

Title: Asymptotically Safe Amplitudes from the Quantum Effective Action

Speakers: Frank Saueressig

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Abstract: This talk will feature a brief introduction to the gravitational asymptotic safety program before reviewing the current status of the field. Motivated by recent developments, I will introduce the form factor formulation of the quantum effective action and explain how various quantum gravity programs can be embedded into this framework. Finally, I will discuss a novel gravity-matter model whose scattering amplitudes exhibit all features expected from Asymptotic Safety.

# Asymptotically Safe Amplitudes from the Quantum Effective Action

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T. Draper, B. Knorr, C. Ripken and F.S., arXiv:2007.00733

T. Draper, B. Knorr, C. Ripken and F.S., arXiv:2007.04396

Quantum Gravity 2020, Perimeter Institute, July 14<sup>th</sup>, 2020

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## Acknowledgments

J.F. Donoghue, A Critique of the Asymptotic Safety Program, arXiv:1911.02967

- Is the running of  $G, \Lambda$  in conflict with EFT scattering amplitudes?
- Is there a Lorentzian version of Asymptotic Safety?

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- Is the running of  $G, \Lambda$  in conflict with EFT scattering amplitudes?
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Bonanno et. al., [Critical reflections on asymptotically safe gravity](#), arXiv:2004.06810

- surveys open questions of both technical and conceptual nature

The “Tuesday Night Club” collaboration

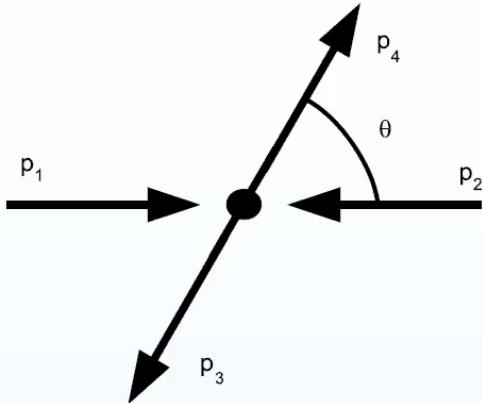


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## Outline

- Setting the scene: scattering amplitudes in General Relativity
- Asymptotic Safety: introduction and recent developments
- Form Factors in Quantum Gravity
  - Parameterizing the quantum effective action
  - Form Factors for Lorentzian Asymptotic Safety
- Conclusions

## scattering of particles



quantum mechanics: amplitude  $\mathcal{A}$  relating in- and out-states:

- parameterized by Mandelstam variables

$$s = (p_1 + p_2)^2 \quad t = (p_1 + p_3)^2 \quad u = (p_1 + p_4)^2$$

- $s$  = center-of-mass energy
- massless scalar fields:  $t = -s(1 + \cos \theta)/2 \quad u = -s(1 - \cos \theta)/2$

- $\mathcal{A}$  computed via Feynman diagrams

## Gravity-mediated particle scattering in general relativity

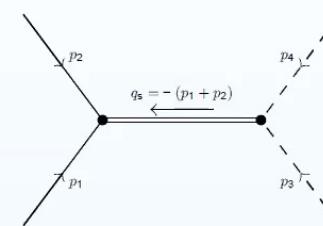
Starting point: Einstein-Hilbert action with minimally coupled scalar fields:

$$\Gamma = -\frac{1}{16\pi G_N} \int d^4x \sqrt{-g} R + \frac{1}{2} \int d^4x \sqrt{-g} [\phi \Delta \phi + \chi \Delta \chi]$$

### 1. amplitude for $\phi\phi \rightarrow \chi\chi$ -scattering

$$\mathcal{A}_s = 8\pi G_N \frac{tu}{s} \sim G_N s$$

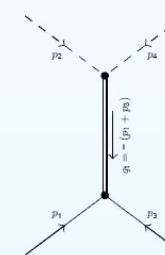
- grows with the center-of-mass energy  $s$ !



### 2. amplitude for $\phi\chi \rightarrow \chi\phi$ -scattering

$$\mathcal{A}_t = 8\pi G_N \frac{su}{t} \sim G_N \frac{s^2}{t}$$

- quadratic divergence in the forward-scattering limit



## perturbative quantization of general relativity: conclusions

a) treat gravity as **effective field theory** [J. Donoghue, gr-qc/9405057]

- compute corrections in  $\mathfrak{s} G_N \ll 1$  (independent of UV-completion)
- predictivity breaks down at  $\mathfrak{s} G_N \gtrsim 1$

b) quantizing gravity may require **new physics concepts**:

- String Theory (Veneziano and Virasoro-Shapiro-amplitudes)
- Loop Quantum Gravity

c) gravity makes sense as quantum field theory:

- Asymptotic Safety
- Causal Dynamical Triangulations
- Infinite Derivative Gravity
- ...

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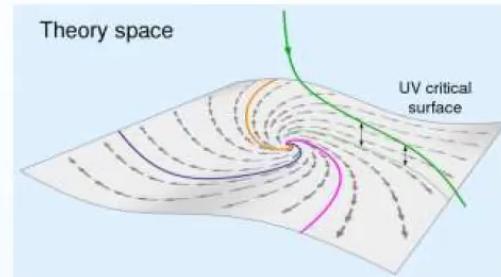
# Asymptotic Safety

## introduction and recent developments

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## Asymptotic Safety: key idea

high-energy completion of gravity by Reuter fixed point



- 2 classes of renormalization group trajectories:
  - relevant = end at fixed point in UV
  - irrelevant = go somewhere else...
- predictive power:
  - hitting the fixed point at high energy  $\iff$  relations between couplings

## The Asymptotic Safety package

### Reuter fixed point

- ensures the absence of UV-divergences

fixed point has finite-dimensional UV-critical surface (predictivity)

- select a RG trajectory with a finite number of measurements

RG flow ends with a effective action compatible with observations (falsifiability)

- tests of general relativity
  - solar system tests, cosmological signatures, gravitational waves, ...
- compatibility with standard model of particle physics at 1 TeV

structural demands

- resolution of spacetime singularities
- unitarity

# Exploring Asymptotic Safety via Functional Renormalization

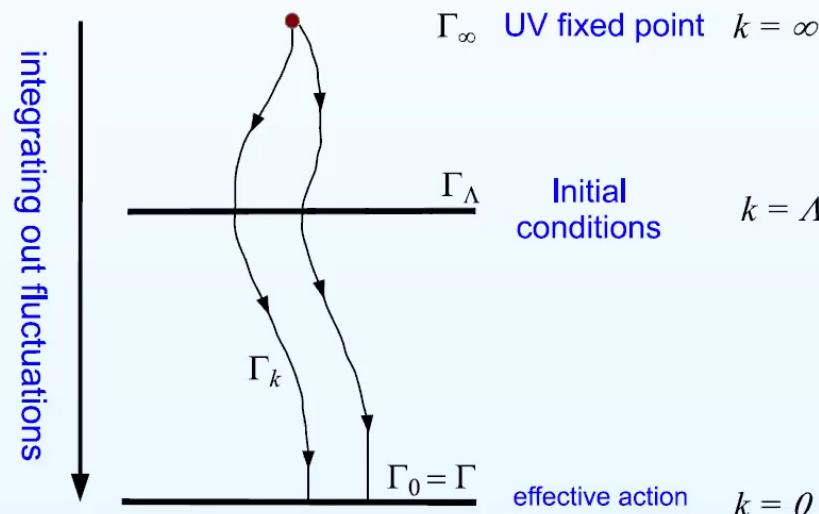
M. Reuter, Phys. Rev. D 57 (1998) 971

idea: integrate out quantum fluctuations shell-by-shell in momentum-space



implementation: Wetterich equation for effective average action  $\Gamma_k$ :

$$k\partial_k \Gamma_k[h_{\mu\nu}, \bar{g}_{\mu\nu}] = \frac{1}{2} \text{Tr} \left[ \left( \Gamma_k^{(2)} + \mathcal{R}_k \right)^{-1} k\partial_k \mathcal{R}_k \right]$$



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conceptual remarks

1. fixed points are a property of the solutions of the Wetterich equation
2. physics is in the effective action  $\Gamma \equiv \Gamma_{k=0}$
3. the fixed point relates couplings in the effective action

## Solving the Wetterich equation: derivative expansion of $\Gamma_k^{\text{grav}}[g]$

⋮

⋮

$$R^8$$

$$C_{\mu\nu\rho\sigma}\Delta^6 C^{\mu\nu\rho\sigma}$$

Einstein-Hilbert truncation

$$R^7$$

$$C_{\mu\nu\rho\sigma}\Delta^5 C^{\mu\nu\rho\sigma}$$

$$R^6$$

$$C_{\mu\nu\rho\sigma}\Delta^4 C^{\mu\nu\rho\sigma}$$

$$R^5$$

$$C_{\mu\nu\rho\sigma}\Delta^3 C^{\mu\nu\rho\sigma}$$

$$R^4$$

$$C_{\mu\nu\rho\sigma}\Delta^2 C^{\mu\nu\rho\sigma}$$

$$R^3$$

$$C_{\mu\nu\rho\sigma}\Delta C^{\mu\nu\rho\sigma}$$

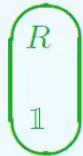
$$R \Delta R$$

+ 5 more

$$R^2$$

$$C_{\mu\nu\rho\sigma} C^{\mu\nu\rho\sigma}$$

$$R_{\mu\nu} R^{\mu\nu}$$



## Asymptotic Safety: selected results

pure gravity:

- non-Gaussian fixed point persists once 2-loop counterterm is included  
[H. Gies, B. Knorr, S. Lipholdt, and F. Saueressig, arXiv:1601.01800]
- low number of free parameters ( $\simeq 3$ )  
[ R. Percacci and A. Codello, arXiv:0705.1769]  
[ D. Benedetti, P.F. Machado and F. Saueressig, arXiv:0901.2984]  
[ K. Falls, D. F. Litim, . . . , arXiv:1801.00162; arXiv:1810.08550]
- non-Gaussian fixed point in  $d = 2$  is a unitary CFT  
[ A. Nink and M. Reuter, arXiv:1512.06805]
- cosmology (Starobinsky inflation) may fix all free parameters  
[G. Gubitosi, R. Ooijer, C. Ripken, F. Saueressig, arXiv:1806.10147]

gravity coupled to matter:

- non-Gaussian fixed point compatible with standard-model matter  
[ P. Dona, A. Eichhorn and R. Percacci, arXiv:1311.2898]  
[ J. Meibohm, J. M. Pawłowski and M. Reichert, arXiv:1510.07018]
- asymptotically safety could predict parameters of the standard model  
[ M. Shaposhnikov and C. Wetterich, arXiv:0912.0208]  
[ A. Eichhorn, A. Held, arXiv:1803.04027]  
[ A. Eichhorn, arXiv:1810.07615]

## Open Questions

- number of free parameters?
- low-energy physics compatible with observations?
- exact matter content supporting asymptotic safety?
- characterization of the quantum geometry?

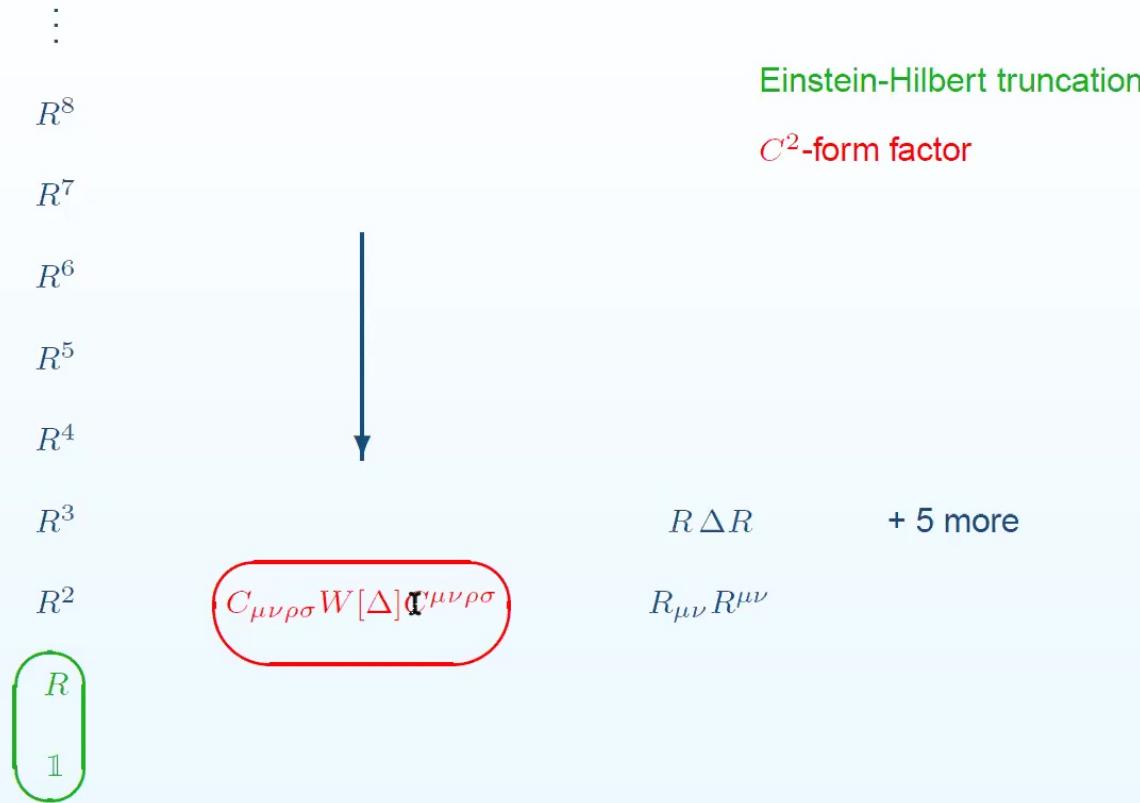
[talk by R. Loll]

- degrees of freedom associated with the Reuter fixed point?
- unitarity?

### investigation method

- derivative expansion of  $\Gamma_k$
- derivative expansion fails by construction

## Derivative expansion $\implies$ curvature expansion of $\Gamma_k^{\text{grav}}[g]$



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## Open Questions

- number of free parameters?
- low-energy physics compatible with observations?
- exact matter content supporting asymptotic safety?
- characterization of the quantum geometry associated with NGFP?
- degrees of freedom associated with the NGFP?
- unitarity?

investigation method

- derivative expansion of  $\Gamma_k$
- derivative expansion fails by construction

derivative expansion  $\implies$  curvature expansion keeping derivatives

## Parameterizing the Quantum Effective action

building blocks for gravity-scalar theory:

$$\Gamma = \Gamma_{\text{grav}} + \Gamma_{\text{gf}} + \Gamma_{\text{matter}}$$

curvature expansion in the gravitational sector

$$\Gamma_{\text{grav}} = \frac{1}{16\pi G_N} \int \sqrt{-g} \left[ -R - \frac{1}{6} R f_R(\Delta) R + \frac{1}{2} C_{\mu\nu\rho\sigma} f_C(\Delta) C^{\mu\nu\rho\sigma} + \dots \right]$$

- form factors  $f_R(\Delta)$  and  $f_C(\Delta)$  fix flat space graviton propagators  $G_2, G_0$

supplemented by two-parameter family of gauge-fixing conditions:

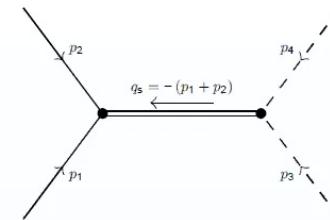
$$\Gamma_{\text{gf}} = \frac{1}{32\pi\alpha} \int \sqrt{-\eta} \left( \partial^\mu h_{\mu\nu} - \frac{1+\beta}{4} \partial_\nu h \right) \left( \partial_\rho h^{\rho\nu} - \frac{1+\beta}{4} \partial^\nu h \right).$$

fluctuations around flat background:  $g_{\mu\nu} = \eta_{\mu\nu} + h_{\mu\nu}$

## Gravity mediated scattering $\phi\phi \rightarrow \chi\chi$

on-shell amplitude

$$\begin{aligned} \mathcal{A}_{\mathfrak{s}}^{\phi\chi} = \frac{4\pi G_N}{3} & \left[ - \left( 1 + \mathfrak{s} f_{Ric\phi\phi}(\mathfrak{s}, m_\phi^2, m_\phi^2) \right) \left( 1 + \mathfrak{s} f_{Ric\chi\chi}(\mathfrak{s}, m_\chi^2, m_\chi^2) \right) \right. \\ & \times G_2(\mathfrak{s}) \times \left( \mathfrak{t}^2 - 4\mathfrak{t}\mathfrak{u} + \mathfrak{u}^2 + 2(m_\phi^2 - m_\chi^2)^2 \right) \\ & + \left( (\mathfrak{s} + 2m_\phi^2)(1 + \mathfrak{s} f_{Ric\phi\phi}(\mathfrak{s}, m_\phi^2, m_\phi^2)) - 12\mathfrak{s} f_{R\phi\phi}(\mathfrak{s}, m_\phi^2, m_\phi^2) \right) \\ & \times \left. \left( (\mathfrak{s} + 2m_\chi^2)(1 + \mathfrak{s} f_{Ric\chi\chi}(\mathfrak{s}, m_\chi^2, m_\chi^2)) - 12\mathfrak{s} f_{R\chi\chi}(\mathfrak{s}, m_\chi^2, m_\chi^2) \right) G_0(\mathfrak{s}) \right] \end{aligned}$$



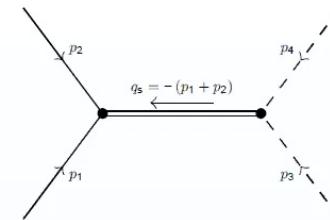
properties:

- gauge invariant
- invariant under a momentum-dependent rescaling of the graviton

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properties:

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- invariant under a momentum-dependent rescaling of the graviton

## Gravity mediated scattering $\phi\phi \rightarrow \chi\chi$ : Stelle gravity

define partial wave amplitudes to separate angular dependence of  $\mathcal{A}_s^{\phi\chi}$ :

$$a_j^{\phi\chi}(s) \equiv \frac{1}{32\pi} \int_{-1}^1 d(\cos\theta) P_j(\cos\theta) \mathcal{A}_s^{\phi\chi}(s, \cos\theta),$$

- $P_j(x)$ : Legendre polynomial of order  $j$

Example: classical Stelle gravity with minimally coupled scalars

$$\Gamma_{\text{grav}}^{\text{Stelle}} = \frac{1}{16\pi G_N} \int \sqrt{-g} \left[ -R - \frac{1}{6} R f_R(\Delta) R + \frac{1}{2} C_{\mu\nu\rho\sigma} f_C(\Delta) C^{\mu\nu\rho\sigma} + \dots \right]$$

- non-zero, constant form factors

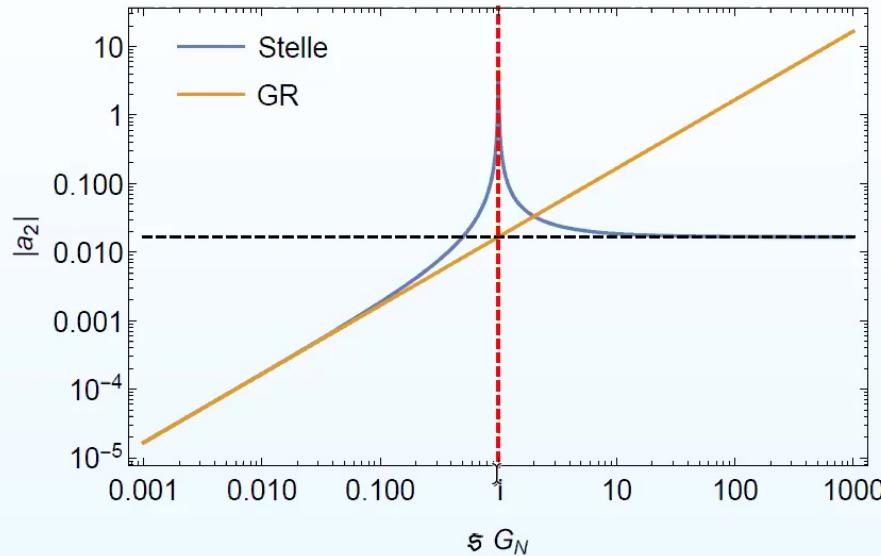
$$f_R = -c_R G_N, \quad f_C = -c_C G_N$$

- partial wave amplitudes

$$a_0^{\text{Stelle}} = \frac{G_N}{12} s^2 \left[ \frac{1}{s} - \frac{1}{s - (c_R G_N)^{-1}} \right], \quad a_2^{\text{Stelle}} = -\frac{G_N}{60} s^2 \left[ \frac{1}{s} - \frac{1}{s - (c_C G_N)^{-1}} \right]$$

## Partial-wave amplitude for classical Stelle gravity

$$a_2^{\text{Stelle}} = -\frac{G_N}{60} \mathfrak{s}^2 \left[ \frac{1}{\mathfrak{s}} - \frac{1}{\mathfrak{s} - (c_C G_N)^{-1}} \right]$$



growth of amplitude tamed by massive spin-two poltergeist

## Form factors: a unifying perspective on quantum gravity

theory	gravity form factors	pole structure	UV behavior
effective field theory	log	pole at cutoff scale	n/a
Stelle gravity	const	massive spin-two ghost d.o.f.	const scale-free
infinite derivative gravity	exp	essential singularity at $\infty$	exp-falloff ( $s$ ), exp-divergent ( $t/u$ )

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powerful framework to study quantum gravity programs in a unifying way

## Blueprint for Asymptotically Safe Amplitudes

requirements

- amplitudes become scale-free at high energy  
     $\implies$  unitarity, Froissart bound
- no new massive degrees of freedom (ghosts)  
     $\implies$  unitarity, causality
- locality at high energy

form factors for Lorentzian Asymptotic Safety

- graviton propagators

$$f_R(\mathfrak{s}) = c_R G_N \tanh(c_R G_N \mathfrak{s}) \quad f_C(\mathfrak{s}) = c_C G_N \tanh(c_C G_N \mathfrak{s})$$

- scalar self-interactions  $f_{\phi^2 \chi^2}$  contributes

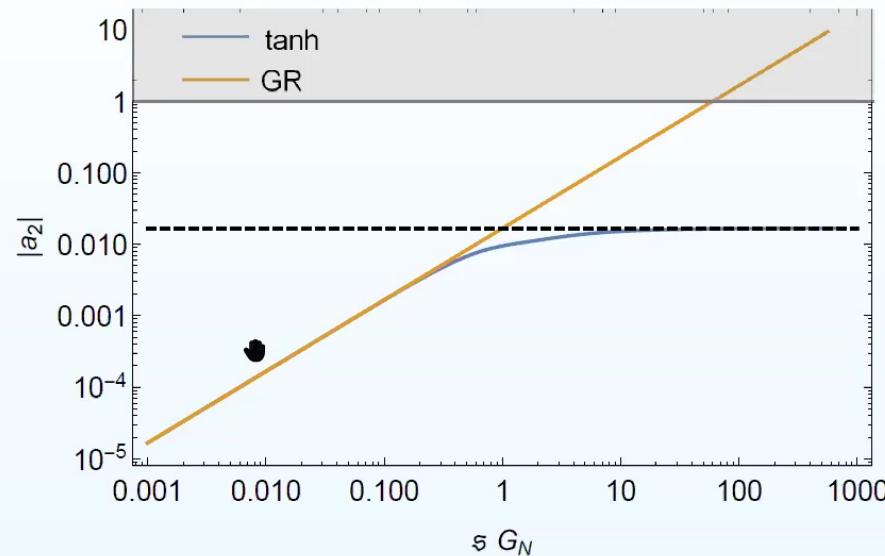
$$\mathcal{A}_4^{\phi\chi}(\mathfrak{s}|\mathfrak{t}, \mathfrak{u}) = 4\pi G_N G_2(\mathfrak{s})(\mathfrak{t}^2 + \mathfrak{u}^2) f^{\text{int}}(\mathfrak{s}^2 + \mathfrak{t}^2 + \mathfrak{u}^2)$$

interpolation function  $f^{\text{int}}(x)$  suppresses self-interaction at low energy

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## Partial-wave amplitude for Lorentzian Asymptotic Safety

$$a_2 = -\frac{G_N}{60} \frac{\varsigma}{1 + c_C G_N \varsigma \tanh(c_C G_N \varsigma)}$$

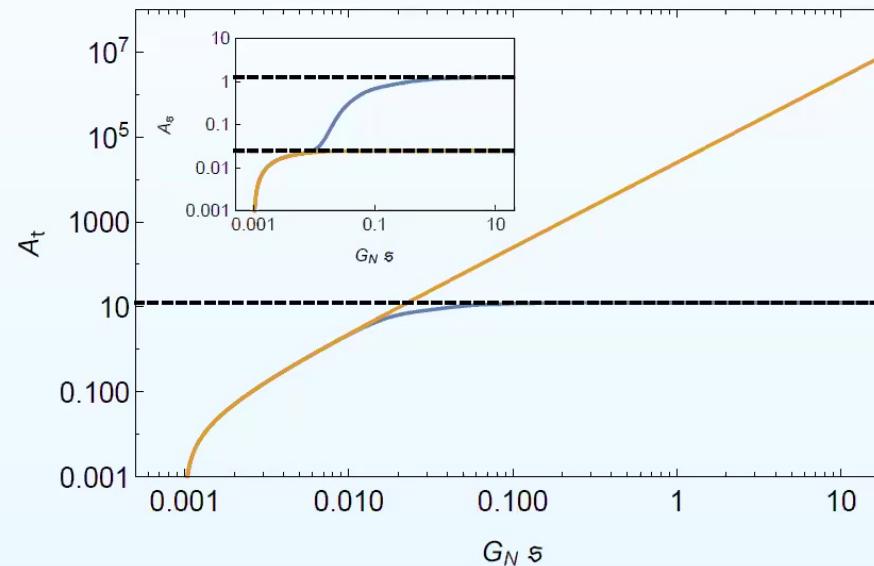
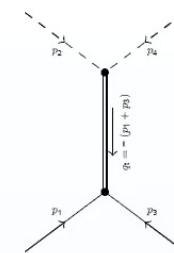


amplitude becomes scale-free at large center-of-mass energy  $G_N \varsigma \gg 1$

## Gravity-mediated scattering $\phi\chi \rightarrow \phi\chi$

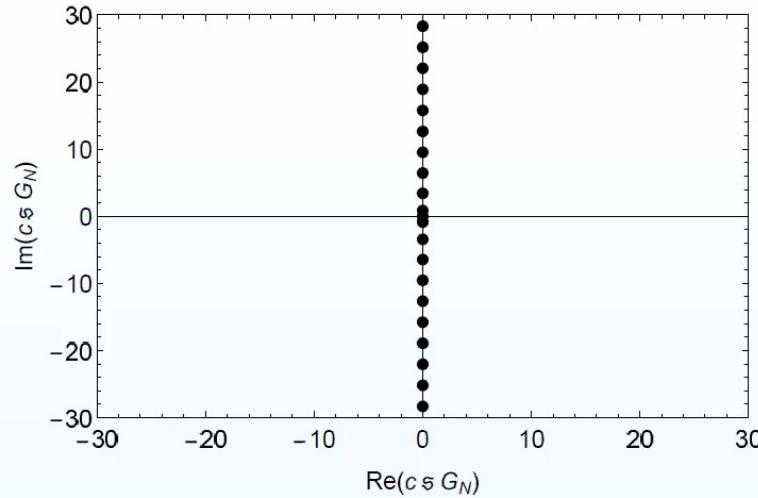
related to  $\phi\phi \rightarrow \chi\chi$  by crossing symmetry  $s \leftrightarrow t$

- graviton-contribution to  $t$ -channel amplitude  $A_t \sim s^2 G_C(t)$
- divergence canceled by scalar self-interaction



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## Mechanism underlying Lorentzian Asymptotic Safety



- propagators exhibit an infinite tower of poles at imaginary squared momentum

$$G_2(\varsigma) \propto \frac{1}{\varsigma [1 + c_C G_N \varsigma \tanh(c_C G_N \varsigma)]}$$

- poles exhibit Regge-type scaling asymptotically  $\hat{\Gamma}_n \propto i n$
- avoids causality violation typical for Lee-Wick models

## Concluding comments

scale-free amplitudes require relations among propagators and vertices  
provided by Reuter fixed point

- evidence from first-principle computations (Wetterich equation)
  - form factor computations

[L. Bosma, B. Knorr, F. Saueressig, arXiv:1904.04845]  
[B. Knorr, C. Ripken, F. Saueressig, arXiv:1907.02903]

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[N. Christiansen, D.F. Litim, J.M. Pawłowski, M. Reichert, arXiv:1710.04669]  
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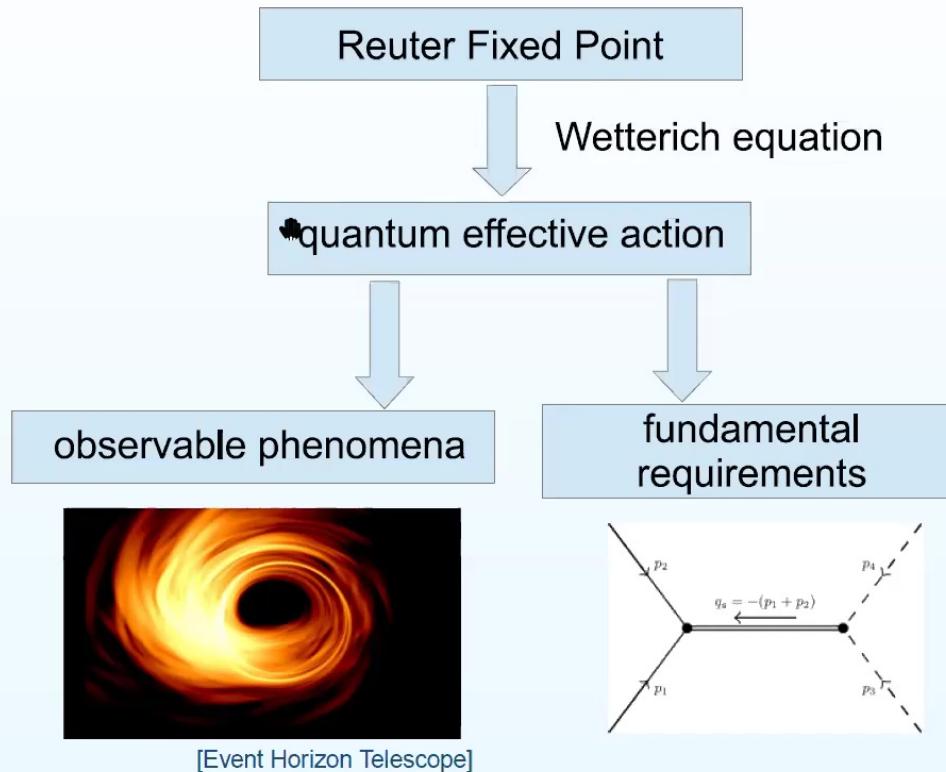
getting quantum-corrected amplitudes via “RG-improvement” fails

- RG-information insufficient for reproducing all amplitudes:

$$G(k^2) \longrightarrow f_C(p^2), f_R(p^2), f_{R\phi\phi}(p_1^2, p_2^2, p_3^2), \dots$$

## Conclusions ...

Structure of the Asymptotic Safety Program



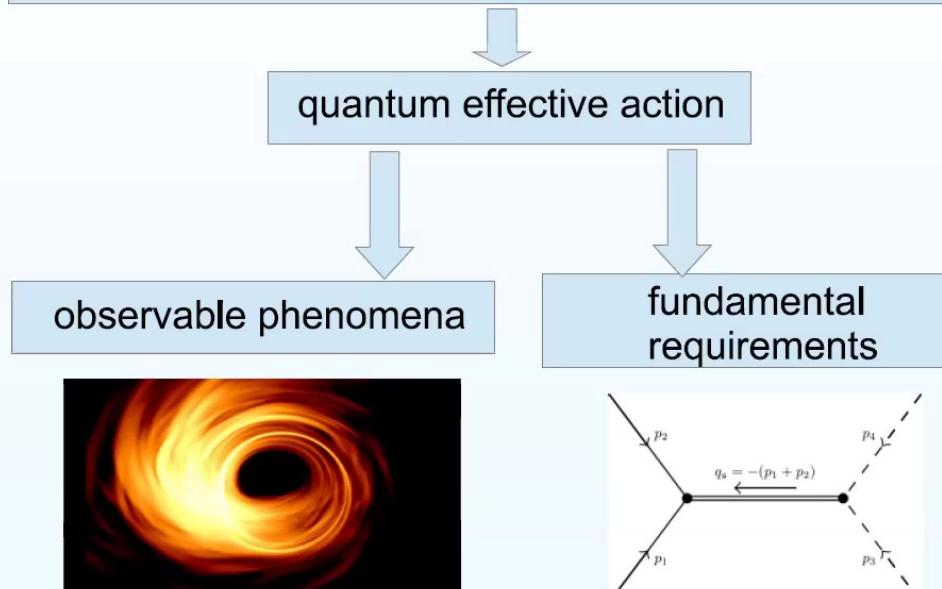
[Event Horizon Telescope]

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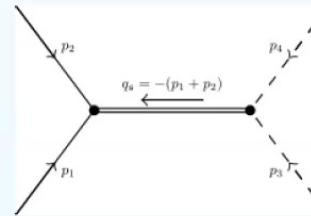
## Conclusions and Outlook

[discussion session: Daniele Oriti, Thursday]

⌚ Asymptotic Safety, Causal Dynamical Triangulations, Causal Set Theory,  
Loop Quantum Gravity, Group Field Theory, Infinite Derivative Gravity,  
weakly non-local gravity, ...



[Event Horizon Telescope]

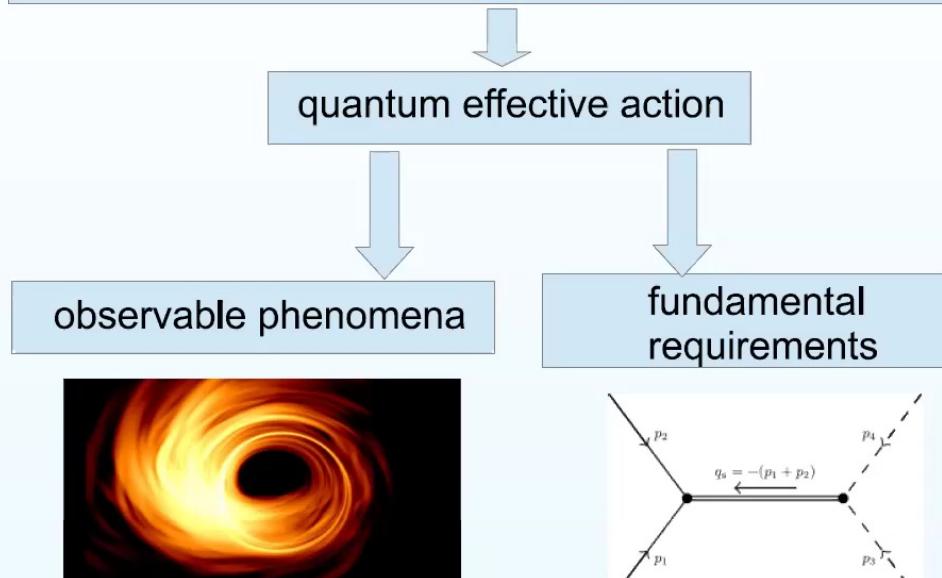


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[Event Horizon Telescope]

Thank You!

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