

Title: How does a dark compact object ringdown?

Speakers: Elisa Maggio

Series: Strong Gravity

Date: October 01, 2020 - 1:00 PM

URL: <http://pirsa.org/20100005>

Abstract: Gravitational waves from the coalescence of compact binaries provide a unique opportunity to test gravity in strong field regime. In particular, the postmerger phase of the gravitational signal is a proxy for the nature of the remnant.

This is of particular interest in view of some quantum-gravity models which predict the existence of horizonless dark compact objects that overcome the paradoxes associated to black holes. Such dark compact objects can emit a modified ringdown with respect to the black hole case and late-time gravitational wave echoes as characteristic fingerprints.

In this talk, I develop a generic framework to the study of the ringdown of dark compact objects and provide a gravitational-wave template for the echo signal. Finally, I assess the detectability of dark compact objects with current and future gravitational-wave detectors.

Strong Gravity Seminar @Perimeter Institute  
October 1<sup>st</sup>, 2020

## How does a dark compact object ringdown?

Elisa Maggio  
Sapienza University of Rome  
<https://web.uniroma1.it/gmunu>



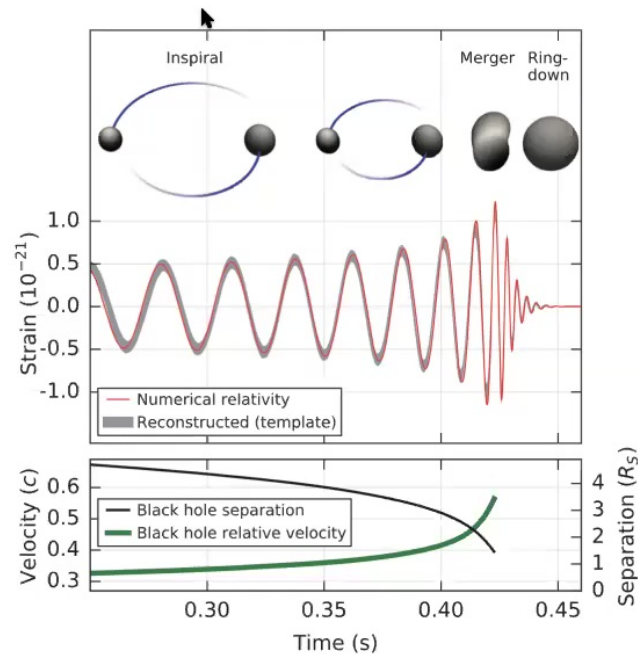
SAPIENZA  
UNIVERSITÀ DI ROMA



DarkGRA 



# Gravitational waves from binary mergers



Abbott+, PRL **116**, 061102 (2016)

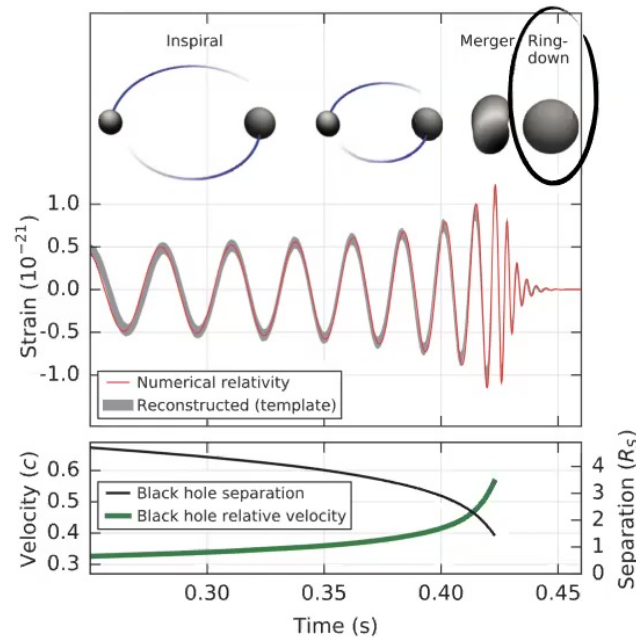
Up to date 15 gravitational wave events from the coalescence of compact binaries have been detected.

The signal has 3 stages:

- Inspiral
- Merger
- Ringdown



# Gravitational waves from binary mergers



Abbott+, PRL **116**, 061102 (2016)

Up to date 15 gravitational wave events from the coalescence of compact binaries have been detected.

The signal has 3 stages:

- Inspiral
- Merger
- Ringdown

**What is the nature of the compact remnant?**

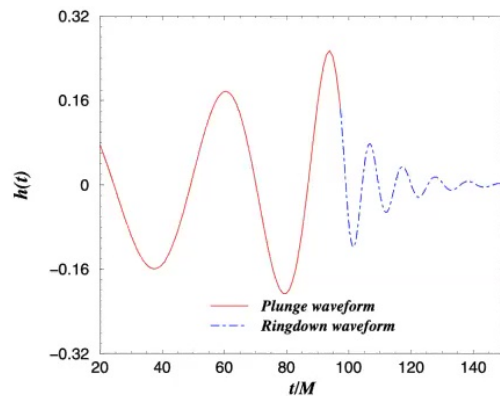


## Ringdown stage

The ringdown stage is dominated by the **quasi-normal modes** of the remnant which describe the response of the compact object to a perturbation.

$$\omega = \omega_R + i\omega_I$$

The signal can be modeled as a sum of exponentially damped sinusoids:



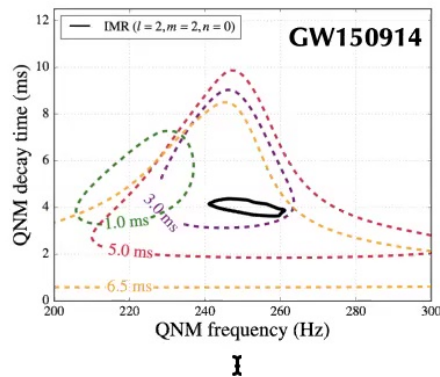
Buonanno, CQG **19**, 1267-1278 (2002)

$$f_{\text{GW}|ringdown} = \frac{\omega_R}{2\pi}$$
$$\tau_{\text{damping}} = -\frac{1}{\omega_I}$$

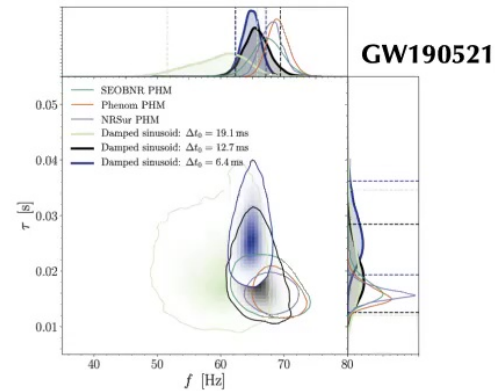


# Ringdown detections

The least-damped quasi-normal mode has been observed in the ringdown of the events GW150914 and GW190521.



Abbott+, PRL **116**, 221101 (2016)



Abbott+, ApJ **900**, L13 (2020)

Both detections are compatible with **Kerr black hole remnants**.  
However the characterization of the remnant is still an open problem.



# Test of the black hole paradigm

- A test of the no-hair theorem requires the identification of **at least two quasi-normal mode frequencies** in the ringdown.

Dreyer+, CQG **21**, 787 (2004)

Kerr black holes are *uniquely* determined by 2 parameters:

- Mass
- Angular momentum

Carter, PRL **26**, 331 (1971); Robinson, PRL **34**, 905 (1975)

- Louder gravitational wave events and improvements of the **detector sensitivity** will allow to test the black hole paradigm.

Berti, Cardoso, Will, PRD **73**, 064030 (2006)

I



## Some questions

- I. Are there alternative models to black holes?
- II. What is the gravitational wave signal that they produce?
- III. Is their gravitational wave signal detectable?





# Alternatives to black holes

There is a zoo of theoretical compact objects **without horizon** which:

- are solutions to **modified gravity** and can overcome paradoxes of BHs, e.g., curvature singularity and Hawking information loss.  
Mazur, Mottola, PNAS **101**, 9545-9550 (2004); Mathur, Fortsch. Phys. **53**, 793-827 (2005)
- are solutions to **GR** in the presence of dark matter/exotic fields  
Liebling, Palenzuela, LRR **20**, 5 (2017); Brito+, Phys. Lett. B **752**, 291-295 (2016)



# Alternatives to black holes

There is a zoo of theoretical compact objects **without horizon** which:

- are solutions to **modified gravity** and can overcome paradoxes of BHs, e.g., curvature singularity and Hawking information loss.  
Mazur, Mottola, PNAS **101**, 9545-9550 (2004); Mathur, Fortsch. Phys. **53**, 793-827 (2005)
- are solutions to **GR** in the presence of dark matter/exotic fields  
Liebling, Palenzuela, LRR **20**, 5 (2017); Brito+, Phys. Lett. B **752**, 291-295 (2016)

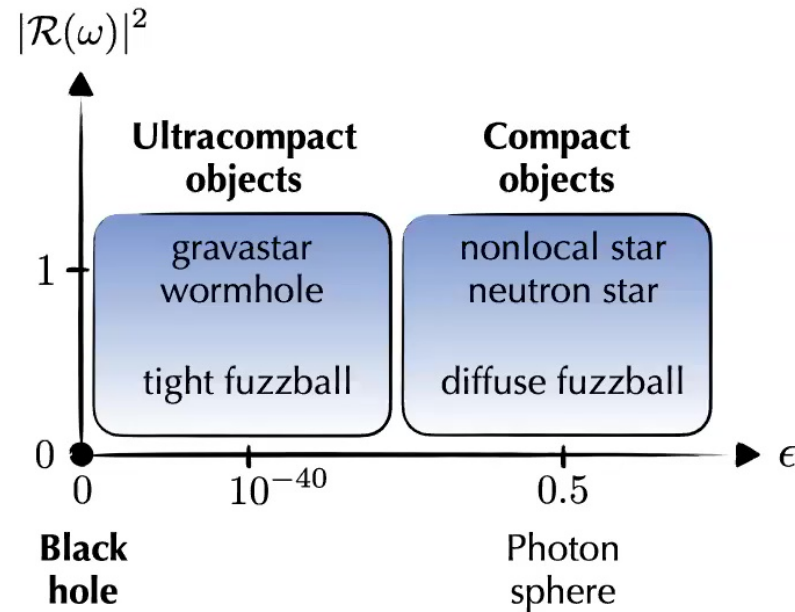
- are not excluded by GW and electromagnetic observations  
Abbott+, ApJ **896**: L44 (2020); Abbott+, PRL **125**, 101102 (2020); Calderón Bustillo+, arXiv: 2009.05376 (2020); EHT, ApJ **875**, L5 (2019)
- quantify the existence of horizons



# Dark compact objects

We analyze a generic model which deviates from a black hole for its:

- **Compactness**  
since the radius of the object is at  $r_0 = r_+(1 + \epsilon)$
- **“Darkness”**  
which is related to the reflectivity of the object  $\mathcal{R}(\omega)$



Cardoso, Pani, Nat. Astron. 1: 586-591 (2017)



# Alternatives to black holes

Model	Formation	Stability	GWs
Fluid stars	✗	✓	✓
Anisotropic stars	✗	✓	✓
Boson stars & oscillatons	✓	✓	✓
→ Gravastars	✗	✓	~
AdS bubbles	✗	✓	✗
→ Wormholes	✗	✓	~
Fuzzballs	✗	✗	~
→ Superspinars	✗	✓	~
2 – 2 holes	✗	✗	~
Collapsed polymers	✗	✗	~
Quantum bounces / black stars	✗	✗	~
Quantum stars*	✗	✗	✗
Fire-walls*	✗	✗	~



Formation



No singularities



No semiclassical paradoxes



Exotic matter

I

Cardoso, Pani, LRR 22:4 (2019)





# What is the gravitational wave signal that dark compact objects produce?

I

# Quasi-normal mode spectrum

We can distinguish dark compact objects from BHs via quasi-normal modes.

We consider a gravitational perturbation

$$\frac{d^2\psi}{dz^2} + V(z, \omega)\psi = 0$$

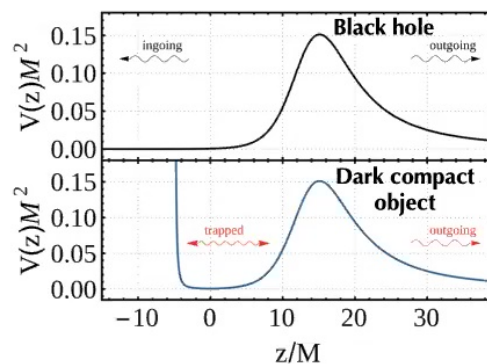
Detweiler, Proc. R. Soc. Lond. A 352 (1977)

+ 2 boundary conditions:

eigenvalue problem for the QNMs

- At infinity: outgoing waves
- At  $r_0$ : ultracompact object with surface reflectivity coefficient  $\mathcal{R}(\omega)$

$$\psi(r_0) \sim A_{\text{in}} e^{-i\tilde{\omega}z_0} + A_{\text{out}} e^{i\tilde{\omega}z_0}$$



Cardoso, Pani, LRR 22:4 (2019)



# Quasi-normal mode spectrum

We can distinguish dark compact objects from BHs via quasi-normal modes.

We consider a gravitational perturbation

$$\frac{d^2\psi}{dz^2} + V(z, \omega)\psi = 0$$

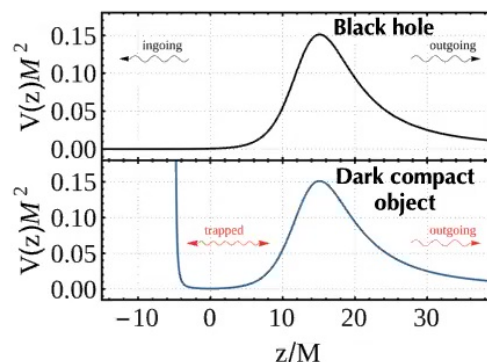
Detweiler, Proc. R. Soc. Lond. A 352 (1977)

+ 2 boundary conditions:

eigenvalue problem for the QNMs

- At infinity: outgoing waves
- At  $r_0$ : ultracompact object with surface reflectivity coefficient  $\mathcal{R}(\omega)$

$$\psi(r_0) \sim A_{\text{in}} e^{-i\tilde{\omega}z_0} + A_{\text{out}} e^{i\tilde{\omega}z_0}$$

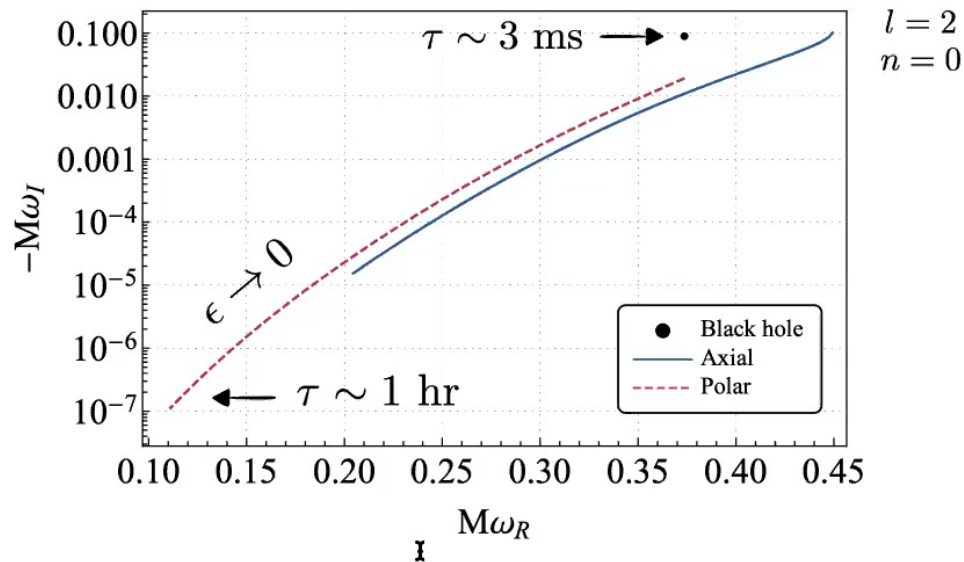


Cardoso, Pani, LRR 22:4 (2019)



# QNMs of ultracompact objects

Static, perfectly reflecting  $|\mathcal{R}(\omega)|^2 = 1$

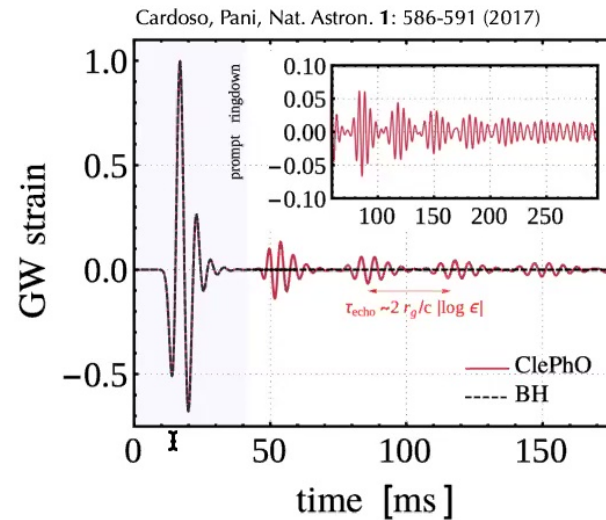


For  $\epsilon \rightarrow 0$ , the deviations from the black hole QNM are arbitrarily large and the QNMs are **low frequencies** and **long-lived**.

Cardoso, Franzin, Pani, PRL **116**, 171101 (2016)



# Gravitational-wave echoes



- The **prompt ringdown** is indistinguishable from that of a black hole since it is excited at the photon sphere.
- A modulated train of **gravitational-wave echoes** appear at late times due to trapped modes.



# Template for echoes

The signal emitted by a dark compact object can be written in terms of the signal emitted by a black hole Mark+, PRD **96**, 084002 (2017); Testa, Pani, PRD **98**, 044018 (2018)

$$\tilde{Z}_{\text{DCO}}^{\infty}(\omega) = \tilde{Z}_{\text{BH}}^{\infty}(\omega) + \mathcal{K} \tilde{Z}_{\text{BH}}^{\text{H}}(\omega)$$

↓  
Transfer function

- Parameters:**
- Standard BH ringdown:  $M, \chi, A_{+,x}, \phi_{+,x}, t_0$
  - 2 extra parameters:  $\epsilon, \mathcal{R}$



# Template for echoes

The signal emitted by a dark compact object can be written in terms of the signal emitted by a black hole Mark+, PRD **96**, 084002 (2017); Testa, Pani, PRD **98**, 044018 (2018)

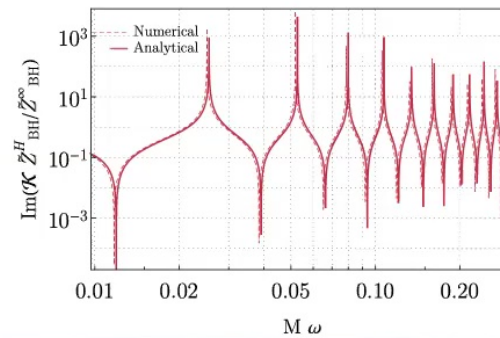
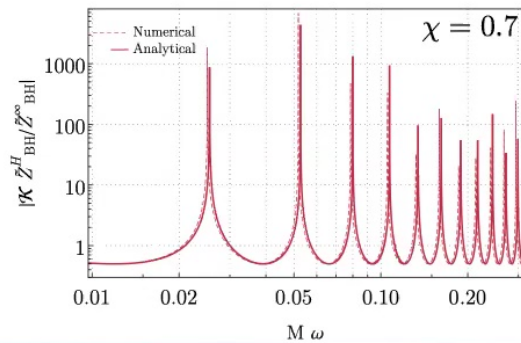
$$\tilde{Z}_{\text{DCO}}^{\infty}(\omega) = \tilde{Z}_{\text{BH}}^{\infty}(\omega) + \mathcal{K} \tilde{Z}_{\text{BH}}^{\text{H}}(\omega)$$

**Analytical  
low-frequency template**

EM, Testa, Bhagwat, Pani, PRD **100**, 064056 (2019)

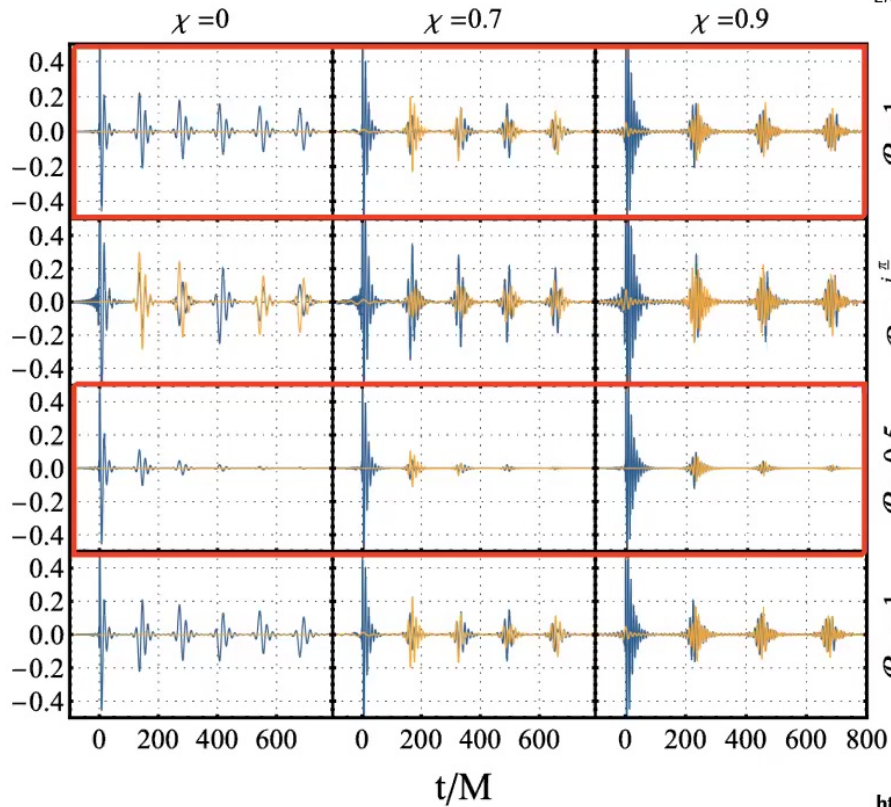
↓  
Transfer function

- Parameters:**
- Standard BH ringdown:  $M, \chi, A_{+, \times}, \phi_{+, \times}, t_0$
  - 2 extra parameters:  $\epsilon, \mathcal{R}$



# Ringdown+echo waveforms

EM, Testa, Bhagwat, Pani, PRD **100**, 064056 (2019)



— Plus polarization  
— Cross polarization

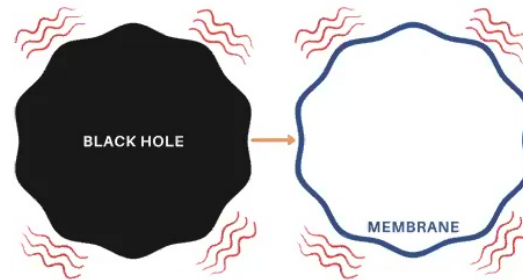
- The mixing of the polarizations occurs for a **spinning object** or a **complex  $\mathcal{R}$** .
- The **phase** in  $\mathcal{R}$  modifies the echo pattern.
- The **amplitude** of the echoes is related to  $\mathcal{R}$ .

<https://www.darkgra.org/gw-echo-catalogue.html>

# Ringdown of dark compact objects

**Membrane paradigm** Damour, PRD **18**, 10 (1978); Price, Thorne, PRD **33**, 4 (1986)

A static observer outside the horizon can replace the interior of a perturbed BH by a perturbed **fictitious** membrane located at the horizon.



The Israel junction conditions  $[[K_{ab} - Kh_{ab}]] = -8\pi T_{ab}$  impose that the membrane is a **viscous fluid** with shear viscosity  $\eta$  and bulk viscosity  $\zeta$ .

**I**

→ We generalize the membrane paradigm to any dark compact object with a Schwarzschild exterior. EM, Buoninfante, Mazumdar, Pani, PRD **102**, 064053 (2020)



# Quasi-normal mode spectrum

The dark compact object is subjected to gravitational perturbations:

$$\frac{d^2\psi}{dz^2} + [\omega^2 - V(r)] \psi = 0$$

Regge, Wheeler, PR **108**, 4 (1957)  
Zerilli, PRL **24**, 13 (1970)

+ 2 boundary conditions:

- At infinity: outgoing waves

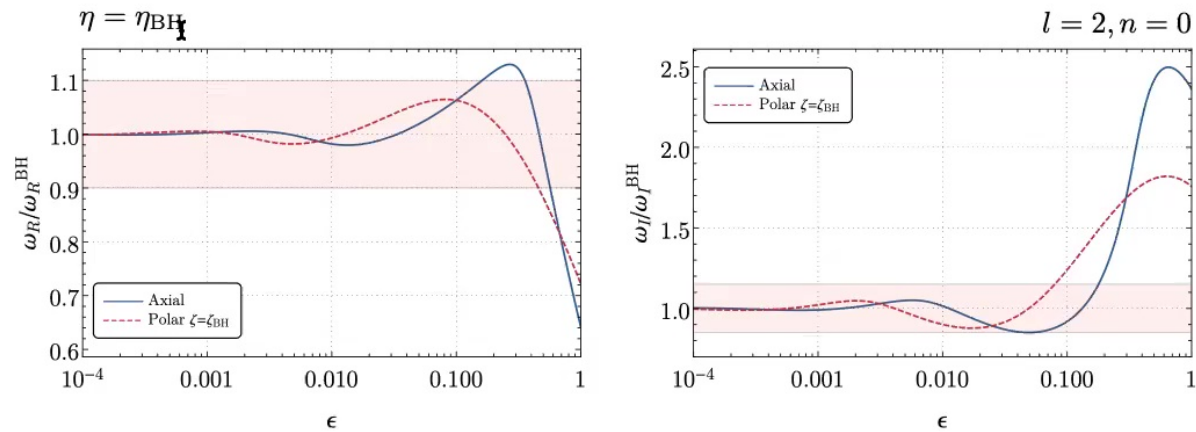
- At  $r_0$ :  $\frac{\psi'(z)}{\psi(z)} = \begin{cases} -\frac{i\omega}{16\pi\eta} - \frac{r_0^2 V(r_0)}{2(r_0 - 3M)} & \text{Axial} \\ -16\pi i\eta\omega + G(r_0, \omega, \eta, \zeta) & \text{Polar} \end{cases}$

**BH limit for**  
 $\eta = 1/(16\pi)$

Price, Thorne, PRD **33**, 4 (1986)



# Quasi-normal modes of dark compact objects

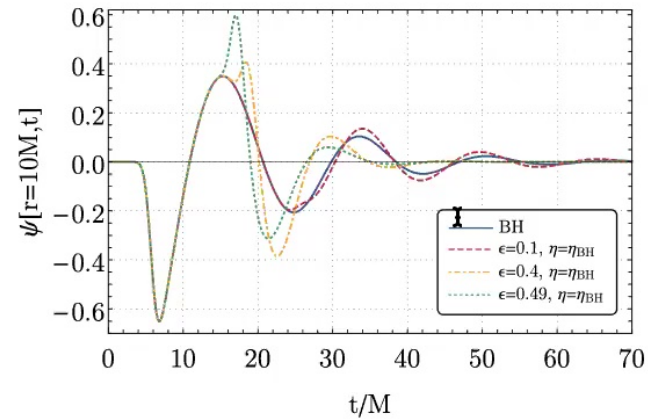


- The isospectrality of axial and polar modes in BHs is broken and the QNMs form a **doublet**.
- The measurement accuracy of the quasi-normal mode of GW150914 agrees with a dark compact object with  $\epsilon \lesssim 0.1$ .

EM, Buoninfante, Mazumdar, Pani, PRD **102**, 064053 (2020)



# Ringdown of dark compact objects



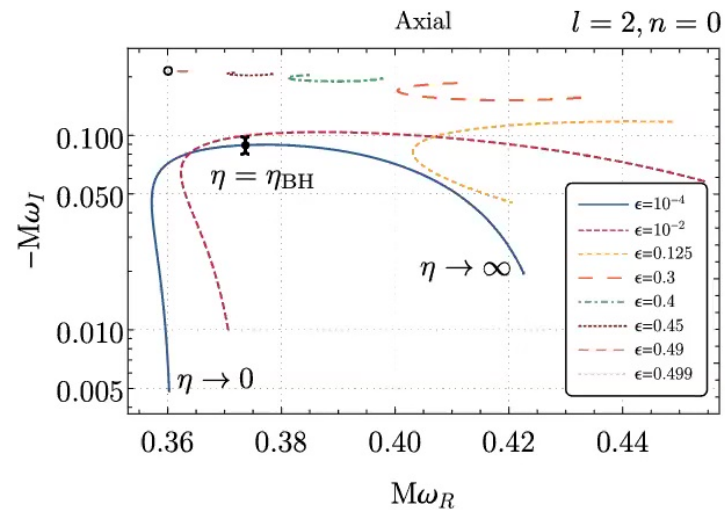
- The **prompt ringdown** is modified with respect to the BH case.
- At late time the prompt ringdown is dominated by the modified QNM of the object.
- Subsequent echoes are absent.

EM, Buoninfante, Mazumdar, Pani, PRD **102**, 064053 (2020)





# Universal axial quasi-normal mode



As the radius of the object approaches the photon sphere, the dark compact object has a universal QNM regardless its interior structure:

$$M\omega_2^{\text{axial}} = 0.3601 - i0.2149, \quad r_0 \rightarrow 3M$$

**Caveat:** gravastars

EM, Buoninfante, Mazumdar, Pani, PRD **102**, 064053 (2020)



# Detectability of GW echoes

- A tentative evidence for echoes in LIGO/Virgo data has been reported

Abadi+, PRD **96**, 082004 (2017); Conklin, Holdom, PRD **98**, 044021 (2018); Abadi, Afshordi, JCAP **11**, 010 (2019)

I

- Independent searches argued that the statistical significance of echoes is low and consistent with noise

Westerweck+, PRD **97**, 124037 (2018); Nielsen+, PRD **99**, 104012 (2019); Uchikata+, PRD **100**, 062006 (2019); Lo+, PRD **99**, 084052 (2019); Tsang+, PRD **101**, 064012 (2020)

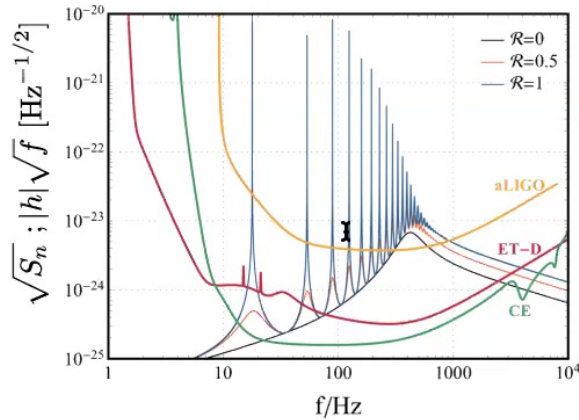
- Reviews

Cardoso, Pani, Living Reviews in Relativity **22:4** (2019)  
Abadi, Afshordi, Oshita, Wang, Universe **6** no. 3, 43 (2020)

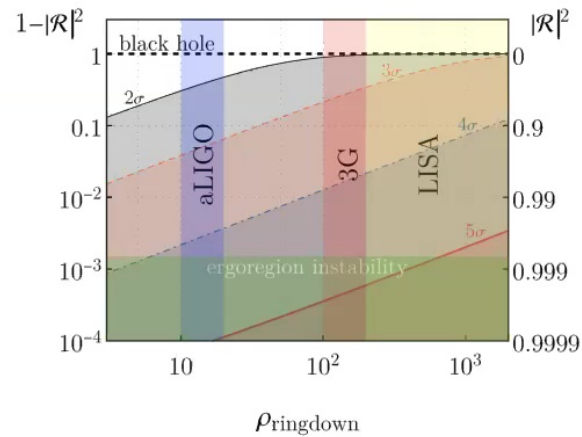


# Constraints on the reflectivity

- For **perfectly reflecting** dark compact objects, the energy emitted in the echoes is larger than the energy emitted in the ringdown.
- LISA will be able to probe values of the **reflectivity close to the BH one**.



Testa, Pani, PRD **98**, 044018 (2018)

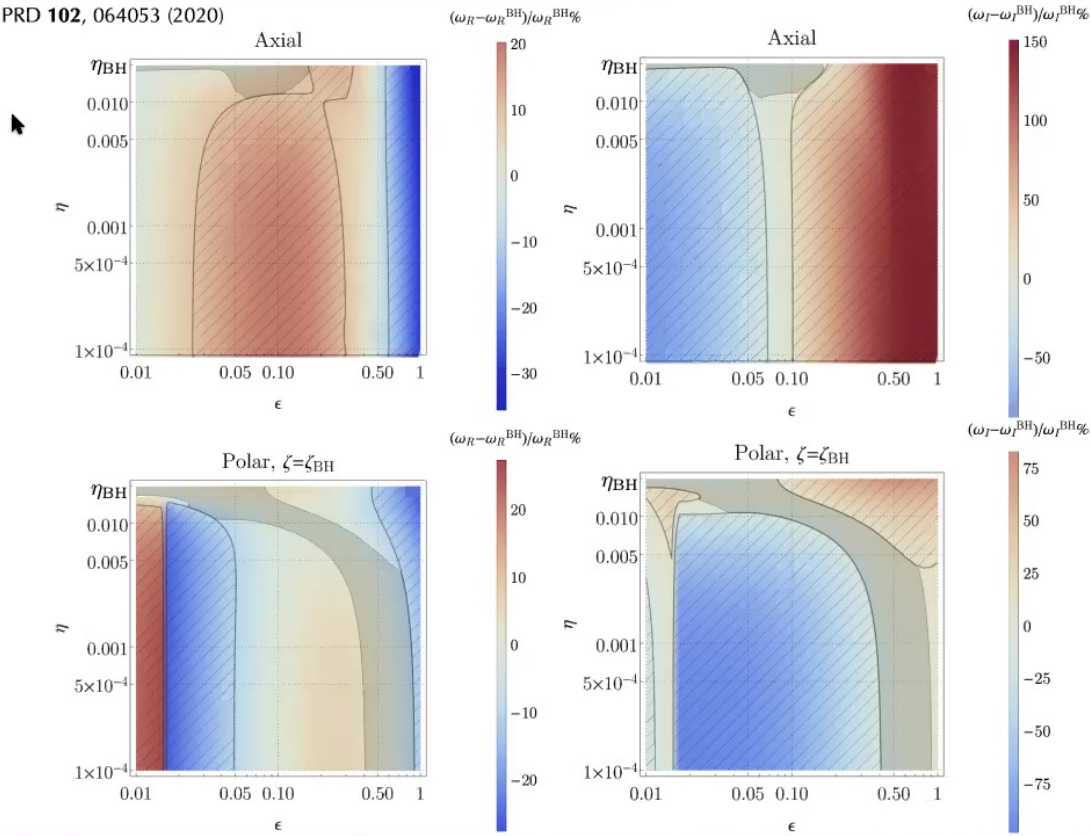


EM, Testa, Bhagwat, Pani, PRD **100**, 064056 (2019)



# Constraints on the compactness

EM+, PRD **102**, 064053 (2020)



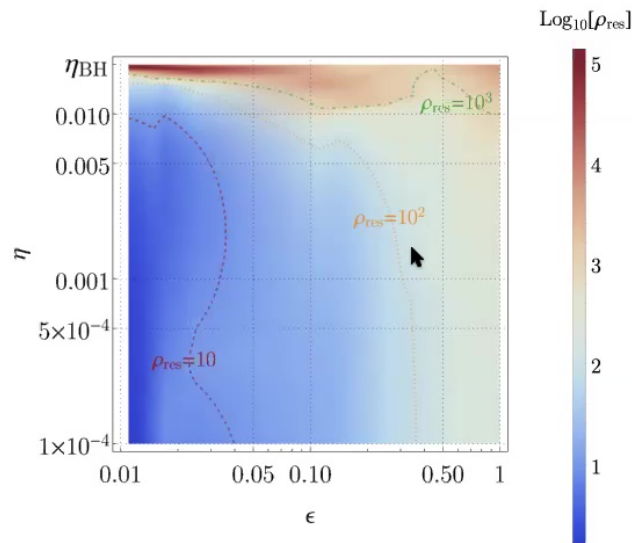
Elisa Maggio - How does a dark compact object ringdown?

Strong Gravity Seminar @Perimeter

21/23



# Resolvability of the doublet



EM, Buoninfante, Mazumdar, Pani, PRD **102**, 064053 (2020)

We use the Rayleigh resolvability criterion:

$$\begin{aligned} \max[\sigma_{f_1}, \sigma_{f_2}] &< |f_1 - f_2| \\ \max[\sigma_{Q_1}, \sigma_{Q_2}] &< |Q_1 - Q_2| \end{aligned}$$

where the uncertainties are computed with a Fisher analysis.

The resolution of the doublet is more challenging than the detection of the deviations from the BH quasi-normal mode.



## Conclusions and future prospects

- We can understand the nature of compact objects and look for new physics at the horizon scale through **gravitational waves**.
- **Dark compact objects** are not excluded by current GW measurements.
- We derived an analytical gravitational-wave template for the **ringdown** and the **echo signal**.
- Accurate echo templates are crucial for matched-filter searches.
- Full inspiral-merger-ringdown waveforms need to be developed.
- **LISA** will allow to perform unprecedented tests of the BH paradigm.

