

Title: Matrix product states and 1d quantum systems

Speakers: Frank Pollmann

Collection: Online School on Ultra Quantum Matter

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Matrix Product States and 1D Quantum Systems

Frank Pollmann (TUM)



$$|\psi\rangle : \dots \begin{array}{ccccccc} & A^{[1]} & A^{[2]} & A^{[3]} & A^{[4]} & A^{[5]} & A^{[6]} & A^{[7]} \\ & \text{---} & \text{---} & \text{---} & \text{---} & \text{---} & \text{---} & \text{---} \\ & \text{Y} & \text{Y} & \text{Y} & \text{Y} & \text{Y} & \text{Y} & \text{Y} \\ & | & | & | & | & | & | & | \\ & & & & & & & & \dots \end{array} \dots$$

Online School on Ultra Quantum Matter

Complexity of a quantum many-body problem

Many-body Hilbert space

$$|\psi\rangle = \sum_{j_1, j_2, \dots, j_L} \psi_{j_1, j_2, \dots, j_L} |j_1\rangle |j_2\rangle \dots |j_L\rangle, \quad j_n = 1 \dots d$$

N spin-1/2 particles: $\text{dim} = 2^L$

10 spins $\text{dim} = 1'024$

20 spins $\text{dim} = 1'048'576$

30 spins $\text{dim} = 1'073'741'824$

40 spins $\text{dim} = 1'099'511'627'776$

- ▶ Full diagonalization up to ~ 20 sites
- ▶ Sparse methods up to ~ 40 sites



Matrix-Product States

Many-body Hilbert space

$$|\psi\rangle = \sum_{j_1, j_2, \dots, j_L} \psi_{j_1, j_2, \dots, j_L} |j_1\rangle |j_2\rangle \dots |j_L\rangle, \quad j_n = 1 \dots d$$

Matrix-Product States: Reduction of #variables $d^L \rightarrow Ld\chi^2$

$$\psi_{j_1, j_2, \dots, j_L} \approx \sum_{\alpha_1, \alpha_2, \dots, \alpha_{L-1}} A_{\alpha_1}^{j_1} A_{\alpha_1, \alpha_2}^{j_2} \dots A_{\alpha_{L-1}}^{j_L} \quad \alpha_j = 1 \dots \chi$$

Once we have an MPS representation, we can calculate (almost) everything exactly!



Matrix-Product States

Many-body Hilbert space

$$|\psi\rangle = \sum_{j_1, j_2, \dots, j_L} \psi_{j_1, j_2, \dots, j_L} |j_1\rangle |j_2\rangle \dots |j_L\rangle, \quad j_n = 1 \dots d$$

Matrix-Product States: Reduction of #variables $d^L \rightarrow Ld\chi^2$

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Once we have an MPS representation, we can calculate (almost) everything exactly!

Why is it a good representation?



Matrix product states and 1d quantum systems

- I) Entanglement and Matrix-Product States
- II) Time Evolving Block Decimation
- III) Density-Matrix Renormalization Group

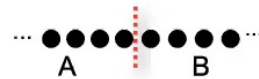


Entanglement

A generic quantum state has a d^L dimensional Hilbert space

$$|\psi\rangle = \sum_{j_1, j_2, \dots, j_L} \psi_{j_1, j_2, \dots, j_L} |j_1\rangle |j_2\rangle \dots |j_L\rangle, \quad j_n = 1 \dots d$$

Decompose a state into a superposition of product states (**Schmidt decomposition**)



$$|\psi\rangle = \sum_{i,j} C_{i,j} |i\rangle_A \otimes |j\rangle_B = \sum_{\alpha} \Lambda_{\alpha} |\alpha\rangle_A \otimes |\alpha\rangle_B, \quad \langle \alpha | \alpha' \rangle = \delta_{\alpha \alpha'}$$

Entanglement entropy as a measure for the

amount of entanglement $S = - \sum_{\alpha} \Lambda_{\alpha}^2 \log \Lambda_{\alpha}^2$

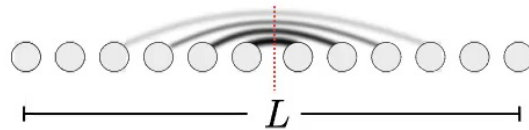
(Equivalent to $S = -\text{Tr} \rho_A \log \rho_A$ with $\rho_A = \text{Tr}_B |\psi\rangle \langle \psi|$)



Entanglement

Area law for ground states of local (gapped) Hamiltonians
in one dimensional systems

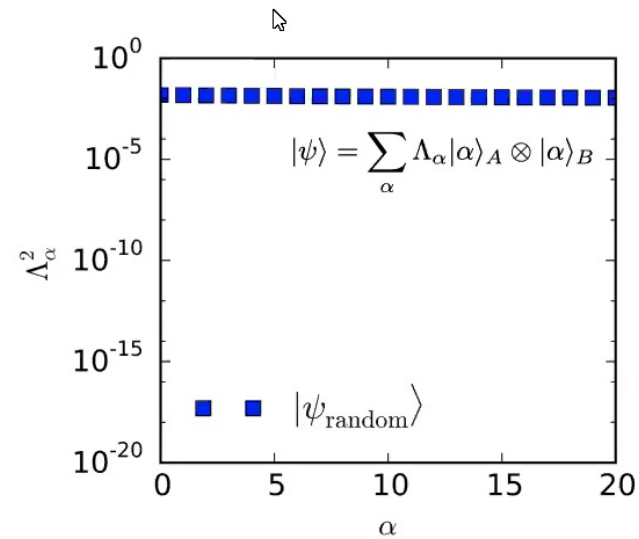
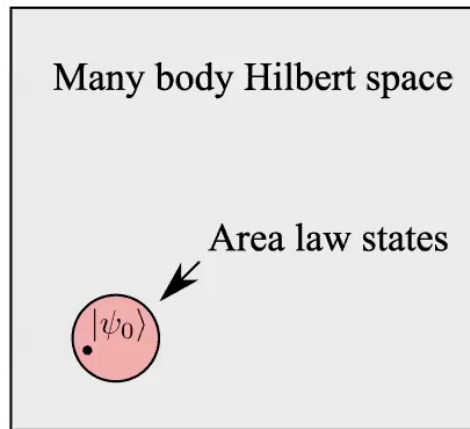
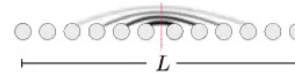
$$S(L) = \text{const.} \quad [\text{Srednicki '93, Hastings '07}]$$



Entanglement

Area law for ground states of local (gapped) Hamiltonians in one dimensional systems

$$S(L) = \text{const.} \quad [\text{Srednicki '93, Hastings '07}]$$



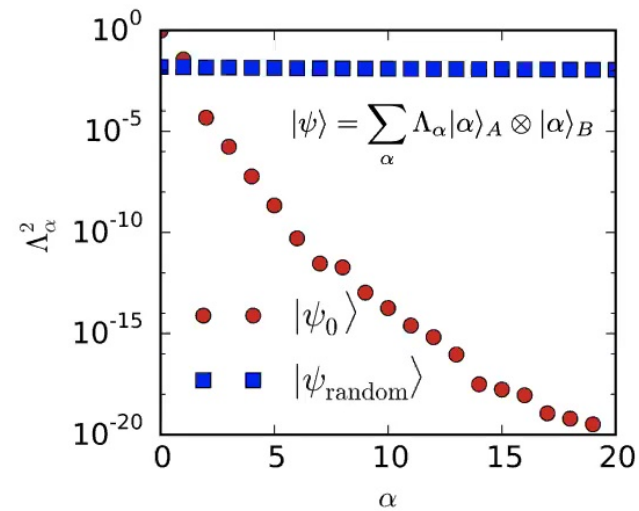
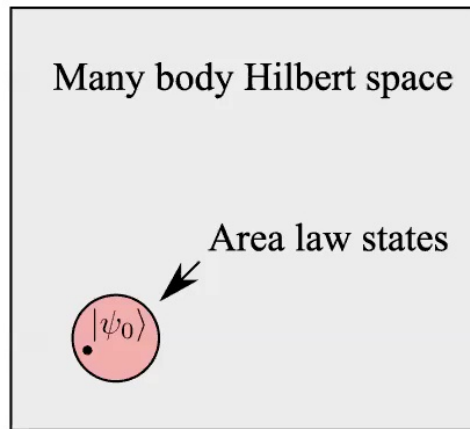
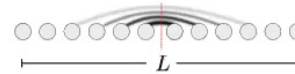
All ground states live in a tiny corner of the Hilbert space!



Entanglement

Area law for ground states of local (gapped) Hamiltonians in one dimensional systems

$$S(L) = \text{const.} \quad [\text{Srednicki '93, Hastings '07}]$$



All ground states live in a tiny corner of the Hilbert space!



SVD Compression

$$\text{Example: } |\psi\rangle = \sum_{i=1}^{N_A} \sum_{j=1}^{N_B} C_{ij} |i\rangle_A |j\rangle_B = \sum_{\gamma} \lambda_{\gamma} |\phi_{\gamma}\rangle_A |\phi_{\gamma}\rangle_B$$

Matrix can represent an image (array of pixel)

$$C = \begin{pmatrix} 0.23 & \dots & 0.56 \\ \vdots & \ddots & \vdots \\ 0.22 & \dots & 0.34 \end{pmatrix} = \left(\text{Image of Golden Gate Bridge} \right)$$

$\chi = 1200$

Reconstruction of the matrix (image) from a small number of Schmidt states (SVD):



Compression of quantum states



Compression of quantum states



Compression of quantum states



Compression of quantum states



Compression of quantum states



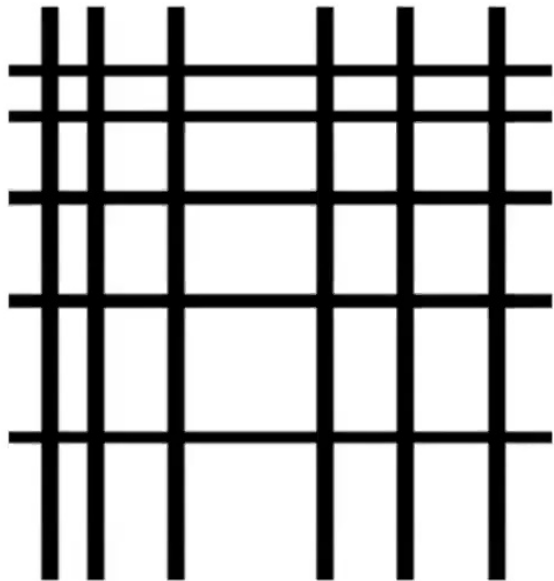
Important features visible already for < 16 states!



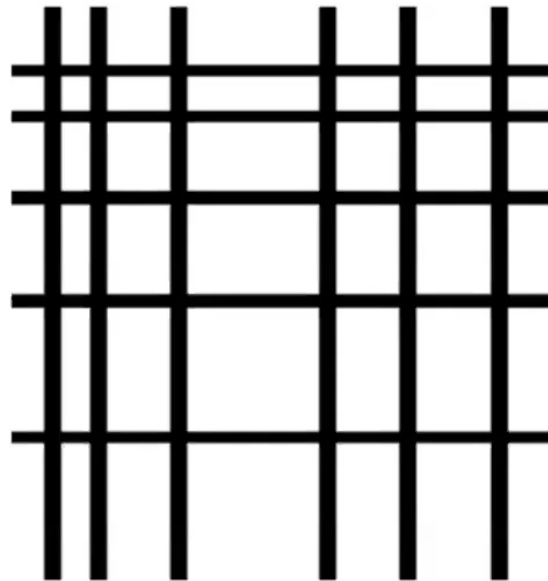
Compression of quantum states



Compression of quantum states



$$\chi = 500$$



$$\chi = 1$$

[Mondrian]



Matrix-Product States

Coefficients in the many-body wave function:

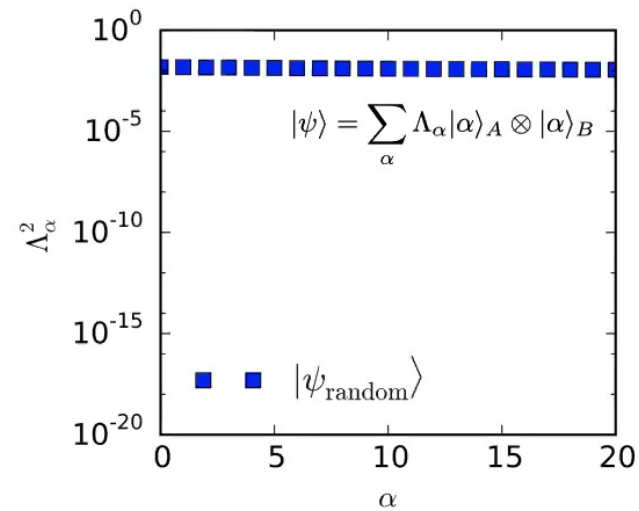
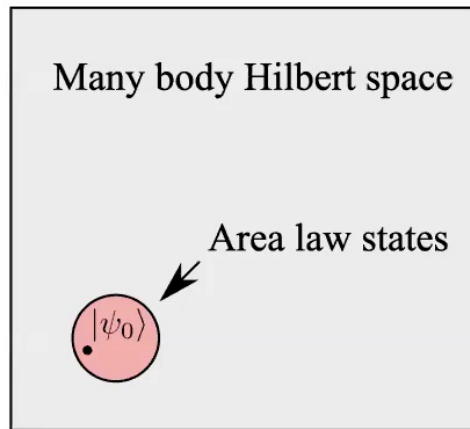
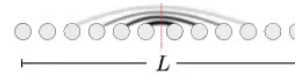
Rank- L tensor: diagrammatic representation

$$\psi_{j_1, j_2, j_3, j_4, j_5} = \begin{array}{c} \psi \\ \hline | \quad | \quad | \quad | \quad | \end{array}$$

Entanglement

Area law for ground states of local (gapped) Hamiltonians in one dimensional systems

$$S(L) = \text{const.} \quad [\text{Srednicki '93, Hastings '07}]$$



All ground states live in a tiny corner of the Hilbert space!



SVD Compression

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Matrix can represent an image (array of pixel)

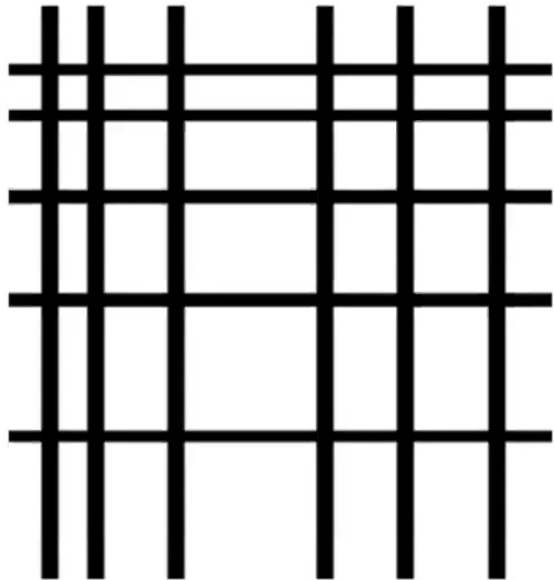
$$C = \begin{pmatrix} 0.23 & \dots & 0.56 \\ \vdots & \ddots & \vdots \\ 0.22 & \dots & 0.34 \end{pmatrix} = \left(\begin{array}{c} \text{Image of Golden Gate Bridge} \\ \chi = 1200 \end{array} \right)$$



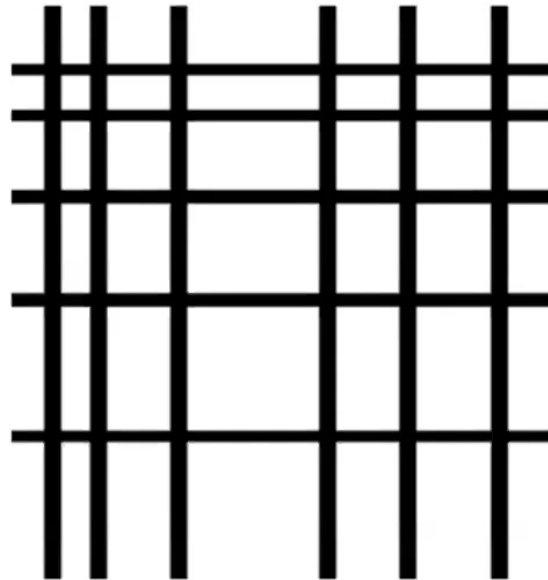
Compression of quantum states



Compression of quantum states



$$\chi = 500$$



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[Mondrian]



Matrix-Product States

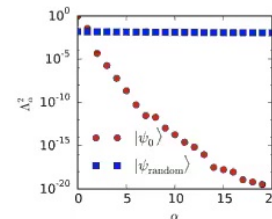
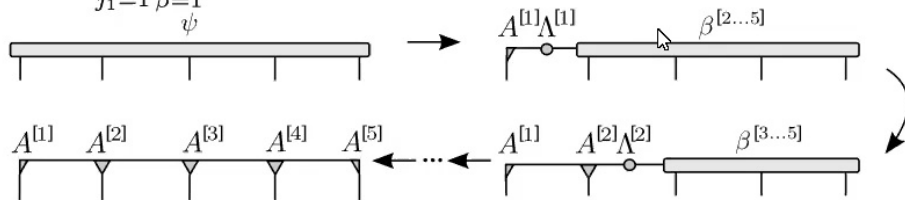
Coefficients in the many-body wave function:

Rank- L tensor: diagrammatic representation

$$\psi_{j_1, j_2, j_3, j_4, j_5} = \text{Diagram of a rank-5 tensor } \psi$$

Successive Schmidt decompositions

$$|\psi\rangle = \sum_{j_1=1}^d \sum_{\beta=1}^d A_{\beta}^{[1]j_1} \Lambda_{\beta}^{[1]} |j_1\rangle |\beta\rangle_{[2, \dots, N]}$$



$$\psi_{j_1, j_2, \dots, j_L} \approx \sum_{\alpha_1, \alpha_2, \dots, \alpha_{L-1}} A_{\alpha_1}^{j_1} A_{\alpha_1, \alpha_2}^{j_2} \dots A_{\alpha_{L-1}}^{j_L}$$





Matrix-Product States

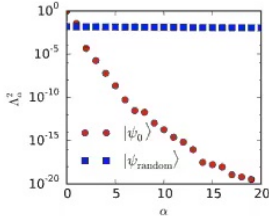
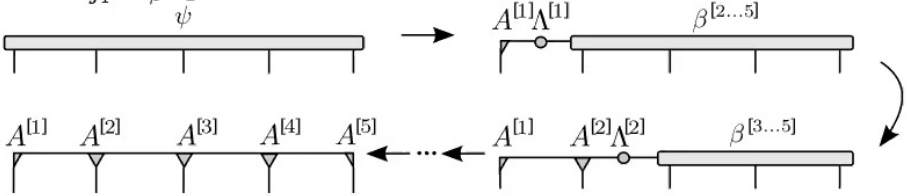
Coefficients in the many-body wave function:

Rank- L tensor: diagrammatic representation

$$\psi_{j_1, j_2, j_3, j_4, j_5} = \text{Diagram of a horizontal bar with 5 legs pointing downwards, labeled } \psi$$

Successive Schmidt decompositions

$$|\psi\rangle = \sum_{j_1=1}^d \sum_{\beta=1}^d A_{\beta}^{[1]j_1} \Lambda_{\beta}^{[1]} |j_1\rangle |\beta\rangle_{[2, \dots, N]}$$



$$\psi_{j_1, j_2, \dots, j_L} \approx \sum_{\alpha_1, \alpha_2, \dots, \alpha_{L-1}} A_{\alpha_1}^{j_1} A_{\alpha_1, \alpha_2}^{j_2} \dots A_{\alpha_{L-1}}^{j_L}$$

MPS are tailored to describe 1D systems with an area law!

MPS and the canonical form

From now on: Leave out site indices!

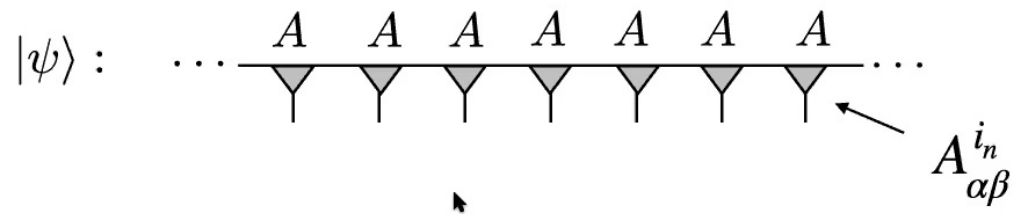
$$|\psi\rangle : \quad \dots \begin{array}{ccccccc} A^{[1]} & A^{[2]} & A^{[3]} & A^{[4]} & A^{[5]} & A^{[6]} & A^{[7]} \\ \text{---} & \text{---} & \text{---} & \text{---} & \text{---} & \text{---} & \text{---} \\ \text{Y} & \text{Y} & \text{Y} & \text{Y} & \text{Y} & \text{Y} & \text{Y} \\ \text{---} & \text{---} & \text{---} & \text{---} & \text{---} & \text{---} & \text{---} \\ \dots & \dots & \dots & \dots & \dots & \dots & \dots \end{array} \dots$$

$A_{\alpha\beta}^{i_n}$



MPS and the canonical form

From now on: Leave out site indices!



MPS is not unique

$$\tilde{A}^{i_n} = X A^{i_n} X^{-1}$$

➔ \tilde{A}^{i_n} describes the same state!



MPS and the canonical form

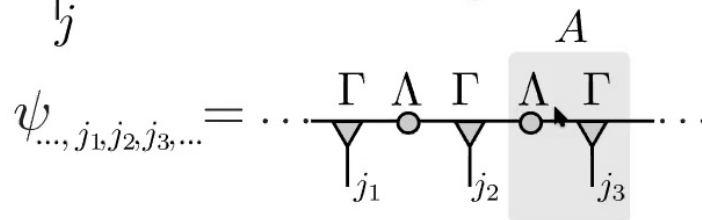
Choose a convenient representation in **Canonical Form**:
 Bond index corresponds to Schmidt decomposition! [Vidal '03]

$$|\psi\rangle = \sum_{\alpha=1}^{\chi} \Lambda_{\alpha} |\alpha\rangle_L \otimes |\alpha\rangle_R \quad \text{with} \quad \langle \alpha | \alpha' \rangle = \delta_{\alpha\alpha'}$$

Write tensor $A_{\alpha\beta}^{i_n}$ as product of

$\Lambda_{\alpha\beta} = \alpha \text{---} \circ \text{---} \beta$: Diagonal matrix with Schmidt values

$\Gamma_{\alpha\beta}^j = \alpha \text{---} \underset{j}{\nabla} \text{---} \beta$: Tensor relating to Schmidt basis



MPS and the canonical form

Schmidt states in terms of the MPS:

$$|\alpha\rangle_L = \dots \text{---} \Lambda \text{---} \Gamma \text{---} \Lambda \text{---} \Gamma \text{---} \alpha$$

$$|\alpha\rangle_R = \alpha \text{---} \Gamma \text{---} \Lambda \text{---} \Gamma \text{---} \Lambda \text{---} \dots$$

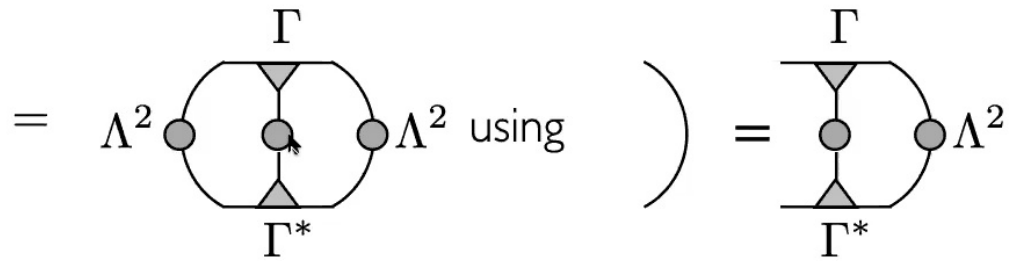
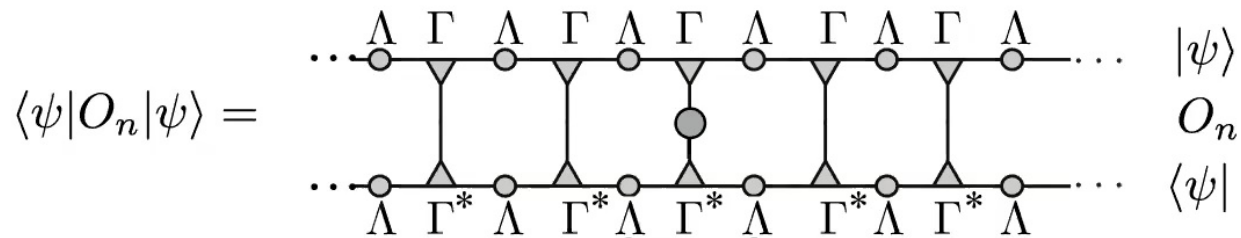
Orthogonality:

$$\langle \alpha' | \alpha \rangle_R = \delta_{\alpha' \alpha} \equiv$$



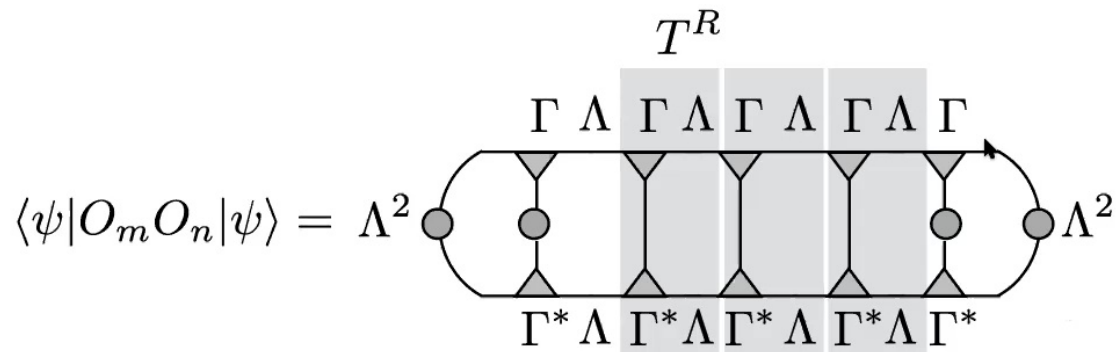
MPS and the canonical form

Efficient evaluation of **expectation values**:



MPS and the canonical form

Efficient evaluation of **correlation functions**:



Matrix Product States and 1D Quantum Systems

- I) Entanglement and Matrix-Product States
- II) Time Evolving Block Decimation
- III) Density-Matrix Renormalization Group



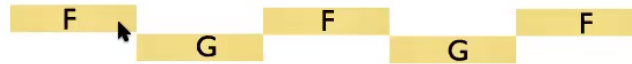
Time evolving block decimation

Consider a Hamiltonian $H = \sum_j h^{[j,j+1]}$ [Vidal '03]

Decompose the Hamiltonian as $H=F+G$

$$F \equiv \sum_{\text{even } j} F^{[j]} \equiv \sum_{\text{even } j} h^{[j,j+1]}$$

$$G \equiv \sum_{\text{odd } j} G^{[j]} \equiv \sum_{\text{odd } j} h^{[j,j+1]}$$



We observe $[F^{[r]}, F^{[r']}] = 0$ ($[G^{[r]}, G^{[r']}] = 0$)
but $[G, F] \neq 0$



Time evolving block decimation

Apply Suzuki-Trotter decomposition of order p

$$\exp(-i(F + G)\delta t) \approx f_p[\exp(-F\delta t), \exp(-G\delta t)]$$

with $f_1(x, y) = xy$, $f_2(x, y) = x^{1/2}yx^{1/2}$, etc.

Two chains of two-site gates

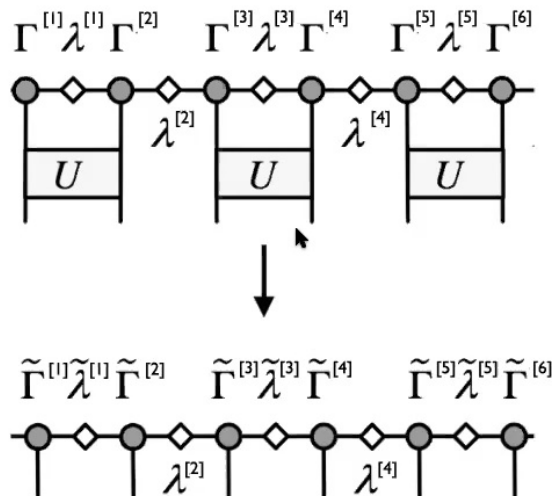
$$U_F = \prod_{\text{even } r} \exp(-iF^{[r]}\delta t)$$

$$U_G = \prod_{\text{odd } r} \exp(-iG^{[r]}\delta t)$$



Time evolving block decimation

Time Evolving Block Decimation algorithm (TEBD)

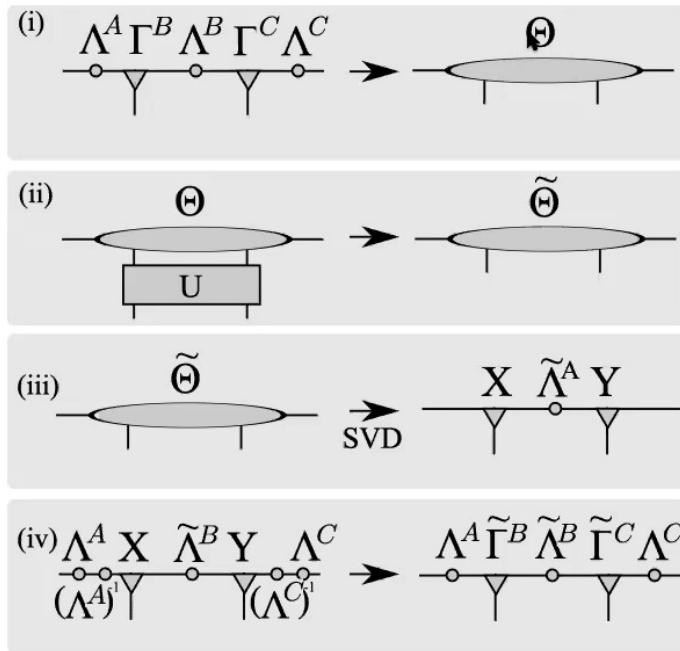


How do we get the original form back?



Time evolving block decimation

Time Evolving Block Decimation (TEBD) algorithm [Vidal '03]

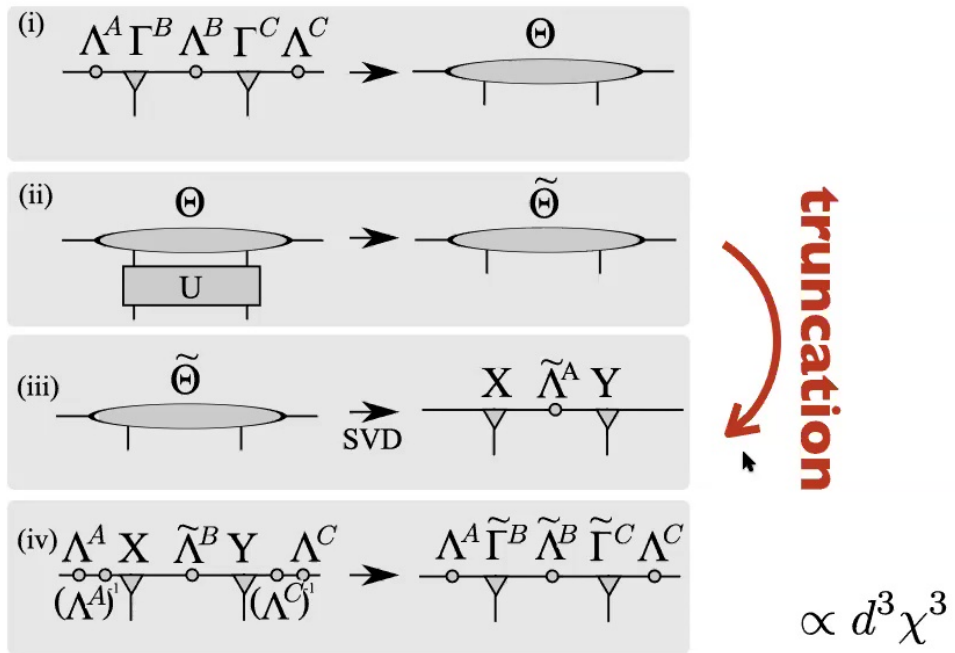


$$\propto d^3 \chi^3$$



Time evolving block decimation

Time Evolving Block Decimation (TEBD) algorithm [Vidal '03]

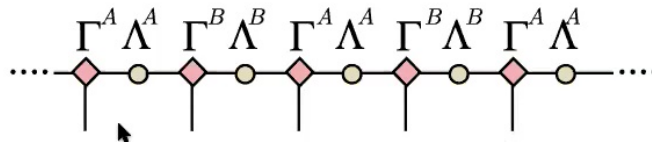


Translationally invariant, infinite systems!

Assume that $|\psi\rangle$ is translational invariant and $L = \infty$:
infinite TEBD (**iTEBD**)

Partially break translational symmetry to simulate
the action of the gates

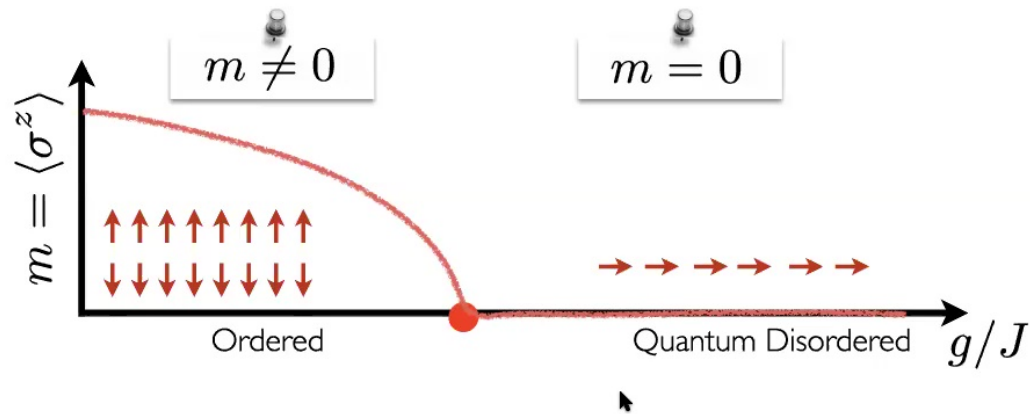
$$\Gamma^{[2r]} = \Gamma^A, \lambda^{[2r]} = \lambda^A, \Gamma^{[2r+1]} = \Gamma^B, \lambda^{[2r+1]} = \lambda^B$$



Quantum Ising Model

Quantum phases at $T = 0$: Transverse field Ising model
with \mathbb{Z}_2 symmetry [Elliott et al. '70]

$$H = - \sum_j (J \sigma_j^z \sigma_{j+1}^z + g \sigma_j^x) \quad \uparrow \uparrow \uparrow \downarrow \downarrow \downarrow \uparrow \uparrow$$



Matrix Product States and 1D Quantum Systems

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Density matrix renormalization group (DMRG)

Matrix-Product operators: Generalization of MPS to the space of linear operators

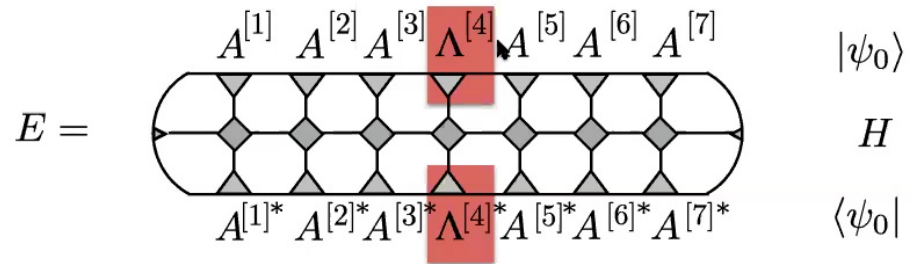
$$\mathcal{O}_{i_1 i_2 \dots i_L, i'_1, i'_2 \dots i'_L} =$$

The diagram shows three layers of tensors. The top layer has four diamond-shaped tensors labeled 'M' connected horizontally. The middle layer has four diamond-shaped tensors labeled 'M' connected horizontally, with a mouse cursor pointing to the first one. Above each 'M' tensor in the middle layer is a triangle pointing down, and below each is a triangle pointing up. The bottom layer has four diamond-shaped tensors labeled 'M' connected horizontally, with a triangle pointing down above each and a triangle pointing up below each. The top layer is connected to the middle layer by vertical lines, and the middle layer is connected to the bottom layer by vertical lines. Labels Γ and Λ are placed above and below the top and bottom layers respectively, with asterisks on the bottom layer labels.



Density matrix renormalization group (DMRG)

Find the **ground state** iteratively



by locally minimizing energy of $H_{\alpha i \beta; \alpha' i' \beta'}$ (e.g., Lanczos)



Density matrix renormalization group (DMRG)

Find the **ground state** iteratively

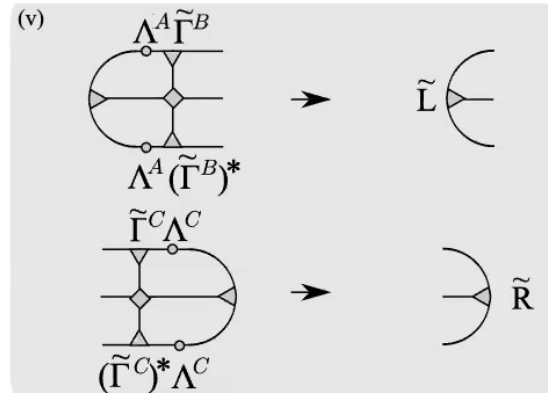
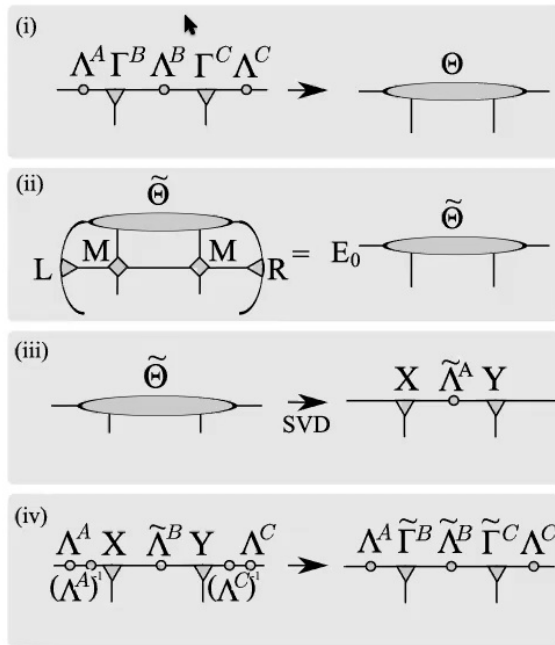
$$H_{\alpha i \beta; \alpha' i' \beta'} = \begin{array}{c} \begin{array}{ccccccc} A^{[1]} & A^{[2]} & A^{[3]} & & A^{[5]} & A^{[6]} & A^{[7]} \\ \alpha & & & \beta & & & \\ \alpha' & & & i & & & \\ \beta' & & & i' & & & \\ A^{[1]*} & A^{[2]*} & A^{[3]*} & & A^{[5]*} & A^{[6]*} & A^{[7]*} \end{array} \\ \langle \psi_0 | \\ H \\ | \psi_0 \rangle \end{array}$$

by locally minimizing energy of $H_{\alpha i \beta; \alpha' i' \beta'}$ (e.g., Lanczos)

Much faster convergence than TEBD + allows for long range interactions!

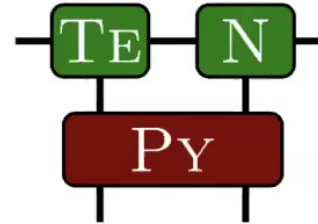


Density matrix renormalization group (DMRG)



Tensor Network Python (TeNPy)

What is TeNPy



- ▶ Python 3 library for simulations with tensor network
<https://github.com/tenpy/tenpy>
- ▶ Object oriented, modular structure, and easy to install
- ▶ HTML documentation
<https://tenpy.github.io>
- ▶ (in)finite DMRG, TEBD, TDVP, ...

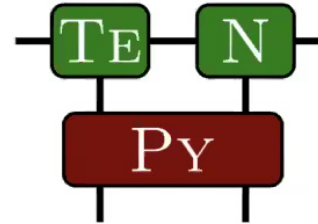


Johannes Hauschild, Berkeley



Tensor Network Python (TeNPy)

Example: DMRG



Example

```
from tenpy.networks.mps import MPS
from tenpy.models.tf_ising import TFChain
from tenpy.algorithms import dmrg

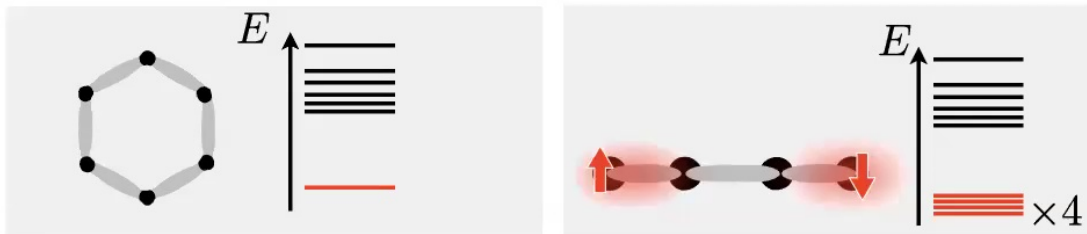
M = TFChain({'L': 16, 'J': 1., 'g': 1.5})
psi = MPS.from_product_state(M.lat.mps_sites(),
                             ['up']*16, 'finite')
dmrg_params = {'trunc_params': {'chi_max': 30,
                                'svd_min': 1.e-10}}
dmrg.run(psi, M, dmrg_params) # find ground state
print("E =", sum(M.bond_energies(psi)))
print("final bond dimensions: ", psi.chi)
```



Symmetry protected topological phases

Spin-1 Heisenberg chain $H = \sum_j \vec{S}_j \cdot \vec{S}_{j+1} \dots \bullet \bullet \bullet \bullet \bullet \bullet \bullet \bullet \dots$

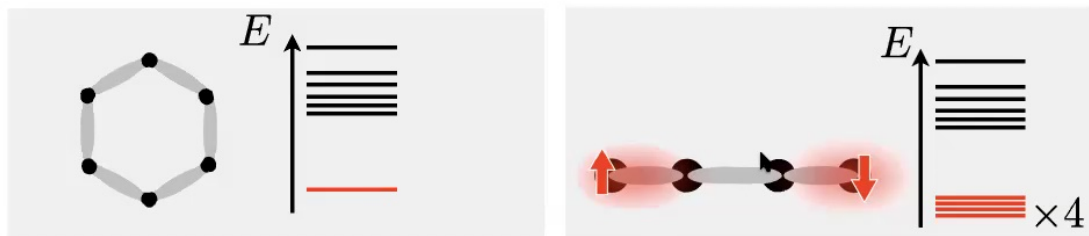
- **Haldane phase:** Gapped and no symmetry breaking [Haldane '83]
- **Spin-1/2** excitations at the edges: Protected by symmetry [Affleck et al '87]



Symmetry protected topological phases

Spin-1 Heisenberg chain $H = \sum_j \vec{S}_j \cdot \vec{S}_{j+1} \dots \bullet \bullet \bullet \bullet \bullet \bullet \bullet \bullet \dots$

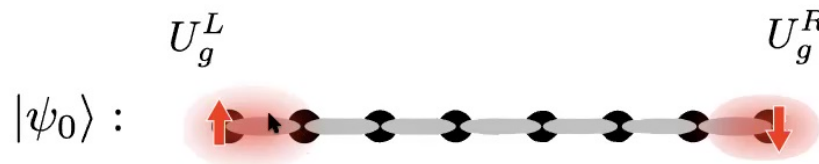
- **Haldane phase:** Gapped and no symmetry breaking [Haldane '83]
- **Spin-1/2** excitations at the edges: Protected by symmetry [Affleck et al '87]



Edge spins have been observed in the NMR profile close to the chain ends of Mg-doped Y_2BaNiO_5 [S.H. Glarum, et al., Tedoldi et al. '99]

Symmetry protected topological phases

Local Hamiltonian and gapped ground state $|\psi_0\rangle$:
Symmetric under $g, h \in G$



Bulk: Linear on-site representation $u_g u_h = u_{gh}$
(e.g., spin-1)

Boundary: Projective representations $U_g U_h = e^{i\phi(g,h)} U_{gh}$
(e.g., spin-1/2)

Classified by the **second cohomology** $H^2[G, U(1)]$ [Schur 1911]

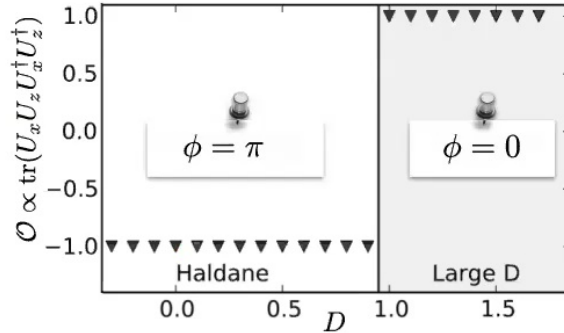
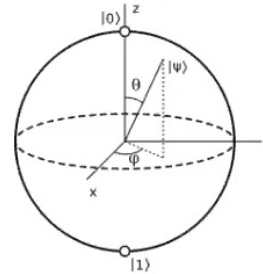
FP, A.M. Turner, E. Berg, and M. Oshikawa, Phys. Rev. B **81**, 064439 (2010).



Symmetry protected topological phases

$\mathbb{Z}_2 \times \mathbb{Z}_2$ symmetry protects the Haldane phase

$$H = \sum_j \vec{S}_j \cdot \vec{S}_{j+1} + D \sum_j (S_j^z)^2$$



$$U_x U_z = e^{i\phi} U_z U_x, \quad \phi = 0, \pi$$

(spin-1 vs. spin)

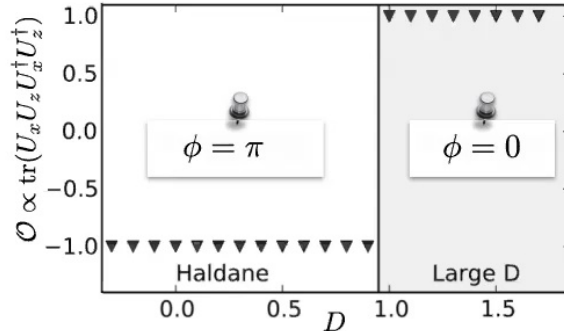
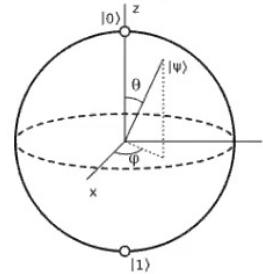
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(spin-1 vs. spin)

- ▶ Non-local order parameters
- ▶ Degeneracies in the entanglement spectrum

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Matrix Product States and 1D Quantum Systems

- I) Entanglement and Matrix-Product States
- II) Time Evolving Block Decimation
- III) Density-Matrix Renormalization Group

Tutorials:

SVD Compression https://colab.research.google.com/drive/1_4V66KqRXHMR-CQ118eqkQBu9EZqw2KM

TEBD Algorithm https://colab.research.google.com/drive/1G_8r4lfiCYKHLmhrFKH4NT4Q2ZWdTNRT

DMRG Algorithm (TeNPy): Ising model https://colab.research.google.com/drive/1V9VAop9_37p2FuqVANHIGGCThtZDD_GM

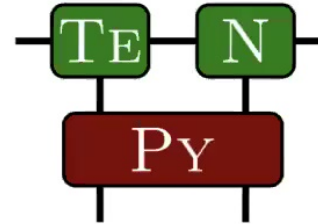
DMRG Algorithm (TeNPy): S-I SPT model https://colab.research.google.com/drive/1A1V7_y75BtmLmUUoDPOw26y41eCGbl3b

(Documentation of TeNPy: <https://tenpy.readthedocs.io/en/latest/>)



Tensor Network Python (TeNPy)

What is TeNPy



- ▶ Python 3 library for simulations with tensor network
<https://github.com/tenpy/tenpy>
- ▶ Object oriented, modular structure, and easy to install
- ▶ HTML documentation
<https://tenpy.github.io>
- ▶ (in)finite DMRG, TEBD, TDVP, ...



Johannes Hauschild, Berkeley



Density matrix renormalization group (DMRG)

Matrix-Product operators: Generalization of MPS to the space of linear operators

$$\mathcal{O}_{i_1 i_2 \dots i_L, i'_1, i'_2 \dots i'_L} =$$

The diagram illustrates the contraction of three tensors to form an MPO. The top row consists of four M tensors connected horizontally. The middle row consists of four M tensors connected horizontally, with a top layer of Γ and Λ tensors and a bottom layer of Γ and Λ tensors. The bottom row consists of four M tensors connected horizontally, with a top layer of Γ and Λ tensors and a bottom layer of Γ^* and Λ tensors.



Matrix Product States and ID Quantum Systems

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