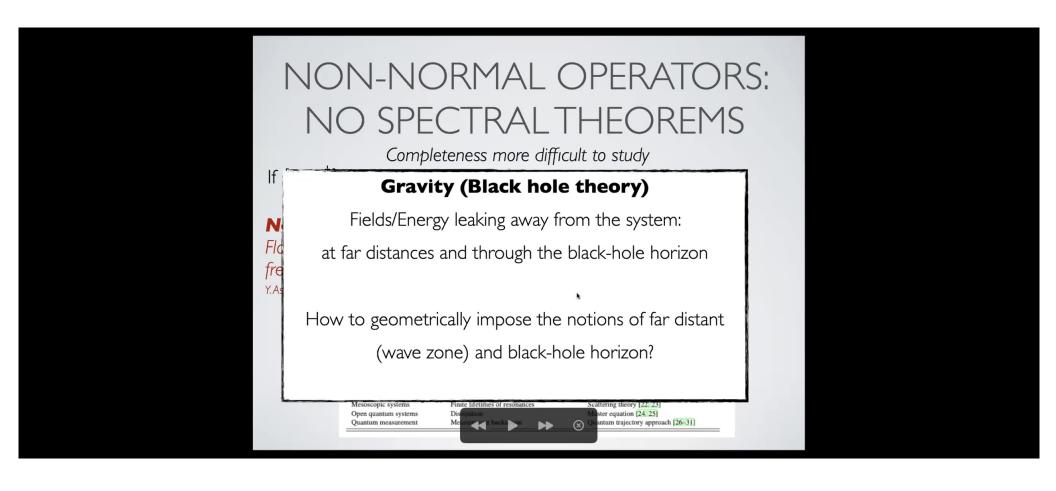
Flow of energy, particles and information to external degrees of freedom out of the Hilbert space

Y. Ashida, Z. Gong, M. Ueda 2006.01827 (2020)

Table 1. A wide variety of classical and quantum systems described by non-Hermitian matrices/operators together with their physical origins of non-Hermiticity, presented in order of appearance in the present review.

Systems / Processes	Physical origin of non-Hermiticity	Theoretical methods
Photonics	Gain and loss of photons	Maxwell equations [12, 13]
Mechanics	Friction	Newton equation [14, 15]
Electrical circuits	Joule heating	Circuit equation [16]
Stochastic processes	Nonreciprocity of state transitions	Fokker-Planck equation [17] 18]
Soft matter and fluid	Nonlinear instability	Linearized hydrodynamics [19-21]
Nuclear reactions	Radiative decays	Projection methods [4-6]
Mesoscopic systems	Finite lifetimes of resonances	Scattering theory [22, 23]
Open quantum systems	Dissipation	Master equation [24, 25]
Quantum measurement	Measure of backs on	Quantum trajectory approach [26-31]

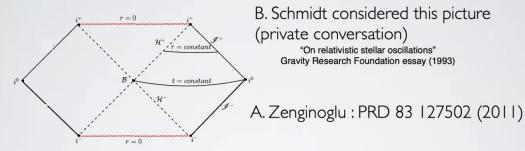
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HYPERBOLOIDAL SLICES

•New coordinate system: spacelike hypersurfaces of constant time extending between black-hole horizon \mathcal{H}^+ (horizon-penetrating) and future null infinity \mathscr{I}^+



"Black hole perturbation theory is typically studied on time surfaces that extend between the bifurcation sphere and spatial infinity. From a physical point of view, however, it may be favourable to employ time surfaces that extend between the future event horizon and future null infinity. This framework resolves problems regarding the representation of quasinormal mode eigenfunctions (...)"

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New co constar (horizo

Boundary Conditions X Regularity

Ingoing/Outgoing external boundary conditions in the original problem are now geometrically taken into account via hyperboloidal slices.

 $n \mathcal{H}^+$

2011)

At the differential equation level, the physical scenario (Black Hole+Wave Zone) is described by the equations' regular solutions

spectral problem of a non-self adjoint operator"

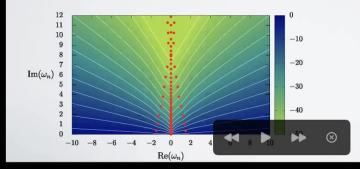
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SPECTRAL INSTABILITY

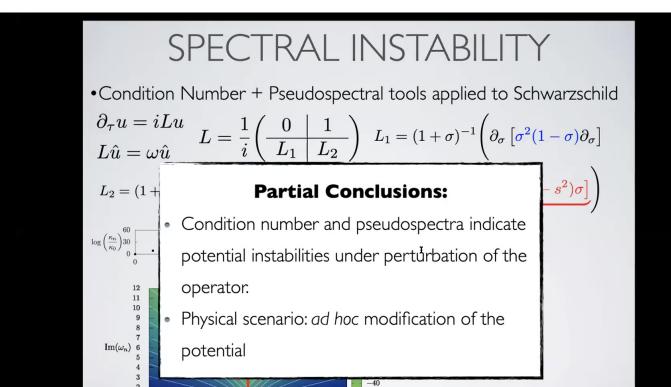
•Condition Number + Pseudospectral tools applied to Schwarzschild

$$L_2 = (1+\sigma)^{-1} \Biggl((1-2\sigma^2) \partial_\sigma - 2\sigma \Biggr) \qquad \sigma \in [0,1] \qquad - \underbrace{ \left[\ell(\ell+1) + (1-s^2)\sigma \right]}_{q_\ell} \Biggr)$$





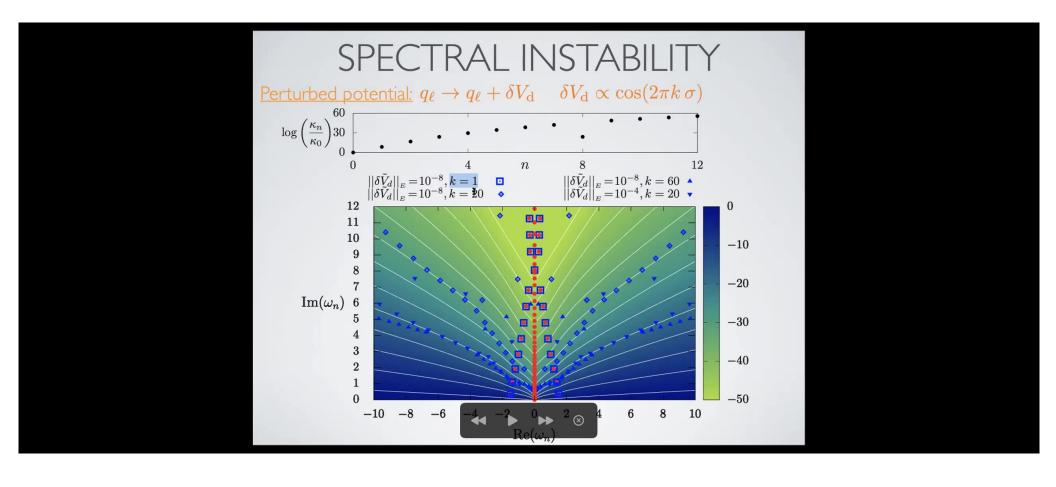
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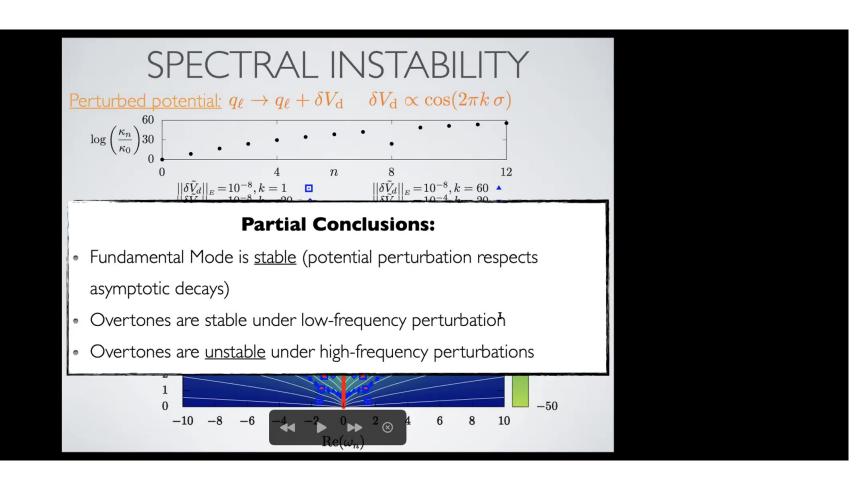
-10 -8 -6 -4 -2 0 2

 $\operatorname{Re}(\omega_n)$

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CONCLUSION

Hyperboloidal Framework for perturbation theory:

Field has reached a mature stage to offer a rich set of theoretical and numerical tools to expand studies in gravitational wave physics (mathematical relativity, theoretical physics, astrophysics and gravitational wave astronomy)

Fundamental aspects of black-hole physics:

Use of tools from theory of scattering resonances and non-self adjoint operators — pseudospectra — to assess structural aspects of QNM spectra.

Perspectives:

<u>Astrophysics and Cosmology</u>: Black hole spectroscopy; universality of compact object QNM

<u>Fundamental gravitational physics</u>: (Sub)planckian-scale physics; QNMs and strong cosmic censorship; spacetime semiclassical limits

Mathematical relativity: Formal statements of the numerical experiements

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