

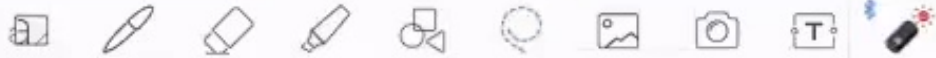
Title: Special Topics in Astrophysics - Numerical Hydrodynamics - Lecture 15

Speakers: Daniel Siegel

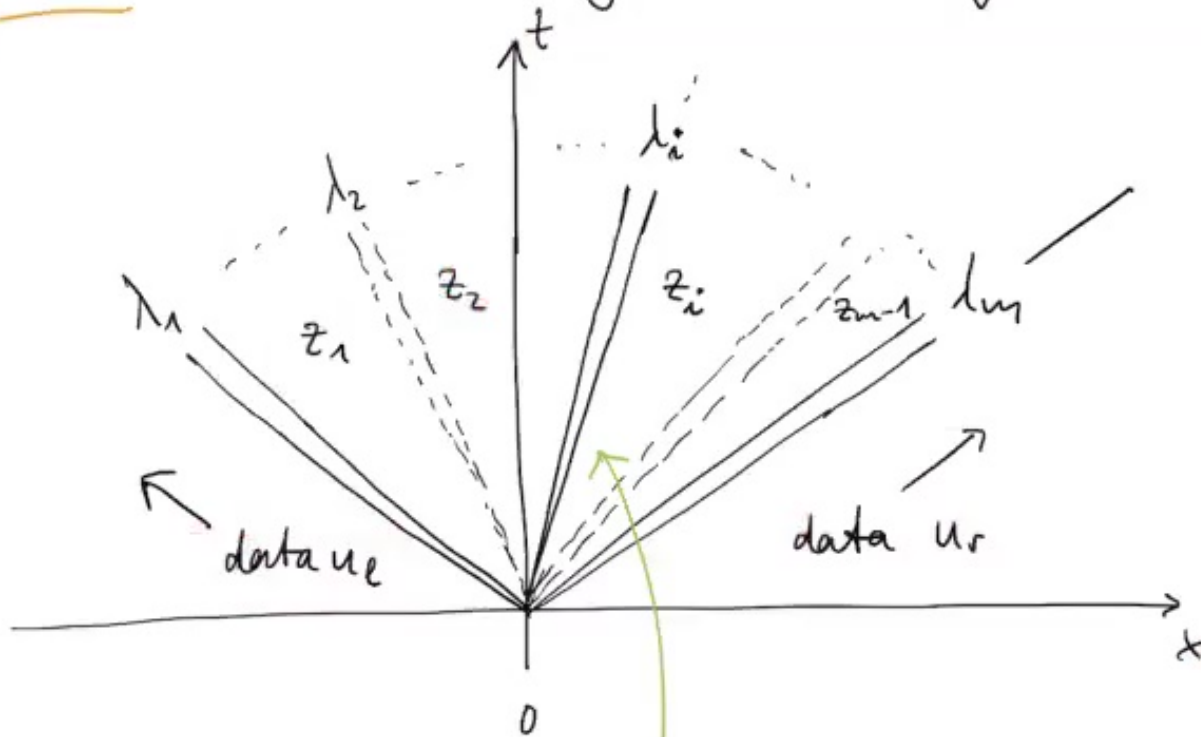
Collection: Special Topics in Astrophysics - Numerical Hydrodynamics

Date: November 05, 2020 - 3:30 PM

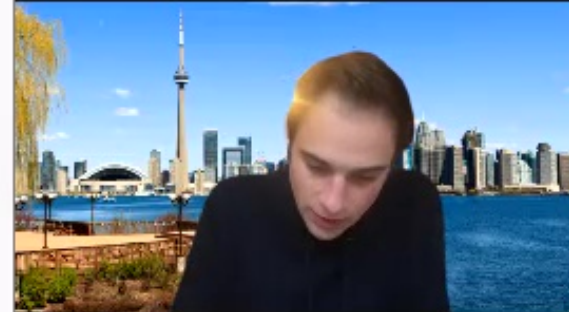
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The global solution u is also unique as it satisfies the entropy condition by construction. □



solution constant on lines through the origin ($\frac{x}{t} = \text{const.}$)





Consider Euler eqns

$$u_t + f(u)_x = 0$$

$$\begin{pmatrix} \rho \\ \rho u \\ E \end{pmatrix}_t + \begin{pmatrix} \rho u \\ \rho u^2 + p \\ (E + p)u \end{pmatrix}_x = 0$$

$$\lambda_1 = u - c \quad \lambda_2 = u \quad \lambda_3 = u + c$$

$$\Gamma_1 = \begin{pmatrix} 1 \\ u - c \\ H - uc \end{pmatrix}, \quad \Gamma_2 = \begin{pmatrix} 1 \\ u \\ \frac{1}{2}u^2 \end{pmatrix}, \quad \Gamma_3 = \begin{pmatrix} 1 \\ u + c \\ H + uc \end{pmatrix}$$

- $(\lambda_1, \Gamma_1), (\lambda_3, \Gamma_3)$: genuinely non-linear
- (λ_2, Γ_2) : linearly degenerate



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$$\begin{pmatrix} \rho \\ \rho u \\ E \end{pmatrix}_t + \begin{pmatrix} \rho u \\ \rho u^2 + p \\ (E + p)u \end{pmatrix}_x = 0$$

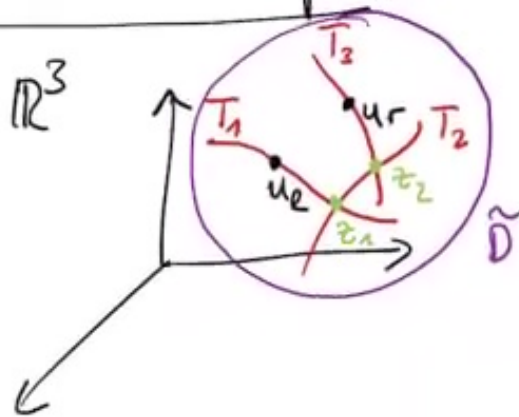
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- $(\lambda_1, \Gamma_1), (\lambda_3, \Gamma_3)$: genuinely non-linear
- (λ_2, Γ_2) : linearly degenerate

no General structure of the RP for the Euler equations

Situation in state space:



$m=3$ equations
 $u = (u_1, u_2, u_3)$
 $= (p, \rho u, E)$

Situation in real space:

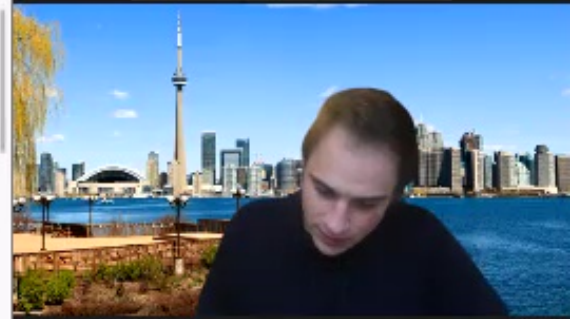
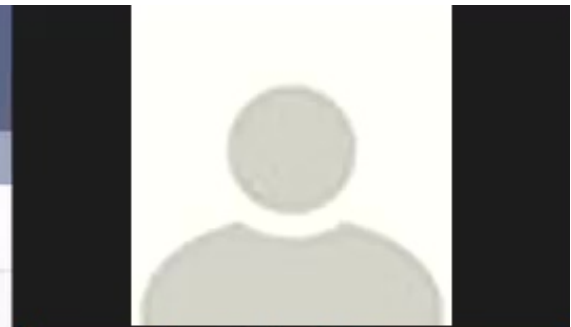
$\lambda_1 = u - c$



contact discontinuity
 $\lambda_2 = u$ ("entropy wave")



$\lambda_3 = u + c$

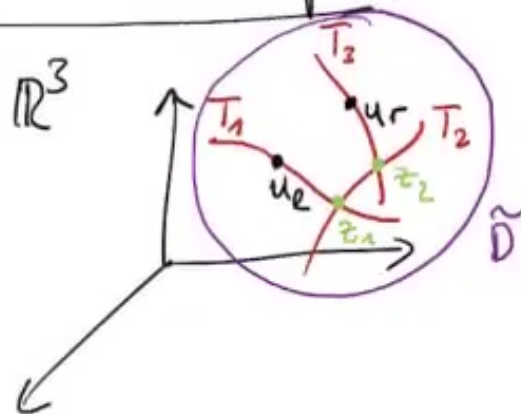


$$2-R1: \quad u, p$$

$$3-R1: \quad s, u - \frac{2c}{\delta-1}$$

no General structure of the RP for the Euler equations

Situation in state space:



$m=3$ equations

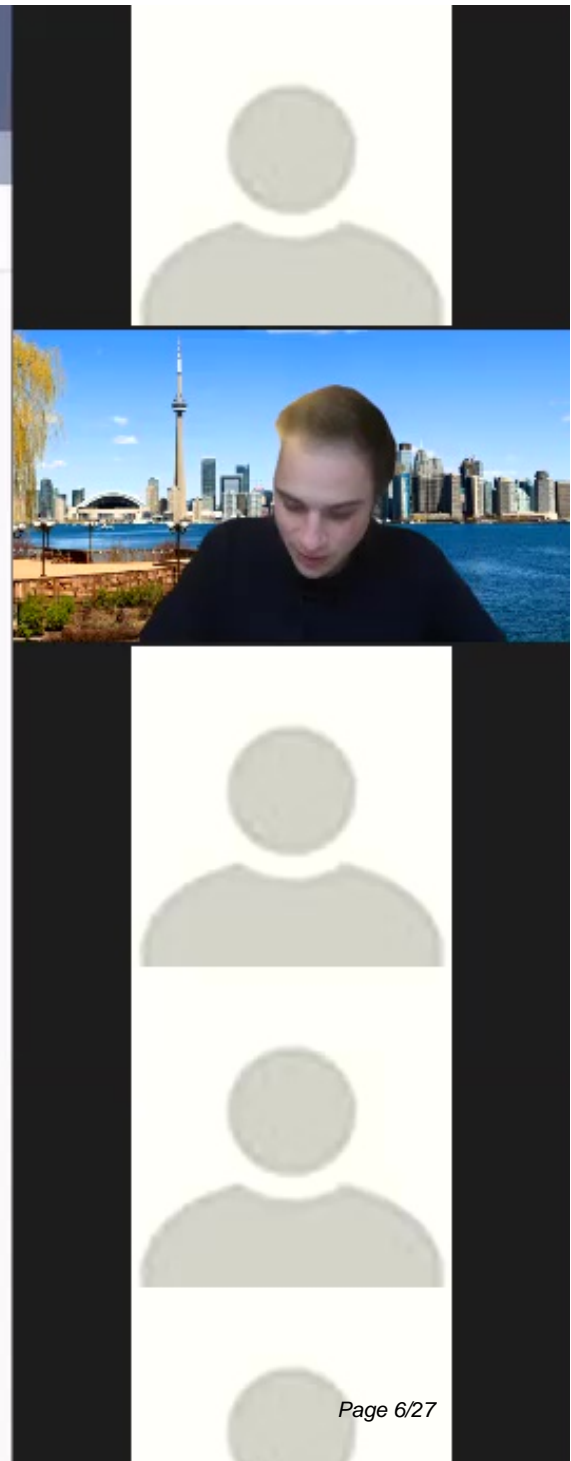
$$u = (u_1, u_2, u_3)$$

$$= (s, pu, E)$$

$$w = (s, u, p)$$

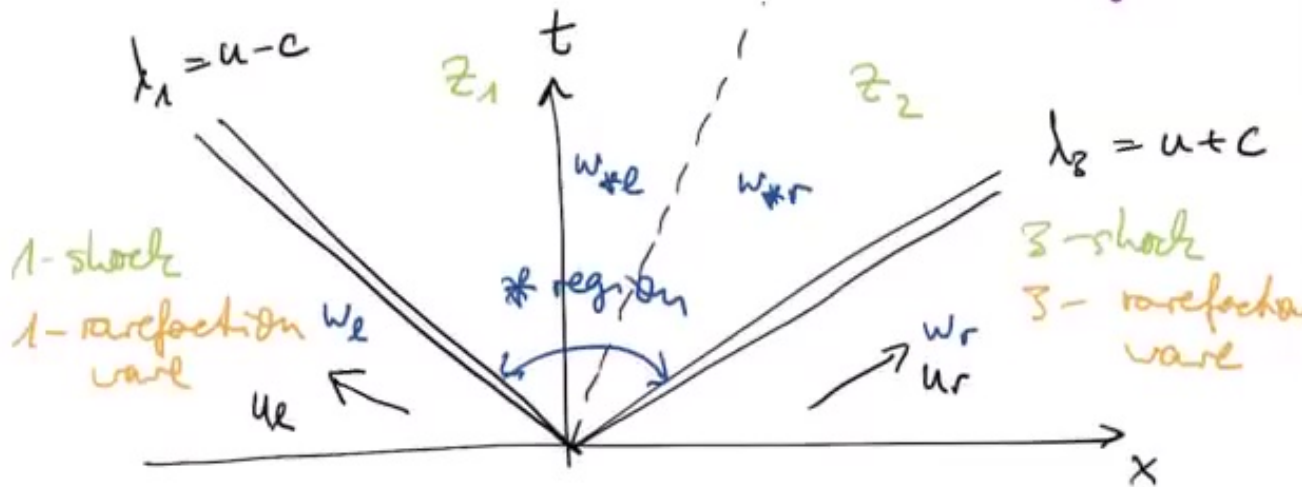
Situation in real space:

contact discontinuity



Situation in real space:

discontinuity
 $\lambda_2 = u$ ("entropy wave")



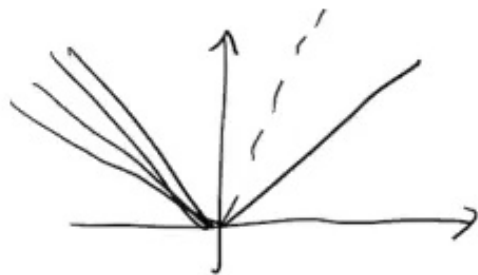
Remarks: 1) Since both u & p are constant across contact discontinuities

it is easier to work with the primitive variables $w = (s, u, p)$

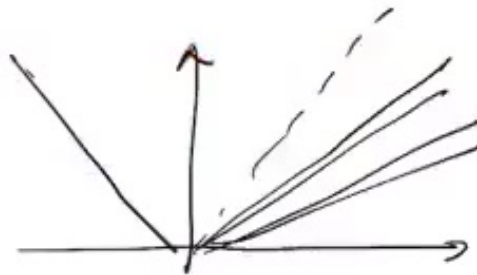


2) There are 4 possible patterns

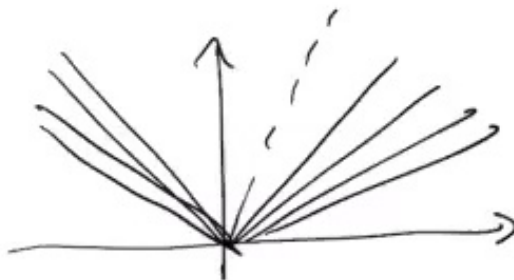
(i)



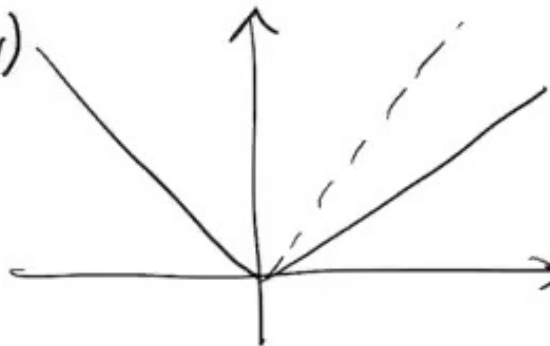
(ii)



(iii)



(iv)



Example:

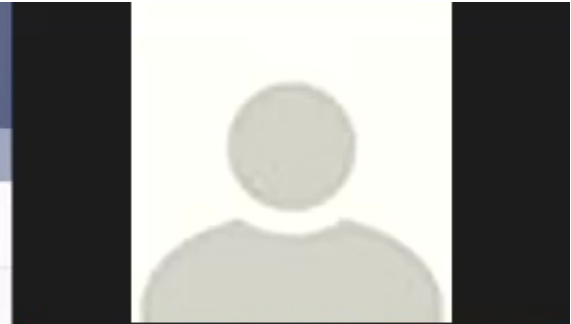
1)



ISM ρ_0

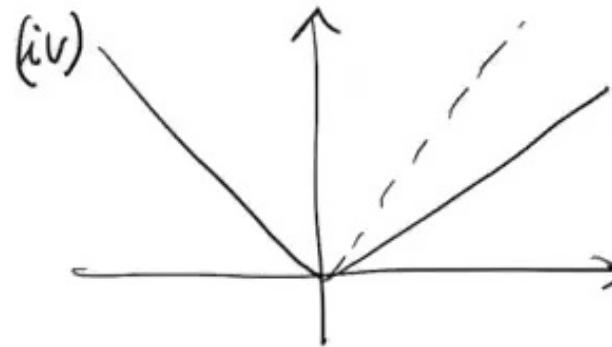
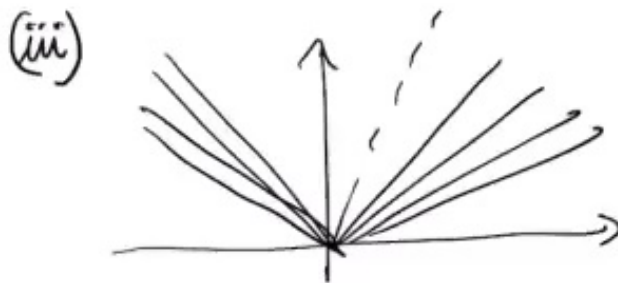
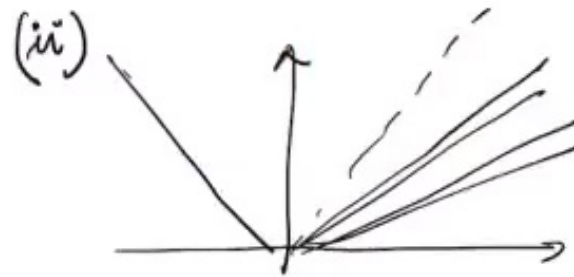
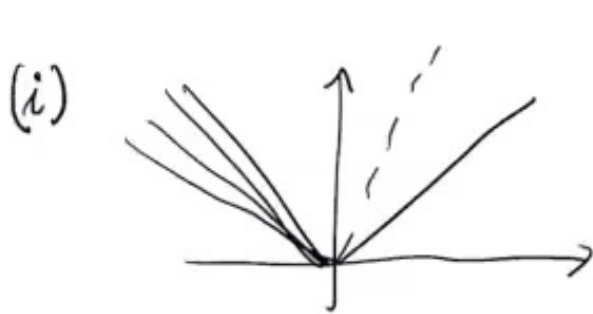


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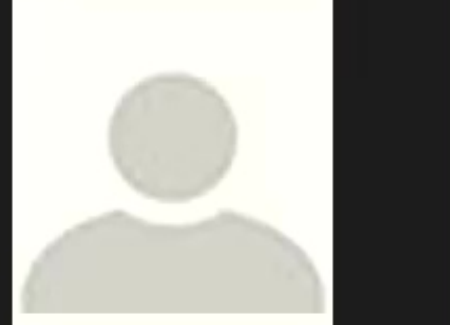
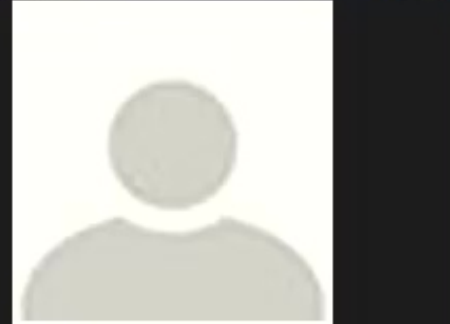
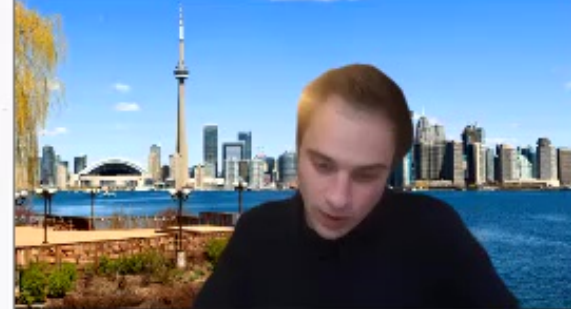
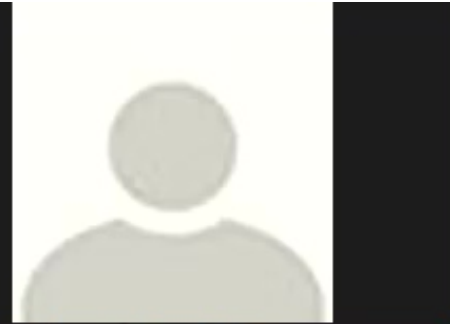
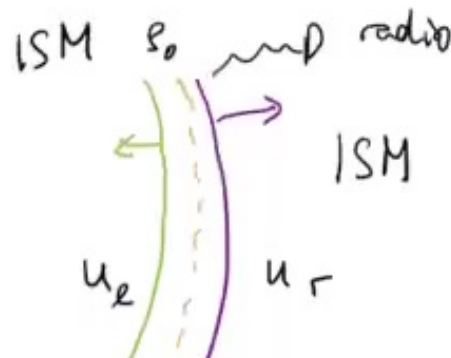
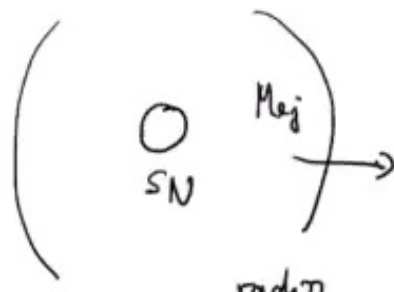


2) There are 4 possible patterns



Example:

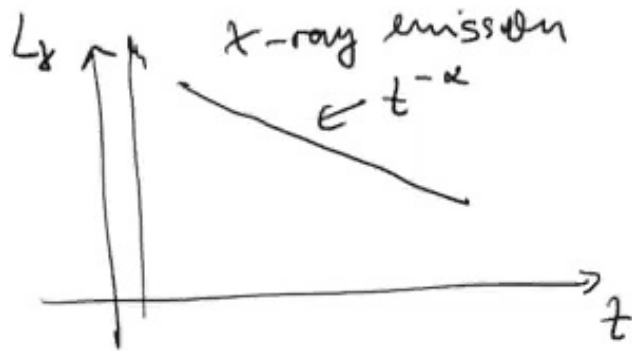
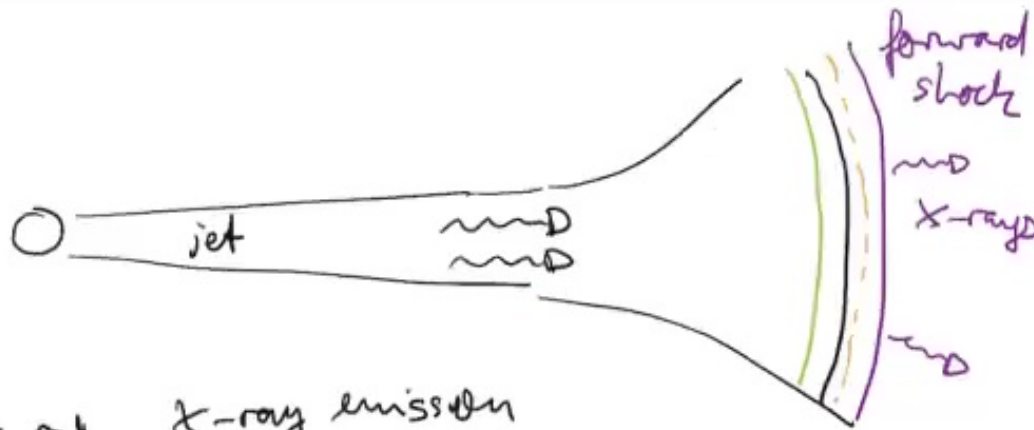
1)



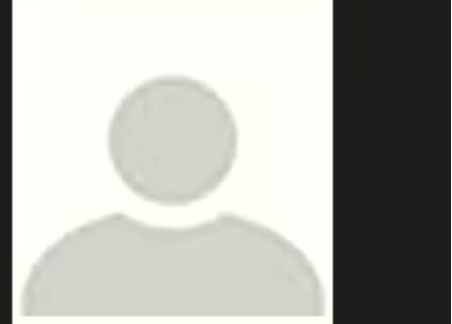
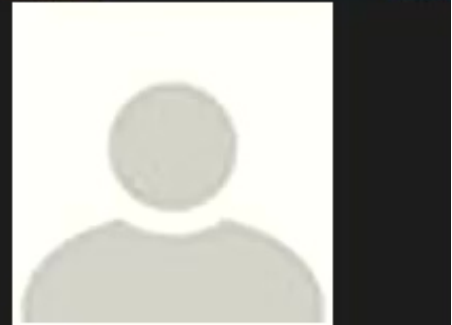
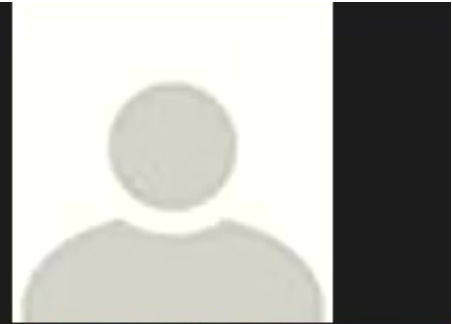


contact

2)



5.2 Solution strate





5.2 Solution strategy

- ① Assume we know u_* , P_* ,
 evaluate the case (i)–(iv) that applies
 to this configuration of (u_*, P_*) and
 write down the combined solution
 $w = w(x, t)$.

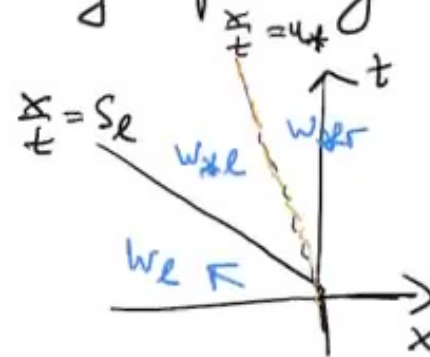
will follow with (1)

step (1): Treat regions left and right of the contact discontinuity separately

(1) $S \equiv \frac{x}{t} \leq u_*$: left side

$p_* > p_l$: left shock
 ($\lambda(u_l) > \lambda(u_*)$) (ii) or (iv)

$$w(x,t)_{\text{left}} = \begin{cases} w_{*l}^{sh} \\ w_l \end{cases}$$



$$\Downarrow S_l \leq \frac{x}{t} \leq u_*$$

$$\Downarrow \frac{x}{t} \leq S_l$$

shock



$P_* > P_L$: left shock
 ($\lambda(u_e) > \lambda(u_*)$) (ii) or (iv)

$$w(x,t)_{\text{left}} = \begin{cases} w_{*L}^{sh} & \text{if } \underline{s_L} \leq \frac{x}{t} \leq u_* \\ w_L & \text{if } \frac{x}{t} \leq \underline{s_L} \end{cases}$$

where $w_{*L}^{sh} = (\underline{s_{*L}^{sh}}, u_*, P_*)$

RH jump conditions

$$s_{*L} = s_L \frac{\frac{P_*}{P_L} + \frac{\delta-1}{\delta+1}}{\frac{\delta-1}{\delta+1} \frac{P_*}{P_L} + 1}$$

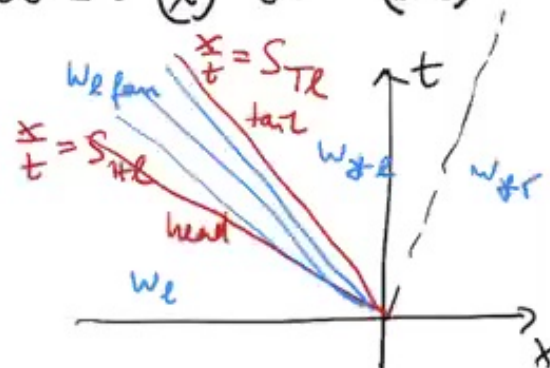
jump conditions

$$S_{*l} = S_l \frac{\frac{p_*}{\rho_l} + \frac{u-1}{u+1}}{\frac{u-1}{u+1} \frac{p_*}{\rho_l} + 1}$$

$$S_l = u_l - c_l \sqrt{\frac{u+1}{2u} \frac{p_*}{\rho_l} + \frac{u-1}{2u}}$$

$\frac{p_*}{\rho_l} \leq p_l : \Rightarrow$ left rarefaction wave: (i) or (iii)

$(\lambda(u_l) < \lambda(u_*))$



Lecture_15

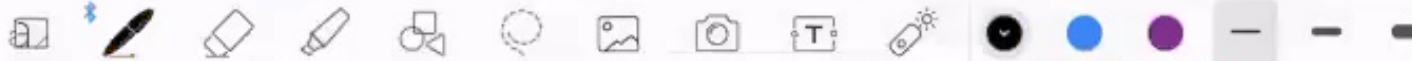
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Chap_5_RiemannPr...

Chap_6_Numerical...



$$w(x,t)_{\text{left}} = \begin{cases} w_L & , \quad \frac{x}{t} \leq S_{HL} \\ w_{\text{fan}} & , \quad S_{HL} \leq \frac{x}{t} \leq S_{TL} \\ w_{*L} & , \quad S_{TL} \leq \frac{x}{t} \leq u_{*} \end{cases}$$

where

$$w_{*L} = (\underbrace{S_{*L}}_{\text{known}}, \underbrace{u_{*}, p_{*}}_{\text{known}})$$

w_{fan} : state "inside" the rarefaction wave

Lecture_15

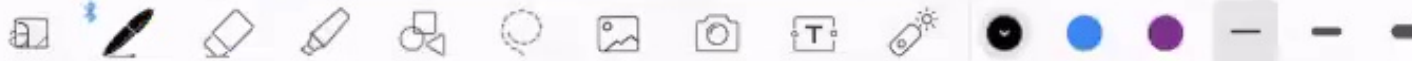
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Chap_6_Numerical...



where

$$w_{*l} = (\underline{s_{*l}}, \underbrace{u_{*l}, p_{*l}}_{\text{known}})$$

w_{lfan}: state "inside" the rarefaction wave

Find: $s_{*l}, s_{*r}, s_{*l}, w_{lfan}$

$$\left. \begin{array}{l} \text{I} \quad s \\ \text{II} \quad u + \frac{2c}{\gamma-1} \end{array} \right\} \begin{array}{l} 1\text{-Riemann} \\ \text{invariants} \end{array}$$

I
we can use isentropic law $p = k s^\gamma$

across wave

$p_0 = k s_0^\gamma \Rightarrow k = \frac{p_0}{s_0^\gamma}$

adiabatic
index
of EOS

Lecture_15

Lecture_15

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$$\left. \begin{array}{l} I \quad S \\ II \quad u + \frac{2c}{\gamma - 1} \end{array} \right\} \text{1-Riemann invariants}$$

I can use isentropic law $p = k s^\gamma$

across wave

$$p_l = k s_l^\gamma \Rightarrow k = \frac{p_l}{s_l^\gamma}$$

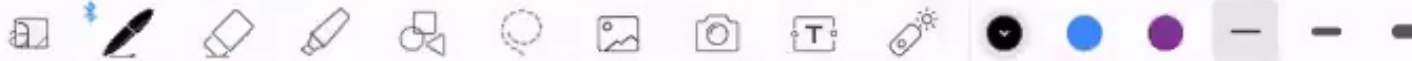
$$s_{*l} = s_l \left(\frac{p_{*l}}{p_l} \right)^{\frac{1}{\gamma}}$$

$$c = \sqrt{\gamma \frac{p}{\rho}}$$

↓

$$c_{*l} = c_l \left(\frac{p_{*l}}{p_l} \right)^{\frac{\gamma-1}{2\gamma}}$$

adiabatic
index
of EOS



$$c_{\#l} = c_l \left(\frac{p_{\#l}}{p_l} \right)^{\frac{\gamma-1}{2\gamma}}$$

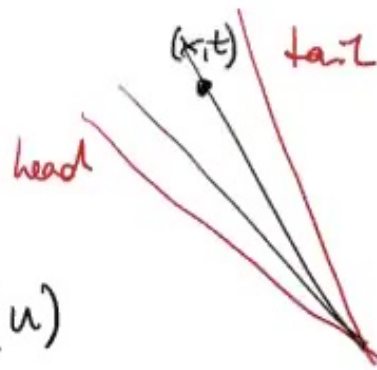
Head & tail speeds:

$$S_{Hl} = \lambda_1(u_l) = u_l - c_l$$

$$S_{Tl} = \lambda_1(u_{\#l}) = u_{\#l} - c_{\#l}$$

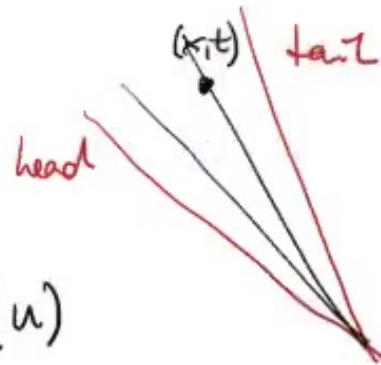
Wave fan:

$$\frac{dx}{dt} = \frac{x}{t} = \lambda_1(u) = u - c \quad (1)$$



$$S_{Tx} = \lambda_1(u) = u - c$$

Wk form:



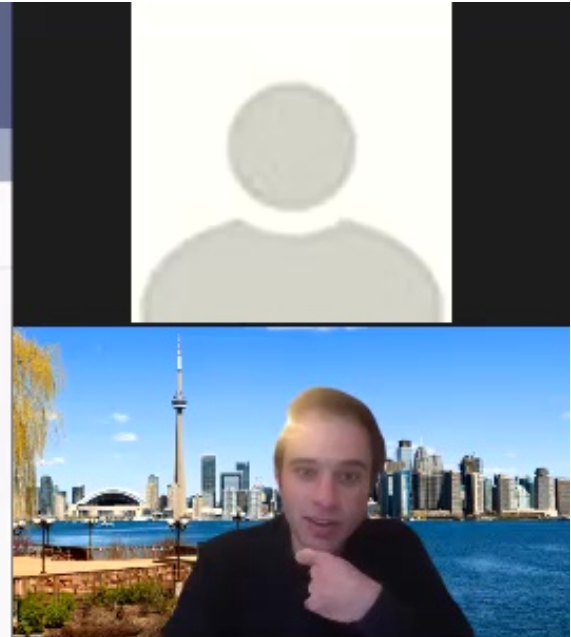
$$\frac{dx}{dt} = \frac{x}{t} = \lambda_1(u) = u - c \quad (1)$$

relate to u_e, c_e via Π : (u, c)

$$u_e + \frac{2c_e}{\gamma-1} = u + \frac{2c}{\gamma-1} \quad (2)$$

combine to:

$$S = S_e \left[\frac{2}{\gamma+1} + \frac{\gamma-1}{(\gamma+1)c_e} \left(u_e - \frac{x}{t} \right) \right]^{\frac{\gamma}{\gamma-1}}$$

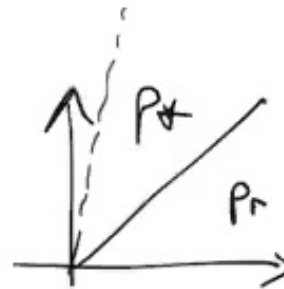


$$w_{\text{left}} = \begin{cases} S = \frac{x}{t} \left[\delta + 1 + (\delta + 1) c e^{-\frac{x}{t}} \right] \\ u = \frac{2}{\delta + 1} \left[c e + \frac{\delta - 1}{2} u_L + \frac{x}{t} \right] \\ p = p_L \left[\frac{2}{\delta + 1} + \frac{\delta - 1}{(\delta + 1) c} \left(u e^{-\frac{x}{t}} \right)^{\delta + 1} \right] \end{cases}$$

(2) $S \equiv \frac{x}{t} \geq u_L$: right side

analogous to above

$\underline{p_L} > p_R$: right shock (i) or (iv)



$$w(x,t)_{\text{right}} = \begin{cases} w_{*r}^{sh} & u_L \leq \frac{x}{t} \leq S_r \\ w_r & \frac{x}{t} \geq S_r \end{cases}$$

$$w(x,t)_{right} = \begin{cases} w_{rfan}, & s_{Tr} \leq \frac{x}{t} \leq s_{Hr} \\ w_r, & \frac{x}{t} \geq s_{Hr} \end{cases}$$

② Find (u_*, p_*) given initial guess:

Proposition: The solution p_* is given by the root of the function (assuming an ideal gas EOS):

$$f(p, w_e, w_r) \equiv f_e(p, w_e) + f_r(p, w_r) + u_r - u_e$$

$$f_r(p, w_r) = \begin{cases} p - p_r \sqrt{\frac{A_r}{p + B_r}}, & p > p_r \\ \frac{2c_r}{\gamma - 1} \left[\left(\frac{p}{p_r}\right)^{\frac{\gamma-1}{2\gamma}} - 1 \right], & p \leq p_r \end{cases}$$

(rarefaction wave)

(shock)

(rarefaction wave)

with given constants

$$A_e \equiv \frac{2}{(\gamma+1)c_e}, \quad B_e \equiv \frac{\gamma-1}{\gamma+1} p_e$$

$$A_r \equiv \frac{2}{(\gamma+1)c_r}, \quad B_r \equiv \frac{\gamma-1}{\gamma+1} p_r$$

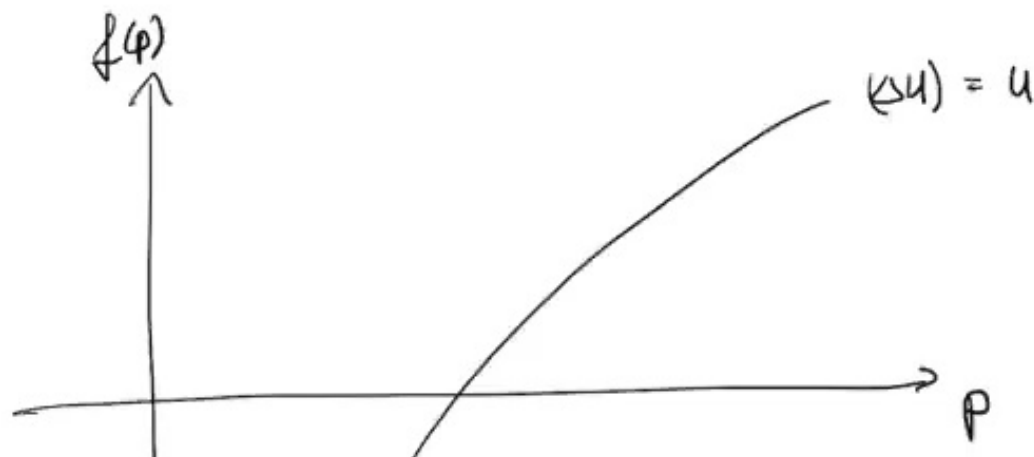


Remarks: 1) Behavior of $f(p)$

Can show: • $f'(p) > 0 \rightarrow$ monotonically increasing

• $f''(p) < 0 \rightarrow$ concave

• $\begin{cases} f'(p) \\ f''(p) \end{cases} \xrightarrow{p \rightarrow \infty} 0$



Lecture_15

Lecture_15

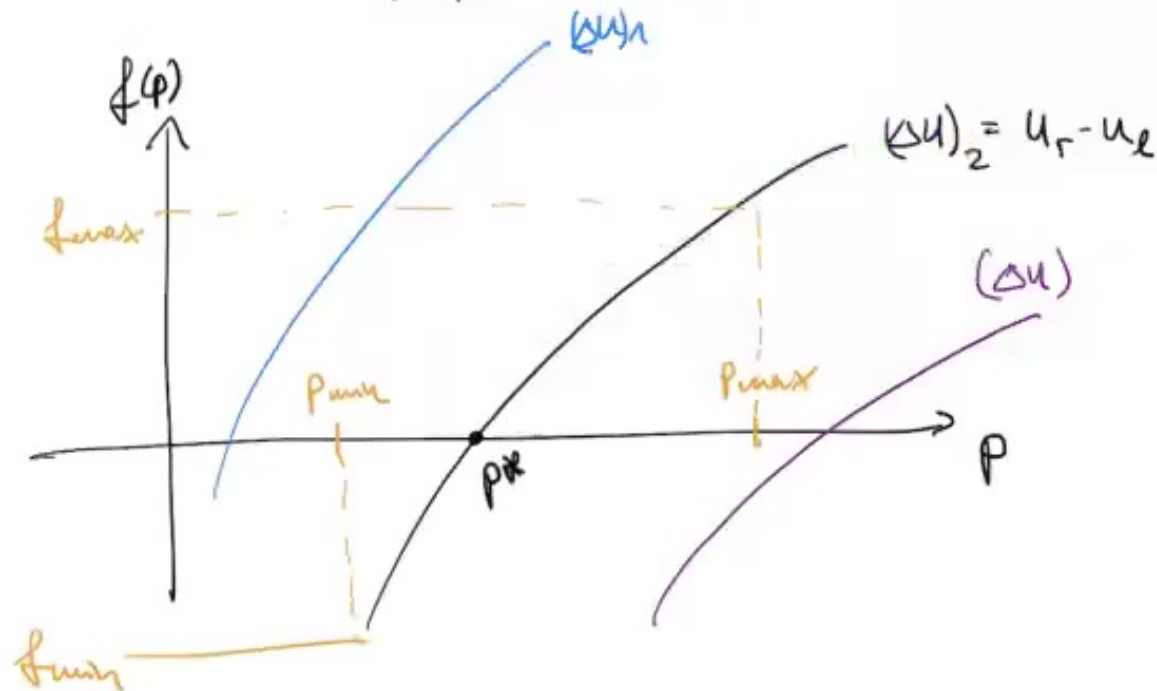
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$$\cdot \begin{cases} f'(p) \\ f''(p) \end{cases} \xrightarrow{p \rightarrow \infty} 0$$



$$p_{\min} = \min(p_l, p_r), \quad f_{\min} = f(p_{\min})$$

$$p_{\max} = \max(p_l, p_r), \quad f_{\max} = f(p_{\max})$$

3 regimes: $f_{min}, f_{max} > 0 \Rightarrow p_* \in I_1$

$f_{min} \leq 0, f_{max} \geq 0 \Rightarrow p_* \in I_2$

$f_{min}, f_{max} < 0 \Rightarrow p_* \in I_3$

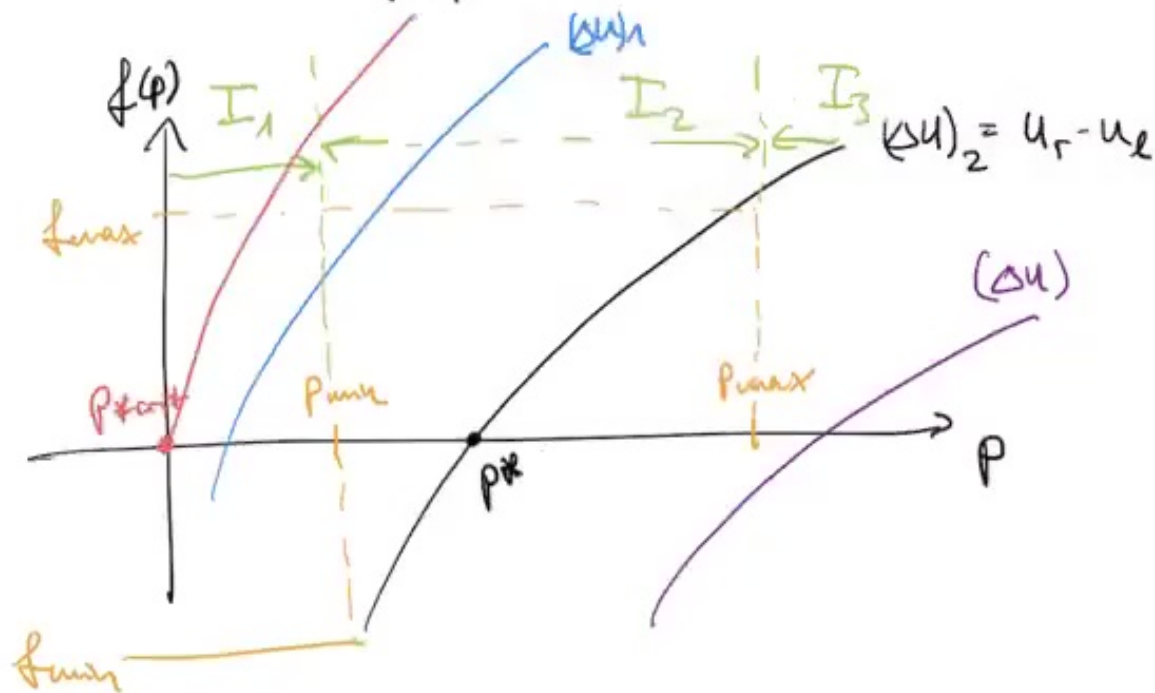
$I_1: p_* < p_l, p_r \Rightarrow 2 \text{ rarefaction waves}$

$I_2: \left. \begin{array}{l} p_l < p_* < p_r \\ p_r < p_* < p_l \end{array} \right\} \text{one rarefaction wave \& one shock}$

$I_3: p_* > p_l, p_r \Rightarrow 2\text{-shock solution}$

- $f''(p) < 0 \rightarrow$ concave

- $\begin{cases} f'(p) \\ f''(p) \end{cases} \xrightarrow{p \rightarrow \infty} 0$



$p_{min} = \min(p_l, p_r), \quad f_{min} = f(p_{min})$



2) Vacuum;

A physical solution to $f(p)$ requires

$$f(0) < 0 \quad (\text{monotonicity property})$$

For $p_* = p_{\text{crit}} = 0$ we have $p_* \leq p_l, p_r$
 \Rightarrow left & right rarefaction

$$\begin{aligned} \text{so } \underline{(u_r - u_l)_{\text{crit}}} &\stackrel{\downarrow f(p_{\text{crit}}) = 0}{=} -f_l(0) - f_r(0) \\ &= \frac{2}{\gamma-1} (c_l + c_r) (> u_r - u_l) \end{aligned}$$