

Title: Special Topics in Astrophysics - Numerical Hydrodynamics - Lecture 7

Speakers: Daniel Siegel

Collection: Special Topics in Astrophysics - Numerical Hydrodynamics

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URL: <http://pirsa.org/20100013>

Chap. 3: Finite Difference  
Methods



### 3.1 Basic notions of discretization

$$L: X \quad Y$$

$$u \quad \longmapsto \quad Lu = S$$

$\uparrow$  differential  
 $\downarrow$  operator

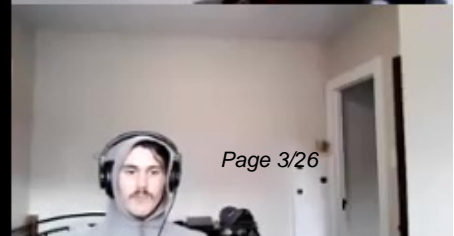
$$\downarrow D_h^x$$

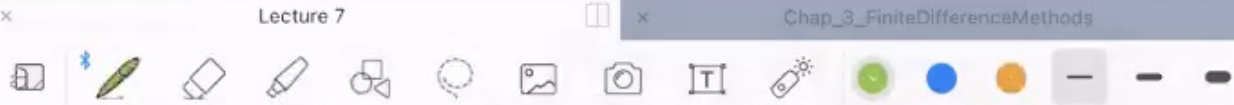
$$\downarrow \text{operator}$$

$$L_h: X_h \quad \longrightarrow \quad Y_h$$

continuum problem:  $Lu = S$       continuum PDE

For example:  $X = \{u \in C^2(\Omega) \mid |Lu| < \infty\}$





$$L: X \rightarrow Y$$

$$u \xrightarrow{\quad} Lu = S$$

$\uparrow$  differential operator  
 $\downarrow$

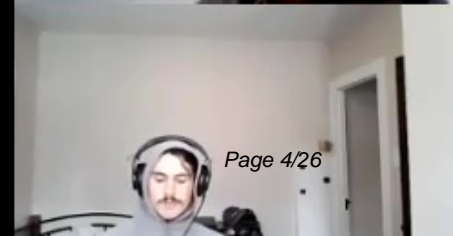
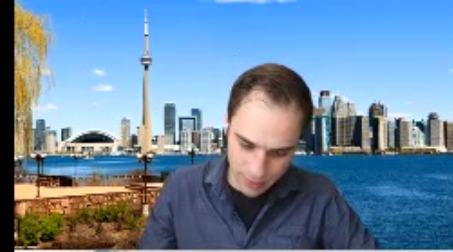
$$L_h: X_h \rightarrow Y_h$$

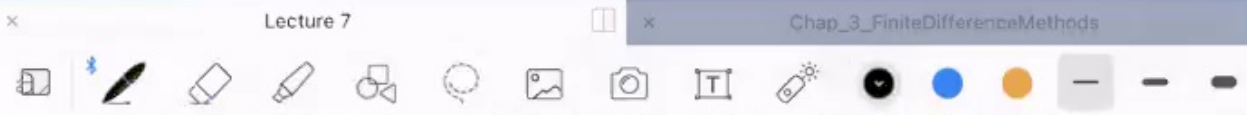
continuum problem:  $Lu = S$       continuum PDE

for example:  $X = \{u \in C^2(\Omega) \mid |u| < \infty\}$

$$L = \Delta, \square, \partial_t - \partial_x f(\cdot), \dots$$

discretized problem:  $L_h u_h = S_h$       discretized problem

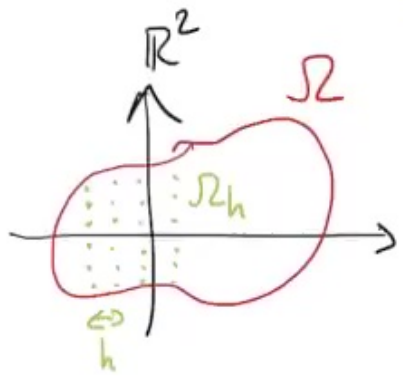




discretized problem:  $L_h u_h = S_h$  discretized problem

Typically:  $X_h = \{\text{grid functions}\}^n$   
 $= \{v: \Omega_h \rightarrow \mathbb{R}^n\}$

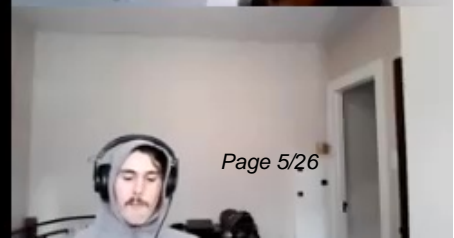
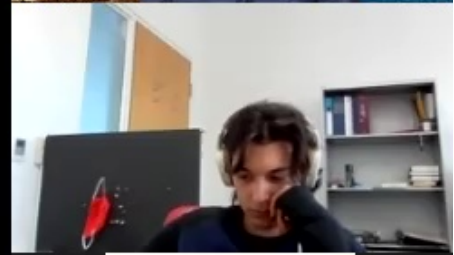
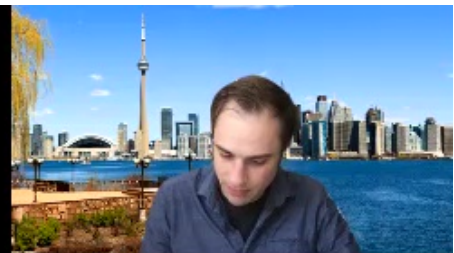
$\Omega_h = \text{"grid"}: \Omega \cap h\mathbb{Z}^n$   
 $h$ : discretization parameter

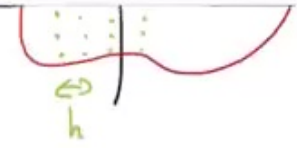
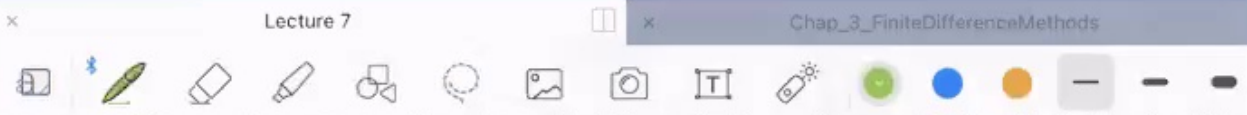


Notation:  $x \rightarrow x_i \in \{x_1, \dots, x_{N_x}\}$   
 $t \mapsto t_n \in \{t_1, \dots, t_{N_t}\}$

$D_h^x$ : discretization operator

$D_h^x: X \rightarrow X_h$





$$t \mapsto t_n \in \{t_0, \dots, t_N\}$$

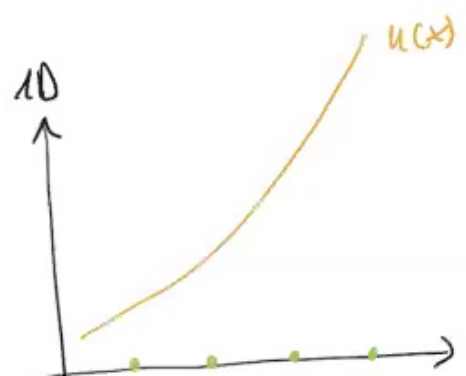
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$$D_h^x: X \rightarrow X_h$$

$$u \mapsto D_h^x u \equiv u_h$$

$$D_h^x u = u_h = \begin{cases} u|_{\Omega_h} & (u_h)_i^n = u(x_i, t^n) \end{cases}$$

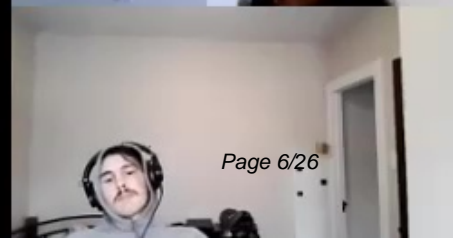
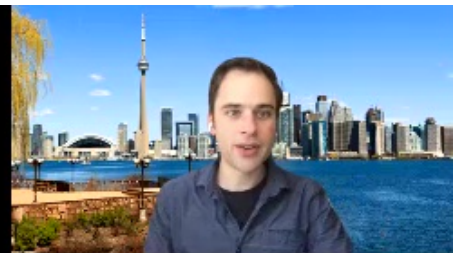
*finite differencing methods*

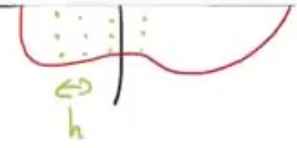
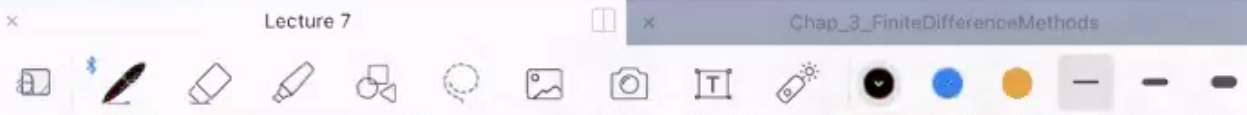


$$(u_h)_i^n = \frac{1}{\Delta x} \int_{x_{i-1/2}}^{x_{i+1/2}} u(x, t^n) dx$$

spatial average across  $\Delta x$

*local volume*





$$t \mapsto t_n \in \{t_0, \dots, t_N\}$$

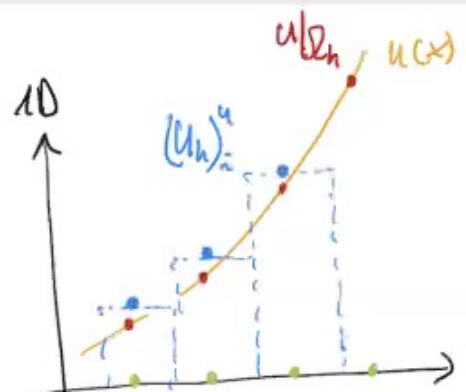
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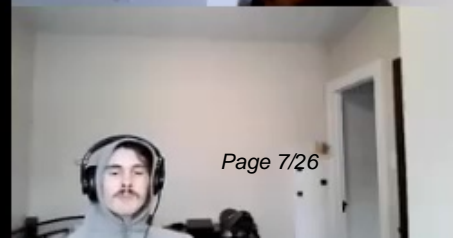
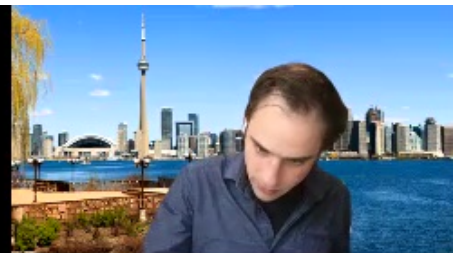
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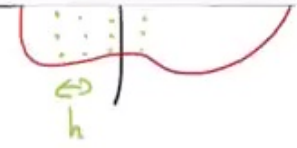
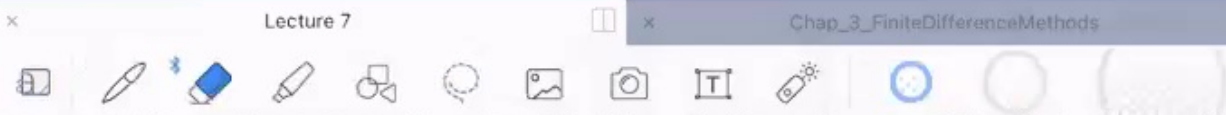


$$(u_h)_i^n = \frac{1}{\Delta x} \int_{x_{i-\frac{1}{2}}}^{x_{i+\frac{1}{2}}} u(x, t^n) dx$$

spatial average across  $\Delta x$

*finite volume*





$$t \mapsto t_n \in \{t_0, \dots, t_N\}$$

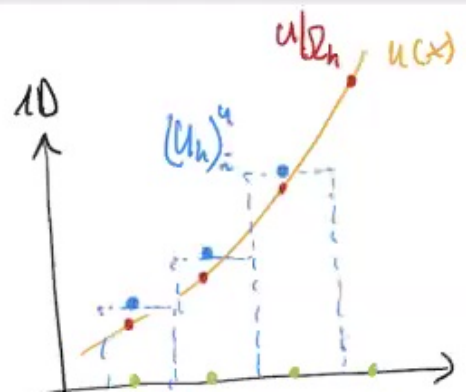
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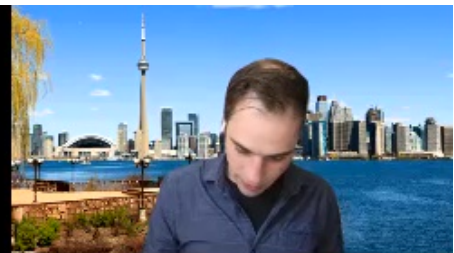
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$$(u_h)_i^n = \frac{1}{\Delta x} \int_{x_{i-1/2}}^{x_{i+1/2}} u(x, t^n) dx$$

spatial average across  $\Delta x$

*finite volume*





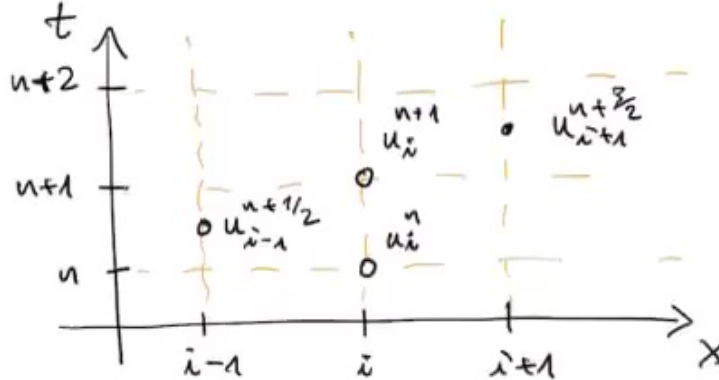


## 3.2 Finite difference approximations

Notation:

$$u(x_i, t^n) \equiv u_i^n$$

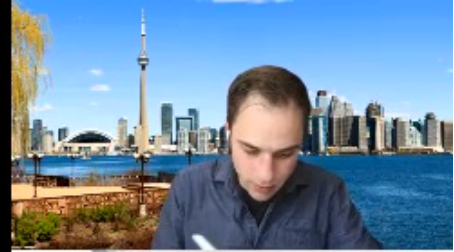
$$u(x_i + \frac{\Delta x}{2}, t^n) \equiv u_{i+\frac{1}{2}}^n$$



### 3.2.1 Partial derivatives

Consider Taylor expansion:

$$u_{i+1}^n$$



### 3.2.1 Partial derivatives

Consider Taylor expansion:

$$u_{i+1}^n = u_i^n + \frac{\partial u}{\partial x} \Big|_i \Delta x + \frac{1}{2} \frac{\partial^2 u}{\partial x^2} \Big|_i (\Delta x)^2 + \mathcal{O}(\Delta x^3)$$

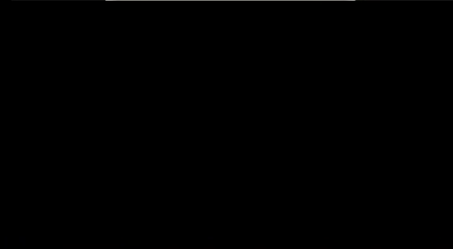
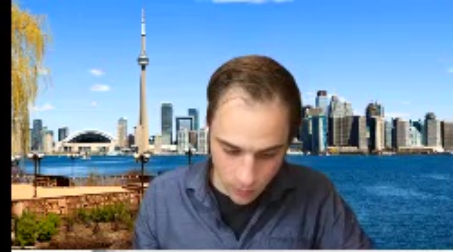
Solve for  $\frac{\partial u}{\partial x}$ :

$$\frac{\partial u}{\partial x} \Big|_{in} = \frac{u_{i+1}^n - u_i^n}{\Delta x} - \left( \frac{1}{2} \frac{\partial^2 u}{\partial x^2} \Big|_i \Delta x + \mathcal{O}(\Delta x^2) \right)$$

$$\boxed{D_1^+ u} = \frac{u_{i+1}^n - u_i^n}{\Delta x}$$

"stencil size"

"forward difference approximation"





Solve for  $\frac{du}{dx}$ :

$$(*) \frac{du}{dx} \Big|_i = \frac{u_{i+1}^n - u_i^n}{\Delta x} - \frac{1}{2} \frac{\partial^2 u^2}{\partial x^2} \Big|_i \Delta x + O(\Delta x^3)$$

$$D_1^+ u = \frac{u_{i+1}^n - u_i^n}{\Delta x}$$

"stencil size"

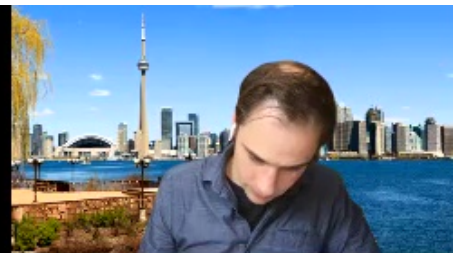
"forward difference approximation"

accuracy / truncation error:

$$D_1^+ u - \frac{du}{dx} \stackrel{(*)}{=} \frac{1}{2} \frac{\partial^2 u}{\partial x^2} \Big|_i \Delta x + O(\Delta x^2)$$

↑  
first-order accurate

Similarly:



Similarly:

$$D_1^- u = \frac{u_i^n - u_{i-1}^n}{\Delta x}$$

first-order accurate

"backward  
difference  
approximation"

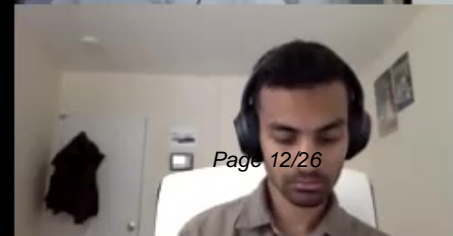
Improve accuracy:

$$\text{forward: } u_{i+1}^n = u_i^n + \frac{\partial u}{\partial x} \Big|_i \frac{\Delta x}{2} + \frac{1}{2} \frac{\partial^2 u}{\partial x^2} \Big|_i \left(\frac{\Delta x}{2}\right)^2 + \dots$$

$$\text{backward: } u_{i-1}^n = u_i^n - \frac{\partial u}{\partial x} \Big|_i \frac{\Delta x}{2} + \frac{1}{2} \frac{\partial^2 u}{\partial x^2} \Big|_i \left(\frac{\Delta x}{2}\right)^2 + \dots$$

↓ subtract

$$\leadsto \frac{\partial u}{\partial x} \Big|_i = \frac{u_{i+1}^n + u_{i-1}^n}{2\Delta x} -$$



Improve accuracy:

$$\text{forward: } u_{i+\frac{1}{2}}^n = u_i^n + \left. \frac{\partial u}{\partial x} \right|_i \frac{\Delta x}{2} + \frac{1}{2} \left. \frac{\partial^2 u}{\partial x^2} \right|_i \left( \frac{\Delta x}{2} \right)^2 + \dots$$

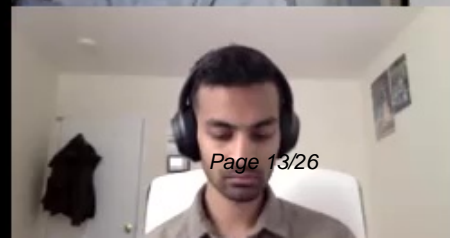
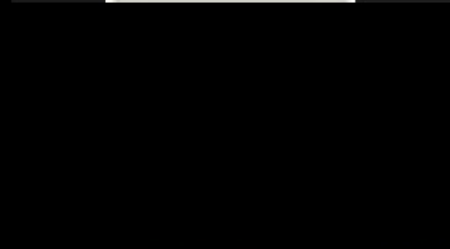
$$\text{backward: } u_{i-\frac{1}{2}}^n = u_i^n - \left. \frac{\partial u}{\partial x} \right|_i \frac{\Delta x}{2} + \frac{1}{2} \left. \frac{\partial^2 u}{\partial x^2} \right|_i \left( \frac{\Delta x}{2} \right)^2 + \dots$$

↓ subtract

$$\leadsto \left. \frac{\partial u}{\partial x} \right|_i = \frac{u_{i+\frac{1}{2}}^n - u_{i-\frac{1}{2}}^n}{\Delta x} + \mathcal{O}(\Delta x)^2$$

$$D_1^0 u = \frac{u_{i+\frac{1}{2}}^n - u_{i-\frac{1}{2}}^n}{\Delta x}$$

"centred  
difference  
approximation"



$$\text{forward: } u_{i+\frac{1}{2}}^n = u_i^n + \left. \frac{\partial u}{\partial x} \right|_i \frac{\Delta x}{2} + \frac{1}{2} \left. \frac{\partial^2 u}{\partial x^2} \right|_i \left( \frac{\Delta x}{2} \right)^2 + \dots$$

$$\text{backward: } u_{i-\frac{1}{2}}^n = u_i^n - \left. \frac{\partial u}{\partial x} \right|_i \frac{\Delta x}{2} + \frac{1}{2} \left. \frac{\partial^2 u}{\partial x^2} \right|_i \left( \frac{\Delta x}{2} \right)^2 + \dots$$

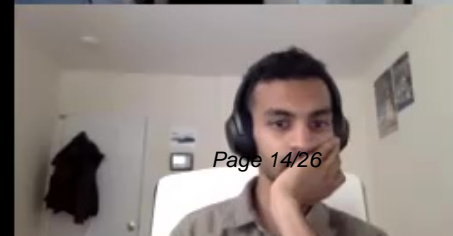
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$$\leadsto \left. \frac{\partial u}{\partial x} \right|_i = \frac{u_{i+\frac{1}{2}}^n - u_{i-\frac{1}{2}}^n}{\Delta x} + \mathcal{O}(\Delta x)^2$$

$$D_1^0 u = \frac{u_{i+\frac{1}{2}}^n - u_{i-\frac{1}{2}}^n}{\Delta x}$$

"centred  
difference  
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Remark: By virtue of Taylor expansion we assume that



the function  $u$  is continuously differentiable.

Higher-order derivatives:

$$\begin{aligned} \text{Example: } \frac{\partial^2 u}{\partial x^2} &\approx D_1^+ D_1^- u = \frac{1}{\Delta x} (D_1^+ u_i^n - D_1^- u_i^n) \\ &= \frac{u_{i+1}^n - 2u_i^n + u_{i-1}^n}{(\Delta x)^2} \\ &\equiv D^2 u \end{aligned}$$

General approach:

Goal: compute appro



Example:  $\frac{\partial^2 u}{\partial x^2} \approx D_1^+ D_1^- u = \frac{1}{\Delta x} (D_1^+ u_i^n - D_1^- u_{i-1}^n)$

$$= \frac{u_{i+1}^n - 2u_i^n + u_{i-1}^n}{(\Delta x)^2}$$

$$\equiv D^2 u$$

General approach:

Goal: compute approximation to

$$u^{(k)} = \frac{\partial^k u}{\partial x^k}$$

based on a stencil of  $n \geq k+1$   
points  $x_1, \dots, x_n$

Example: want one-sided approximation  
to  $\frac{\partial u}{\partial x} \Big|_i$  based on  $u_i^n, u_{i-1}^n, u_{i-2}^n$





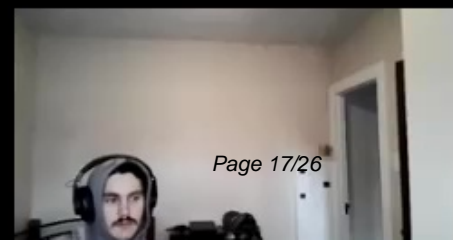
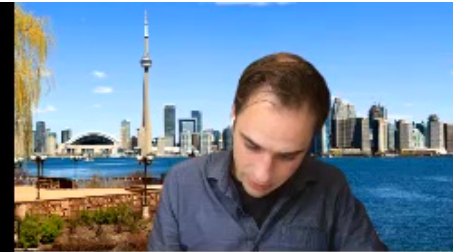
Example: want one-sided approximation  
to  $\frac{\partial u}{\partial x} \Big|_i$  based on  $u_i^n, u_{i-1}^n, u_{i-2}^n$

$$\leadsto \underline{u_{i-1}} = u_i - u^{(1)} \Big|_i \Delta x + \frac{1}{2} u^{(2)} \Big|_i (\Delta x)^2 - \frac{1}{6} u^{(3)} \Big|_i (\Delta x)^3 + O(\Delta x^4)$$

$$\underline{u_{i-2}} = u_i - u^{(1)} \Big|_i (2\Delta x) + \frac{1}{2} u^{(2)} \Big|_i (2\Delta x)^2 - \frac{1}{6} u^{(3)} \Big|_i (2\Delta x)^3 + O(\Delta x^4)$$

Ansatz:

$$\begin{aligned} D_2 u \Big|_i &= c_1 u_i + c_2 \underline{u_{i-1}} + c_3 \underline{u_{i-2}} \\ &= (c_1 + c_2 + c_3) u_i - (c_2 + 2c_3) \Delta x u^{(1)} \\ &\quad + \frac{1}{2} (c_2 + 4c_3) (\Delta x)^2 \end{aligned}$$

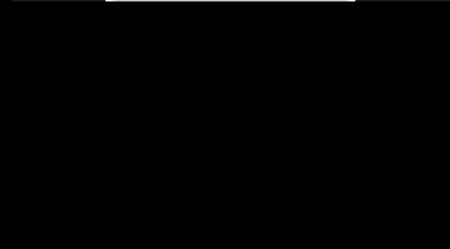


Ansatz:

$$\begin{aligned}
 D_2 u|_i &= c_1 u_i + c_2 u_{i-1} + c_3 u_{i-2} \\
 &= \underbrace{(c_1 + c_2 + c_3)}_{=0} u_i - \underbrace{(c_2 + 2c_3)}_{=1} \Delta x u^{(1)} \\
 &\quad + \underbrace{\frac{1}{2}(c_2 + 4c_3)}_{=0} (\Delta x)^2 u^{(2)} \\
 &\quad - \frac{1}{6}(c_2 + 8c_3) (\Delta x)^3 u^{(3)} + \dots
 \end{aligned}$$

(note: coefficients are of the form  $\frac{1}{(l-1)!} \sum_{m=1}^n c_m (x_m - x_i)^{l-1}$ )

$$\left. \begin{aligned}
 c_1 + c_2 + c_3 &= 0 \\
 -(c_2 + c_3) \Delta x &= 1 \\
 \frac{1}{2}(c_2 + 4c_3) (\Delta x)^2 &= 0
 \end{aligned} \right\} \begin{aligned}
 c_1 &= \frac{3}{2\Delta x} \\
 c_2 &= -\frac{2}{\Delta x} \\
 c_3 &= \frac{1}{2\Delta x}
 \end{aligned}$$



Truncation error:

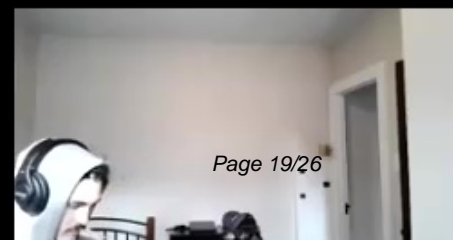
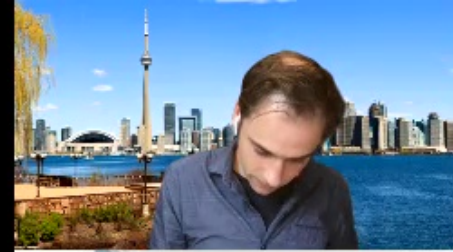
$$\begin{aligned}
 D_2 u|_i - u^{(1)}|_i &= -\frac{1}{6}(c_2 + 8c_3)(\Delta x)^3 u^{(3)}|_i \\
 &= -\frac{1}{3}(\Delta x)^2 \underbrace{u^{(3)}|_i}_{+ \dots} + \dots \\
 &\rightarrow \underline{O(\Delta x^2)}
 \end{aligned}$$

Generalize procedure: Consider stencil

$\{x_\ell\}_{\ell=1, \dots, n}$  around  $x_i$ ,  $n \geq k+1$ .

Consider  $n$  Taylor series  $\ell=1, \dots, n$   
and obtain  $u^{(k)}$  that is

$O(\Delta x^p)$  ( $p = n - k$ ) accurate.





$$\begin{aligned}
 D_2 u|_i - u^{(1)}|_i &= -\frac{1}{6}(c_2 + 8c_3)(\Delta x)^3 u^{(3)}|_i + \dots \\
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 &\rightarrow \underline{O(\Delta x^2)}
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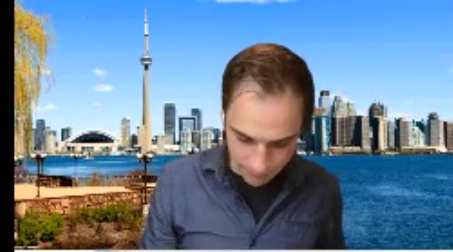
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and obtain  $u^{(k)}$  that is

$O(\Delta x)^p$  ( $p = n - k$ ) accurate.

$p > n - k$  achievable if additional higher order terms cancel ( $\rightarrow$  centered differences)



General: increasing stencil size  
increases accuracy.

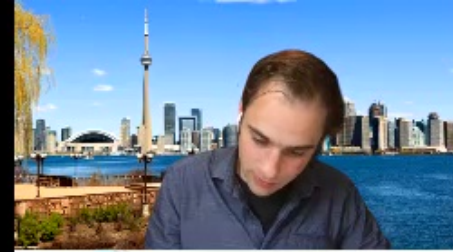
### 3.2.2 Sample Discretizations

1D advection equation

$$u_t + a u_x = 0$$

$$u(x, 0) = u_0(x)$$

$$(u(0, t) = u(1, t) = 0)$$



## 3.2.2 Sample Discretizations

1D advection equation

$$u_t + a u_x = 0 \quad 0 \leq x \leq 1, t \geq 0$$

$$u(x, 0) = u_0(x)$$

$$(u(0, t) = u(1, t) = 0)$$

(i) Consider simple forward in time,  
centred in space FDA:

$$u_x \Big|_i^n \approx \frac{u_{i+1}^n - u_{i-1}^n}{2\Delta x}$$

$$u \Big|_i^n = u_i^n$$



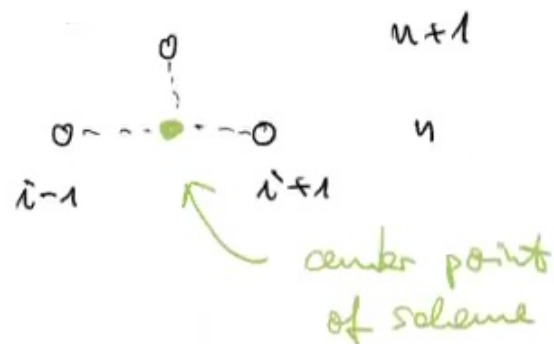
$$(u(0,t) = u(1,t) = 0)$$

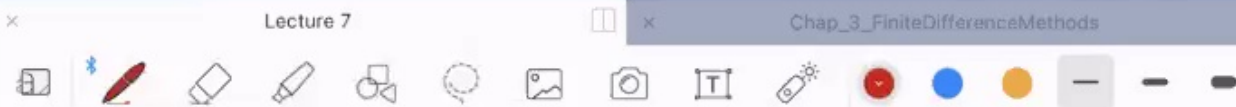
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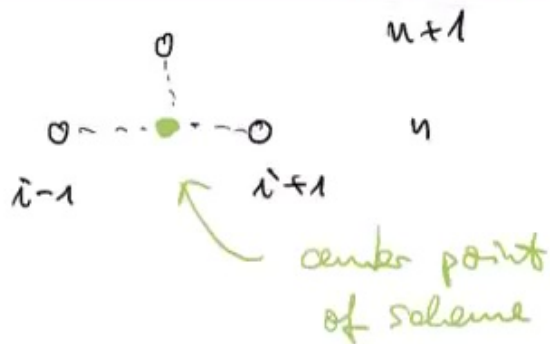
$$u_t \Big|_i^n = \frac{u_i^{n+1} - u_i^n}{\Delta t}$$

stencil:





stencil:



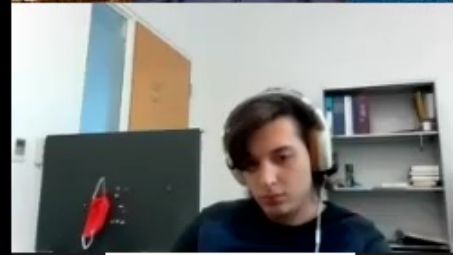
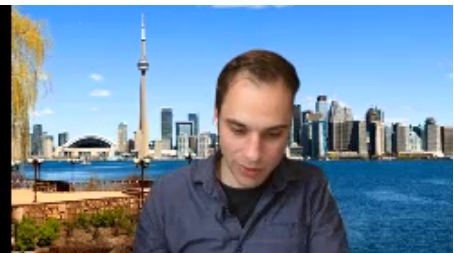
$$\leadsto \frac{u_i^{n+1} - u_i^n}{\Delta t} + \frac{a}{2\Delta x} (u_{i+1}^n - u_{i-1}^n) = 0$$

$$\Leftrightarrow u_i^{n+1} = u_i^n - \frac{a\Delta t}{2\Delta x} (u_{i+1}^n - u_{i-1}^n)$$

linear system for  $u^{n+1}$ :  $A u^{n+1} = b$

$$\begin{pmatrix} 1 & & & 0 \\ & \ddots & & \\ & & \ddots & \\ 0 & & & 1 \end{pmatrix} \begin{pmatrix} u_1^{n+1} \\ \vdots \\ u_i^{n+1} \\ \vdots \\ u_I^{n+1} \end{pmatrix} = \begin{pmatrix} u_i^n - \lambda (u_{i+1}^n - u_{i-1}^n) \end{pmatrix}$$

$\uparrow x_i \in \{x_1, \dots, x_I\}$





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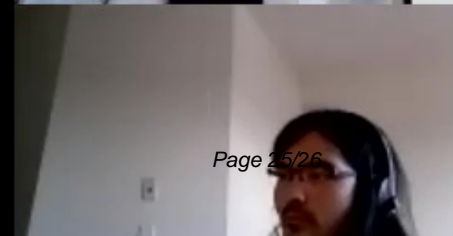
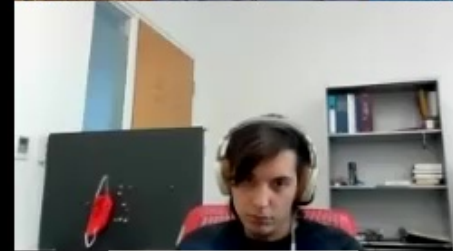
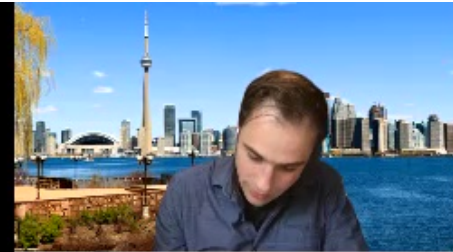
$\mathcal{O}(\Delta t, (\Delta x)^2)$

no explicit scheme

will see: unstable for any fixed  $\frac{\Delta t}{\Delta x}$

(ii) Lax-Friedrichs method

replace:  $u_i^n \rightarrow \frac{1}{2} (u_{i-1}^n + u_{i+1}^n)$





$$2\Delta x (u_{i+1}^{n+1} - u_{i-1}^{n+1})$$

linear system for  $u^{n+1}$ :  $A u^{n+1} = b$

$$\begin{pmatrix} 1 & & & 0 \\ & \ddots & & \\ & & \ddots & \\ 0 & & & 1 \end{pmatrix} \begin{pmatrix} u_1^{n+1} \\ \vdots \\ u_I^{n+1} \end{pmatrix} = \begin{pmatrix} u_i^n - \lambda (u_{i+1}^n - u_{i-1}^n) \\ \vdots \\ \end{pmatrix}$$

$\uparrow x_i \in \{x_1, \dots, x_I\}$

$$O(\Delta t, (\Delta x)^2)$$

no explicit scheme

will see: unstable for any fixed  $\frac{\Delta t}{\Delta x}$

(ii) Lax-Friedrichs method

$$\text{replace: } u_i^n \rightarrow \frac{1}{2} (u_{i-1}^n + u_{i+1}^n)$$

$$\text{no } u_i^{n+1} = \frac{1}{2} (u_{i+1}^n + u_{i-1}^n) - \frac{a \Delta t}{2\Delta x} (u_{i+1}^n - u_{i-1}^n)$$

