

Title: Path-integral Optimization and AdS/CFT

Speakers: Tadashi Takayanagi

Collection: Tensor Networks: from Simulations to Holography III

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Abstract: In this talk we will start with a review of path-integral optimization, which provides a useful description of non-unitary tensor networks for Euclidean path-integrals in CFTs. We will explain an emergence of AdS geometry in this method and an interpretation as a computational complexity. Next we will give its application to analytical calculations of entanglement of purification, which was quite recently reproduced by numerical calculations. Finally, we would like to present a derivation of a path-integral optimization method directly from the AdS/CFT.

TENSOR NETWORKS: FROM SIMULATIONS TO HOLOGRAPHY III
@Perimeter, Nov.16-20, 2020



Path-integral Optimization and AdS/CFT
(A continuous approach to “AdS =Tensor Network”)

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① Introduction: AdS/CFT and Tensor Network

② Path-Integral Optimization in CFTs

* Caputa-Kundu-Miyaji-Watanabe-TT [Phys.Rev.Lett. 119 (2017) 071602]

* Caputa-Kundu-Miyaji-Watanabe-TT [JHEP 11 (2017) 097]

* Bhattacharyya-Caputa-Das-Kundu-Miyaji-Watanabe-TT [JHEP 07 (2018) 086]

③ An Application: Entanglement of Purification

* Caputa-Miyaji-Umemoto-TT [Phys.Rev.Lett. 122 (2019) 11, 111601]

* Camargo-Hackl-Heller-Jahn-TT, arXiv: 2009.11881

④ PI-Opt. from Hartle-Hawking Wave function in AdS

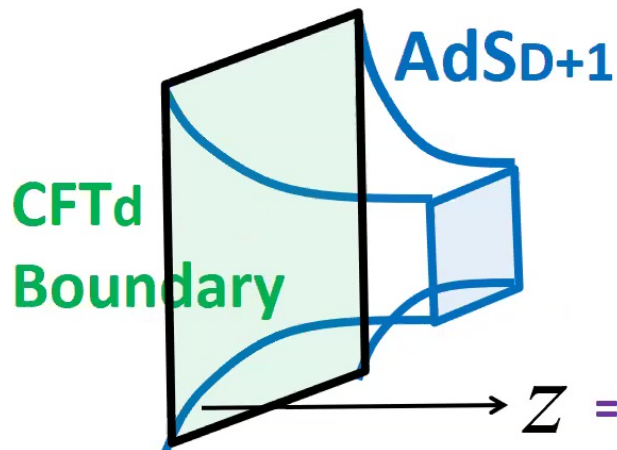
* Boruch-Caputa-TT arXiv:2011.08188

⑤ Conclusions

① Introduction

AdS/CFT [Anti de-Sitter space/Conformal Field Theory] [Maldacena 1997]

Gravity on D+1 dim. AdS = D dim. CFT on bdy



Metric of AdS_{d+1}

$$ds^2 = R^2 \cdot \frac{dz^2 - dt^2 + \sum_{i=1}^{d-1} dx_i^2}{z^2}$$

Z = a length scale of RG flow

⇒ "Geometrization" of Renormalization Group Flow
[AdS = Real Space Renormalization]



An essence of this emergence of geometry from QFTs is

Quantum entanglement

⇒ “Geometry” of Quantum States in many-body system

In AdS/CFT, the entanglement entropy is computed as the area of minimal surface [Ryu-TT 06, Hybeny-Rangamani-TT 07]

Boundary = CFT_d **Bulk = AdS_{d+1}**

$$S_A = -\text{Tr}[\rho_A \log \rho_A], \quad \rho_A = \text{Tr}_B \rho_{tot} .$$

$$S_A = \frac{\text{Area}(\gamma_A)}{4G_N} \approx \frac{\text{Area}(\gamma_A)}{l_{pl}^{d-1}} .$$

Area in the unit of Planck length

γ_A : Minimal Area surface

Planck length ~ 1 qubit

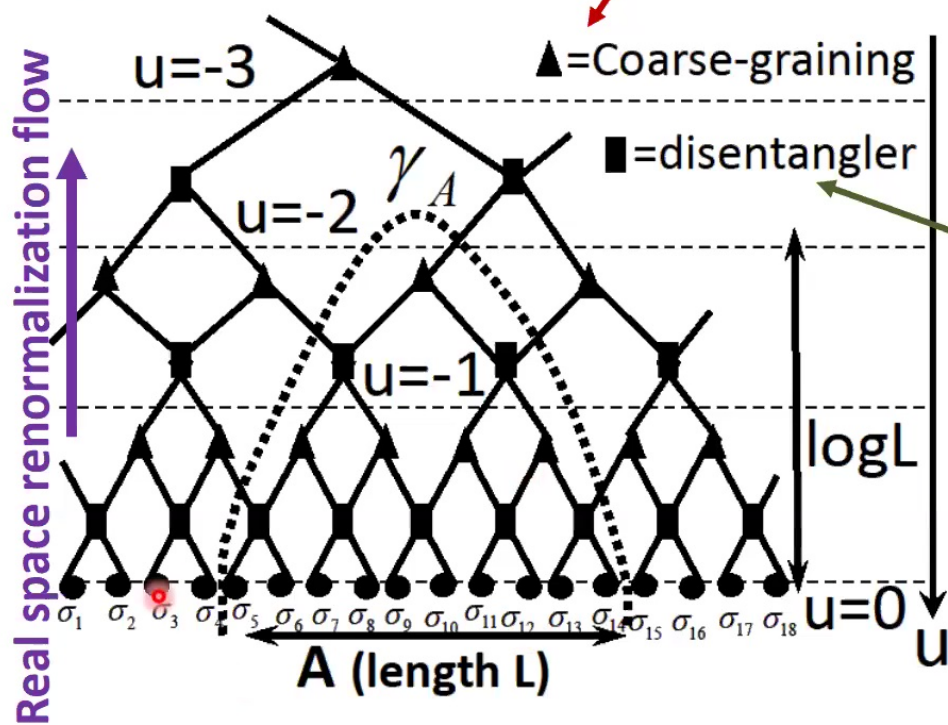


Spacetime in gravity = Collections of bits of entanglement
 = Real space renormalization ?
 ⇒ Tensor Networks !



MERA [Vidal 05]

Coarse-graining = Isometry



$$[T]_{abc}^\dagger [T]_{bcd} = \delta_{ad}$$

$$\begin{array}{c}
 a \quad b \\
 \diagdown \quad / \\
 \square \\
 / \quad \diagdown \\
 c \quad d
 \end{array}
 = a \quad d$$

Disentangler = Unitary trf.

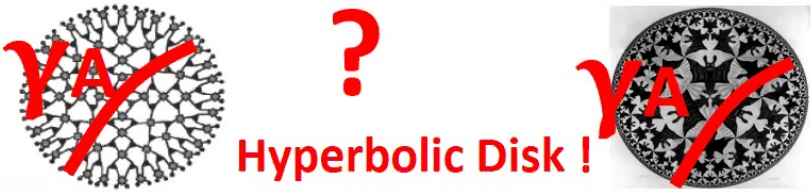
$$S_A \leq \text{Min}[\# \text{links}]$$

$$\propto \log L$$

⇒ agrees with results in 2d CFT !

This suggests
a connection:
[Swingle 2009]

TN of MERA = A time slice of AdS



Hyperbolic Disk !



However, \exists many difficulties: [Beny 2011, Bao et.al. 2015, Czech et.al. 2015]

- The MERA has a causal structure but the time slice does not ?
- Why is the EE bound saturated ?
- How can we find the sub-AdS locality ? (i.e. approximation by GR ?)

Though there are improvements of models/interpretations:

[Perfect TN: Pastawski-Yoshida-Harlow-Preskill 15]

[Kinematical Space: Czech-Lamprou-McCandlish-Sully 15]

[Random TN: Hayden-Nezami-Qi-Thomas-Walter-Yang 16]

[Matchgate TN: Jahn-Gluza-Pastawski-Eisert 17] [Hyper inv. TN: Evenbly 17]

[P-adicTN: Bhattacharyya-Hung-Lei-Li 17] [TN geometry: Misted-Vidal 18],.....

, the true connection between TN and AdS is not still clear.

Our Key Observations towards AdS/TN

[1] A disadvantage of the original TN approach is the presence of *discretized structure*. The genuine AdS/CFT is continuous.

[2] To describe a time slice of AdS, it looks better to regard it as a *Euclidean evolution* rather than Lorentzian one.



A Time slice of AdS

= Continuous TN from Euclidean CFT path-integrals ?

Motivated by this, we give an alternative approach **based on path-integrals**, related to a continuum limit of TNs.

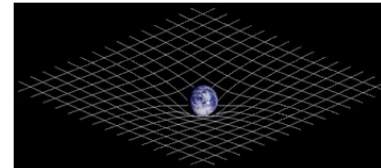
⇒ Path-integral Optimization !





➡ We will describe the discretized lattice structure by **the metric in a continuous way** and remove unnecessary tensors to make the TN efficient.

Our guiding principle



Eliminating unnecessary tensors in TN for a given state
= Creating the most efficient TN (= **Optimization of TN**)

↔ Solving the dynamics of Gravity (i.e. Einstein eq.)

② Path-Integral Optimization and CFT

(2-1) Formulation

A Basic Rule: Simplify a path-integral s.t. it produces the correct UV wave functional.

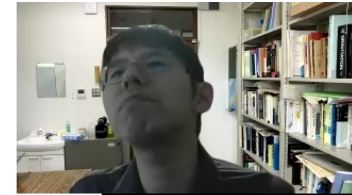
Consider 2D CFTs for simplicity. (z = Euclidean time, x = space)

**Deformation of discretizations in path-integral
= Curved metric such that one cell (bit) = unit length.**

$$\rightarrow ds^2 = e^{2\phi(x,z)} (dx^2 + dz^2).$$

Note: The original flat metric is given by (ϵ is UV cutoff):

$$ds^2 = \epsilon^{-2} \cdot (dx^2 + dz^2).$$



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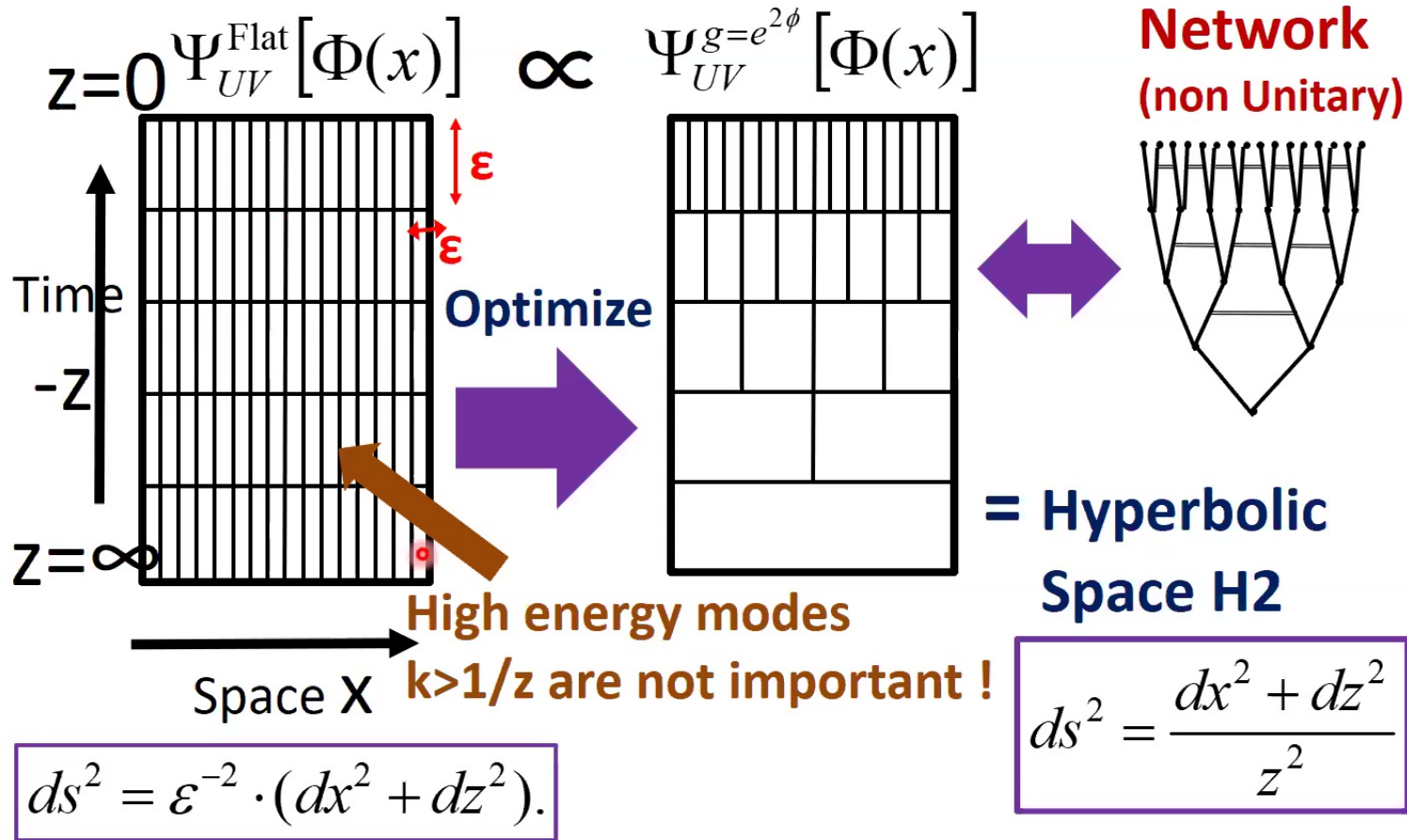
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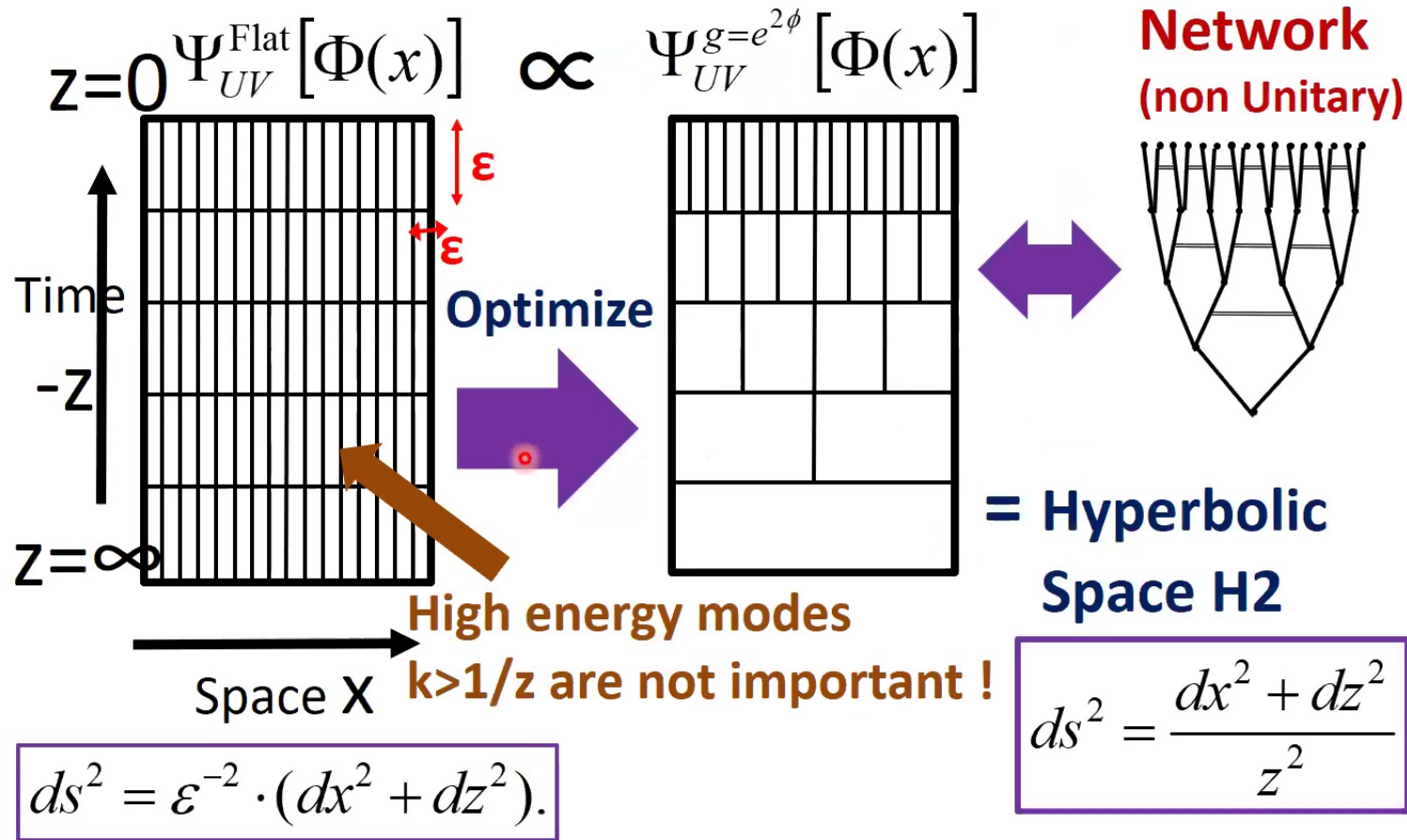
$$ds^2 = \epsilon^{-2} \cdot (dx^2 + dz^2).$$



A Sketch: Optimization of Path-Integral



A Sketch: Optimization of Path-Integral



The wave functional for CFT vacuum is given by

$g_{ab}(x,z)$: background metric

$$\Psi_{UV}^g[\Phi(x)] = \int \prod_{\substack{0 < z < \infty \\ -\infty < x < \infty}} D\Phi(x, z) e^{-S_{CFT}(\Phi)} \cdot \delta(\Phi(x) - \Phi(x, z=0))$$

In CFTs, owing to the Weyl invariance, we have

$$\Psi_{UV}^{g_{ab}=e^{2\phi}\delta_{ab}}[\Phi(x)] = \exp(I[\phi(x, z)]) \cdot \Psi_{UV}^{\text{Flat}}[\Phi(x)].$$

Optimized wf.

Original wf.

Our Proposal (Optimization of Path-integral for CFTs):

Minimize $I[\phi(x, z)]$ w.r.t $\phi(x, z)$
with the boundary condition $e^{2\phi} \big|_{z=\epsilon} = \epsilon^{-2}$.



Motivation of Our Proposal

The normalization N of wave functional estimates repetitions of same operations of path-integration.

→ **Minimize this !**

⇒ **Our conjecture:**

$$\begin{aligned} \text{Min}[I[\varphi]] &\sim \text{Min\# of Tensors in TN} \\ &\approx \text{Computational complexity of TN state} \end{aligned}$$

[For a justification based on circuit complexity, refer to Camargo-Heller-Jefferson-Knaute 2019]

[cf. Holographic Complexity: Susskind 2014, Brown-Roberts-Susskind-Swingle-Zhao 2015, Lehner-Myers-Poisson-Sorkin 2016, Chapman-Marrochio-Myers 2016,...]



For 2D CFTs, $I[\phi]$ is given by the Liouville action !



$$I[\phi] = \text{Log} \left[\frac{\Psi_{g=e^{2\phi} \delta_{ab}}}{\Psi_{g=\delta_{ab}}} \right] = S_L[\phi],$$

Liouville Action

of Isometries
[Czech 17]



of Unitaries



$$S_L[\phi] = \frac{c}{24\pi} \int dx dz \left[(\partial_x \phi)^2 + (\partial_z \phi)^2 + e^{2\phi} \right]$$

$$= \frac{c}{24\pi} \int dx dz \left[(\partial_x \phi)^2 + (\partial_z \phi + e^\phi)^2 \right] + (\text{surface term})$$

$$\Rightarrow \text{Minimum: } e^{2\phi} = \frac{1}{z^2}.$$

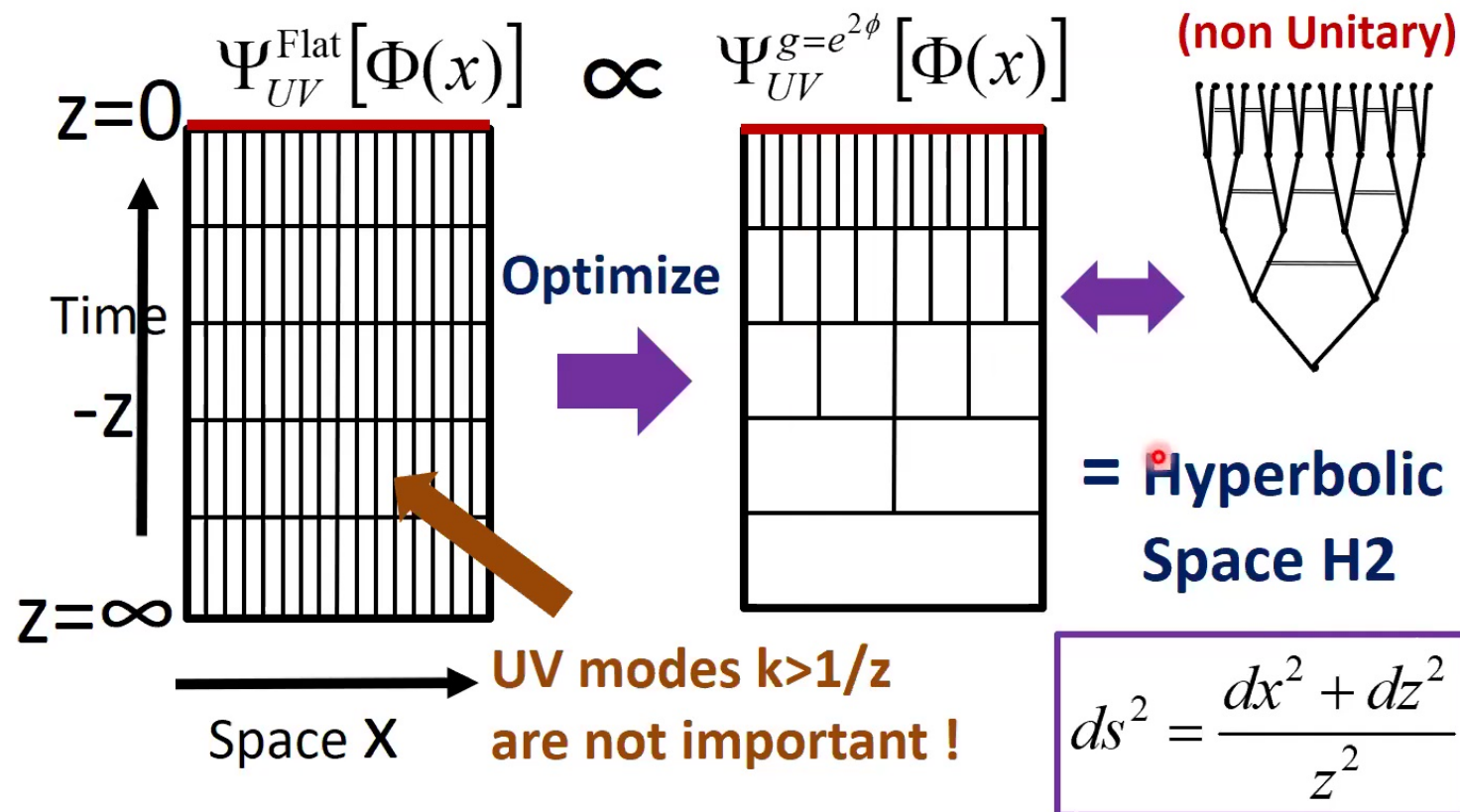


Hyperbolic plane (H₂)

= Time slice of AdS₃

$$ds^2 = (dx^2 + dz^2) / z^2.$$

Summary: Optimization of Path-Integral



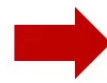
(2-2) Excited States and Back-reactions

Consider the vacuum state of 2d CFT on a circle dual to a global AdS3.

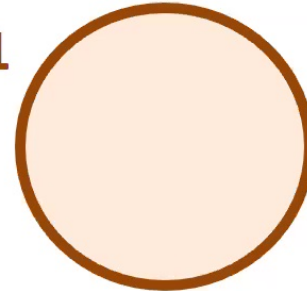
$$|w|=1$$

The optimization of the path-integral of 2d CFT on a disk with leads to

$$ds^2 = \frac{4dw d\bar{w}}{(1-|w|^2)^2}$$



Hyperbolic Disk H2
= time slice of global AdS3

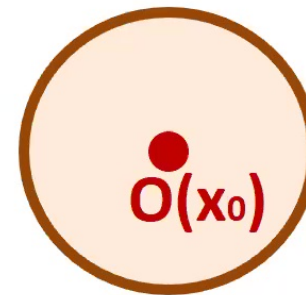


Now we insert an operator $O(x)$ in the center of the disk $x=x_0$.

$O(x)$: conformal dim. $h_L=h_R=h$

$$\Rightarrow O(x) \sim e^{-2h\cdot\varphi}$$

$$|w|=1$$



$$\frac{\Psi_{g=e^{2\varphi}}}{\Psi_{\text{Flat}}} \propto e^{\frac{c}{24\pi} S_L} \cdot e^{-2h\varphi(x_0)}.$$

Minimize →

$$\partial_w \partial_{\bar{w}} \varphi - \frac{1}{4} e^{2\varphi} + \frac{6\pi h}{c} \delta^2(w) =$$

← **Solve**

$$\text{Metric: } ds^2 = \frac{4d\zeta d\bar{\zeta}}{(1-|\zeta|^2)^2}, \quad \zeta = re^{i\theta}$$

⇒ Deficit angle geometry $\theta \sim \theta + 2\pi\alpha$.

This agrees with the expected gravity dual if $h/c \ll 1$.

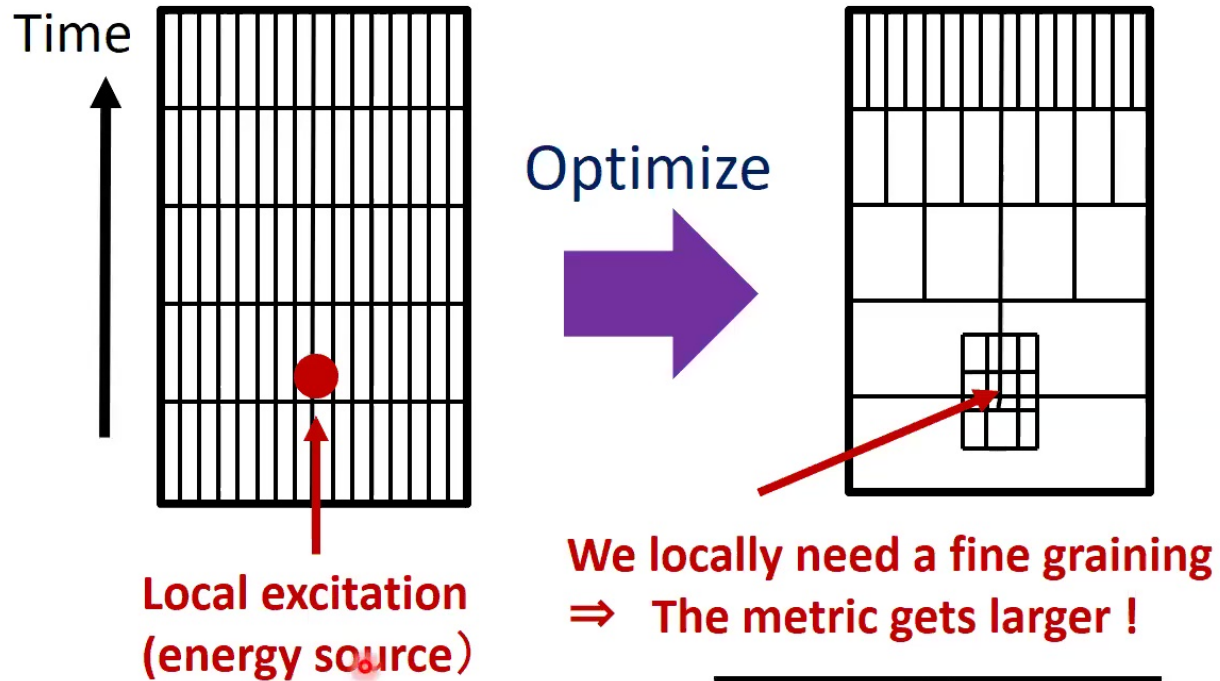
Note: the AdS/CFT predicts $a = \sqrt{1 - 24h/c}$.

Interestingly, if we consider **the quantum Liouville CFT**,
 then $h = \frac{\gamma\alpha}{4}(Q - \alpha\gamma/2)$, $c = 1 + 3Q^2$, ($Q \equiv 2/\gamma + \gamma$).

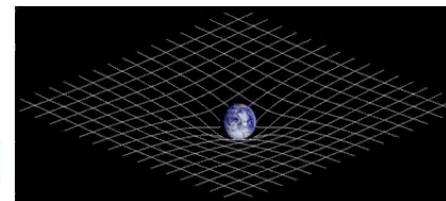
⇒ We get $a = \sqrt{1 - 24h/c}$.



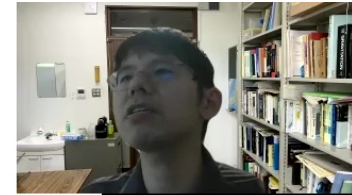
Heuristic Summary



This agrees with
general relativity !



Information Source = Back-reactions in GR



③ An Application: Entanglement of Purification

(3-1) Entanglement of Purification

When we divide a total system into A and B: $H_{tot} = H_A \otimes H_B$ the entanglement entropy (EE) is defined by

$$S_A = -\text{Tr}[\rho_A \log \rho_A]. \quad \rho_A = \text{Tr}_B[|\Psi\rangle\langle\Psi|]$$

The entanglement entropy is known to be a measure of quantum entanglement only for pure states.

We know that mixed states can always be purified:

$$\rho_C = \sum_i \lambda_i |i\rangle_C \langle i| \xrightarrow{H_C \rightarrow H_C \otimes H_D} |\Psi\rangle_{CD} = \sum_i \sqrt{\lambda_i} |i\rangle_C |i\rangle_D$$

$\rho_C = \text{Tr}_D[|\Psi\rangle\langle\Psi|]$



Definition of EoP [Terhal-Horodecki-Leung-Divincenzo 2002]

Consider all purifications $|\Psi\rangle_{A\tilde{A}B\tilde{B}}$ of ρ_{AB} in the extended Hilbert space: $H_A \otimes H_B \rightarrow H_A \otimes H_B \otimes H_{\tilde{A}} \otimes H_{\tilde{B}}$.

Then, **Entanglement of Purification (EoP)** is defined by

$$E_P(\rho_{AB}) = \underset{\text{All purifications } |\Psi\rangle \text{ of } \rho_{AB}}{\text{Min}} S_{A\tilde{A}}(|\Psi\rangle_{A\tilde{A}B\tilde{B}})$$

$$\rho_{AB} = \text{Tr}_{\tilde{A}\tilde{B}} [|\Psi\rangle\langle\Psi|]$$

Entanglement Entropy

Note: $E_p(\rho_{AB}) \geq 0$ and $E_p(\rho_{AB}) = 0 \Leftrightarrow \rho_{AB} = \rho_A \otimes \rho_B$.



(3-2) Holographic EoP

[Umemoto-TT 17, Nguyen-Dev
-Halbasch-Zaletel-Swingle 17]

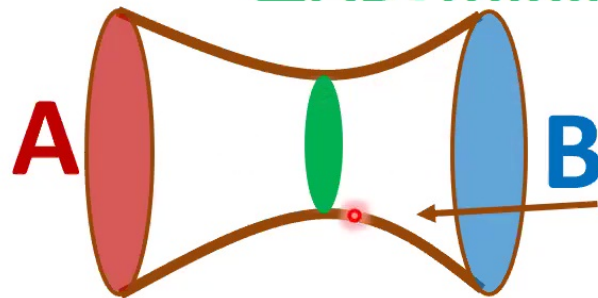


Conjecture

$$E_P(\rho_{AB}) = \frac{\text{Area}(\Sigma_{AB})}{4G_N}$$

Entanglement Wedge
Cross Section

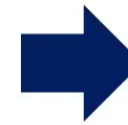
Σ_{AB} : Minimal Surface of M_{AB}



M_{AB} : Minimal surface
connecting ∂A and ∂B

Note: When ρ_{AB} is pure,

we simply have $E_P(\rho_{AB}) = S_A = S_B$.



Standard
Hol. EE

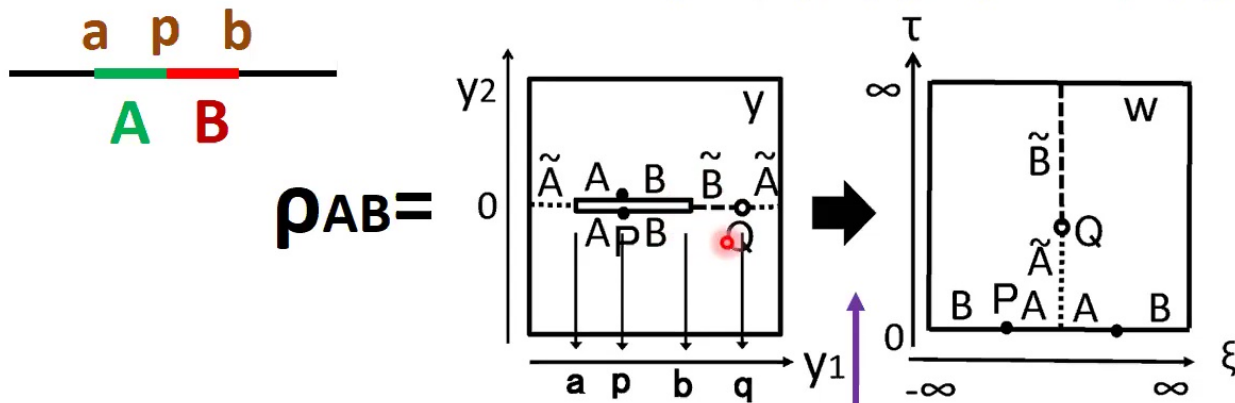
(3-3) EoP from Path-Integral Optimization [Caputa-Miyaji -Umemoto-TT 18]



Key Idea: Min over all purifications is realized by PI-opt !

Consider the setup when $AB =$ a single interval:

$$A=[a,p] \quad B=[p,b] : \Rightarrow \tilde{A}=[-\infty,a] \cup [q,\infty] \quad \tilde{B}=[b,q]$$



Conformal Map

$$w = \sqrt{\frac{y-a}{b-y}}$$

The final optimized metric looks like

$$ds^2 = \frac{\epsilon^2}{\tau^2} \cdot dwd\bar{w} = \frac{\epsilon^2}{\tau^2} \cdot \frac{(b-a)^2}{4|b-y|^3|y-a|} \cdot dyd\bar{y} \equiv e^{2\tilde{\phi}} \cdot dyd\bar{y}.$$

In this setup, we obtain $e^{2\tilde{\phi}_P} = 1$ and $e^{2\tilde{\phi}_Q} = \frac{\epsilon^2(b-a)^2}{4(q-a)^2(q-b)^2}$.

The entanglement entropy $S_{A\tilde{A}} = S_{B\tilde{B}}$ is found to be

$$\begin{aligned} S_{A\tilde{A}} &= \frac{c}{3} \log \left(\frac{q-p}{\epsilon} \right) + \frac{c}{6} \tilde{\phi}_P + \frac{c}{6} \tilde{\phi}_Q \\ &= \frac{c}{6} \log \left[\frac{(b-a)(q-p)^2}{2\epsilon(q-a)(q-b)} \right], \end{aligned}$$

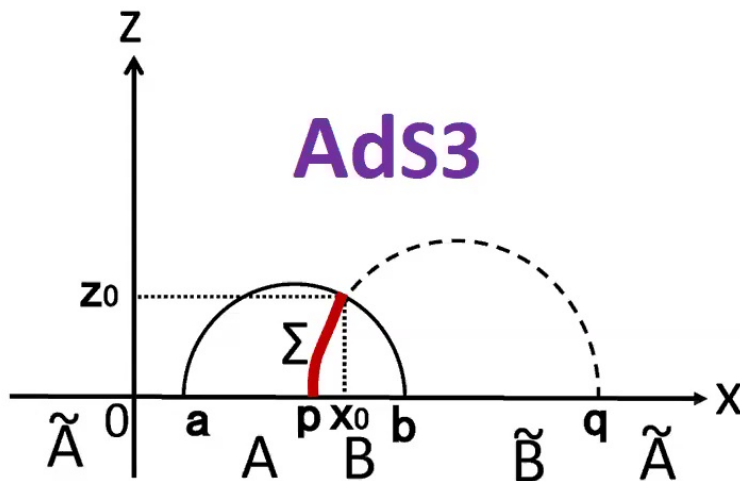
Minimize
w.r.t q



$$\begin{aligned} S_{A\tilde{A}}^{min} &= \frac{c}{6} \log \left[\frac{2(p-a)(b-p)}{\epsilon(b-a)} \right], \\ \text{at } q &= \frac{2ab - (a+b)p}{a+b-2p}. \end{aligned}$$



Calculations of Hol EoP in AdS3/CFT2



$$ds^2 = \frac{dx^2 + dz^2}{z^2}$$

$$L(\Sigma) = \frac{q-p}{2} \int_{\epsilon}^{z_0} \frac{dz}{z \sqrt{\frac{(q-p)^2}{4} - z^2}} = \log \left[\frac{2(p-a)(p-b)}{\epsilon(b-a)} \right].$$

➡ The Hol EoP $L(\Sigma)/4G$ agrees with the CFT result !

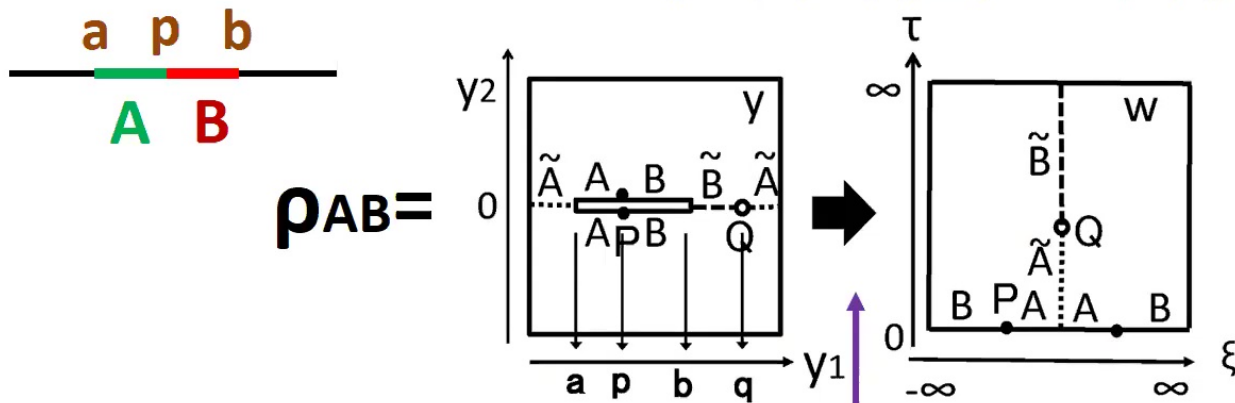
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Minimize
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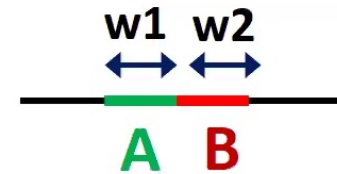


Holographic EoP

Path-integral Optimization



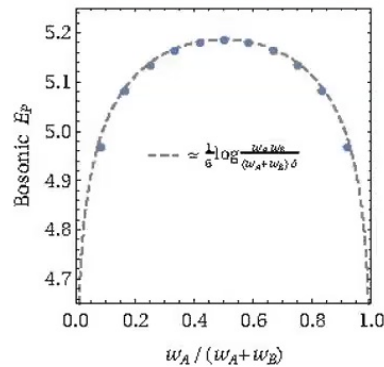
$$E_p(\rho_{AB}) = \frac{c}{6} \log \frac{w_1 w_2}{(w_1 + w_2) \delta}$$



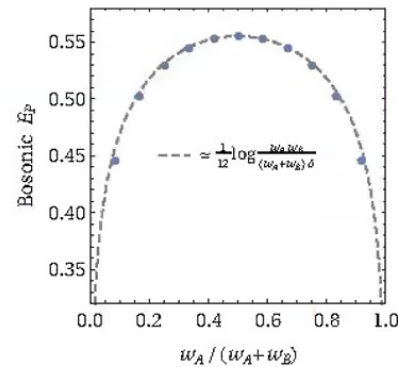
Confirmed by direct Gaussian numerical calculations in free scalar/fermion CFT !

[Camargo-Hackl-Heller-Jahn-TT, 2020]

Free Scalar CFT



Free Fermion CFT



This provides a further support of path-integral optimization as well as Hol. EoP conjecture !

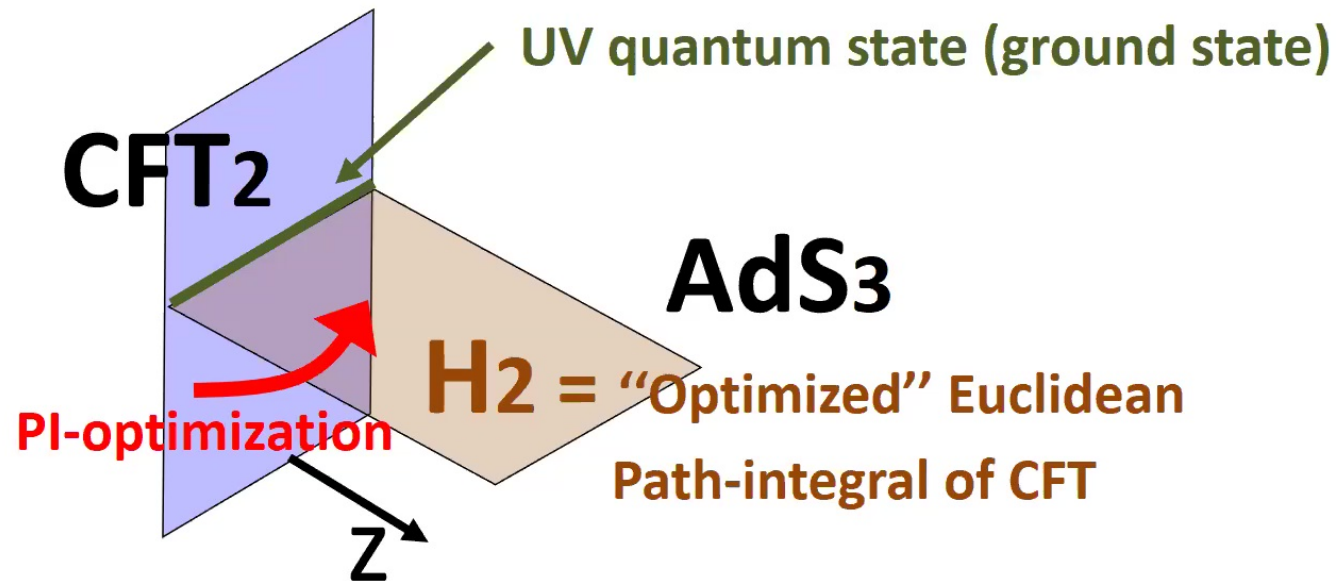
④ Connection between PI-Optimization and AdS/CFT

[Boruch-Caputa-TT, 2020]

We observed that a time slice of AdS emerges from our path-integral optimization. This raises the question:

Q. What does optimization mean in AdS/CFT ?

Heuristic Expectation




$$\text{PI-optimization} = \text{Min}_{\phi} \left[\Psi_{UV}^{g_{ab} = e^{2\phi} \delta_{ab}} [\Phi(x)] \right]$$

One might think the optimization corresponds to *a saddle point approximation* of some quantum gravity.

➔ However, a saddle point approximation does not give a minimum but a maximum !

This is important when we would like to take into account quantum gravity (=1/c) corrections in AdS/CFT.



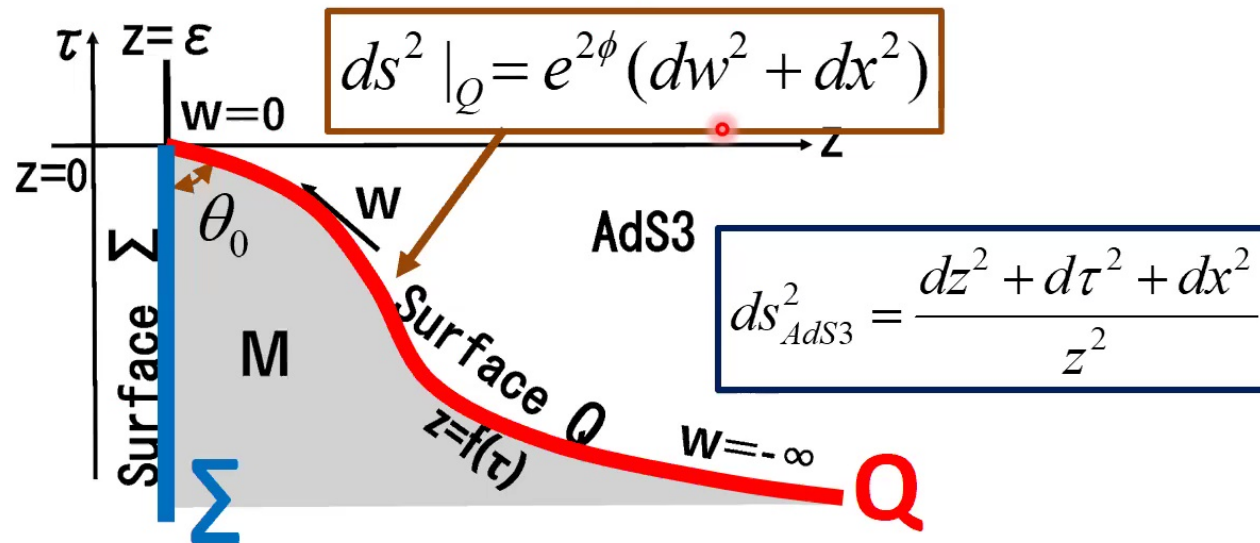


Consider **Hartle-Hawking wave function** on a time slice \mathcal{Q} in AdS3 gravity, with an initial condition on the AdS bdy Σ .

$$\Psi_{HH}^{(T)}[e^{2\phi} \delta_{ab}] = \int Dg_{\mu\nu} e^{-I_G[g] - T \int_{\mathcal{Q}} e^{2\phi}} \delta(g|_{\mathcal{Q}} - e^{2\phi} \delta_{ab})$$

“Tension” term

Metric on \mathcal{Q} \sim bdy cosmological const.



By evaluating the HH wave function using classical gravity action, we obtain (we assume $\partial_x \phi = 0$ and write $\dot{\phi} \equiv \partial_w \phi$)



$$\Psi_{HH}^{(T)}[e^{2\phi} \delta_{ab}] \approx e^{-I^{(T)}[\phi] + I^{(T)}[\phi_0]}, \quad (e^{\phi_0} = 1/\epsilon)$$

↑ Minus sign

$$I^{(T)}[\phi] = \frac{c}{12\pi} \int dx dw \left[e^{2\phi} \sqrt{1 - \dot{\phi}^2 e^{-2\phi}} + \dot{\phi} \text{ArcSin}[\dot{\phi} e^{-\phi}] \right] - \frac{c}{12\pi\epsilon} \int dx \theta_0$$

When the UV cut off is dominant $\dot{\phi}^2 \ll e^{2\phi}$, we can trust the continuum limit of field theory. In this case, we have

$$I^{(T)}[\phi] \approx \frac{c}{24\pi} \int dx dw [\dot{\phi}^2 + 2e^{2\phi}] - \frac{c}{12\pi} \int dx \text{ArcSin} \sqrt{1 - T^2}$$



Indeed we reproduce the Liouville action !
The HH wave function gives finite cut off corrections.



Therefore, we can interpret the PI-optimization as an optimization of HH wave function $\Psi_{HH}^{(T)}[e^{2\phi} \delta_{ab}]$:

$$\langle O_1 O_2 \cdots O_n \rangle = \int D\phi \left| \Psi_{HH}^{(T)}[\phi] \right|^2 O_1 O_2 \cdots O_n \Rightarrow \text{Min}[I^{(T)}[\phi]]$$

$$\Psi_{HH}^{(T)} \propto e^{-I^{(T)}[\phi]}$$

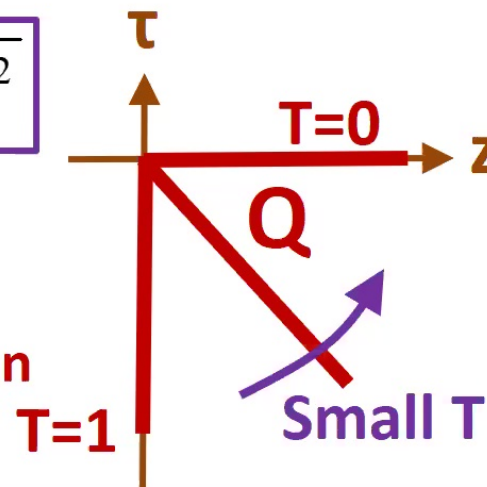
Indeed, its saddle point approximation lead to the maximization of $\Psi_{HH}^{(T)}[e^{2\phi} \delta_{ab}]$, leading to the EOM:

$$\ddot{\phi} + \dot{\phi}^2 - 2e^{2\phi} = 2Te^{2\phi} \sqrt{1 - e^{-2\phi} \dot{\phi}^2}$$

Solution

$$e^{2\phi} = \frac{1}{(1 - T^2)w^2}$$

Agree with PI-optimization





Therefore, we can interpret the PI-optimization as an optimization of HH wave function $\Psi_{HH}^{(T)}[e^{2\phi} \delta_{ab}]$:

$$\langle O_1 O_2 \cdots O_n \rangle = \int D\phi \left| \Psi_{HH}^{(T)}[\phi] \right|^2 O_1 O_2 \cdots O_n \Rightarrow \text{Min}[I^{(T)}[\phi]]$$

$$\Psi_{HH}^{(T)} \propto e^{-I^{(T)}[\phi]}$$

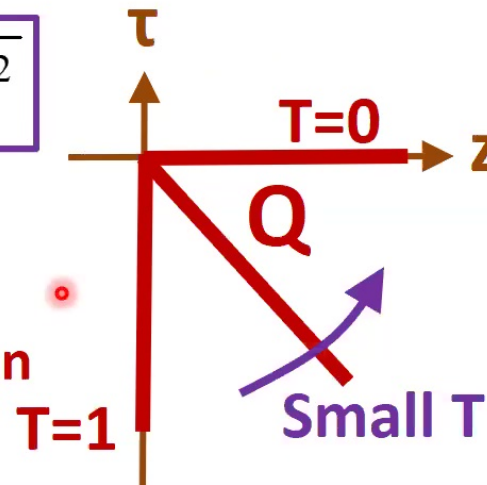
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The Liouville potential and the tension of the brane Q is related via $\mu = 1 - T^2$.

This allows us to rewrite the AdS3 metric as follows:

$$ds^2 = \frac{d\mu^2}{4\mu^2(1-\mu)} + e^{2\phi(w)}(dw^2 + dx^2)$$

Emergent Time direction
= the energy scale (μ) of
path-integral optimization

**Optimized CFT metric
at the scale μ**

$$e^{2\phi} = \frac{1}{\mu w^2}$$

Generalizations



We can also generalize the HH-wave function derivation of PI-optimization to

- Higher Dim. CFTs
- BTZ black holes
- JT gravity



CFT analysis looks difficult because we need more than the Weyl transformation to optimize PI. But for Holographic CFTs, HH-wave function gives an answer !

For details refer to [arXiv:2011.08188](https://arxiv.org/abs/2011.08188) [Boruch-Caputa-TT]



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⑤ Conclusions

- We expect that the AdS/CFT provides a geometrization of real space renormalization. \Rightarrow Conjectured connection to TN.
But, the direct TN approaches suffer from lattice artifacts etc.
- We argued that the path-integral optimization gives a definite framework which covers also a continuum limit.
- We found : An PI-optimized geometry = a time slice of AdS
- We gave the explicit connection between the AdS/CFT and PI-optimization by considering the Hartle-Hawking wave function.
 \Rightarrow PI-optimization = saddle point calculation of HH wave function.
 \Rightarrow We can consider quantum fluctuations !



Future problems

- Meaning of Emergence of Time ? Metric in the time direction ?
- Time-dependent b.g. ?
- Implications of large N strongly coupled CFTs and Bulk locality ?
- dS/CFT version ? (Formally, $\mu < 0$)

⋮
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