

Title: Classical algorithms, correlation decay, and complex zeros of partition functions of quantum many-body systems

Speakers: Mehdi Soleimanifar

Series: Perimeter Institute Quantum Discussions

Date: May 13, 2020 - 3:30 PM

URL: <http://pirsa.org/20050017>

Abstract: Basic statistical properties of quantum many-body systems in thermal equilibrium can be obtained from their partition function. In this talk, I will present a quasi-polynomial time classical algorithm that estimates the partition function of quantum many-body systems at temperatures above the thermal phase transition point. It is known that in the worst case, the same problem is NP-hard below this temperature. This shows that the transition in the phase of a quantum system is also accompanied by a transition in the computational hardness of estimating its statistical properties. The key to this result is a characterization of the phase transition and the critical behavior of the system in terms of the complex zeros of the partition function. I will also discuss the relation between these complex zeros and another signature of the thermal phase transition, namely, the exponential decay of correlations. I will show that in a system of  $n$  particles above the phase transition point, where the complex zeros are far from the real axis, the correlation between two observables whose distance is at least  $\log(n)$  decays exponentially. This is based on joint work with Aram Harrow and Saeed Mehraban.

# Classical algorithms, correlation decay, and complex zeros of quantum partition functions

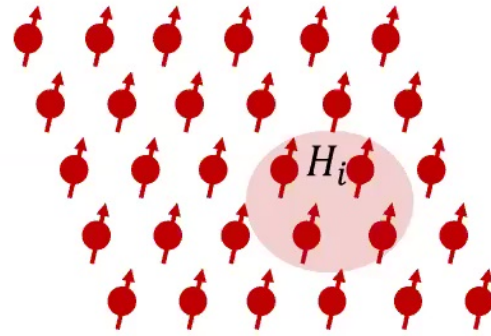
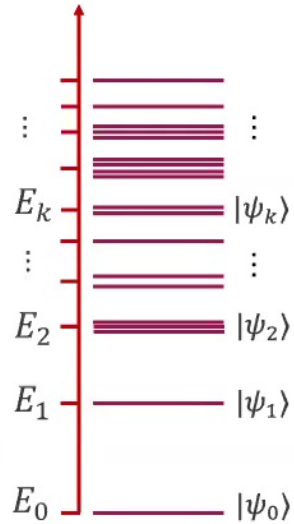
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Joint work with  
Aram Harrow and Saeed Mehraban

arXiv:1910.09071, STOC'20



Energy

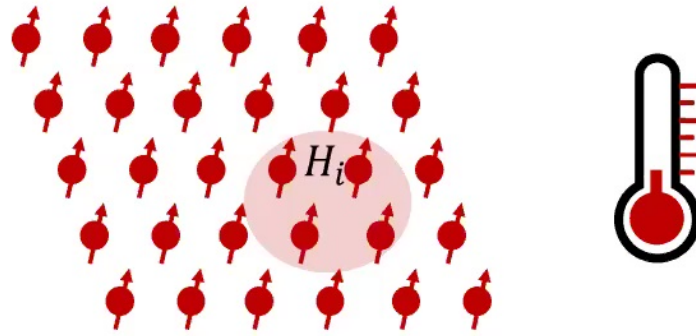
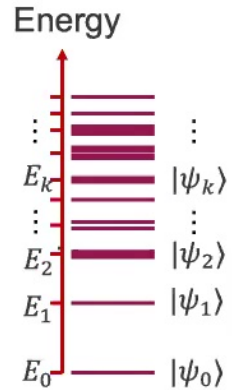


$n$  qudits (or spins) on a **lattice**

Interactions described by  $H = \sum_i H_i$

$H_i$  acting on a **spatially-local** region





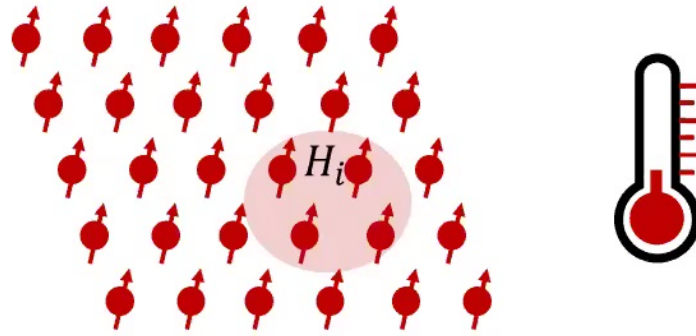
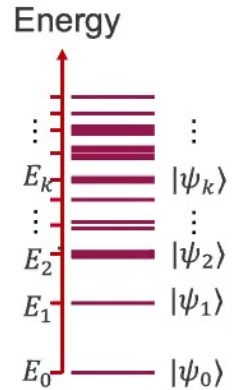
How hard is it to approximate  $E_0$  in the worst case?

It is **QMA-complete** [K'02, KKR'05, ...]

**Efficient** algorithms in special cases [ALVV'16, BH'13, ...]

1D Lattice   
 $H$  gapped



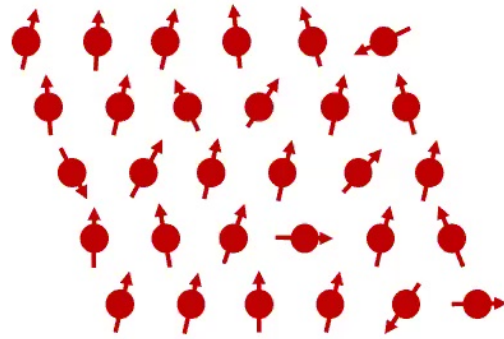
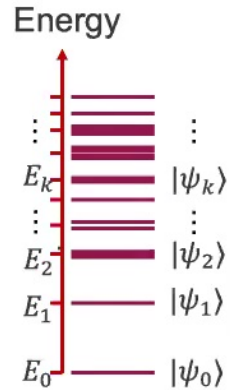


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At **non-zero** temperature  $T = 1/\beta$ , system in state

**Gibbs state**

$$\rho = \frac{1}{Z} \sum_k e^{-\beta E_k} |\psi_k\rangle \langle \psi_k|$$

**Partition function**

$$Z(\beta) = \sum_k e^{-\beta E_k} = \text{Tr}[e^{-\beta H}]$$

$$\beta \rightarrow \infty, \rho \rightarrow |\psi_0\rangle \langle \psi_0|$$

$$\beta \rightarrow 0, \rho \rightarrow I/d^n$$

**Free energy**

$$\beta \rightarrow \infty, F \rightarrow E_0$$

$$F(\beta) = -\frac{1}{\beta} \log Z = \text{Tr}[H\rho] - \frac{1}{\beta} S(\rho)$$



Finding  $Z(\beta)$  exactly is **#P-hard** [Val'79]

**Goal:** design an **approximation** algorithm that

**Input:**  $H, \beta, \varepsilon$

**Output:** estimate  $\hat{Z}(\beta)$  such that

$$|\log Z(\beta) - \log \hat{Z}(\beta)| \leq \varepsilon$$

*Hardness depends on temperature*

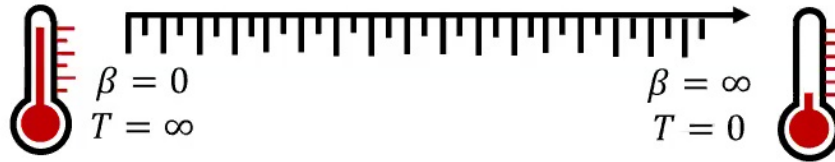


Maximally mixed state

Trivial

Ground state

QMA-hard



Transition in the **hardness**

but we also know about

Transition in the **phase**



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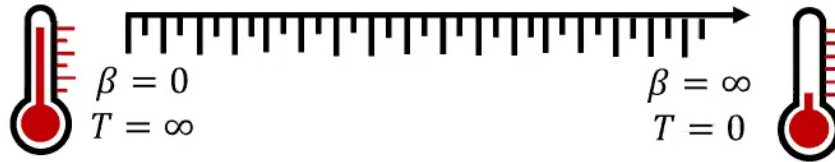


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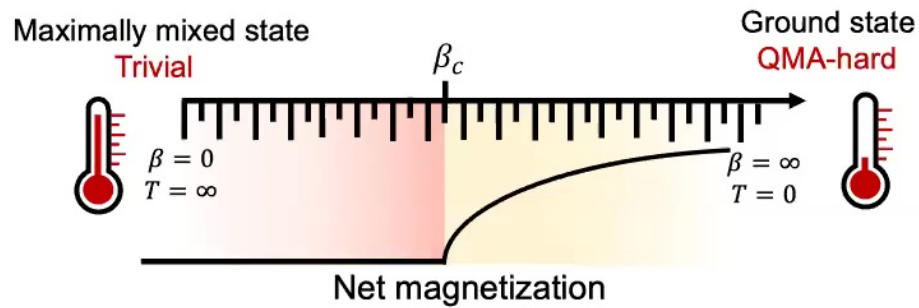
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**Transition** in the **phase**

*Physical properties abruptly change*





**Transition in the hardness**

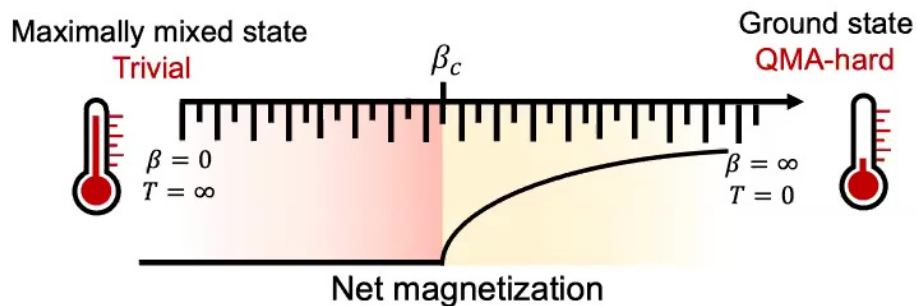
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**Transition** in the **hardness**

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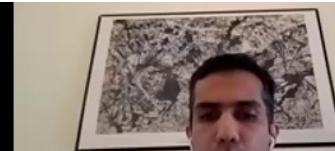
**Transition** in the **phase**

*Physical properties abruptly change*

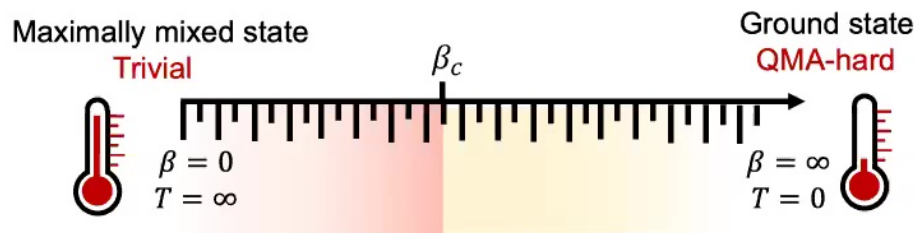
*Hardness of approximating  $Z(\beta)$*

VS

*Physical phase transition?*



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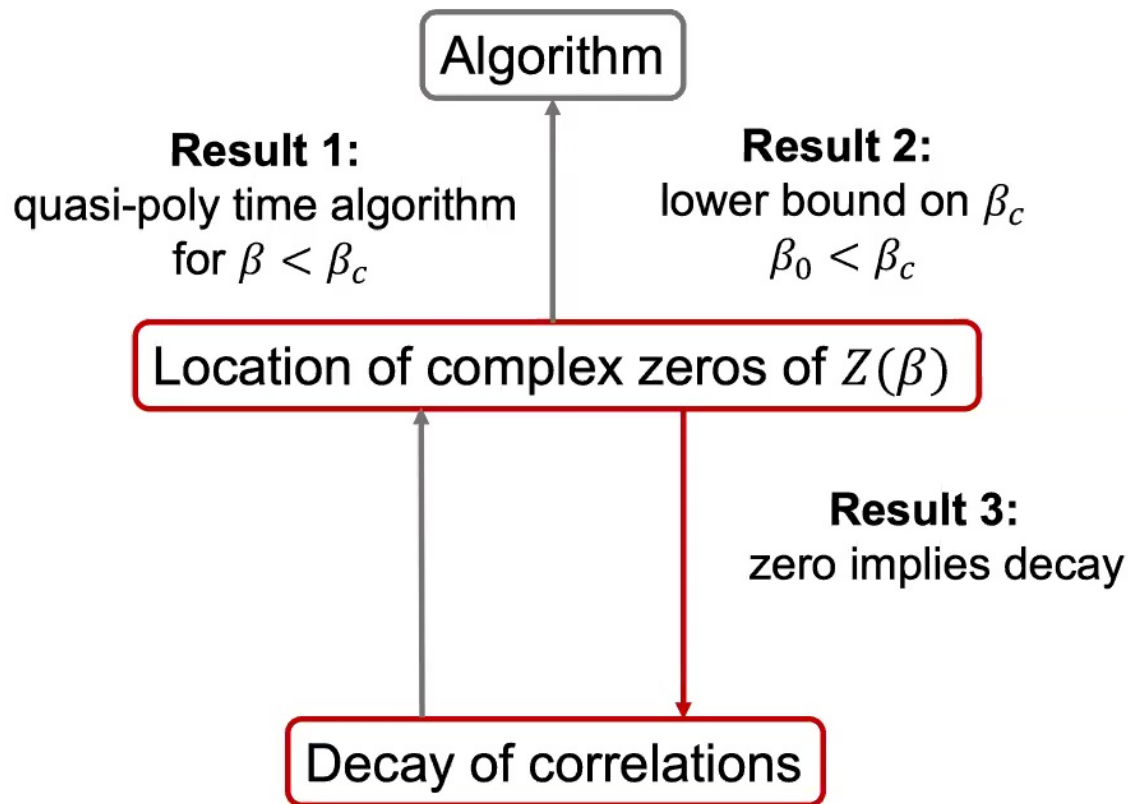
For **classical** Ising model

- Above  $T_c$  there is a polynomial time algorithm
- Below  $T_c$  no efficient algorithm unless  $NP = RP$  [Sly'10]

Computational phase transition  
matches  
Physical phase transition

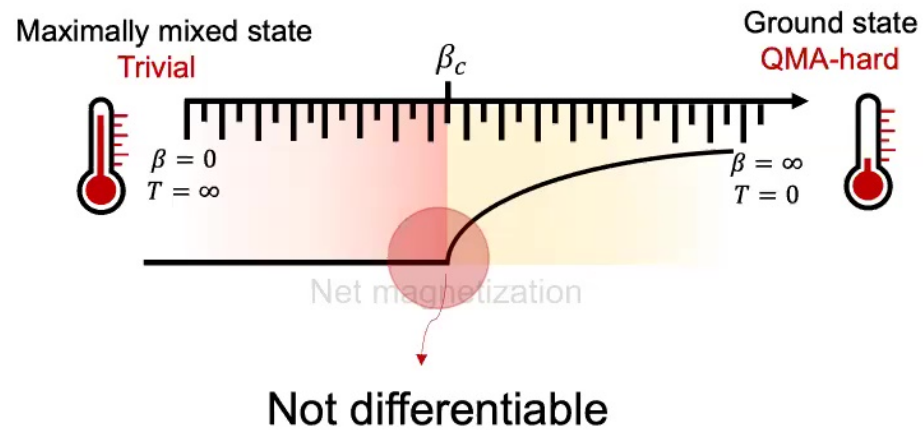


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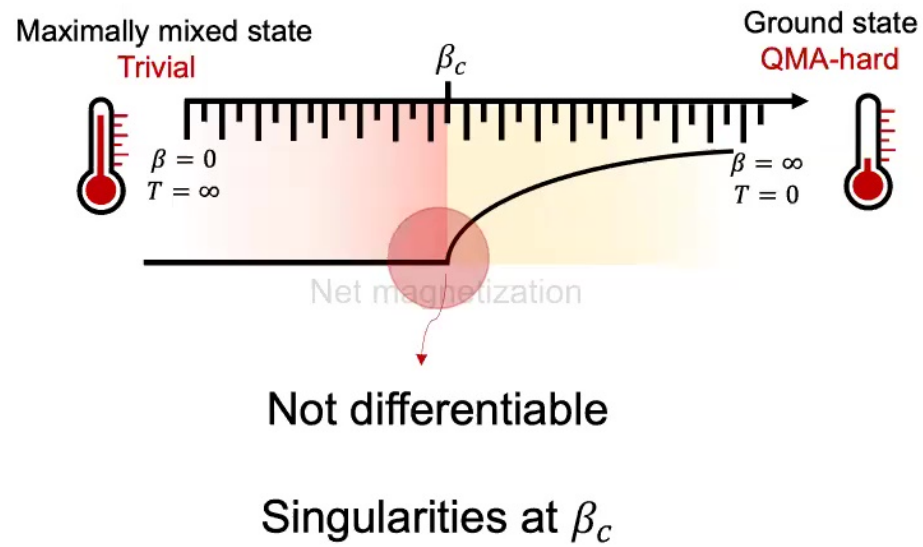
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# Complex zeros vs phase transition

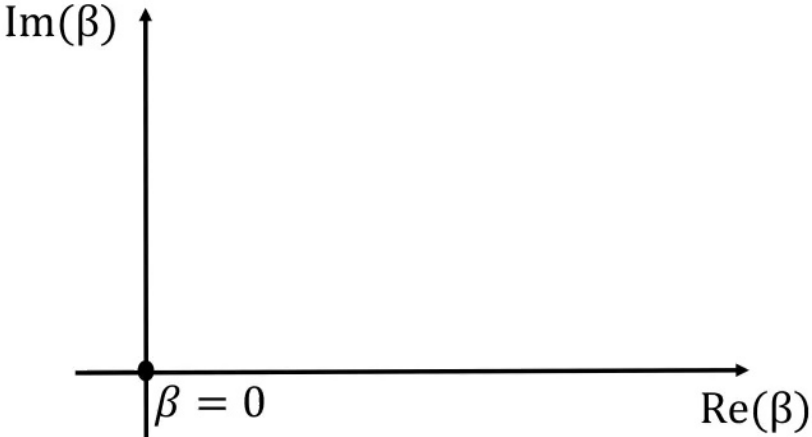


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# Complex zeros vs phase transition



# Complex zeros vs phase transition



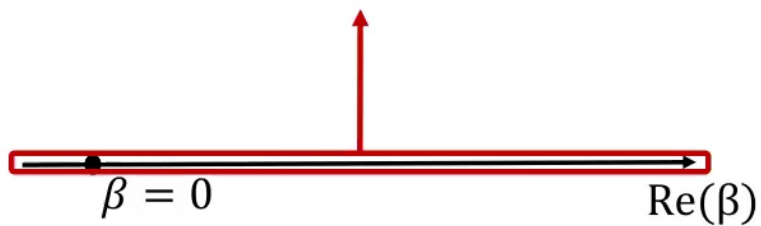
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## Complex zeros vs phase transition

$$Z(\beta) = \sum_k e^{-\beta E_k}$$

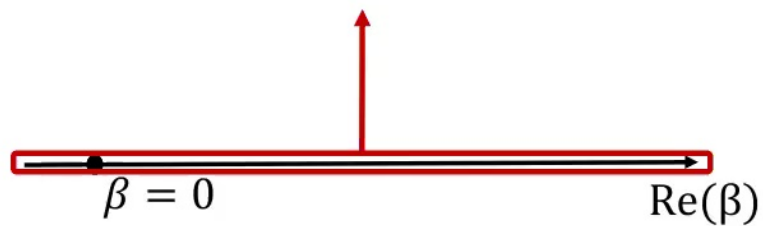
Sum of **strictly positive** terms



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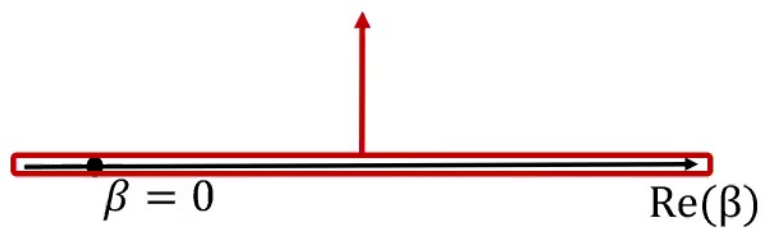


## Complex zeros vs phase transition

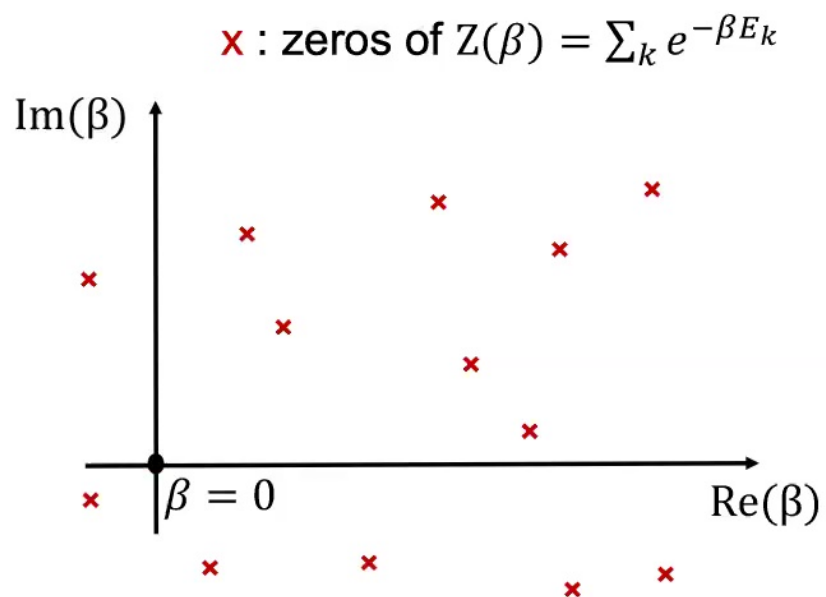
$$Z(\beta) = \sum_k e^{-\beta E_k}$$

Sum of **strictly positive** terms

No **singularities** for free energy



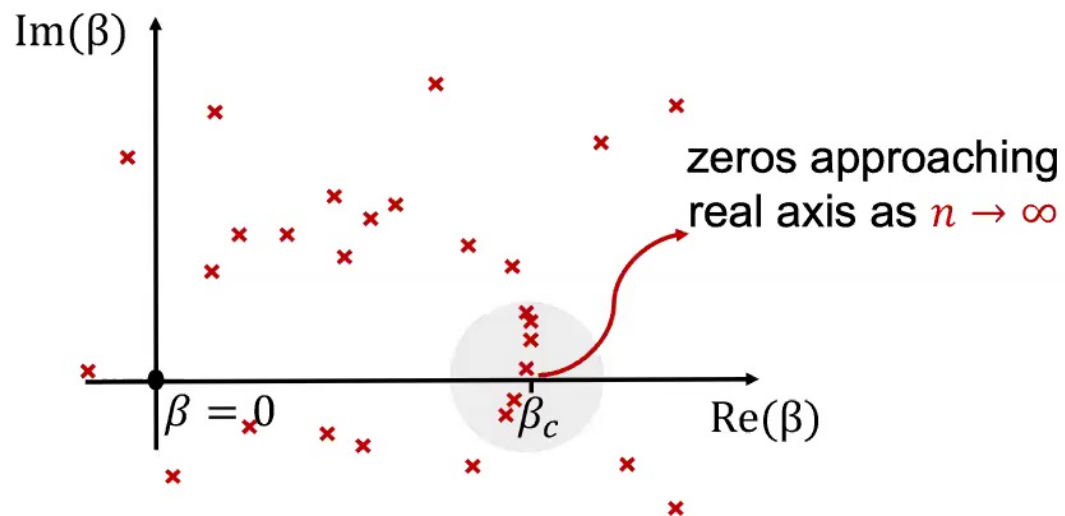
## Complex zeros vs phase transition



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## Complex zeros vs phase transition

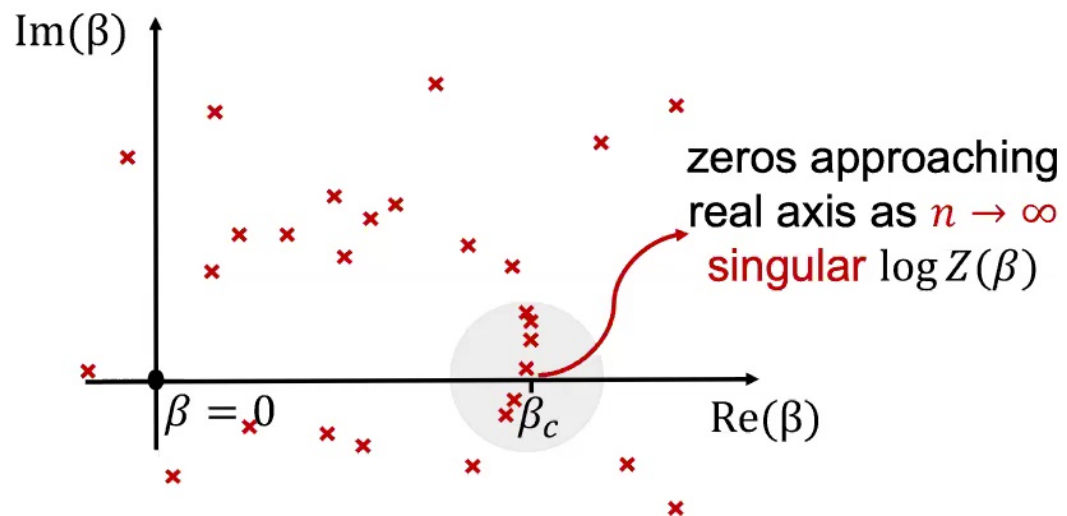
x : zeros of  $Z(\beta) = \sum_k e^{-\beta E_k}$



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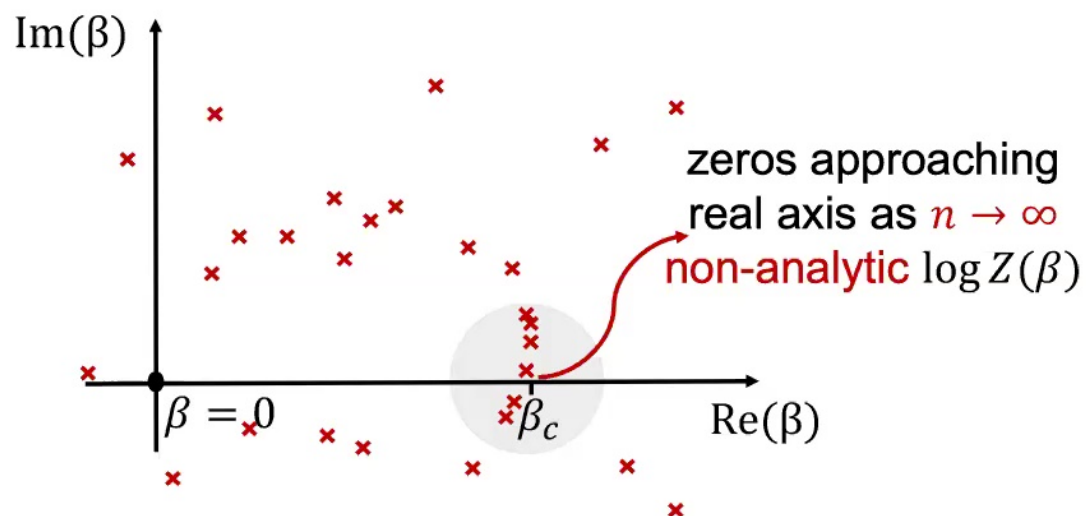


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## Complex zeros vs phase transition

Location of **complex zero of  $Z(\beta)$**   
tells us about  
**phase transition point**

$$x : \text{zeros of } Z(\beta) = \sum_k e^{-\beta E_k}$$



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- Studying zeros initiated by Lee and Yang [LY'52]

Phase transition in **classical Ising model** by  
changing the **external magnetic field**

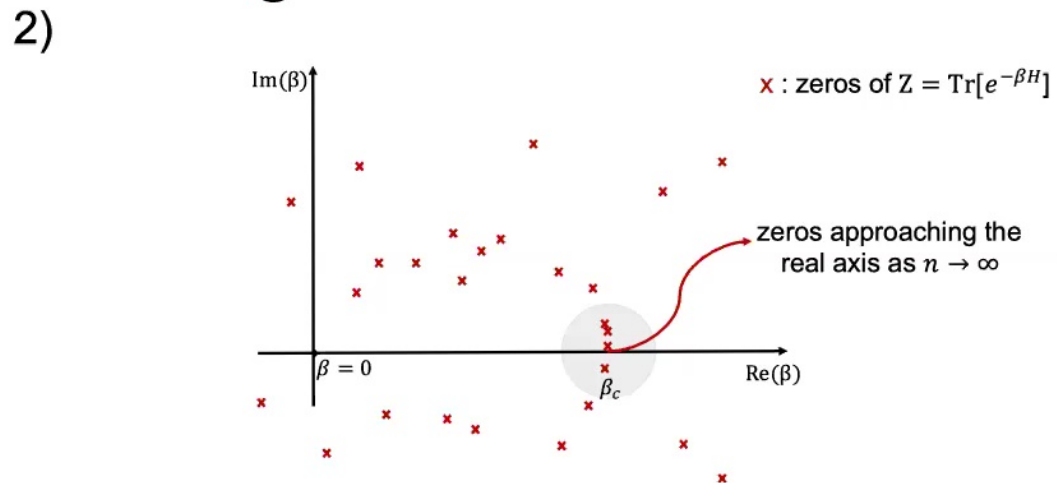
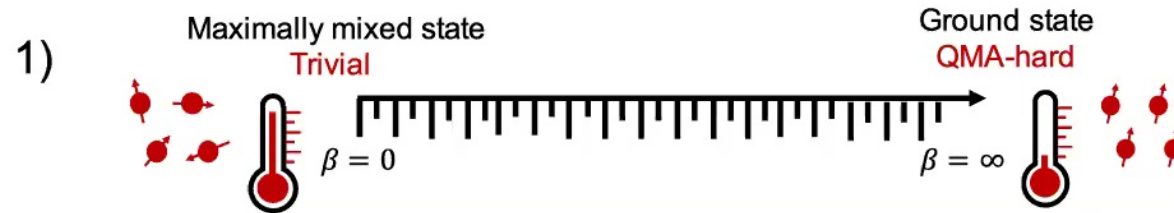
- Extended to **thermal** phase transition by Fisher [F'65]

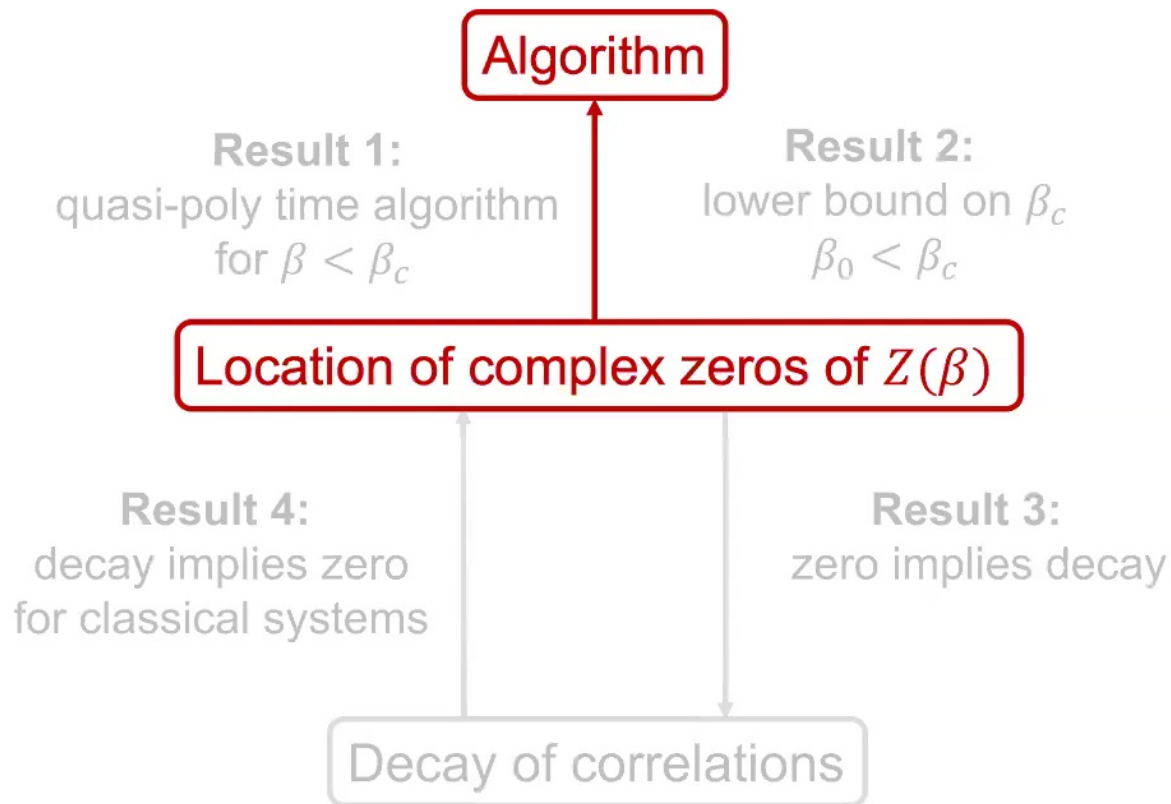




## To Recap

We made two observations





**Goal:** design an algorithm that

**Input:**  $H, \beta, \varepsilon$

**Output:** estimate  $\hat{Z}(\beta)$  such that

$$|\log Z(\beta) - \log \hat{Z}(\beta)| \leq \varepsilon$$



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**Result 1:**

*An approximation algorithm for  $Z(\beta)$  with running time  $n^{O(\log(n/\varepsilon))}$  that works **above the phase transition point***



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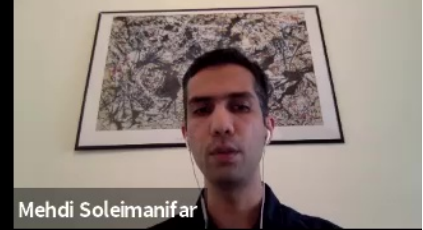


## Analysis:

The idea is to

**Extrapolate** the solution in the **easy regime** (high T)

to find the solution in the **harder regime** (lower T)



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Taylor expanding  $\log Z(\beta)$



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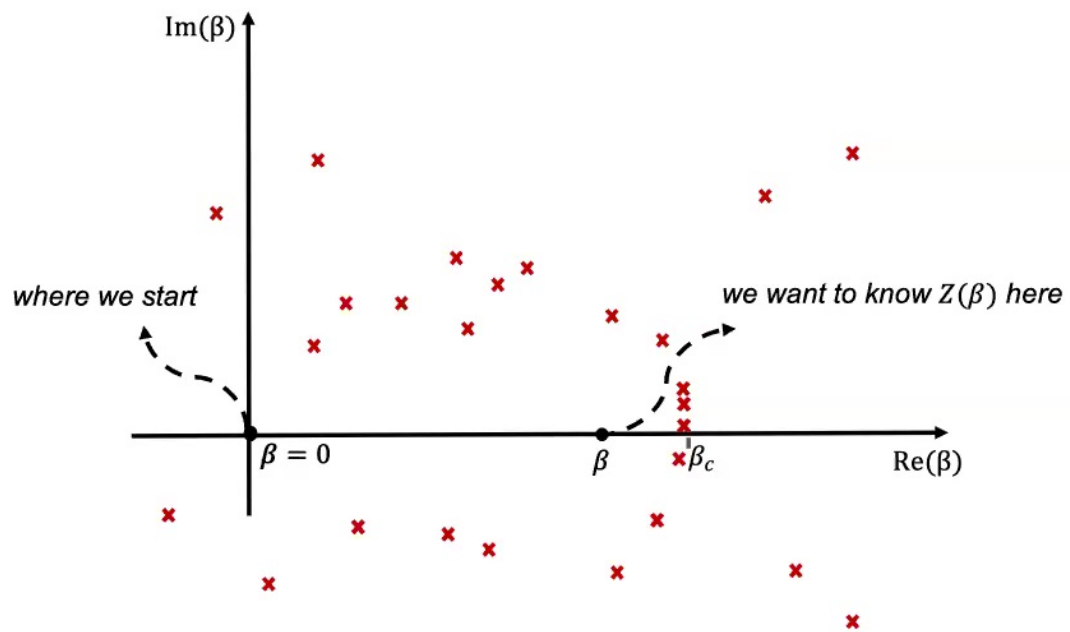
*Need to compute the low order derivatives*

*Need to make sure Taylor expansion converges*



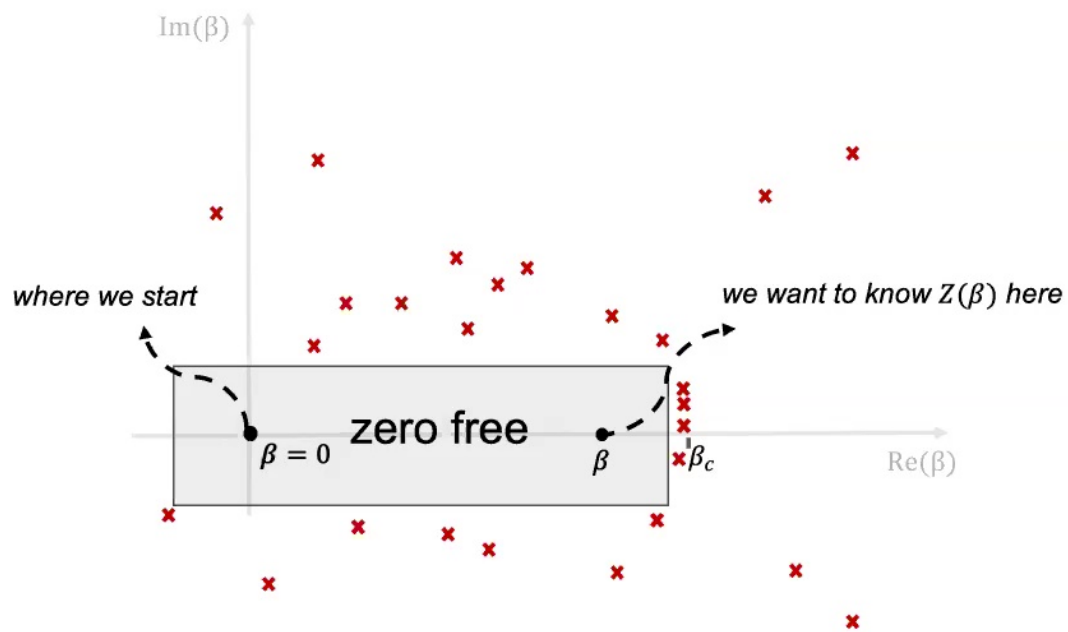


# Extrapolating $\log Z(\beta)$



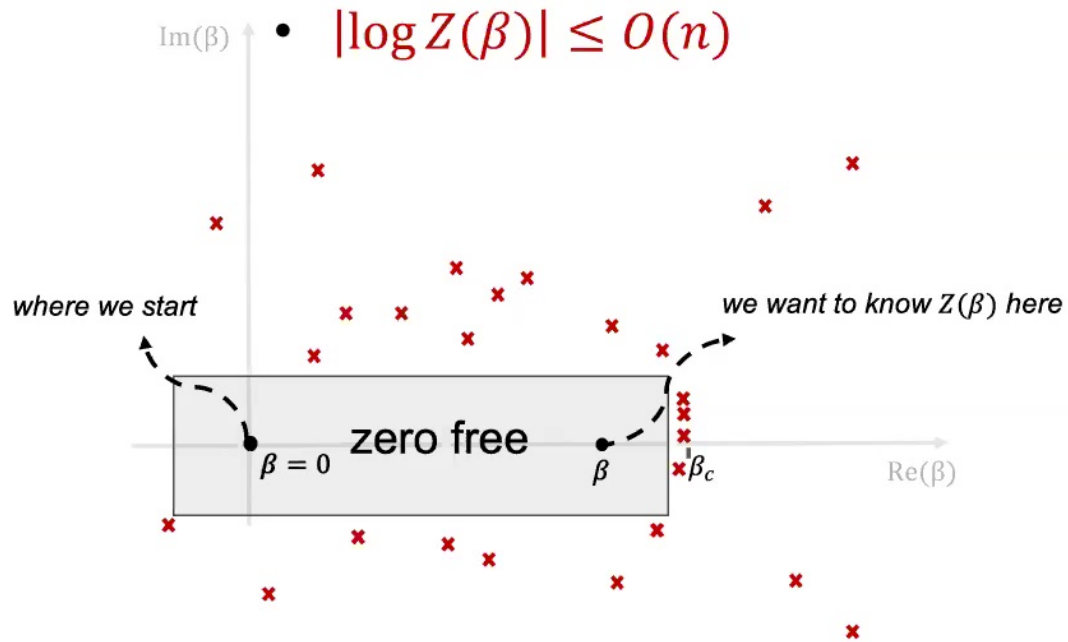
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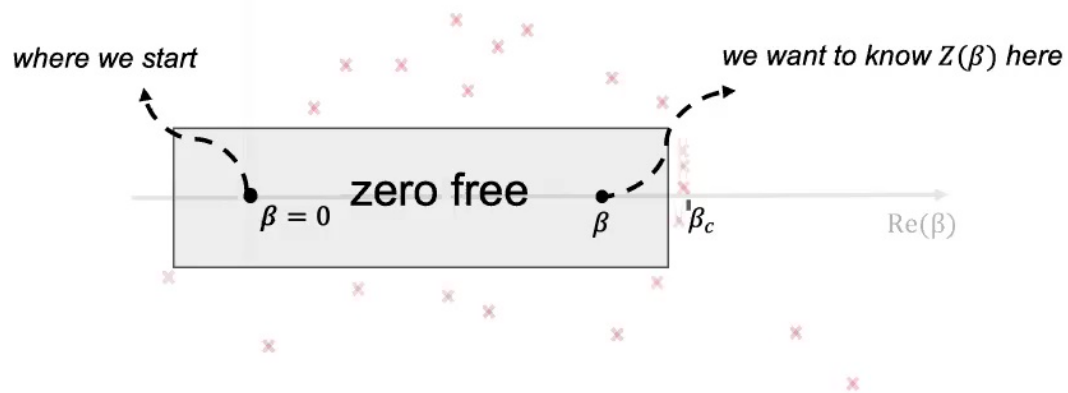


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## Extrapolating $\log Z(\beta)$

- $\log Z(\beta)$  is **analytic** in zero free region
  - $|\log Z(\beta)| \leq O(n)$

$$\left| \log Z(\beta) - \sum_{\ell=0}^K \frac{1}{\ell!} \frac{d^\ell \log Z(0)}{d^\ell \beta} \beta^\ell \right| \leq O(n) e^{-\alpha K}$$



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↓

$1/\text{poly}(n)$



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↓  
 $1/\text{poly}(n)$

$K = O(\log n)$  derivatives needed



How to compute  $O(\log n)$  derivatives of  $\log Z(\beta)$  ?

Sufficient to find  $O(\log n)$  derivative of  $Z(\beta)$

$$\frac{d^k Z(0)}{d^k \beta} \propto \text{Tr}[H^k]$$

$H$  is sum of  $\text{poly}(n)$  many local terms

$$H = \sum_{i=1}^{\text{poly}(n)} H_i$$

So takes **time**  $n^{O(K)}$  to find  $\text{Tr}[H^K]$



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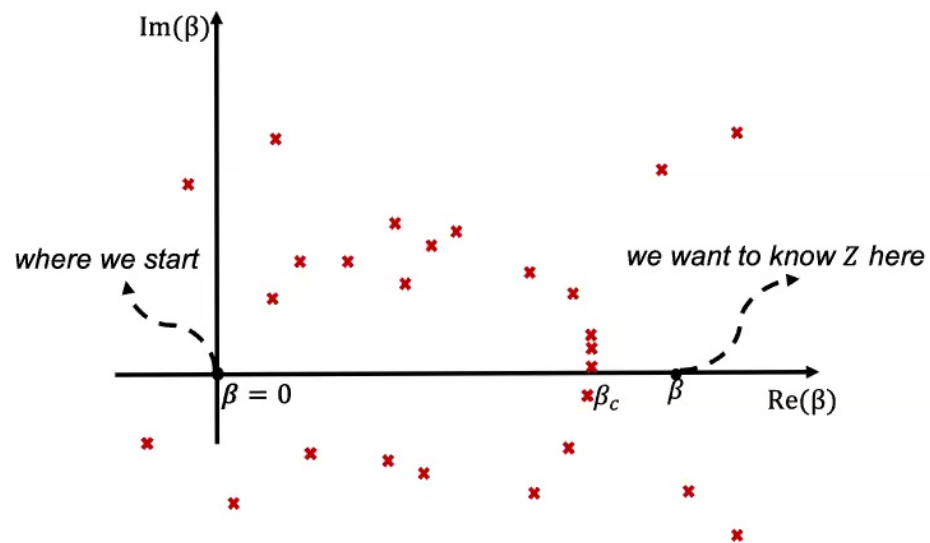
$$H = \sum_{i=1}^{\text{poly}(n)} H_i$$

So takes **time**  $n^{O(\log n)}$  to find all the derivatives

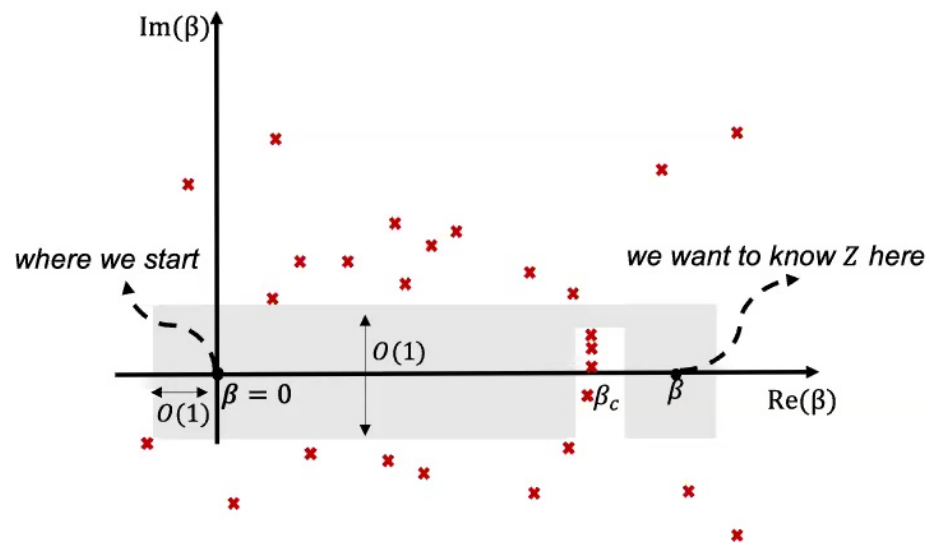




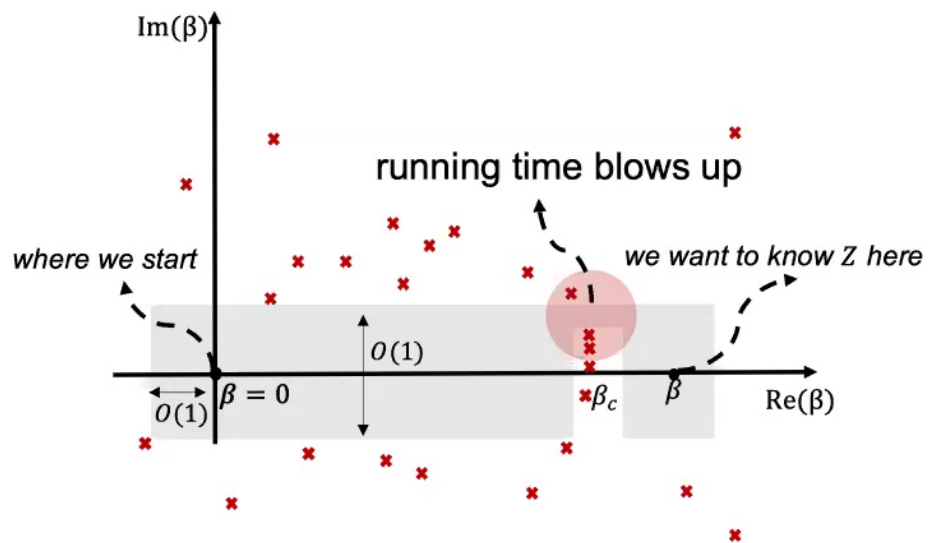
Why not go **beyond** the **phase transition point**?



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Why not go **beyond** the **phase transition point**?



## Previous work

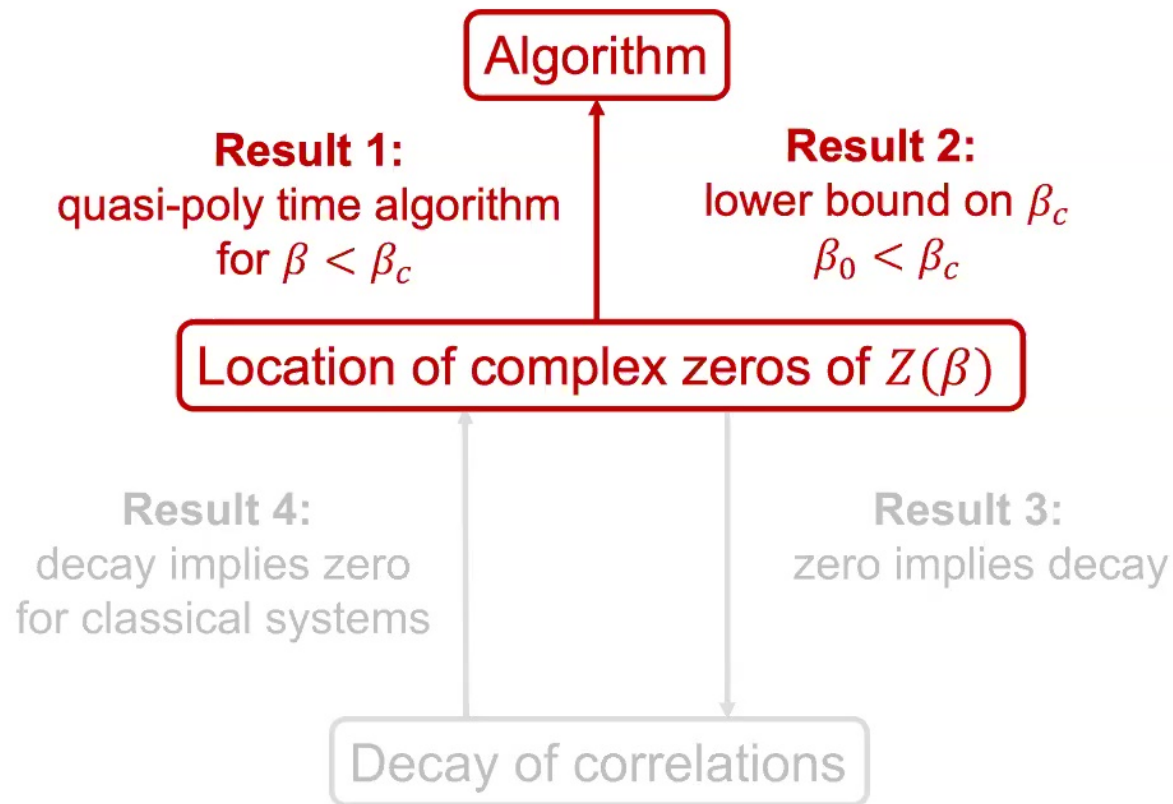
- Introduced by Barvinok to compute permanent of matrices [Bar'16]
- Used in many new algorithms for old **counting** problems [LSS'18, PR'16, EM'18,...]
- Running times are usually **quasi-polynomial**



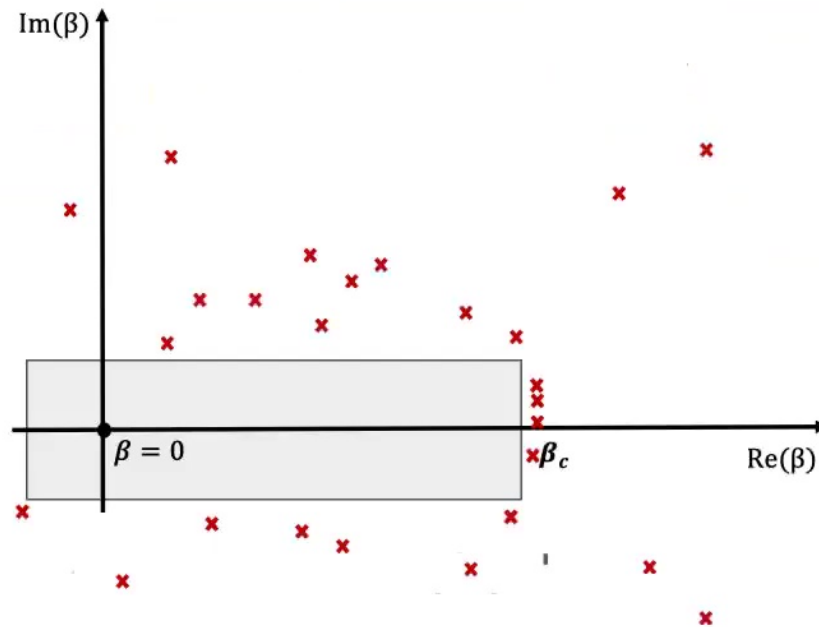
## Previous work

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- Used in many new algorithms for old **counting** problems [LSS'18, PR'16, EM'18,...]
- Running times are usually **quasi-polynomial**
- More recent results [KKB'19, MH'20] improve our running time to **polynomial** time at high enough temperatures



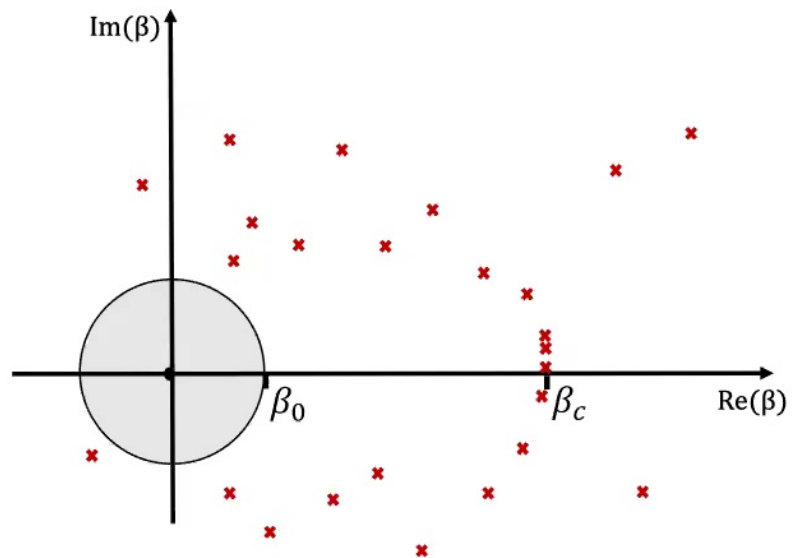


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## Result 2:

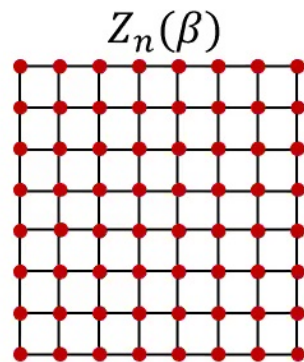
There is a *constant*  $\beta_0$  such that no zeros in *disk*  $|\beta| \leq \beta_0$





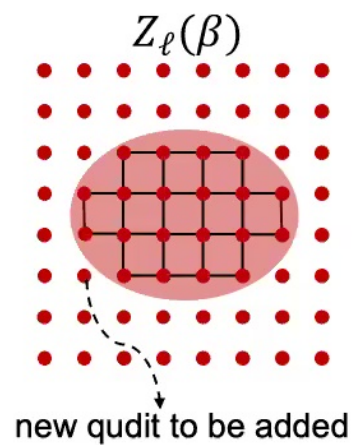
## Proof:

We lower bound  $Z(\beta)$  **recursively**  
using “**cluster expansions**” [Hastings'05, KGK+'14]



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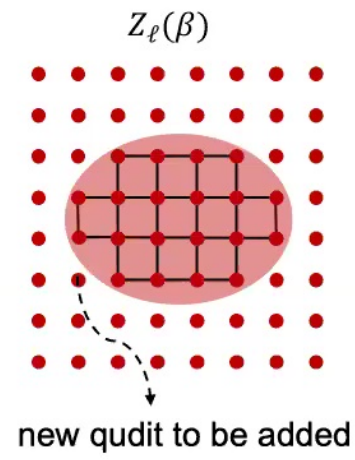
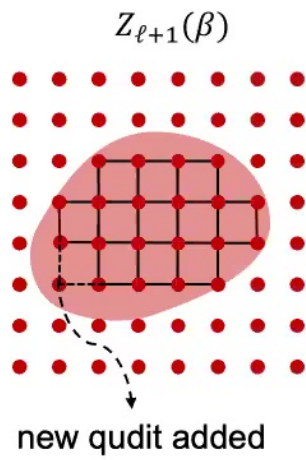


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**Proof:**

First show

$$|Z_{\ell+1}(\beta)| \geq c^{-1} |Z_{\ell}(\beta)|$$



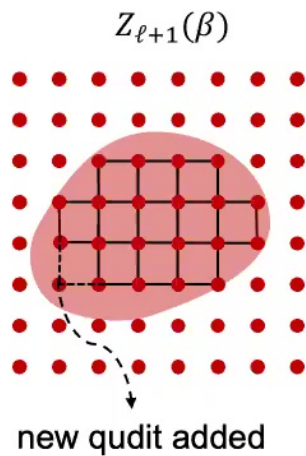
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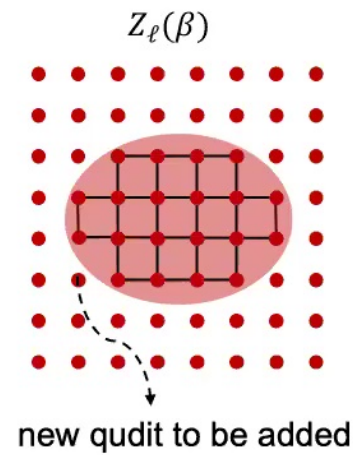
$$|Z_{\ell+1}(\beta)| \geq c^{-1} |Z_{\ell}(\beta)|$$

Repeat to get

$$|Z_n(\beta)| \geq c^{-n} |Z_1(\beta)|$$



$$|\log Z_n(\beta)| \leq O(n)$$

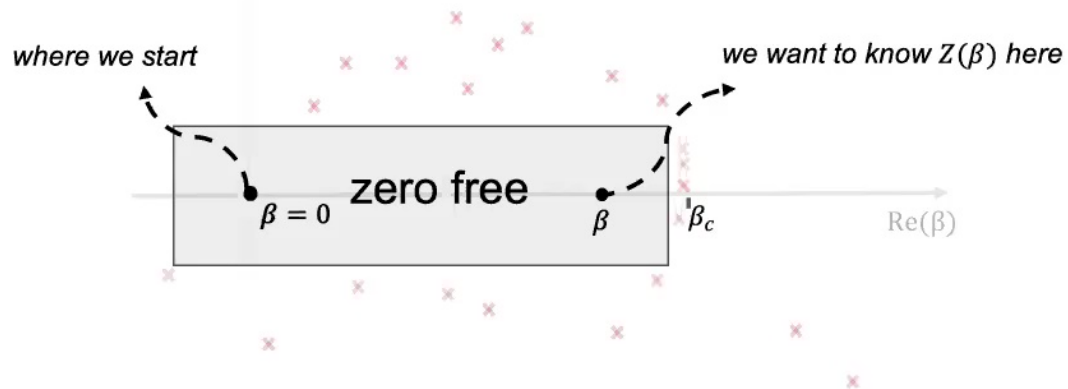


## Extrapolating $\log Z(\beta)$

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**Proof:**

How to show?

$$|Z_{\ell+1}(\beta)| \geq c^{-1} |Z_{\ell}(\beta)|$$

Cluster expansion: for  $|\beta| \leq \beta_0$ ,

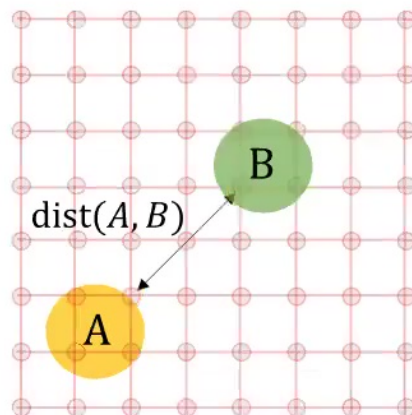
$$e^{-\beta H} \approx \sum \text{product of } H_i \text{'s}$$

$$Z_{\ell+1}(\beta) = Z_{\ell}(\beta) + \text{corrections}$$

$$|\text{corrections}/Z_{\ell}(\beta)| \leq O(1)$$



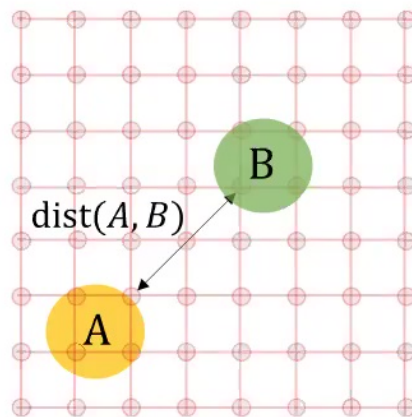
## Decay of correlations vs phase transition



$$| \text{Tr}[AB\rho] - \text{Tr}[A\rho]\text{Tr}[B\rho] | \leq c e^{-\text{dist}(A,B)/\xi}$$



## Decay of correlations vs phase transition



$$| \text{Tr}[AB\rho] - \text{Tr}[A\rho]\text{Tr}[B\rho] | \leq c e^{-\text{dist}(A,B)/\xi}$$



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## Decay of correlations vs phase transition

Decay of correlations is  
a signature of phase transition

- Above  $T_c$  exponential decay of correlations
- Below  $T_c$  there is long range order



# Decay of correlations vs phase transition

## *Algorithmic implications?*

- Classical spin systems [Weitz'99,...]



# Decay of correlations vs phase transition

*Algorithmic implications?*

- Classical spin systems [Weitz'99,...]

Mixing in time

=

Mixing in space



efficient sampling algorithm



exponential decay of correlations

- General Hamiltonians [BK'16]

Mixing in time

←

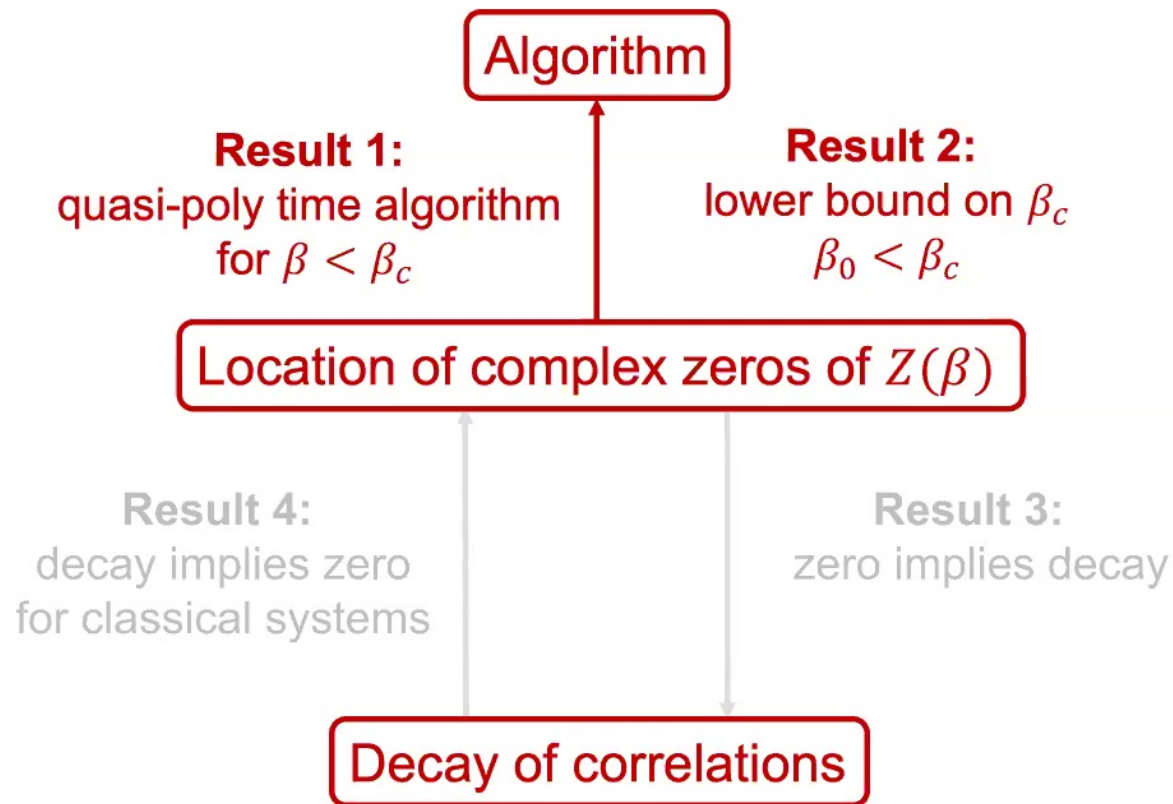
Mixing in space

+

Decay of quantum CMI  
"Markov property"



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*What is the relation between  
zeros of  $Z(\beta)$  and decay of correlations?*

For **translationally-invariant classical**  
system proved to be **equivalent** [DS'85]

*How about quantum systems?*



## Result 3

*We show absence of zeros near real axis implies exponential decay of correlations*

When

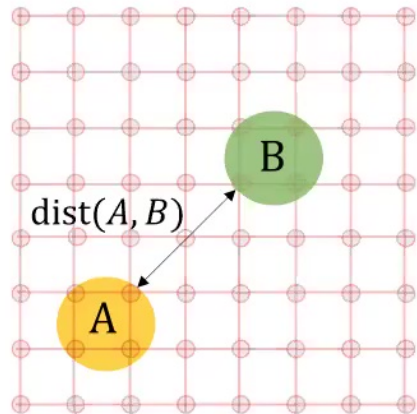
- For any quantum system if  $\text{dist}(A, B) = \Omega(\log n)$
- $H$  consists of **commuting** terms  $H = \sum_i H_i, [H_i, H_j] = 0$
- General  $H$  on a **one-dimensional chain**



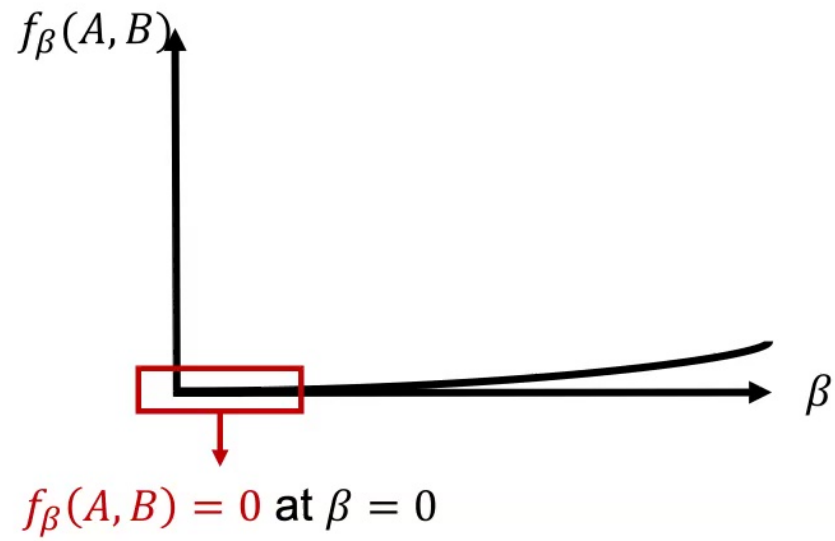
**Proof:** Similar to the extrapolation idea in our algorithm and the one used in [DS'85]

- Define a function that **measures correlations btw  $A, B$**

$$f_{\beta}(A, B)$$

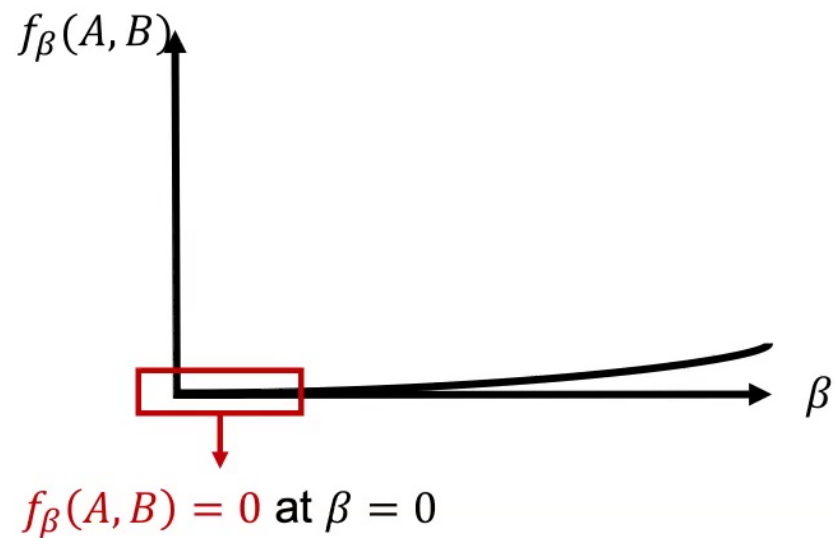


We show





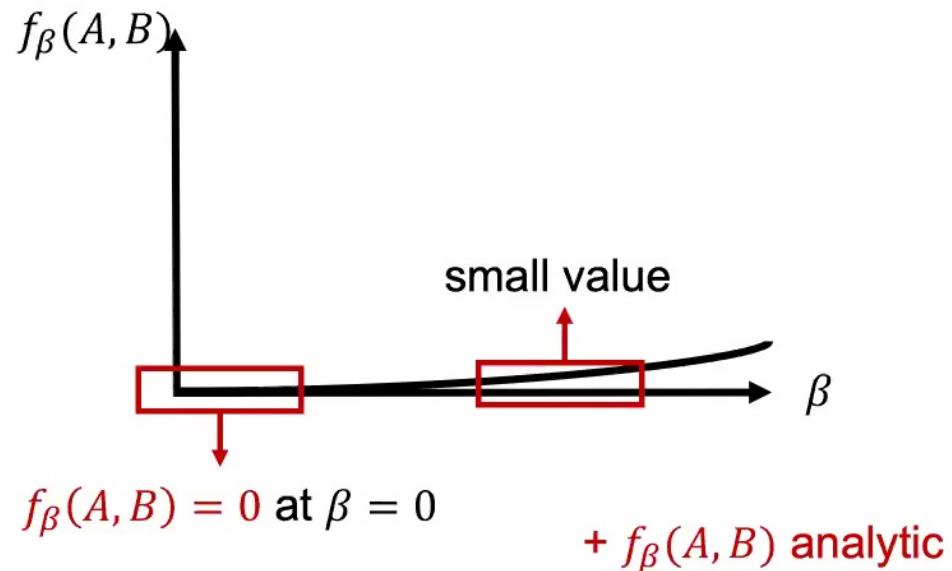
We show



$$\frac{d^m}{d\beta^m} f_{\beta}(A, B) = 0 \text{ for } m < O(\text{dist}(A, B))$$

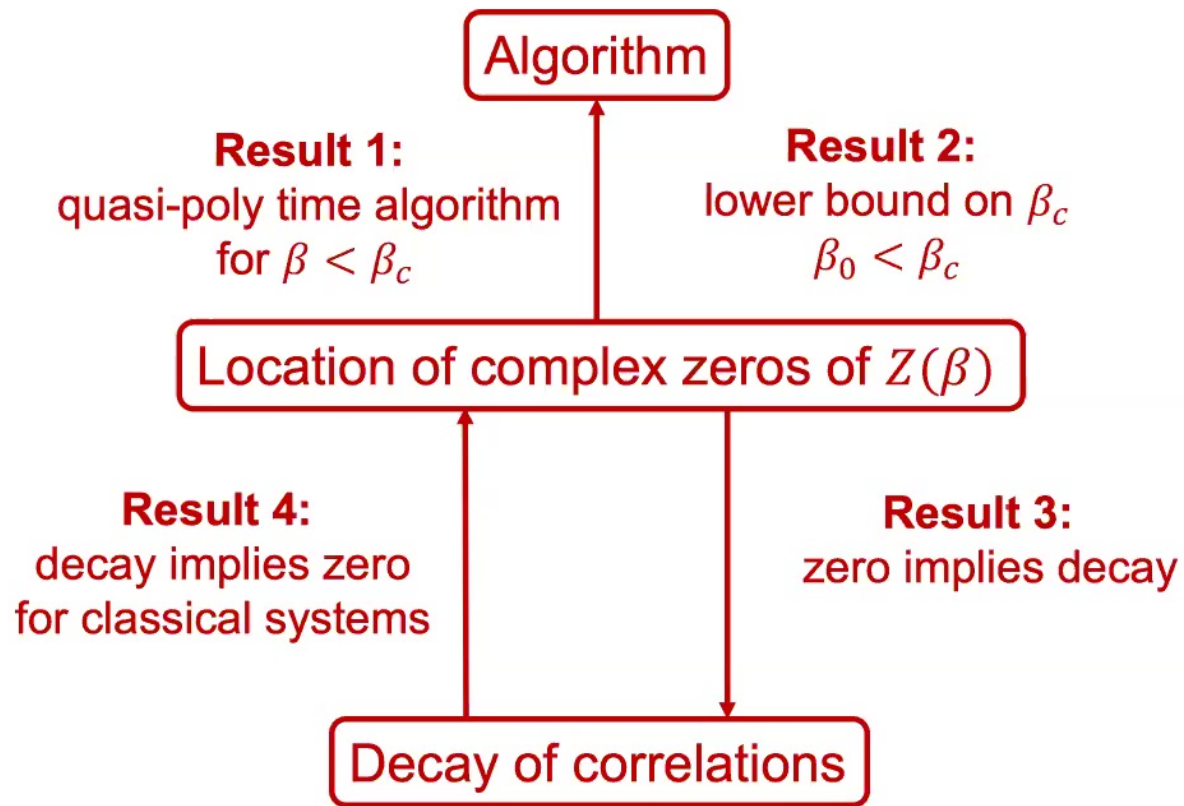


We show



$$\frac{d^m}{d\beta^m} f_{\beta}(A, B) = 0 \text{ for } m < O(\text{dist}(A, B))$$





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## Result 4:

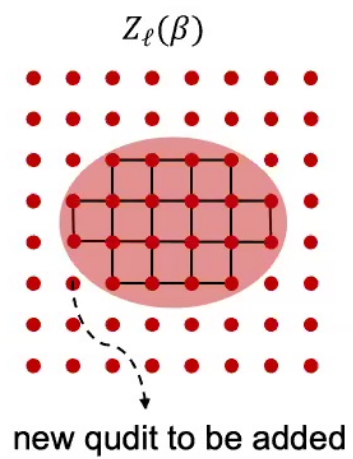
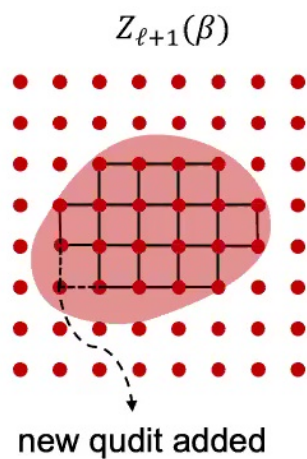
*Decay of correlations at real temperature  $\beta$*   
*implies no zeros close to real axis at  $\beta$*

- Proved for **translationally-invariant classical** systems [DS'85]
- We can extend their proof for **general classical** systems



## Rough high level idea

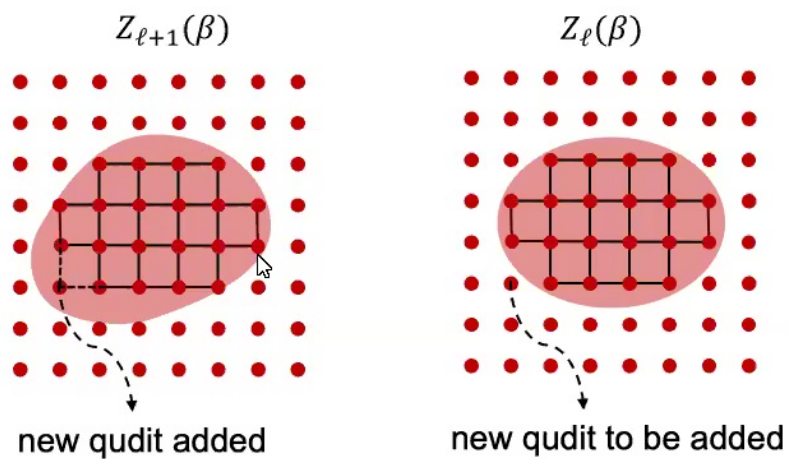
Similar to **Result 2**

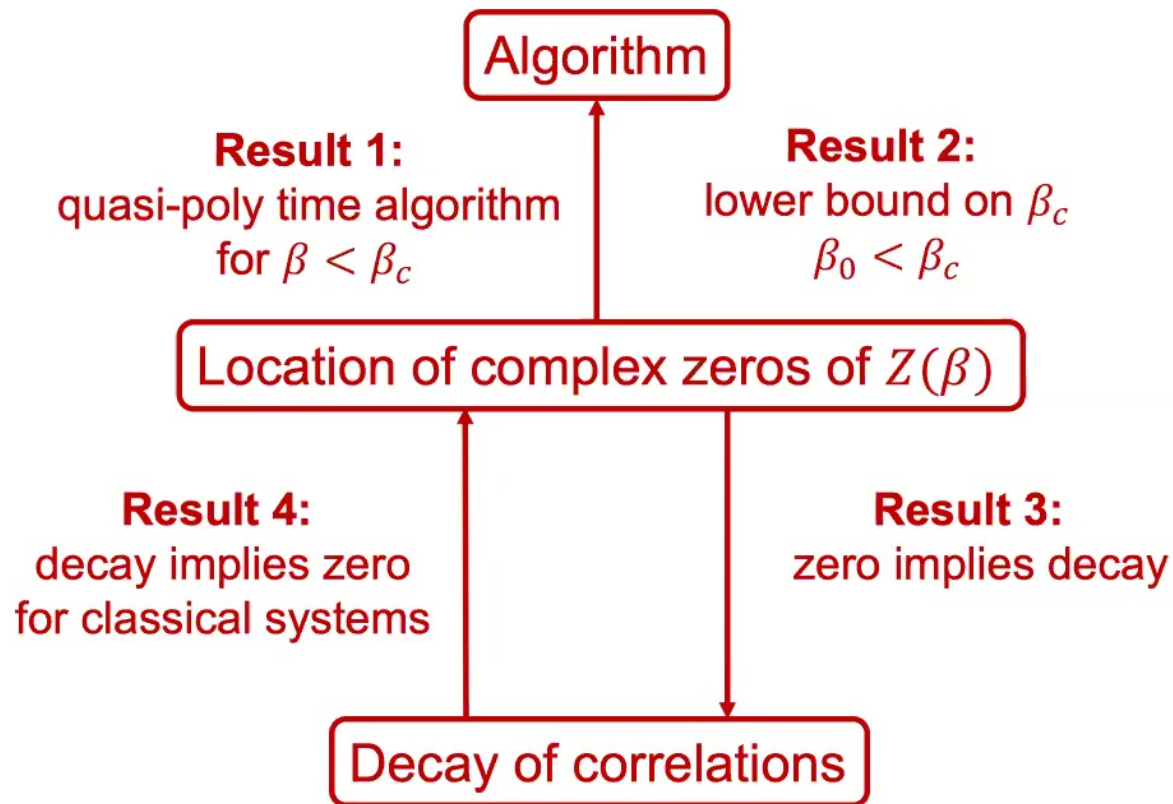


## Rough high level idea

Similar to **Result 2**

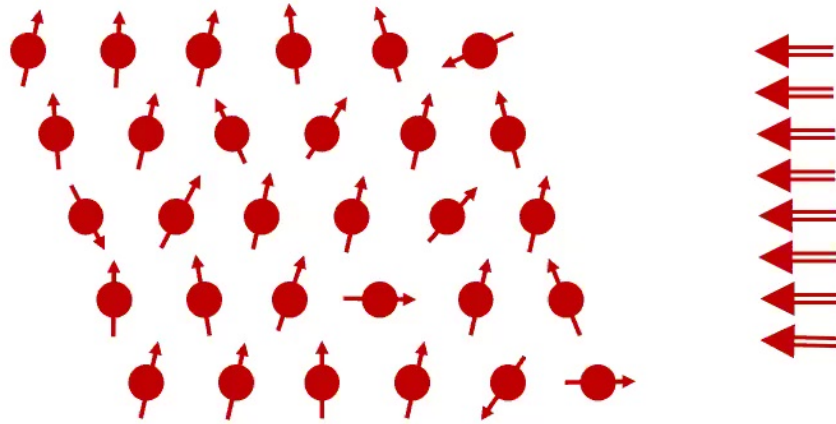
Instead of **cluster expansions** use **decay of correlations**





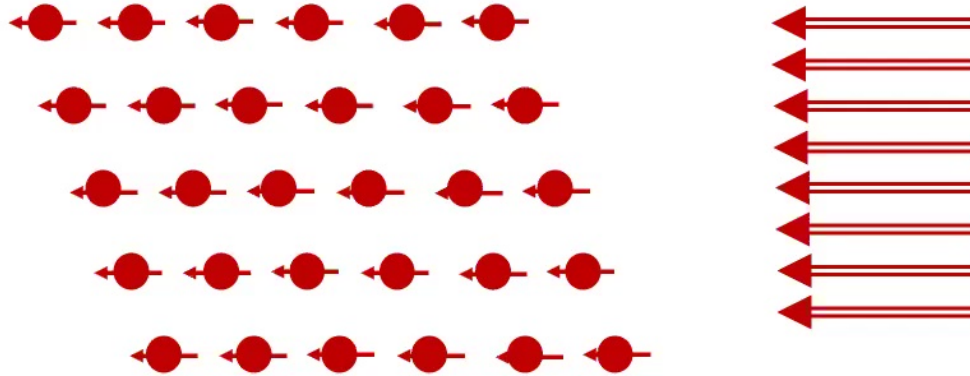
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## Extrapolating in external field





## Extrapolating in external field



Easy regime



## Result 5:

### Ferromagnetic Heisenberg model

$$H = -\sum_{ij} K_{ij} Z_i Z_j - \sum_{ij} J_{ij} (X_i X_j + Y_i Y_j) - \sum_i J_i Z_i$$

$$K_{ij} \geq |J_{ij}|$$

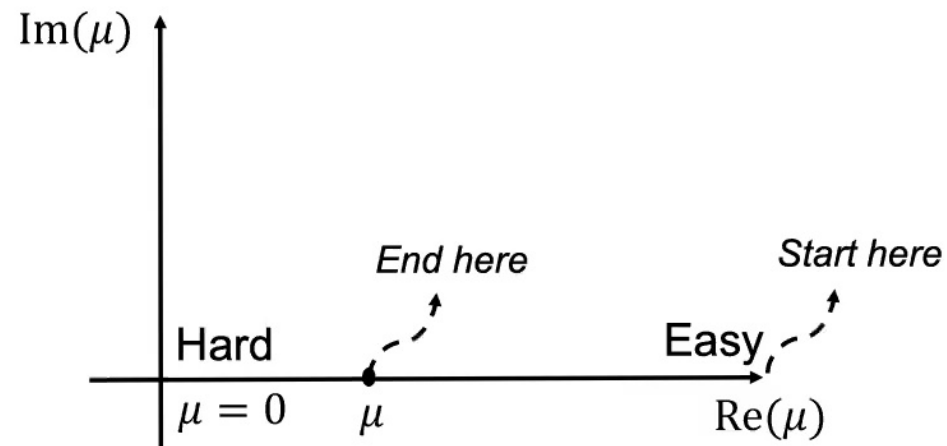


## Result 5:

### Ferromagnetic Heisenberg model

$$H = -\sum_{ij} K_{ij} Z_i Z_j - \sum_{ij} J_{ij} (X_i X_j + Y_i Y_j) - \mu \sum_i J_i Z_i$$

$$K_{ij} \geq |J_{ij}|$$

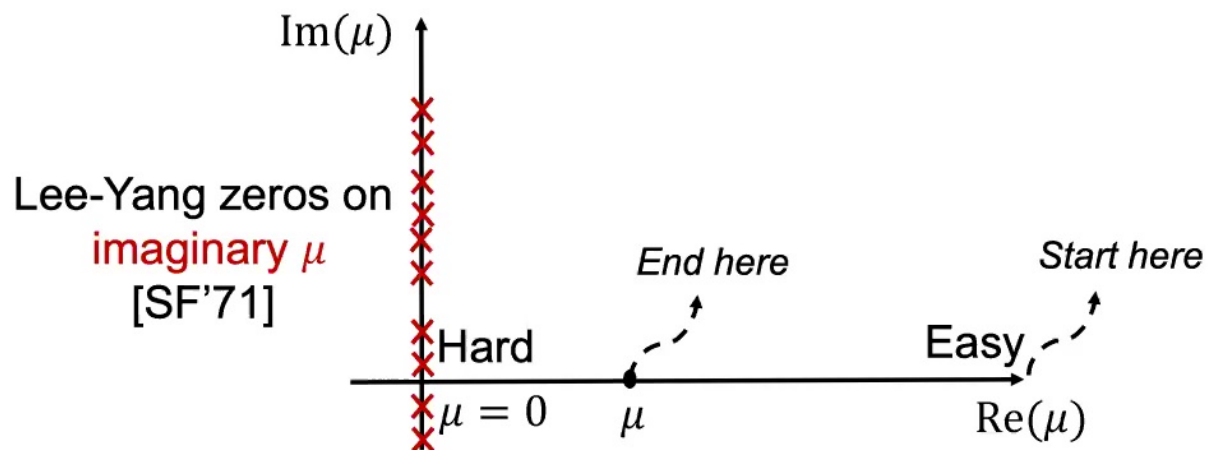


## Result 5:

### Ferromagnetic Heisenberg model

$$H = -\sum_{ij} K_{ij} Z_i Z_j - \sum_{ij} J_{ij} (X_i X_j + Y_i Y_j) - \mu \sum_i J_i Z_i$$

$$K_{ij} \geq |J_{ij}|$$



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Previously, polynomial time algorithm for  
**ferromagnetic XY model** [BG'17]

$$H = -\sum_{ij} b_{ij} X_i X_j - \sum_{ij} c_{ij} Y_i Y_j - \sum_i d_i Z_i$$

$$b_{ij} \geq |c_{ij}|$$

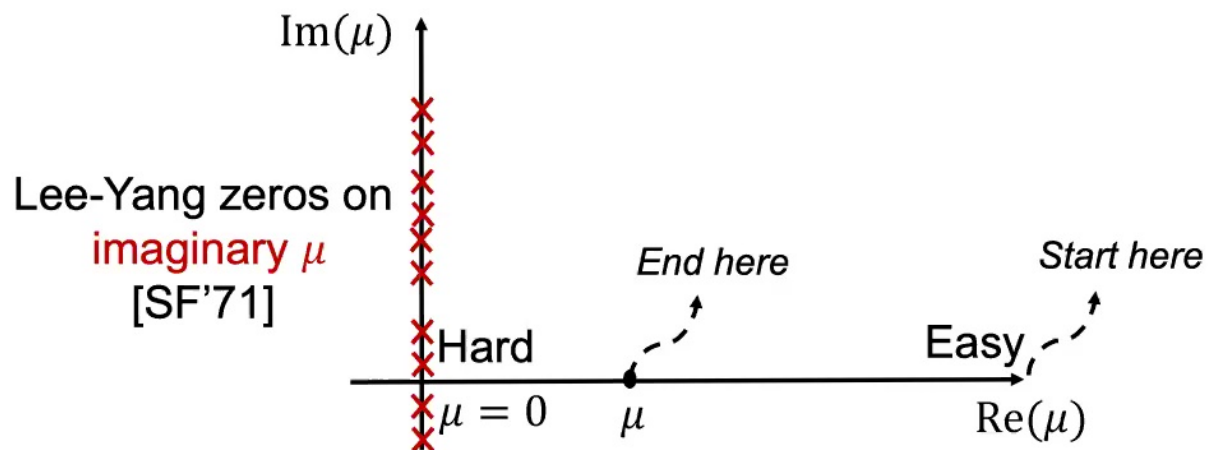
Solved by reducing to counting **perfect matchings**

## Result 5:

### Ferromagnetic Heisenberg model

$$H = -\sum_{ij} K_{ij} Z_i Z_j - \sum_{ij} J_{ij} (X_i X_j + Y_i Y_j) - \mu \sum_i J_i Z_i$$

$$K_{ij} \geq |J_{ij}|$$



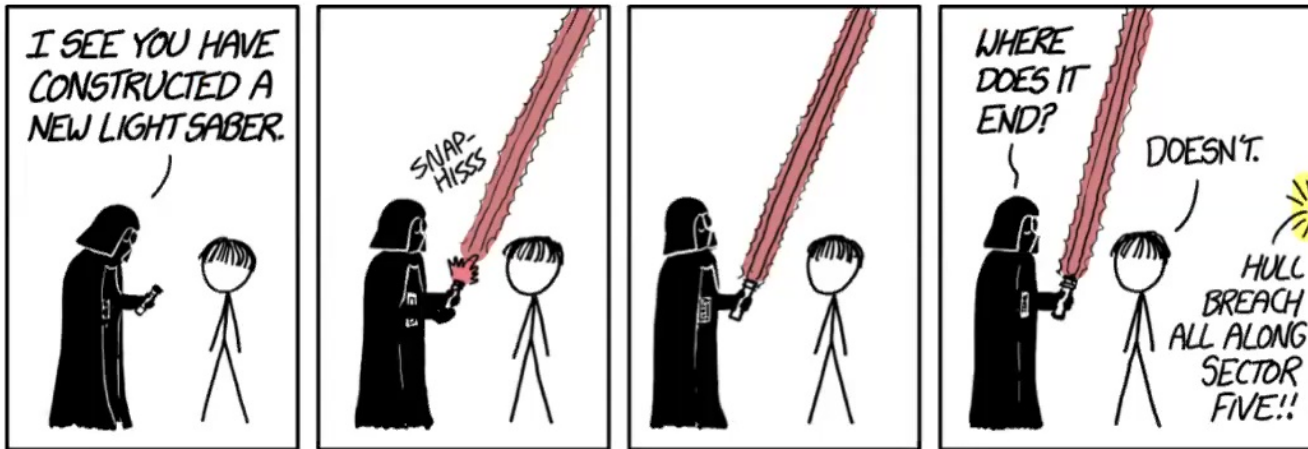
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## Open questions

- Fully establish the equivalence between absence of **complex zeros** and **decay of correlations**
- **Absence of zeros** implies **decay of quantum CMI**?
- A regime where quantum computer can't sample but extrapolation works?
- Other applications for extrapolation (avoiding sign problem, adiabatic algorithms,...)
- Other algorithms for  $Z(\beta)$  like convex relaxations?



Thanks!



xkcd.com