

Title: Special Topics in Astrophysics - Numerical Hydrodynamics - Lecture 2

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Collection: Special Topics in Astrophysics - Numerical Hydrodynamics

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Chap 1_PDEs Lecture 2

1.2 First-order PDEs

Def (classification): $(x \in \mathbb{R}^n, \Omega_1, \Omega_2 \subseteq \mathbb{R}^n)$

$$L: \mathcal{F}_1(\Omega_1) \rightarrow \mathcal{F}_2(\Omega_2)$$

(i) non-linear:

$$Lu = L(x, u, \{u_{x_k}\})$$

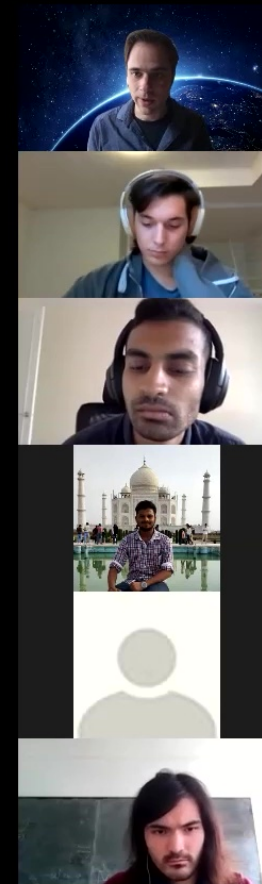
Example: $u_{x_1} - u_{x_2}^2 = 0$

$u_{x_2} \equiv \partial_{x_2} u = \frac{\partial u}{\partial x_2}$

(ii) quasi-linear:

$$Lu = \sum_{k=1}^n a_k(x, u) u_{x_k} + b(x, u) = 0$$

Example: $u_{x_1} + u u_{x_2} = 0$ inviscid
 $\partial_t u + u \partial_x u = 0$ Burger's equation



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$$Lu = L(x_i, u, \{u_{x_k}\}) \quad u_{x_i} \equiv \partial_{x_i} u$$

Example: $u_{x_1} - u_{x_2}^2 = 0 \quad = \frac{\partial u}{\partial x_2}$

(ii) quasi-linear:

$$Lu = \sum_{k=1}^n a_k(x_i, u) u_{x_k} + b(x_i, u) = 0$$

Example: $u_{x_1} + u u_{x_2} = 0$ inviscid
 $\partial_t u + u \partial_x u = 0$ Burger's equation

(iii) semi-linear:

$$Lu = \sum_{k=1}^n a_k(x) u_{x_k} + b(x, u) = 0$$

Example: $u_{x_1} = u$

(iv) linear:

$$Lu = \sum_{k=1}^n a_k(x) u_{x_k} + b(x) = 0$$

Example: $u_{x_1} + a u_{x_2} = 0$ linear advection



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equation

Def: let $U = (u_1, \dots, u_m) : I \subseteq \mathbb{R} \times \Omega \subseteq \mathbb{R}^n \rightarrow \mathbb{R}^m$
 $x_i, \zeta \in \Omega$, $A_i(t, x, U)$ $m \times m$ matrices $\forall i \in \{1, \dots, n\}$
 and

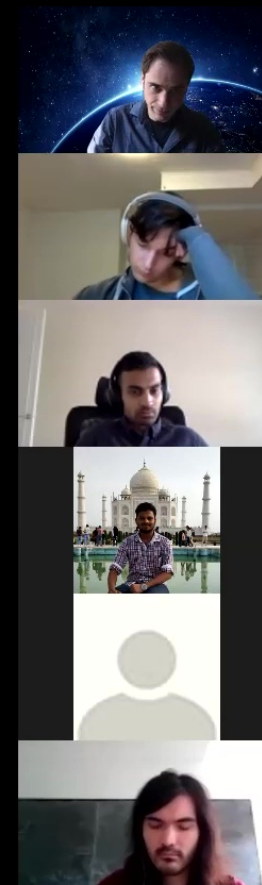
$$\mathcal{A}(t, x, U; \zeta) \equiv \sum_{i=1}^n A_i(t, x, U) \zeta_i.$$

The quasi-linear system

$$(*) \quad U_t + \sum_{i=1}^n A_i(t, x, U) U_{x_i} + B(t, x, U) = 0$$

is called hyperbolic if $\mathcal{A}(t, x, U; \zeta)$ is diagonalizable for $x, \zeta \in \Omega$, $t \in I$, $u \in \mathbb{R}^m$.
 In particular, $(*)$ is hyperbolic

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$\forall x_i \in \Omega, t \in I$ A has real eigenvalues

$$\lambda_1(t, x_i, u; \beta) \leq \dots \leq \lambda_m(t, x_i, u; \beta)$$

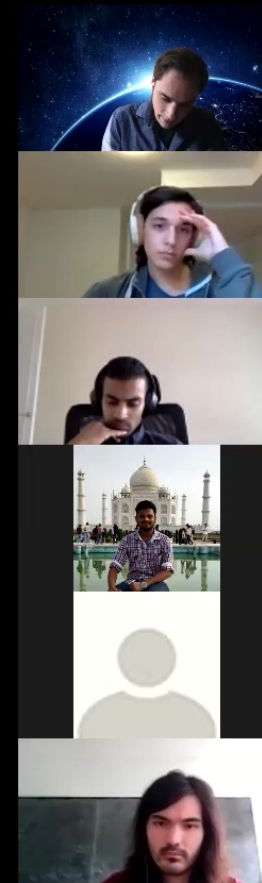
and corresponding eigenvectors

$$\{r_i(x_i, t, u; \beta)\}_{i=1, \dots, m}$$

which form a basis of \mathbb{R}^m . Two special cases:

- 1) $A_i(t, x_i, u)$ symmetric $\forall i \in \{1, \dots, n\}$
 - $\Rightarrow A$ symmetric
 - $\Rightarrow A$ diagonalizable

"(*)

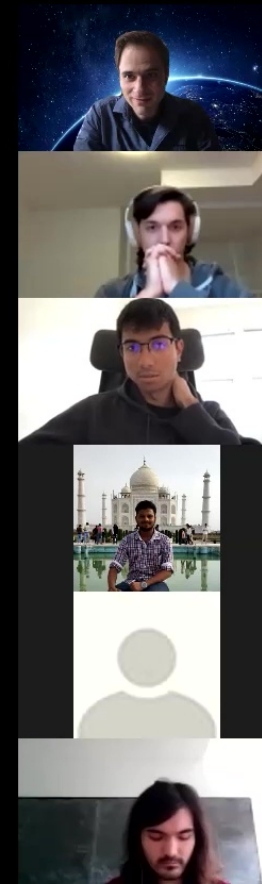


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$\hookrightarrow A(t, x, u; \zeta)$ has real distinct
eigenvalues
 $\lambda_1(t, x, u; \zeta) < \lambda_2(\dots) < \dots < \lambda_m(\dots)$
 $\forall x, \zeta \in \Omega, t \in I$

$\Rightarrow A$ diagonalizable
“(*) strictly hyperbolic”

Examples:



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$$\lambda_1(t, x, \xi) < \lambda_2(\dots) < \dots < \lambda_n(\dots)$$

$$\forall x, \xi \in \Omega, t \in I$$

\Rightarrow A diagonalizable
 “ $(*)$ strictly hyperbolic”

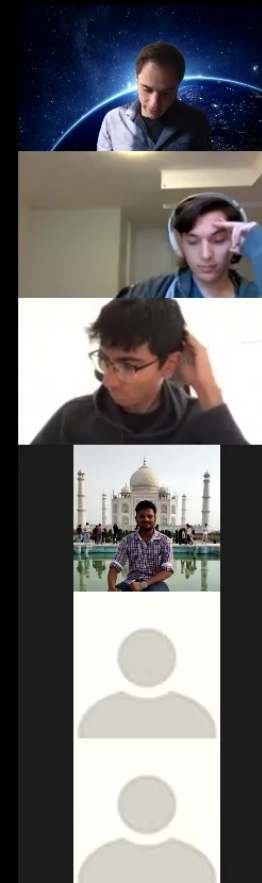
Examples:

$$\begin{pmatrix} u_1 \\ u_2 \end{pmatrix}_t + \begin{pmatrix} 1 & 2 \\ 3 & 4 \end{pmatrix} \begin{pmatrix} u_1 \\ u_2 \end{pmatrix}_x = \begin{pmatrix} 0 \\ 0 \end{pmatrix} \begin{matrix} \text{strictly} \\ \text{hyperbolic} \end{matrix}$$

$$\begin{pmatrix} u_1 \\ u_2 \end{pmatrix}_t + \begin{pmatrix} 1 & 2 \\ 2 & 3 \end{pmatrix} \begin{pmatrix} u_1 \\ u_2 \end{pmatrix}_x = \begin{pmatrix} b_1 \\ b_2 \end{pmatrix} \begin{matrix} \text{symmetric} \\ \text{hyperbolic} \end{matrix}$$

Remarks: 1) The system $(*)$ is elliptic if none of the eigenvalues is real.

2) “hyperbolic”: 2nd order hyperbolic PDEs can be recast as systems of 1st order hyperbolic PDEs



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can be reduced to systems of hyperbolic PDEs.

hyperbolic PDEs.

Def (Conservation laws): Conservation laws are systems of 1st order PDEs that can be written as

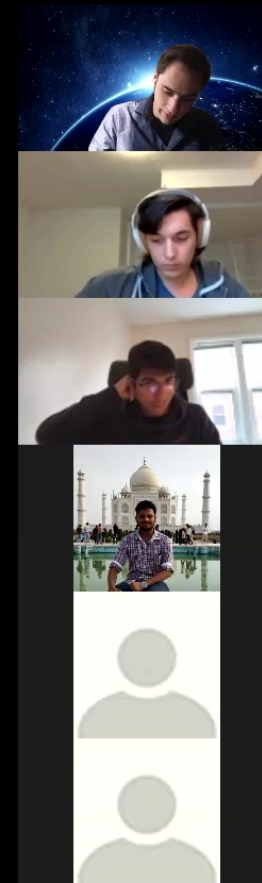
$$u_t + \sum_{i=1}^n \underbrace{F^i(u)}_{x_i} = S(t, x, u)$$

$$= \partial_{x_i} F^i(u)$$

$u = (u_1, \dots, u_m)$ "conserved variables"

$F^i(u) = (f_1^i(u), \dots, f_m^i(u))$, $i = 1, \dots, n$
"vector of fluxes"

$S(t, x, u)$: "sources terms"



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"vector of fluxes"

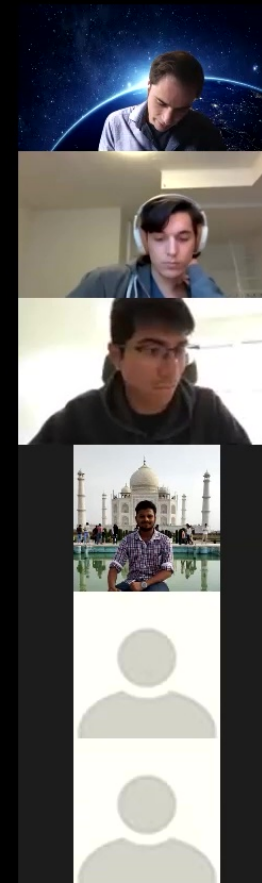
$S(t_i, x_i, U)$: "sources terms"

The Jacobians of the flux functions $F^i(U)$ is

$$A_i(U) = \frac{\partial F^i}{\partial U} = \begin{pmatrix} \frac{\partial f_1^i}{\partial u_1} & \dots & \frac{\partial f_n^i}{\partial u_n} \\ \vdots & & \vdots \\ \frac{\partial f_m^i}{\partial u_1} & \dots & \frac{\partial f_m^i}{\partial u_n} \end{pmatrix}$$

$S \neq 0$: often called "balance law"

Remarks: 1) conservation laws (CL) can be as system of quasi-linear PDEs of 1st order by applying the chain rule:

$$\partial_{x_i} F$$


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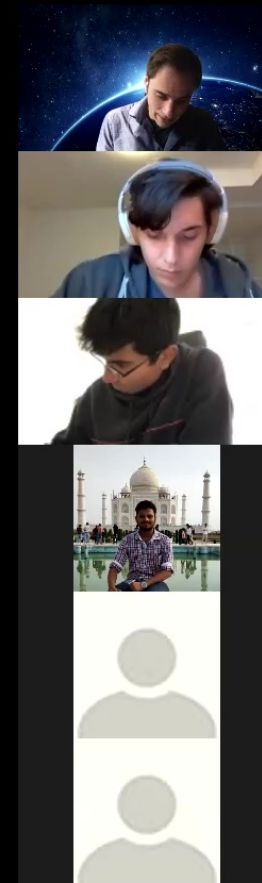
Remarks: 1) conservation laws (CL) can be as system of quasi-linear PDEs of 1st order by applying the chain rule:

$$\partial_{x_i} F^i(U) = \frac{\partial F^i}{\partial U} \frac{\partial U}{\partial x_i} = A_i(U) U_{x_i}$$

2) S : arise due to

- physical sources or sinks of otherwise conserved quantities
- due to changes of coordinates
"geometric source terms"

3)



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3) Motivation: consider integral form

$$\frac{d}{dt} \int_V U \, dV = - \int_{\partial V} F^i(U) \, dV \quad V: \text{control volume}$$

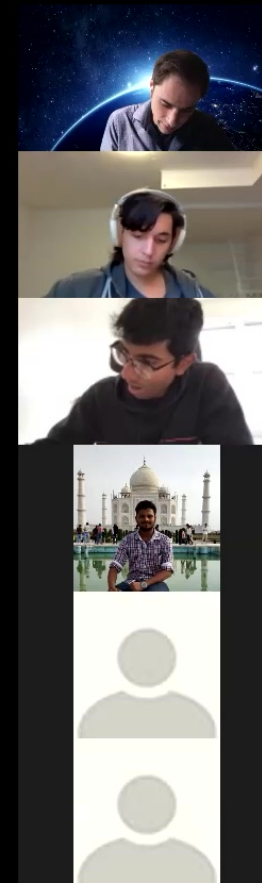
divergence theorem = $-\int_{\partial V} F^i(U) \, n_i \, dS$

↑ unit normal pointing outward

"time-rate change of U inside V depends only on the total flux through the surface S "

no " U is conserved"

Example



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no "u is conserved"

Examples:

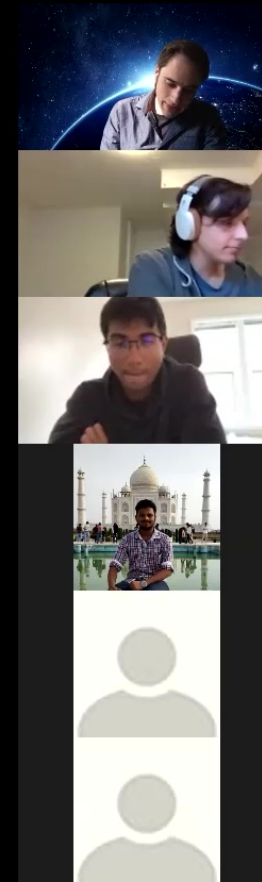
$$\frac{\partial u}{\partial t} + \frac{\partial f(u)}{\partial x} = 0, \quad f(u) = au$$

no linear advection equation

$$f(u) = \frac{1}{2}u^2 \Rightarrow \frac{\partial f}{\partial x} = uu_x$$

no inviscid Burger's equation

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1.5 some properties of hyperbolic systems

1.3.1 characteristics

1) linear hyperbolic systems

Consider hyperbolic system of the form

$$U_t + AU_x = 0 \quad n=1$$

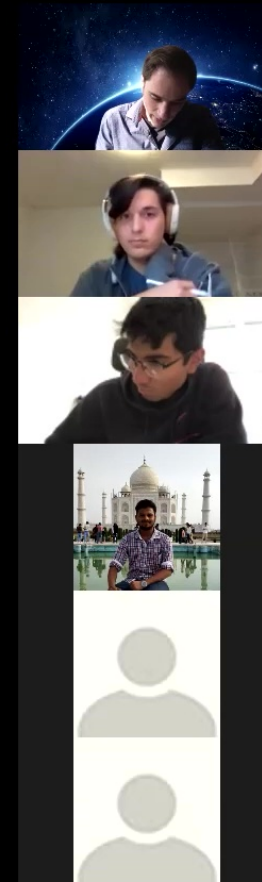
↑ const. coefficients

hyperbolicity $\Rightarrow \exists Q$ with $A = Q\Lambda Q^{-1}$

$$\Lambda = \begin{pmatrix} \lambda_1 & & 0 \\ & \dots & \\ 0 & & \lambda_m \end{pmatrix}$$

↑ eigenvalues of A

Q



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1.5 some properties of hyperbolic systems

1.3.1 Characteristics

1) linear hyperbolic systems

Consider hyperbolic system of the form

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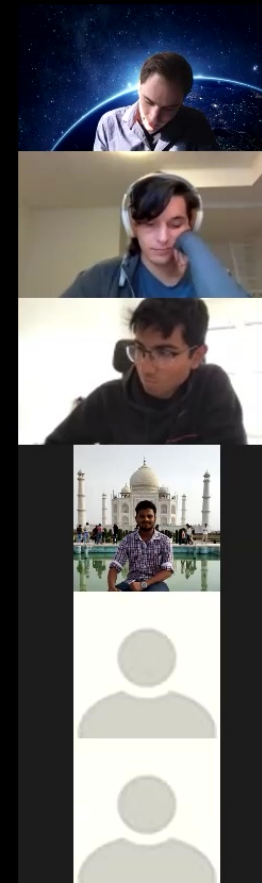
$$\Lambda = \begin{pmatrix} \lambda_1 & & 0 \\ & \dots & \\ 0 & & \lambda_m \end{pmatrix}$$

↑ eigenvalues of A

$$Q = (r_1, \dots, r_m)$$

↑ columns

$\Rightarrow u_t + Q\Lambda Q^{-1}u_x = 0$ |



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hyperbolicity $\Rightarrow \exists Q$ with $A = Q\Delta Q^{-1}$

$$\Delta = \begin{pmatrix} \lambda_1 & & 0 \\ & \dots & \\ 0 & & \lambda_m \end{pmatrix}$$

λ_i eigenvalues of A

$$Q = (r_1, \dots, r_m)$$

columns

$$\Rightarrow u_t + Q\Delta Q^{-1}u_x = 0 \quad | \times Q^{-1}$$

define $V = Q^{-1}u$ "characteristic variables"

$$\Rightarrow \boxed{V_t + \Delta V_x = 0}$$

deco



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define $V = Q^{-1}U$ "characteristic variables"

$\Rightarrow V_t + \Delta V_x = 0$ decoupled system

characteristics: curves $\gamma(t)$ along which the PDE becomes an ODE

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$\Rightarrow V_t + \Delta V_x = 0$ variables⁴
decoupled system

Characteristics: curves $\gamma(t)$ along which the PDE becomes an ODE

Consider $\gamma(t) = x$, $v_i = v_i(\underbrace{\gamma(t)}_x, t)$

$\Rightarrow \frac{dv_i}{dt} = \frac{\partial v_i}{\partial t} + \frac{dx}{dt} \frac{\partial v_i}{\partial x} = \frac{\partial v_i}{\partial t} + \gamma'(t) \frac{\partial v_i}{\partial x}$

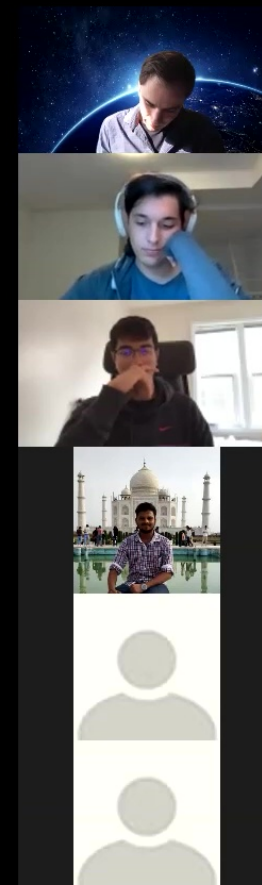
choose $\gamma'(t) = \lambda_i$ "characteristic speeds"

$\Rightarrow \frac{dv_i}{dt} = 0 \Rightarrow v_i$ are constant along characteristics

Therefore: given $v_i(x, 0) = v_i^0(x)$ at $t=0$

$v_i(x, t) = v_i^0(x_0) = v_i^0(x - \lambda_i t)$

where $x(t) = x_0 + \lambda_i t$ is the characteristic



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$x(t)$ that passes through x at t .

no v_x at (x,t) is entirely determined by initial data v_x^0 at x_0

no "advection equation"

characteristic curve $x = x_0 + h_x t$

$v_x(x,t)$

$v_x^0(x)$

x_0

x

t

v_x

→ advection of v_x^0 with speed h_x

General initial value problem:

Given $u_t + Au_x = 0$ (I)

$u^0 = (u_1^0, \dots, u_m^0)$



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→ advection of v_i^0 with speed λ_i

General initial value problem:

Given $U_t + AU_x = 0$ (I)

$U^0 = (u_1^0, \dots, u_m^0)$

↓ introduce characteristic variables $V = Q^{-1}U$

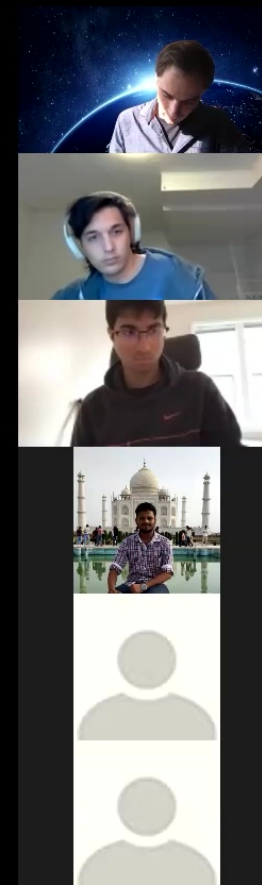
$$\begin{cases} V_t + \Delta V_x = 0 \\ V^0 = Q^{-1}U^0 = (v_1^0, \dots, v_m^0) \end{cases}$$

↓

$V = \{v_i(x,t)\} = \{v_i^0(x - \lambda_i t)\}_{i=1, \dots, m}$

↓ transform back

$U(x,t) = QV(x,t)$



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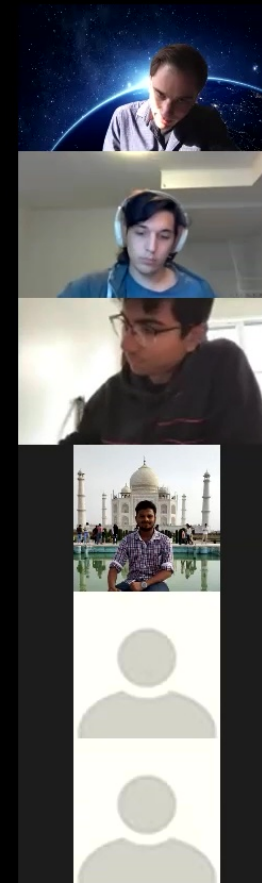
$$V = \{v_i(x,t)\} = \{v_i^0(x - \lambda_i t)\}_{i=1, \dots, m}$$

$$\downarrow \text{transform basis}$$

$$\begin{aligned}
 u(x,t) &= QV(x,t) \\
 &= \sum_{i=1}^m v_i(x,t) r_i \\
 &= \sum_{i=1}^m v_i^0(x - \lambda_i t) r_i
 \end{aligned}$$

solution to (I)

Remarks: • $u(x,t)$ is superposition of eigenvectors
i.e. waves propagating at



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$$= \sum_{i=1}^n v_i^0(x - \lambda_i t) r_i$$

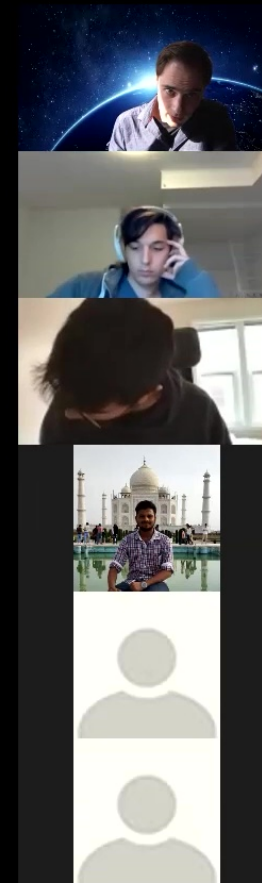
Remarks:

- $U(x,t)$ is superposition of eigenvectors
i.e. waves propagating at finite
speed $\lambda_i \neq 0$
- $U(x,t)$ entirely determined by
initial data v_i^0 at points $x_0^i = x - \lambda_i t$

2) Systems of CLs
considers hyperbolic system

$$U_t + F(U)_x = 0$$

no can locally transform into decoupled
system



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$$u_t + F(u)_x = 0$$

↳ can locally transform into decoupled system

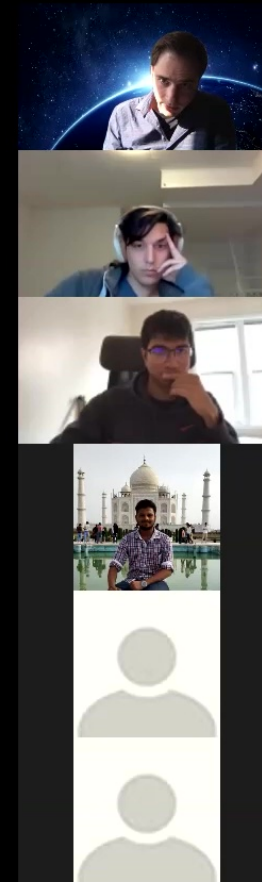
hyperbolicity $\Rightarrow Q = Q(x, t, u), A \equiv \frac{\partial F}{\partial u}$
 $= Q \Delta Q^{-1}$
 \uparrow
 diagonal matrix

locally: $V_t + \Delta V_x = 0$
 meaning of V depends on u

Δ_i characteristics depend on state vector u

↳ the state/solution is "self-propagating"

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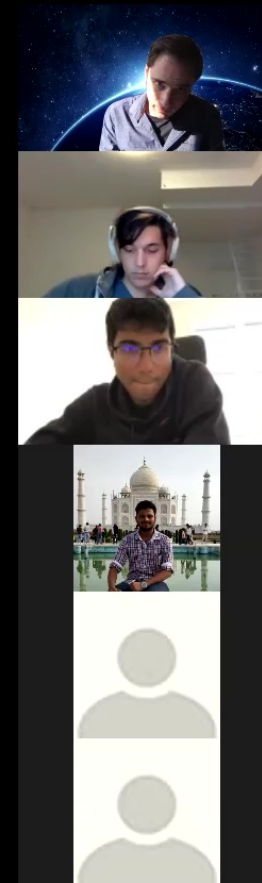
1.3.2 Domain of dependence & range of influence

Def: The domain of dependence $\mathcal{D}(x,t) \subseteq \mathbb{R}^n$

is the domain of the initial data $u^0(x)$ that entirely determines the solution $U(x,t)$ of a hyperbolic system

$$U_t + \sum_{i=1}^n A_i U_{x_i} = 0$$

Theorem: Let U be solution



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that entirely determines the solution $U(x,t)$ of a hyperbolic system

$$U_t + \sum_{i=1}^n A_i U_{x_i} = 0$$

Theorem: let U be solution of symmetric hyperbolic system

$$U_t + \sum_{i=1}^n A_i U_{x_i} = 0$$

\uparrow const. coeff.

Consider (x_0, t_0) , $0 \leq t_1 \leq t_0$ and

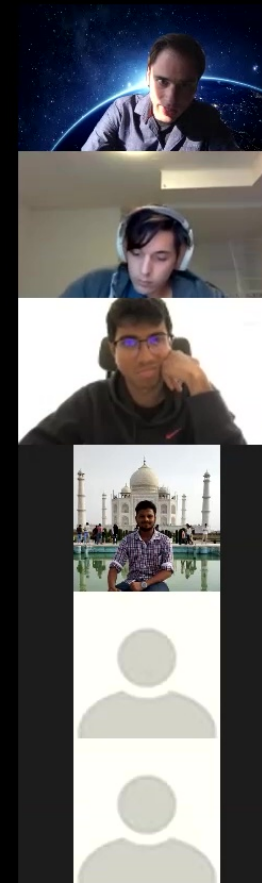
$$B \equiv \{x \in \mathbb{R}^n \mid |x - x_0| \leq M(t_0 - t_1)\}$$

$$C \equiv \{x \in \mathbb{R}^n \mid |x - x_0| \leq M(t_0 - t_1)\}$$

with

$$M = \max \{ \lambda_i \}$$

λ_i are eigenvalues of $A(\xi) = \sum_{i=1}^n A_i \xi_i$.



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$$C \equiv \{x \in \mathbb{R}^n \mid |x - x_0| \leq M(b_0 - t_1)\} \quad M b_0$$

with

$$M = \max \{\lambda_i\}$$

λ_i are eigenvalues of $A(s) = \sum_{i=1}^n A_i s_i$.

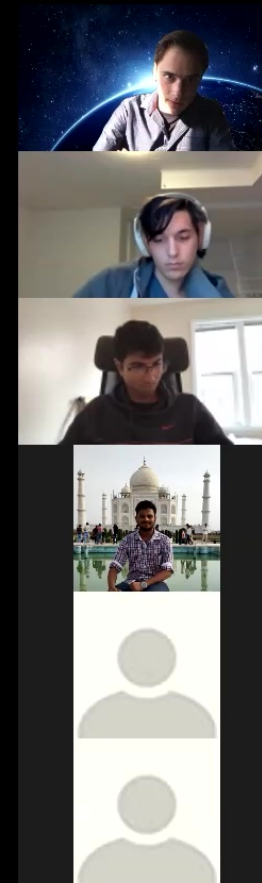
If $|U| \equiv 0$ on B_1 , then $|U| \equiv C$.

Remarks: 1) D is determined by fastest and slowest characteristic speed

2) D is bounded (λ_i are finite)

Boundedness also applies to non-linear systems

"Information propagates at finite speed"



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Remarks: 1) \mathcal{D} is determined by fastest and slowest characteristic speed

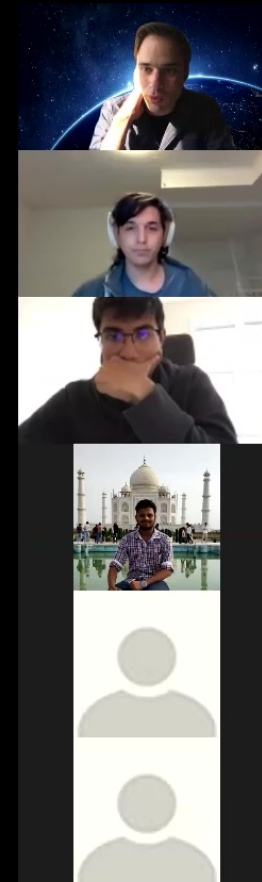
2) \mathcal{D} is bounded (λ_i are finite)

Boundedness also applies to non-linear systems

"Information propagates at finite speed"

Example: 1D: $u_t + \lambda u_x = 0$

$$u(x,t) = \sum_{i=1}^m v_i^0(x - \lambda_i t) r_i$$

$$\rightarrow \mathcal{D}(x,t) = \{x - \lambda_i t \mid i=1, \dots, m\}$$


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Remarks: 1) \mathcal{D} is determined by fastest and slowest characteristic speed

2) \mathcal{D} is bounded (λ_i are finite)

Boundedness also applies to non-linear systems

"Information propagates at finite speed"

Example: 1D: $u_t + \Delta u_x = 0$

$$u(x,t) = \sum_{i=1}^m v_i^0(x - \lambda_i t) r_i$$

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$m=3$

