

Title: Summer Undergrad 2020 - Path Integrals (M) - Lecture 2

Speakers: Dan Wohns

Collection: Summer Undergrad 2020 - Path Integrals

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Abstract: Propagator in real and imaginary time

## Lecture 2

$$K(q_f, t_f, q_i, t_i) = \int \prod_{j=1}^{N-1} dq_j \overbrace{K_{q_{j+1}, q_j}}^{\text{rename } K_{q_{j+1}, q_j}}(q_{j+1}, t_{j+1}, q_j, t_j)$$

$$K_{q_{j+1}, q_j} = \langle q_{j+1} | e^{-iH\Delta t/\hbar} | q_j \rangle$$

$$= \langle q_{j+1} | 1 - i \frac{H\Delta t}{\hbar} + \mathcal{O}(\Delta t^2) | q_j \rangle$$

$$\langle q_{j+1} | q_j \rangle = \delta(q_{j+1} - q_j) \quad \leftarrow \text{breakout rooms}$$

$$= \int \frac{dp_j}{2\pi\hbar} e^{ip_j(q_{j+1} - q_j)/\hbar}$$

$$K_{q_{j+1}, q_j} = \int \frac{dp_j}{2\pi\hbar} e^{i p_j (q_{j+1} - q_j) / \hbar} \left( 1 - \frac{i\Delta t H}{\hbar} + \mathcal{O}(\Delta t^2) \right)$$

$$= \int \frac{dp_j}{2\pi\hbar} e^{i\frac{\Delta t}{\hbar} \left( p_j q_j - \frac{p_j^2}{2m} - V(q) \right)} e^{-\frac{i\Delta t}{\hbar} H}$$

$H = \frac{p_j^2}{2m} + V(q_j)$

$$= \sqrt{\frac{m}{2\pi i \hbar \Delta t}} e^{i\frac{\Delta t}{\hbar} \left( \frac{m}{2} \dot{q}^2 - V(q) \right)}$$

Gaussian int  $\int e^{ax^2+bx+c} dx$

$$K = \left( \frac{m}{2\pi i \hbar} \right)^{N/2} \int \left( \prod_{j=1}^{N-1} dq_{vj} \right) \exp \left[ \frac{i\Delta t}{\hbar} \sum_{j=0}^{N-1} \left( \frac{m \dot{q}_j^2}{2} - V(q_j) \right) \right]$$



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# Lecture 2

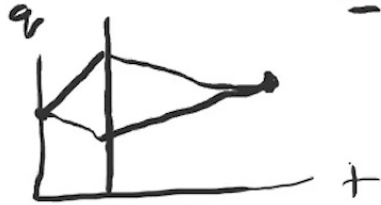
$$K(q_f, t_f, q_i, t_i) = \int \prod_{j=1}^{N-1} dq_j \overbrace{K_{q_{j+1}, q_j}}^{\text{rename } K_{q_{j+1}, q_j}}(q_{j+1}, t_{j+1}, q_j, t_j)$$

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$$\langle q_{j+1} | q_j \rangle = \delta(q_{j+1} - q_j) \quad \leftarrow \text{breakout rooms}$$

$$= \int \frac{dp_j}{2\pi\hbar} e^{ip_j(q_{j+1} - q_j)/\hbar}$$



$$K_{q_{j+\Delta t}, q_j} = \int \frac{dp_j}{2\pi\hbar} e^{i p_j (q_{j+\Delta t} - q_j) / \hbar} \left( 1 - \frac{i\Delta t H}{\hbar} + \mathcal{O}(\Delta t^2) \right)$$

$$= \int \frac{dp_j}{2\pi\hbar} e^{i\frac{\Delta t}{\hbar} (p_j q_j - \frac{p_j^2}{2m} - V(q))} e^{-\frac{i\Delta t}{\hbar} H}$$

$H = \frac{p_j^2}{2m} + V(q_j)$

$$= \sqrt{\frac{m}{2\pi i \hbar \Delta t}} e^{i\frac{\Delta t}{\hbar} \left( \frac{m}{2} \dot{q}^2 - V(q) \right)}$$

Gaussian int  
 $\int e^{ax^2+bx+c} dx$

$$K = \left( \frac{m}{2\pi i \hbar \Delta t} \right)^{N/2} \int \left( \prod_{j=1}^{N-1} dq_{vj} \right) \exp \left[ \frac{i\Delta t}{\hbar} \sum_{j=0}^{N-1} \left( \frac{m}{2} \dot{q}_j^2 - V(q_j) \right) \right]$$

$$S = \int dt \left( \frac{m}{2} \dot{q}^2 - V \right)$$



$$K = \int \mathcal{D}q(t) e^{iS[q(t)]/\hbar}$$

$$T = -iT_E \quad \leftarrow i \text{ leads to oscillations}$$

$$K_E = \langle q_f | e^{-T_E H/\hbar} | q_i \rangle$$

$t = -i\tau$  everywhere

$$iS[q(t)] = i \int_0^T dt \left( \frac{m}{2} \left( \frac{dq}{dt} \right)^2 - V(q) \right)$$

$$\rightarrow - \int_0^{T_E} dz \left( \frac{m}{2} \left( \frac{dq}{dz} \right)^2 + V(q) \right)$$

$$= -S_E[q(z)]/\hbar$$



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$$Z[\beta] = \sum_j e^{-\beta E_j}$$

$$K_E = \int \mathcal{D}q(z) e^{-S_E[q(z)]/\hbar}$$

$$Z = \int dq K_E(q, \beta \hbar, q, 0) \leftarrow \text{breakout rooms}$$

$\uparrow$   
 $T_E = \beta \hbar$

$$E_0 = -\lim_{\beta \rightarrow \infty} \frac{1}{\beta} \log Z$$



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$\uparrow$   
ground state energy



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## Free path integrals

IF  $S$  is at most quadratic in  $q$

→  $N$  coupled Gaussian integrals

→ exact results

$$K_E = \sqrt{\frac{m}{2\pi\beta}} e^{-S[q_c] \hbar}$$

↑  
breakout rooms

↑  
 $q_c = \text{classical trajectory}$



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$$\overbrace{K_E(q_i/T_E, q_j, 0)}$$

$$Z[\beta] = \int dq \langle q | e^{-HT_E/\hbar} | q \rangle$$

$$1 = \sum_j |j\rangle \langle j|$$

$N \rightarrow \infty$  limit  $\rightarrow$  algebra

$N=1$  can understand much of physics

next: perturbation theory

$$y = \sqrt{\frac{m \Delta x}{2}} \quad \checkmark$$

