

Title: The Cosmological Evolution of Light Dark Photon Dark Matter

Speakers: Samuel McDermott

Series: Particle Physics

Date: April 28, 2020 - 2:00 PM

URL: <http://pirsa.org/20040089>

Abstract: Light dark photons are subject to various plasma effects, such as Debye screening and resonant oscillations, which can lead to a more complex cosmological evolution than is experienced by conventional cold dark matter candidates. Maintaining a consistent history of dark photon dark matter requires ensuring that the super-thermal abundance present in the early Universe (i) does not deviate significantly after the formation of the CMB, and (ii) does not excessively leak into the Standard Model plasma after BBN. In this talk, I will clarify the roles of resonant and non-resonant absorption and address some implications of plasma inhomogeneities on resonant transitions.

DARK PHOTONS, COSMOLOGICALLY

PERIMETER INSTITUTE

APRIL 28, 2020

SAMUEL D. MCDERMOTT

PRD & 1911.05086 W/ SAM WITTE
2003.13698 W/ S. WITTE, S. ROSAURO-ALCARAZ, AND V. POULIN

OUTLINE

- I. **What?** $\mathcal{L} = \mathcal{L}_{\text{SM}} - \frac{1}{2}m_{A'}^2 A'_\mu A'^\mu - \frac{1}{4}F'_{\mu\nu} F'^{\mu\nu} + \frac{\epsilon}{2}F'_{\mu\nu} F^{\mu\nu}$
- II. **Why now?** better genesis mechanisms: starting ca 2015
continuing throughout 2019
- III. **What's new?** ~5 orders of magnitude improvement in
constraints over ~5 decades in dark matter mass +
signals for future observatories



DARK MATTER CANDIDATES

THERMAL

- ▶ was once in thermal equilibrium with Standard Model
- ▶ need to get energy & entropy *out of* the dark matter sector
- ▶ “freeze out” paradigm

NONTHERMAL

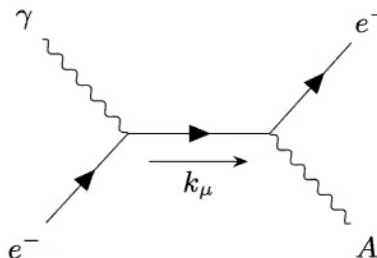
- ▶ never thermally equilibrated
- ▶ need to get (just the right amount of) energy *into* the dark matter sector
- ▶ misalignment, gravitational, ...

“FREEZING IN” DARK PHOTONS AS DARK MATTER?

- ▶ needs to be stable, but if $m_{A'} \gg m_e$ then

$$\Gamma_{A' \rightarrow e^+ e^-} \simeq \epsilon^2 \alpha_{\text{EM}} m_{A'} / 3 \simeq 10^{40} \epsilon^2 \left(\frac{m_{A'}}{\text{GeV}} \right) \tau_U^{-1}$$

- ▶ whereas rate for “semi-Compton” goes like



$$\frac{\Gamma_{e\gamma \rightarrow eA'}}{H} \sim \frac{\epsilon^2 n_e \sigma_T}{T^2 / M_{\text{Pl}}} \sim \frac{\epsilon^2 \alpha_{\text{EM}}^2 M_{\text{Pl}}}{T} \sim 10^{19} \epsilon^2 \left(\frac{\text{GeV}}{T} \right)$$

MISALIGNING DARK PHOTON DARK MATTER?

- ▶ what about misalignment? axion field density goes like

$$\rho_a \simeq \rho_{\text{CDM}} \times \left(\frac{f_a}{10^{12} \text{GeV}} \right)^{7/6} \Theta_i^2 (1+z)^3$$



"just give it the
right kick in the
early universe"

MISALIGNING DARK PHOTON DARK MATTER?

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- ▶ however, vector modes redshift differently $\Rightarrow$
misalignment doesn't work for A'

cf. Arias et al. JCAP 1201.5902
Graham, Mardon, Rajendran PRD 1504.02102

NEW PRODUCTION CHANNELS

- ▶ inflationary production: $m_{A'} \gtrsim 10\mu\text{eV}$



Vector Dark Matter from Inflationary Fluctuations

1504.02102

Peter W. Graham,¹ Jeremy Mardon,^{1,*} and Surjeet Rajendran²

Stanford Institute for Theoretical Physics, Department of Physics, Stanford University, Stanford, CA 94305

*²Berkeley Center for Theoretical Physics, Department of Physics,
University of California, Berkeley, CA 94720*

NEW PRODUCTION CHANNELS

- ▶ misalignment followed by decay or preheating /
parametric resonance

Parametric Resonance Production of Ultralight Vector Dark Matter

1810.07195¹ Jeff A. Dror,^{1,2} Keisuke Harigaya,^{3,1,2} and Vijay Narayan¹

¹Berkeley Center for Theoretical Physics, University of California
²Theoretical Physics Group, Lawrence Berkeley
³School of Natural Sciences, Institute

1810.07188 Relic Abundance of Dark Photon Dark Matter

Prateek Agrawal,¹ Naoya Kitajima,² Matthew Reece,¹ Toyokazu Sekiguchi,³ and Fuminobu Takahashi¹
¹Jefferson Physical Laboratory, Harvard University, 17 Oxford Street, Cambridge, MA 02138, USA
²Department of Physics, Nagoya University, Chikusa, Nagoya 464-8602, Japan
³for the Early Universe (RESCEU),

Dark Photon Dark Matter from a Network of Cosmic Strings

Andrew J. Long
Leinweber Center for Theoretical Physics, University of Michigan, Ann Arbor

1901.03312

Lian-Tao Wang
Institute for Cosmological Physics and Enrico Fermi Institute, University of Chicago
(Dated: March 19, 2019)

We study the production of ultralight dark photons from a network of cosmic strings. We find that dark photons produced in this way are nonrelativistic for dark photon masses as small as $m_A \sim 10^{-22}$ eV.

Vector dark matter production at the end of inflation

1810.07208¹ Mar Bastero-Gil,^{1,*} Jose Santiago,^{1,*} Lorenzo Ubaldi,^{2,3,*} and Roberto Vega-Morales

¹Departamento de Física Teórica y del Cosmos and CAFPE, Universidad de Granada, Campus de Fuentenueva, E-18071 Granada, Spain
²SISSA International School for Advanced Studies, Via Bonomea 265, 34136 Trieste, Italy
³INFN Sezione di Trieste, Via Valerio 2, 34100 Trieste, Italy

Dark Photon Dark Matter Produced by Axion Oscillations

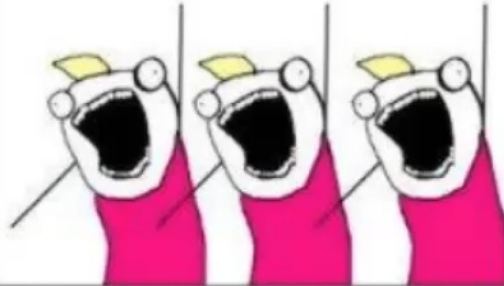
1810.07196¹ Raymond T. Co,¹ Aaron Pierce,¹ Zhengkang Zhang,^{1,2,3} and Yue Zhao^{1,4}

¹Leinweber Center for Theoretical Physics, Department of Physics, University of Michigan, Ann Arbor, MI 48109, USA
²Department of Physics, University of California, Berkeley, CA 94720, USA

WHAT DO WE HAVE?



DARK PHOTON DARK MATTER



**WHAT ARE WE
GOING TO DO WITH IT?**

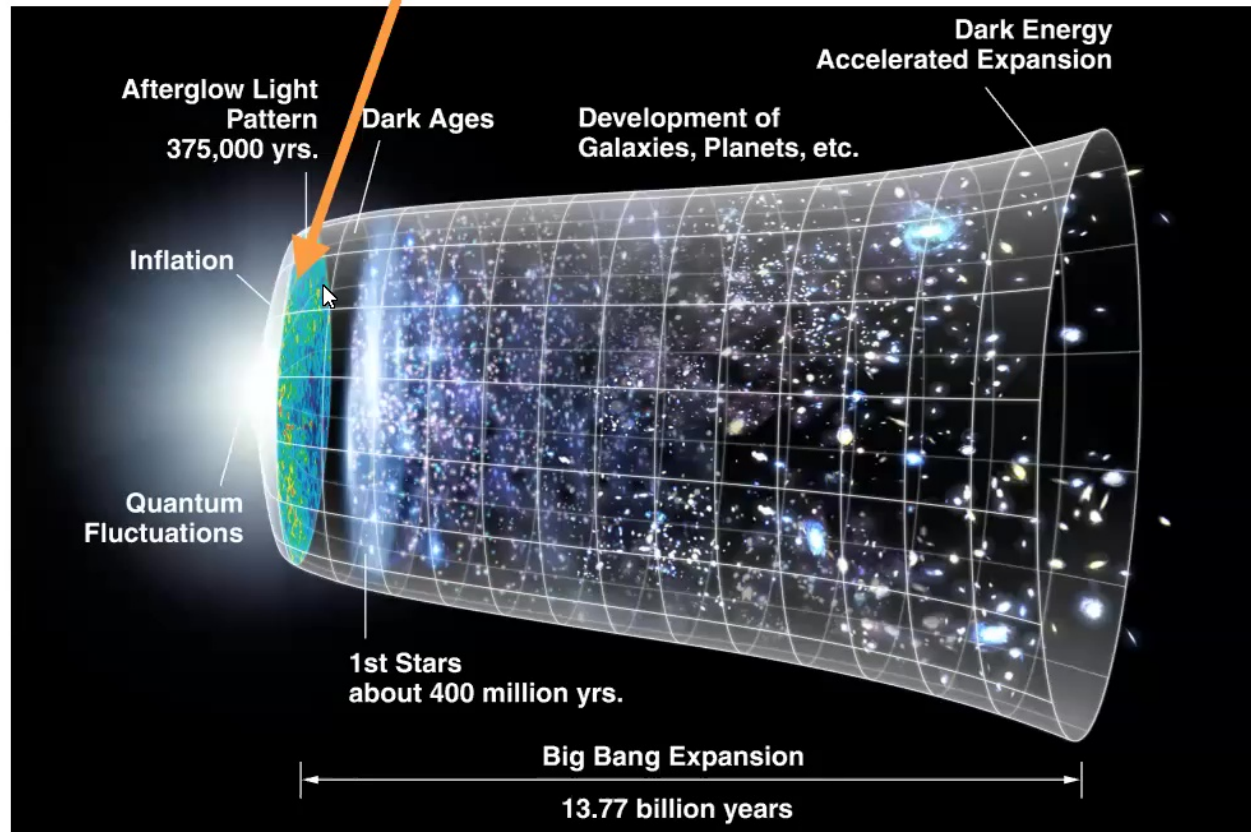


imgflip.com

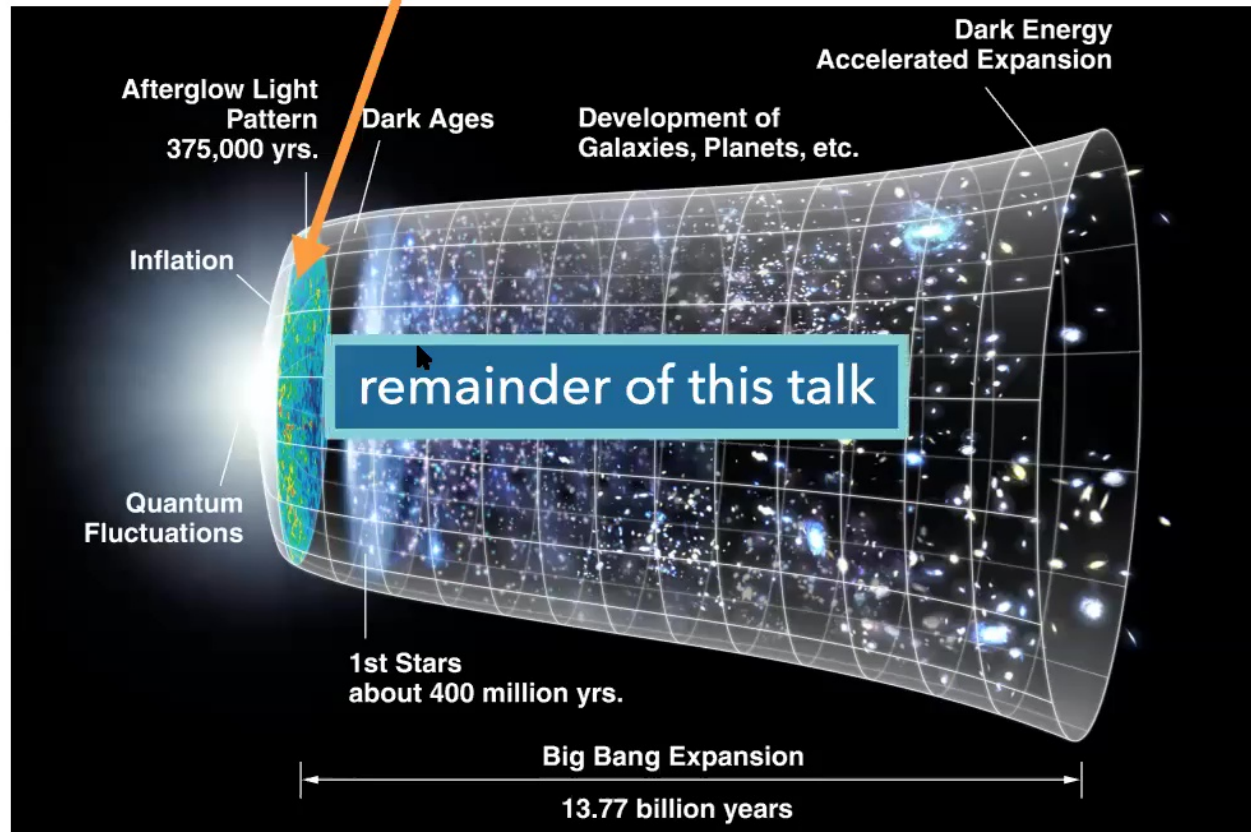
**THINK ABOUT HOW TO
CONSTRAIN AND/OR OBSERVE IT**



"initial conditions"



"initial conditions"



DARK PHOTONS AND REGULAR PHOTONS

unique interaction with Standard Model

$$\mathcal{L} \supset \epsilon F'_{\mu\nu} F^{\mu\nu} / 2 \iff \begin{array}{c} A' \quad \gamma \\ \text{wavy line} \times \text{wavy line} \end{array}$$

$\partial_\mu \partial^\mu \rightarrow K^2$

first Okun 1982 Holdom 1986

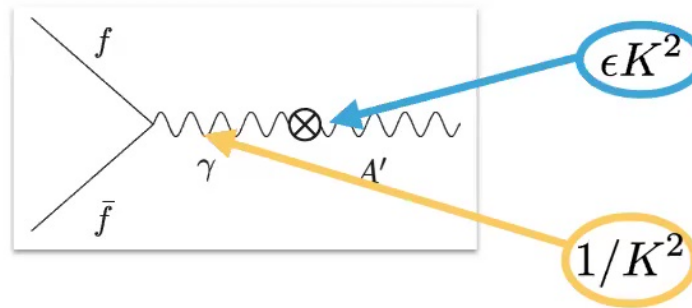
then Jaeckel, Redondo, Ringwald 0804.4157 Pospelov, Ritz, Voloshin 0807.3279

Redondo and Postma 0811.0326 Mirizzi, Redondo, Sigl 0901.0014

COUPLING TO DARK PHOTON

in vacuum:

$$\mathcal{L} \supset \epsilon J_{\mu}^{\text{SM}} A'^{\mu}$$



“PLASMAS GIVE PHOTON A MASS”

nonzero density of free electrons modifies
the SM photon dispersion relation:

$$\boxed{\omega^2 \simeq k^2 + \omega_p^2}^*$$

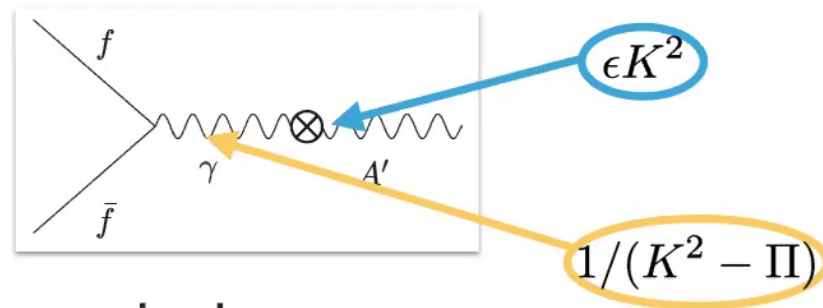
plasma mass: $\omega_p(m_e) \simeq \sqrt{4\pi\alpha_{\text{EM}}n_e/m_e}$

*this is not a real mass (in reality $\omega_p \rightarrow \Pi$, the SM photon self-energy,
which is momentum-dependent)

PHOTON MIXING IN PLASMA

in vacuum:

$$\mathcal{L} \supset \epsilon J_{\mu}^{\text{SM}} A'^{\mu}$$



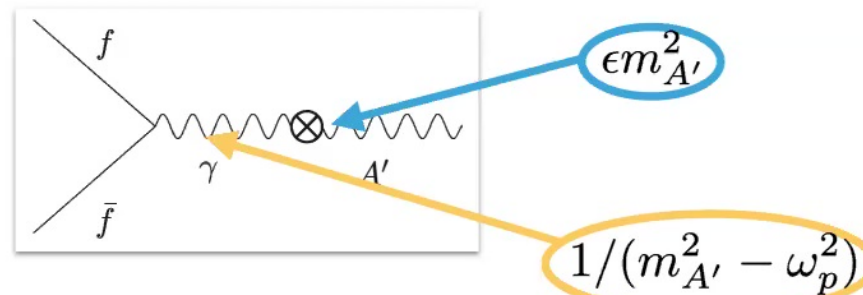
in plasma:

$$\mathcal{L} \supset \frac{\epsilon K^2}{K^2 - \Pi} J_{\mu}^{\text{SM}} A'^{\mu}$$

PHOTON MIXING IN PLASMA

in vacuum:

$$\mathcal{L} \supset \epsilon J_{\mu}^{\text{SM}} A'^{\mu}$$



in plasma, γ & A' on-shell, low-momentum exchange:

$$\mathcal{L} \supset \frac{\epsilon}{1 - \omega_p^2/m_{A'}^2} J_{\mu}^{\text{SM}} A'^{\mu}$$

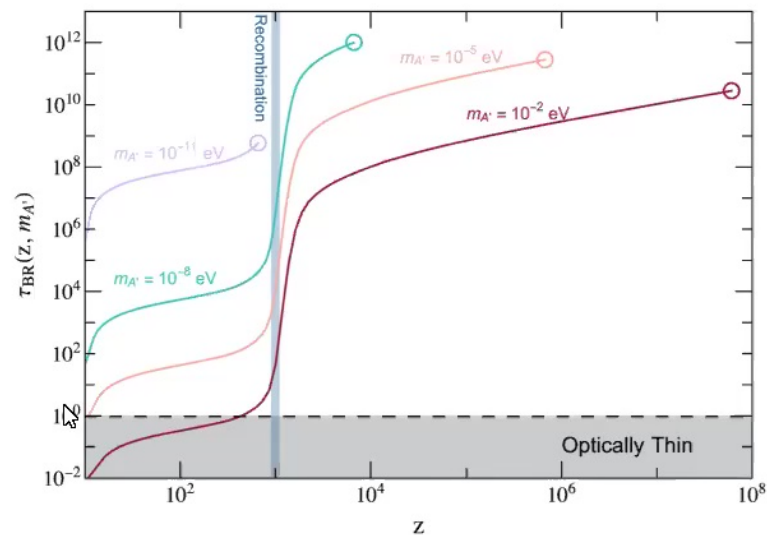
DARK PHOTON COSMOLOGY

- ▶ SM photon: $\omega_p = \sqrt{\frac{4\pi\alpha\eta_b x_e T_{\text{CMB}}^3}{m_e}} \simeq \text{neV} \left(\frac{1+z}{1+z_{\text{CMB}}} \right)^{3/2}$
 $\simeq \text{MHz} \left(\frac{1+z}{1+z_{\text{CMB}}} \right)^{3/2}$
- ▶ Plasma effects are of qualitative importance for cosmological evolution of dark photons below this mass
- ▶ Can experience a level crossing – populations determined by Landau-Zener-type formula

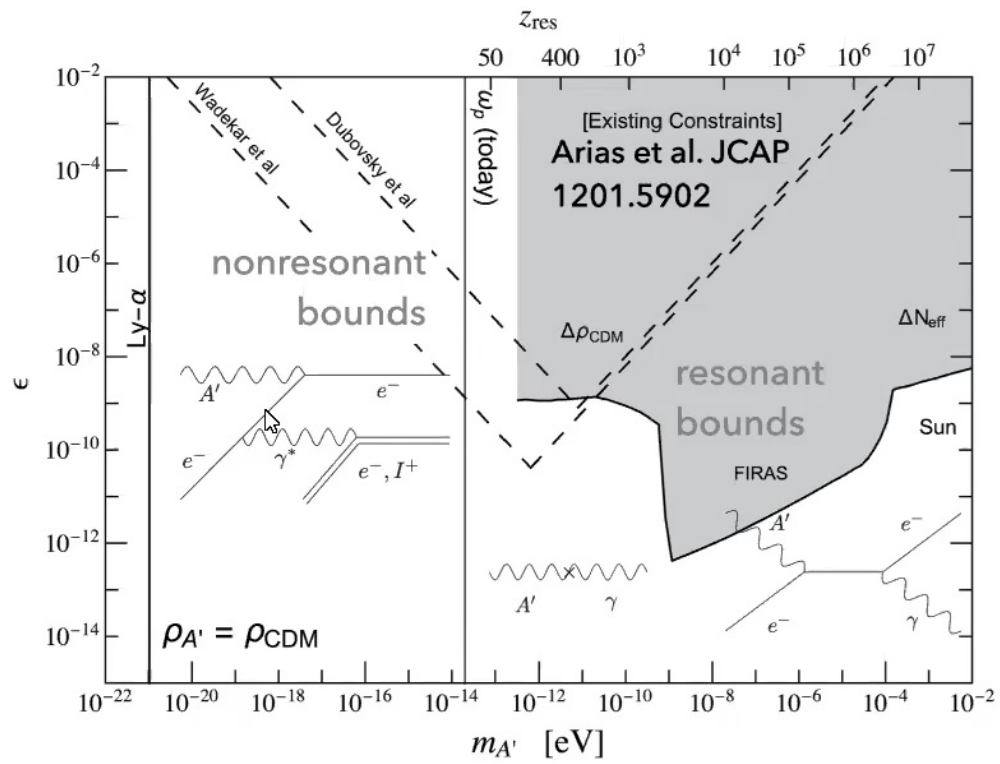
DARK PHOTON COSMOLOGY

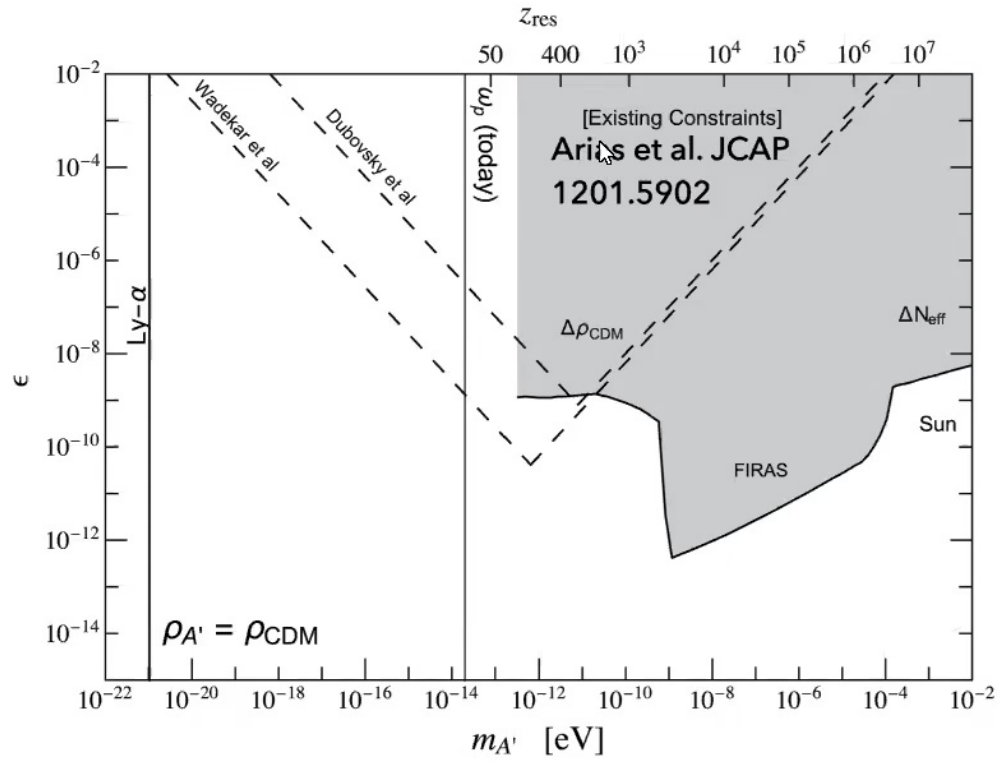
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DARK PHOTON COSMOLOGY

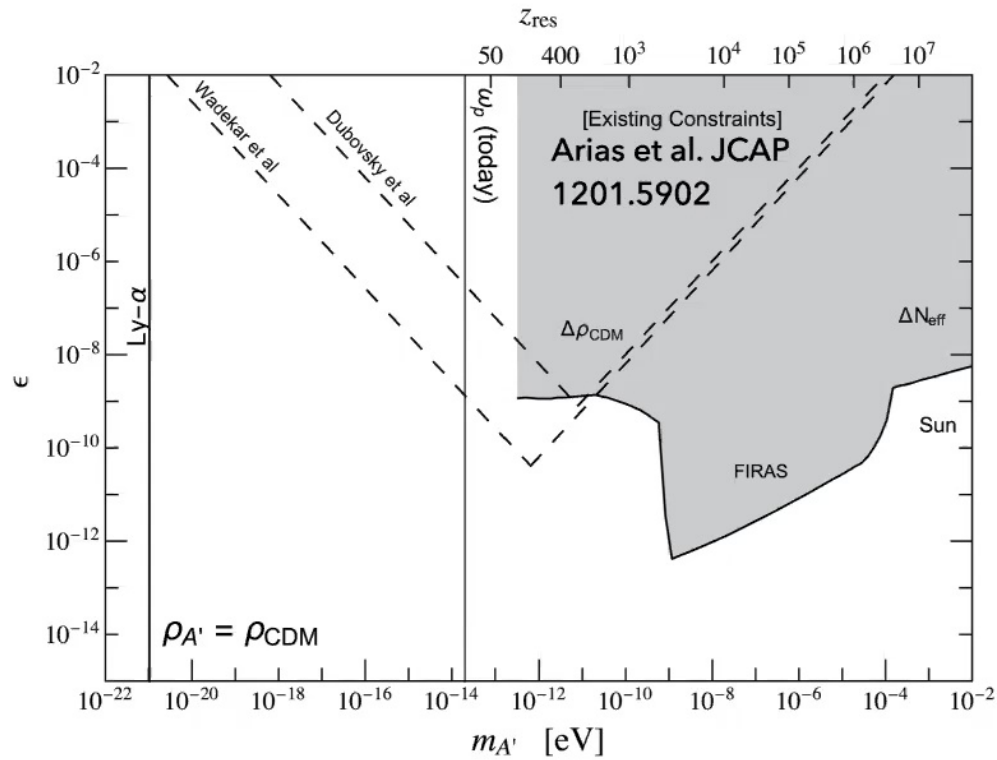


any dark photon in the $\sim \text{neV}$ mass range that converts to a photon "instantly" heats the medium

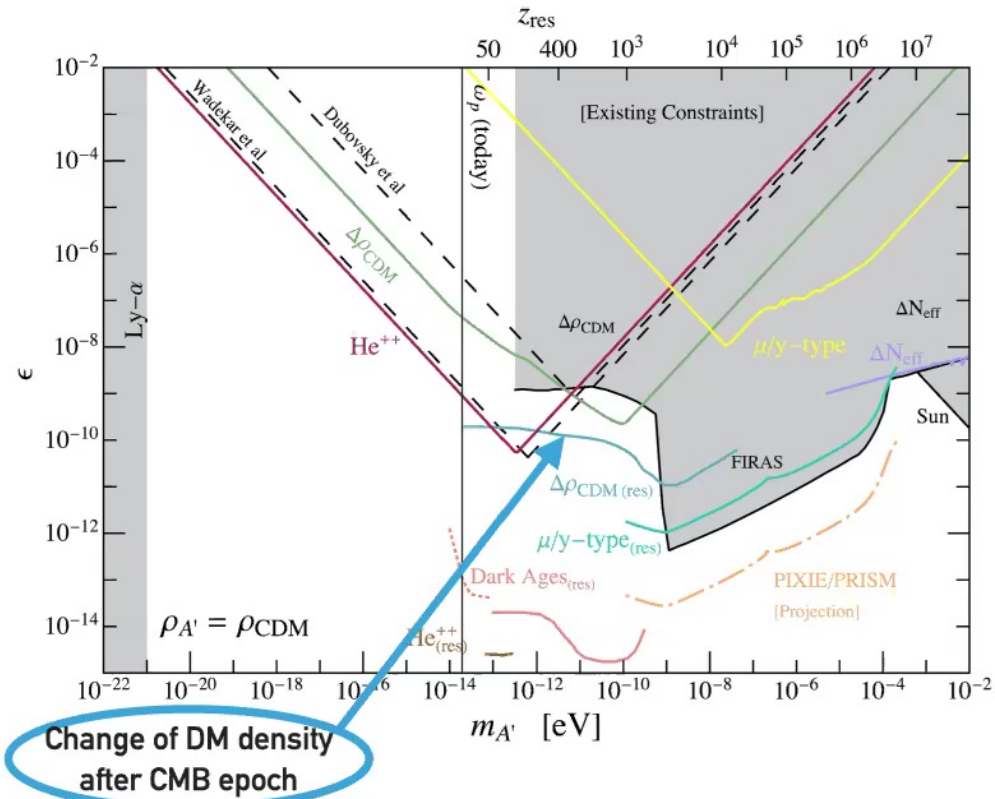




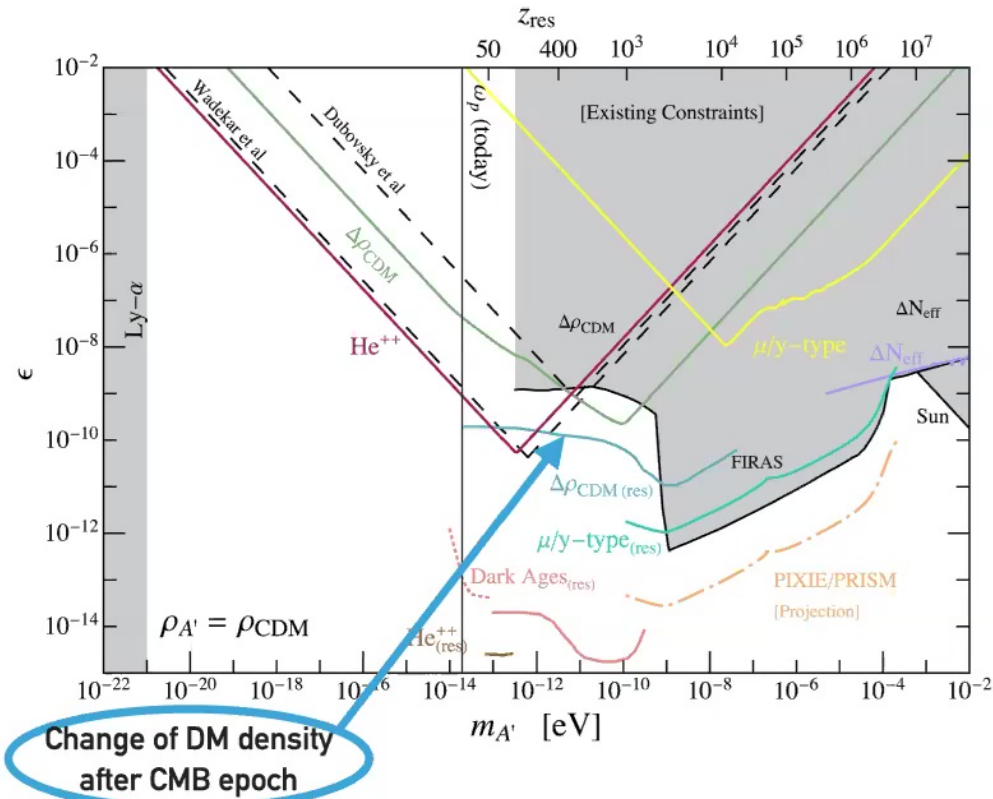
$$\Delta\rho_{\text{CDM}} \iff \text{Prob}_{A' \rightarrow \gamma} = \exp(-1)$$



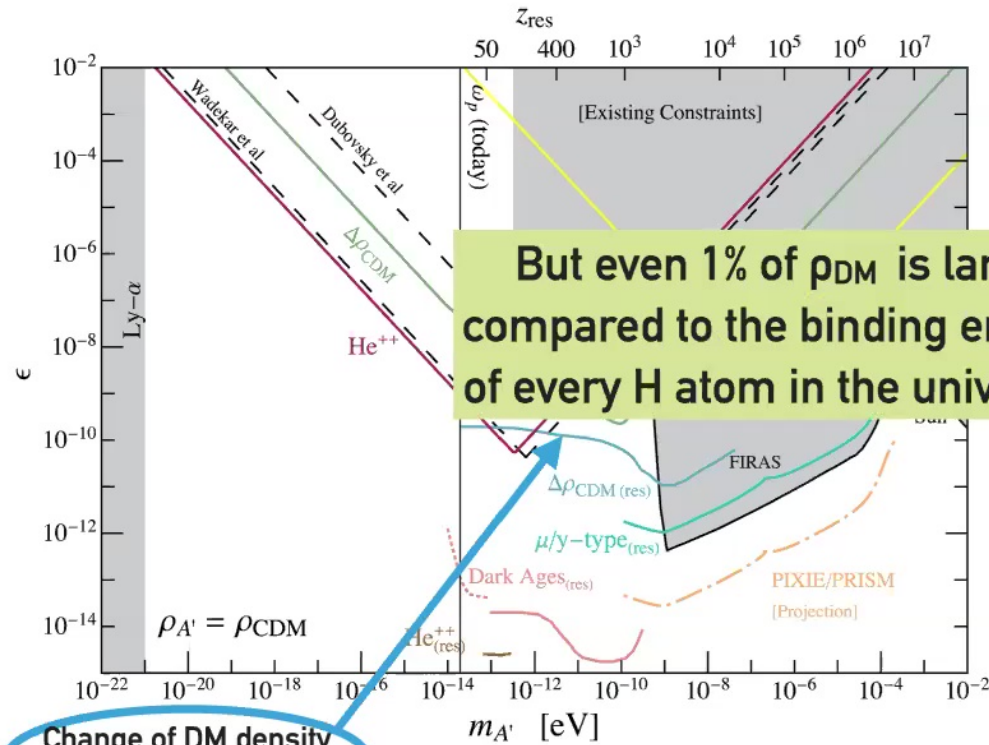
$\Delta\rho_{\text{CDM}} \iff \text{Prob}_{A' \rightarrow \gamma} = \exp(-1) \iff \text{depleting } \Omega_{\text{CDM}} \text{ by 30\%!}$



$$\Delta \rho_{A'}|_{\text{on resonance}} \lesssim 1\% \rho_{A'}$$



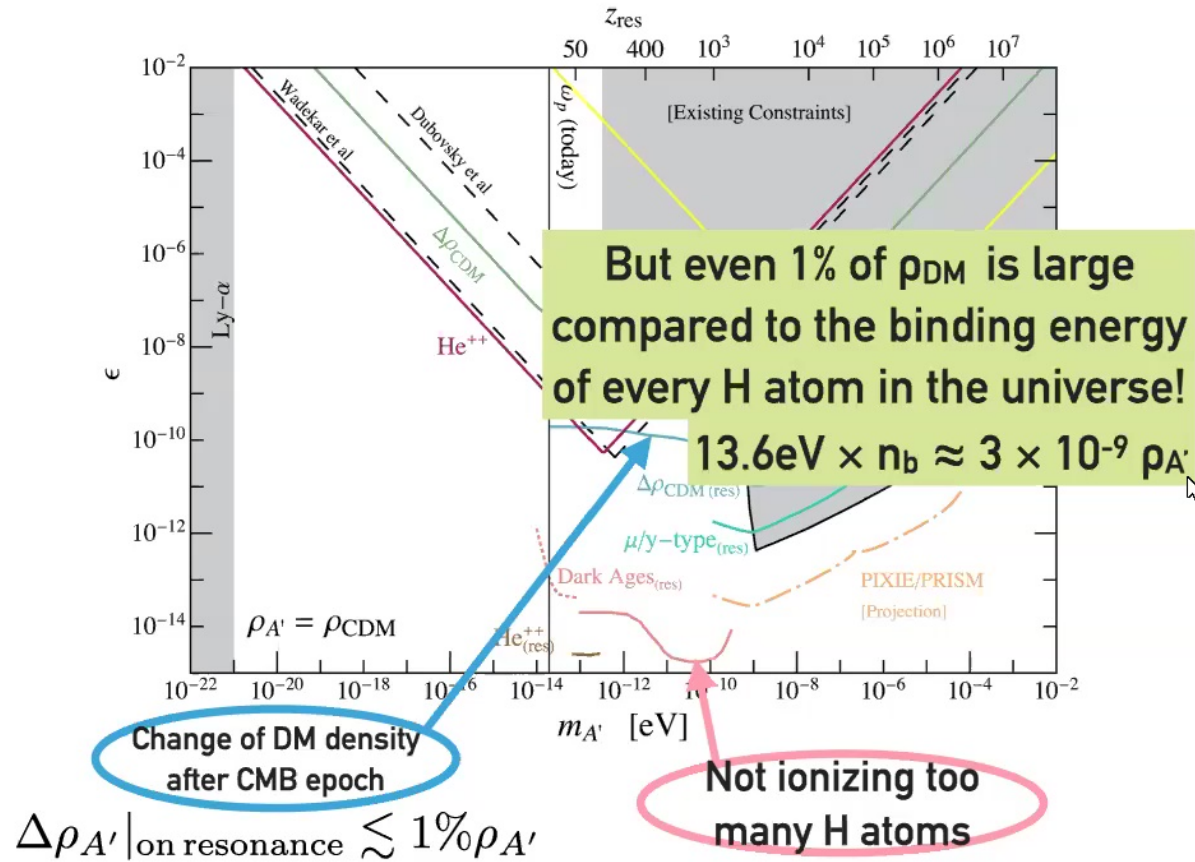
$$\Delta\rho_{A'}|_{\text{on resonance}} \lesssim 1\%_{18}\rho_{A'}$$

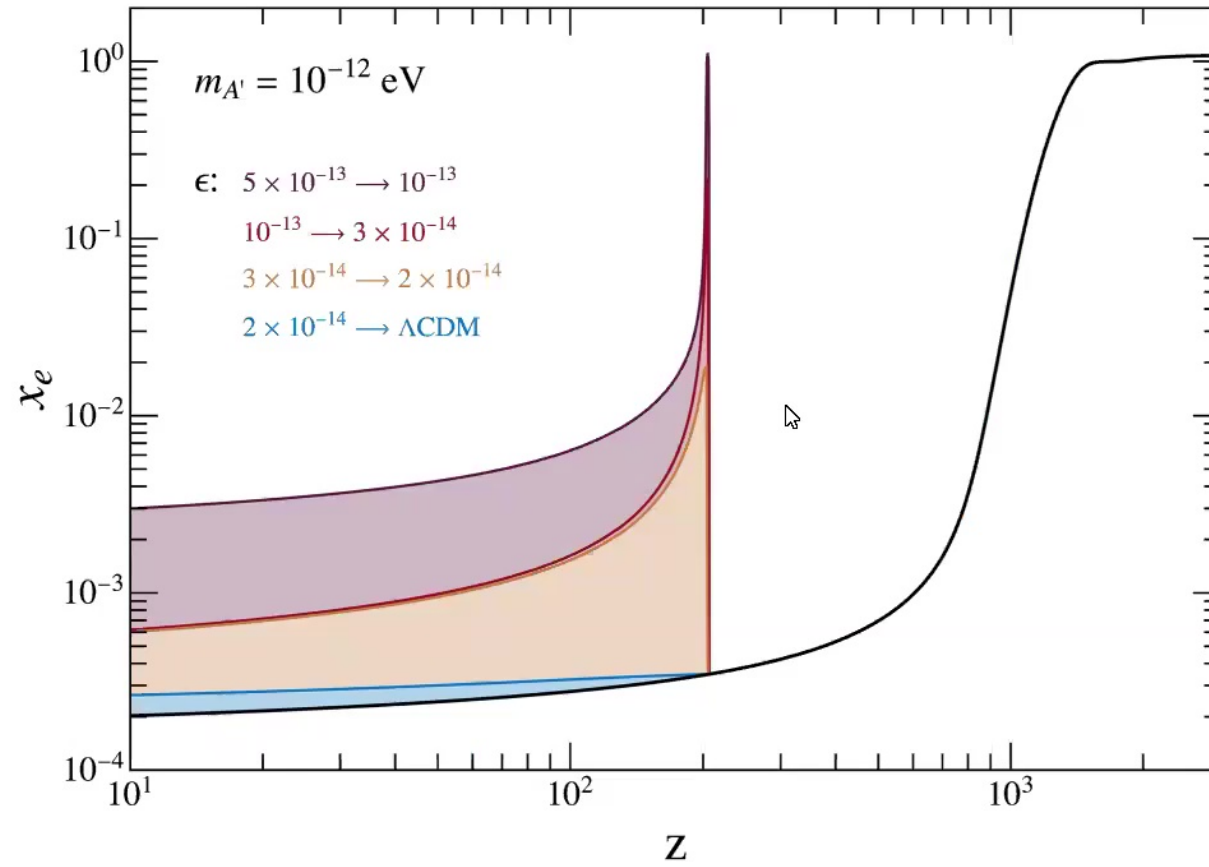


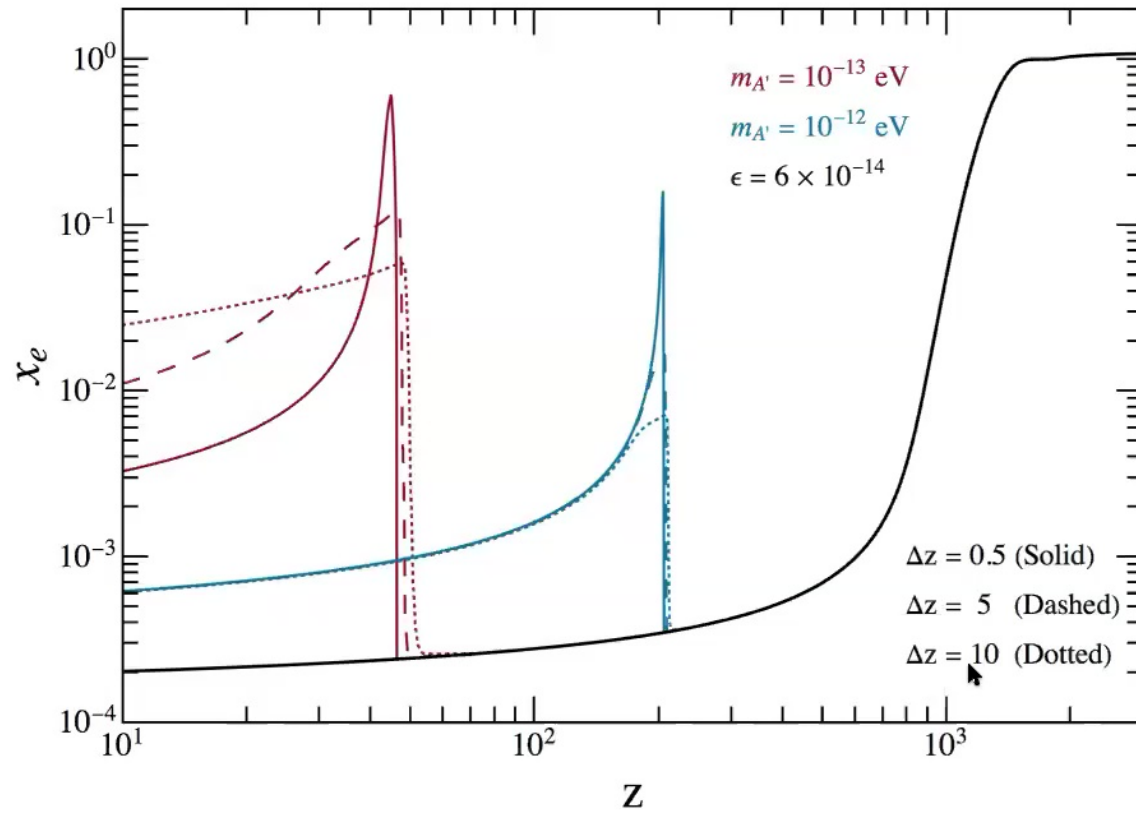
Change of DM density
after CMB epoch

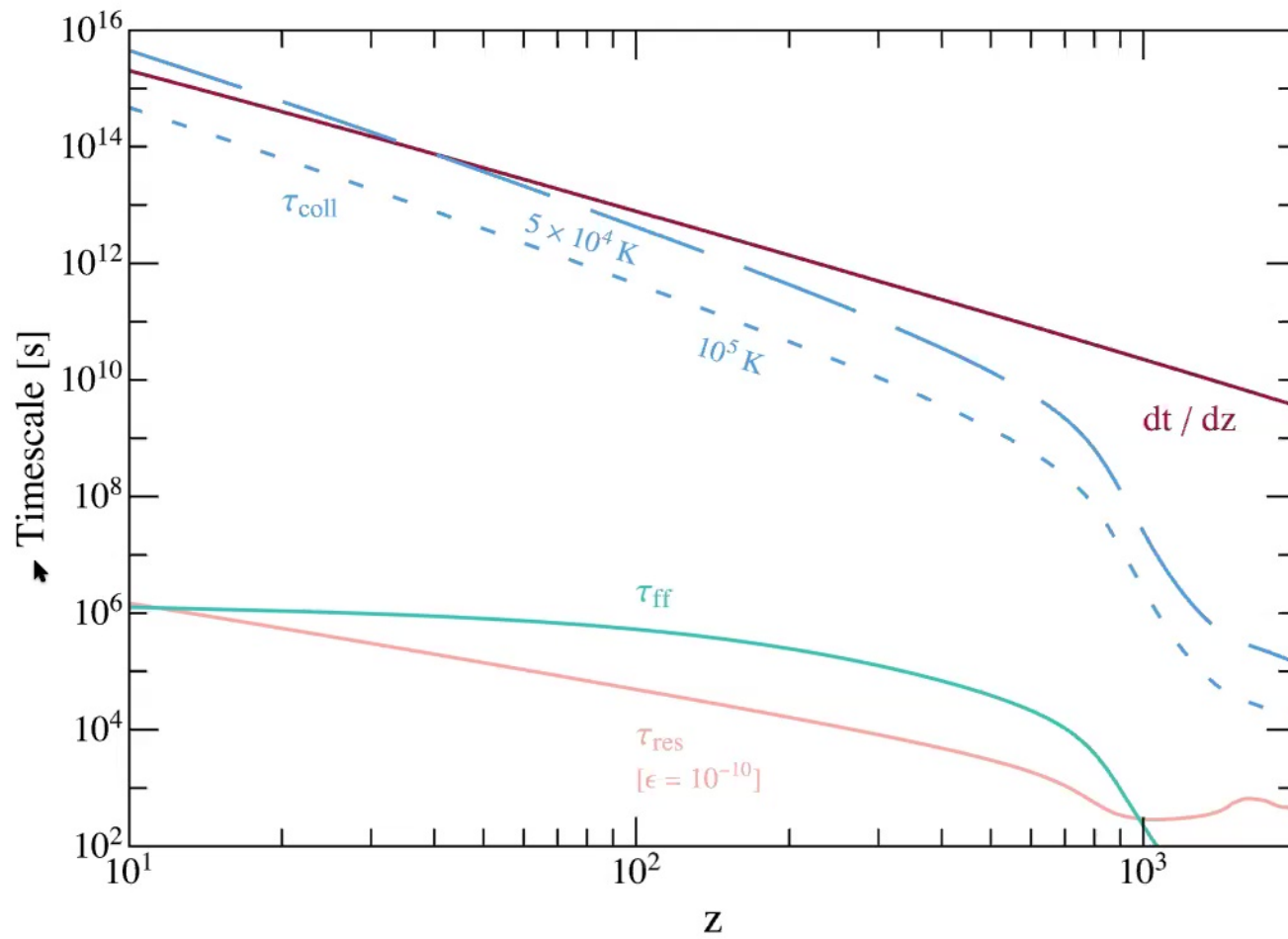
$$\Delta\rho_{A'}|_{\text{on resonance}} \lesssim 1\% \rho_{A'}$$

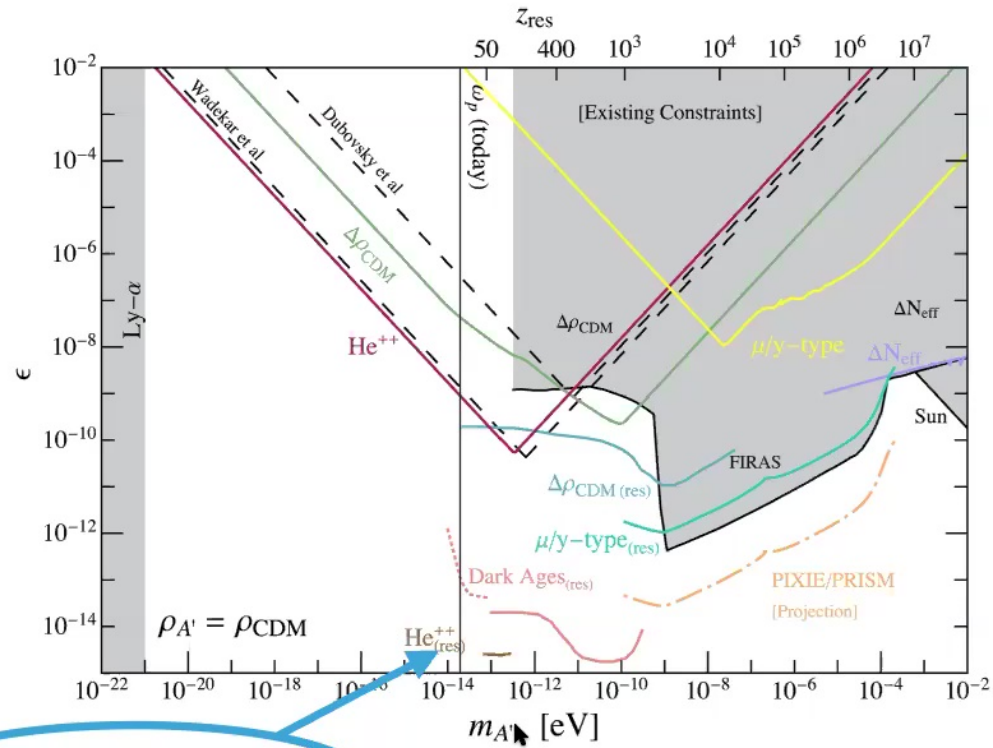




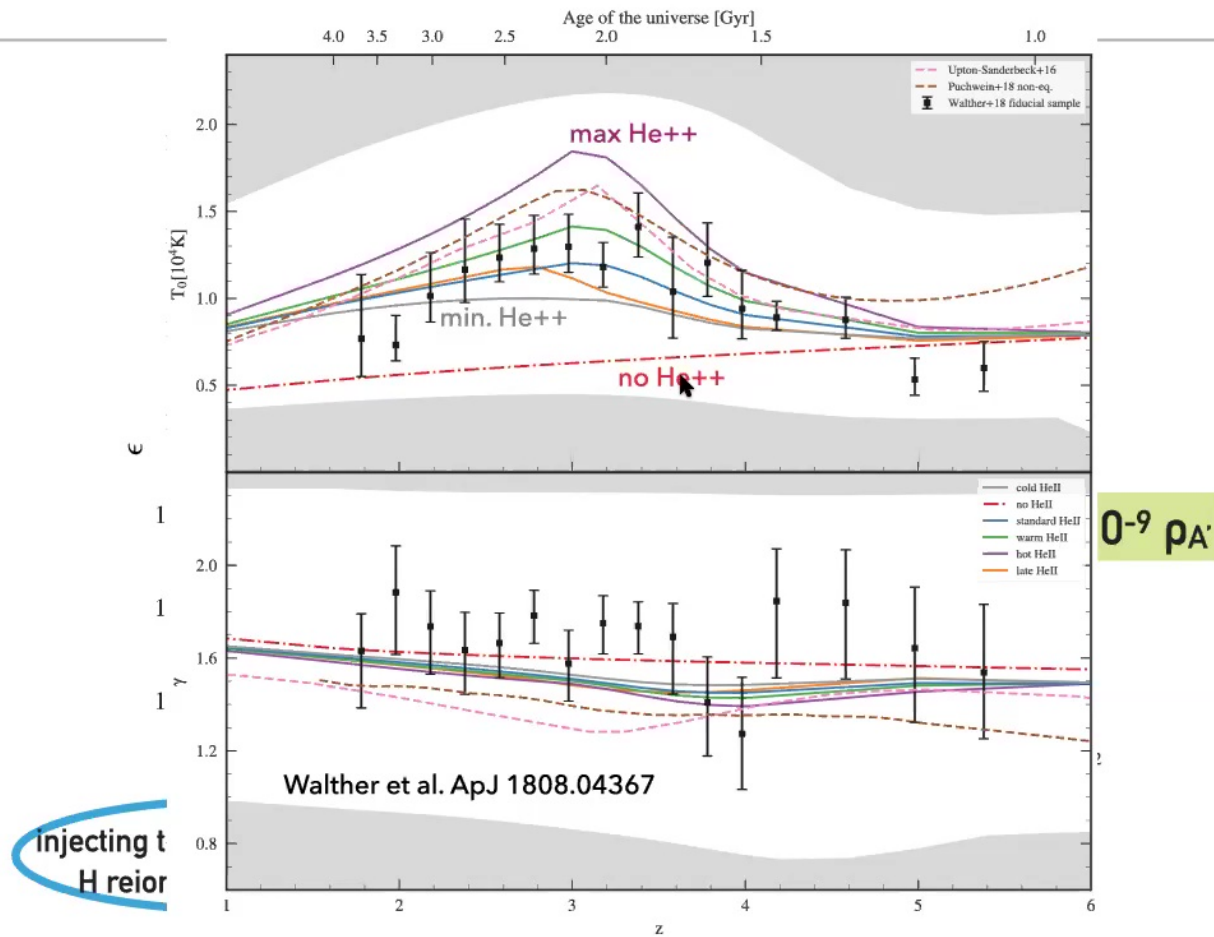


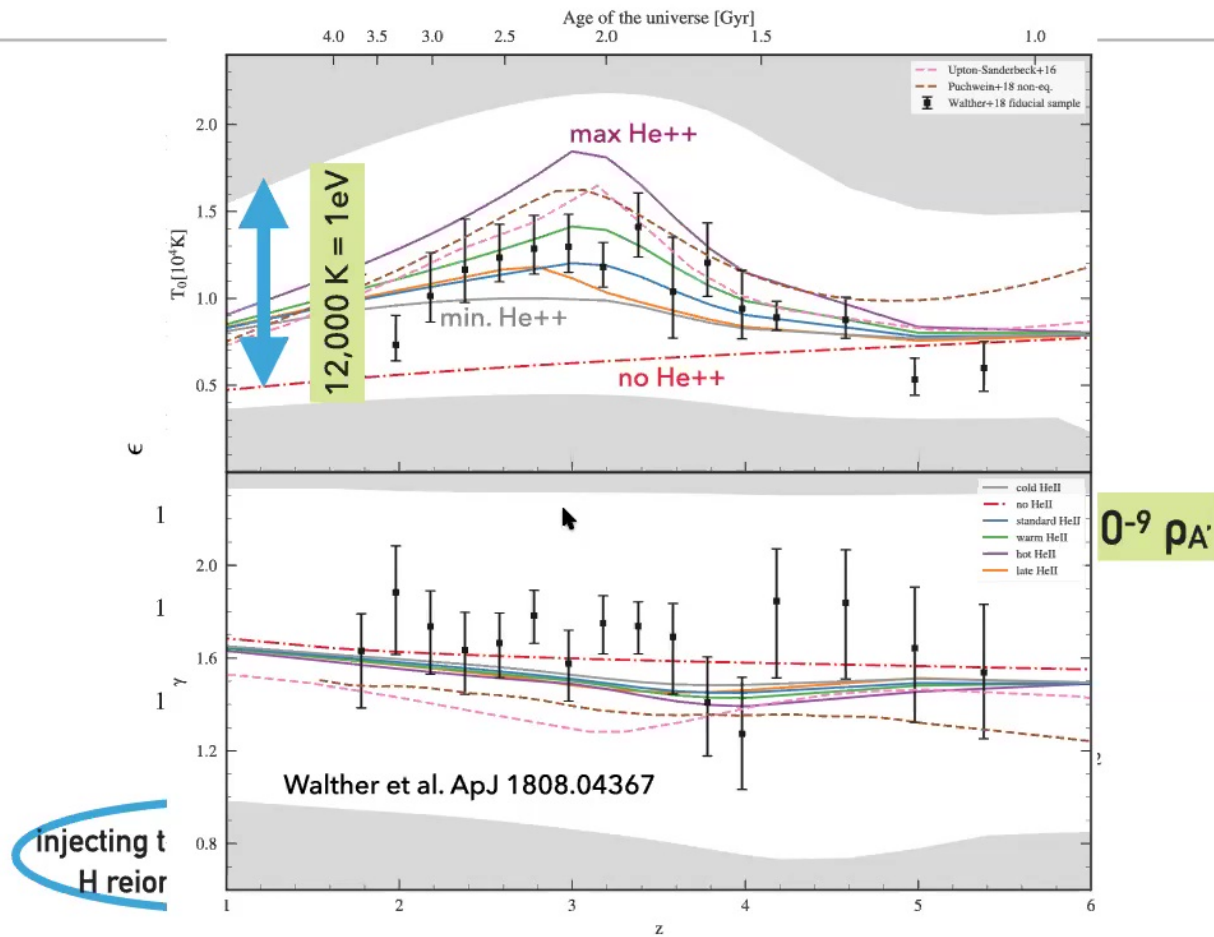


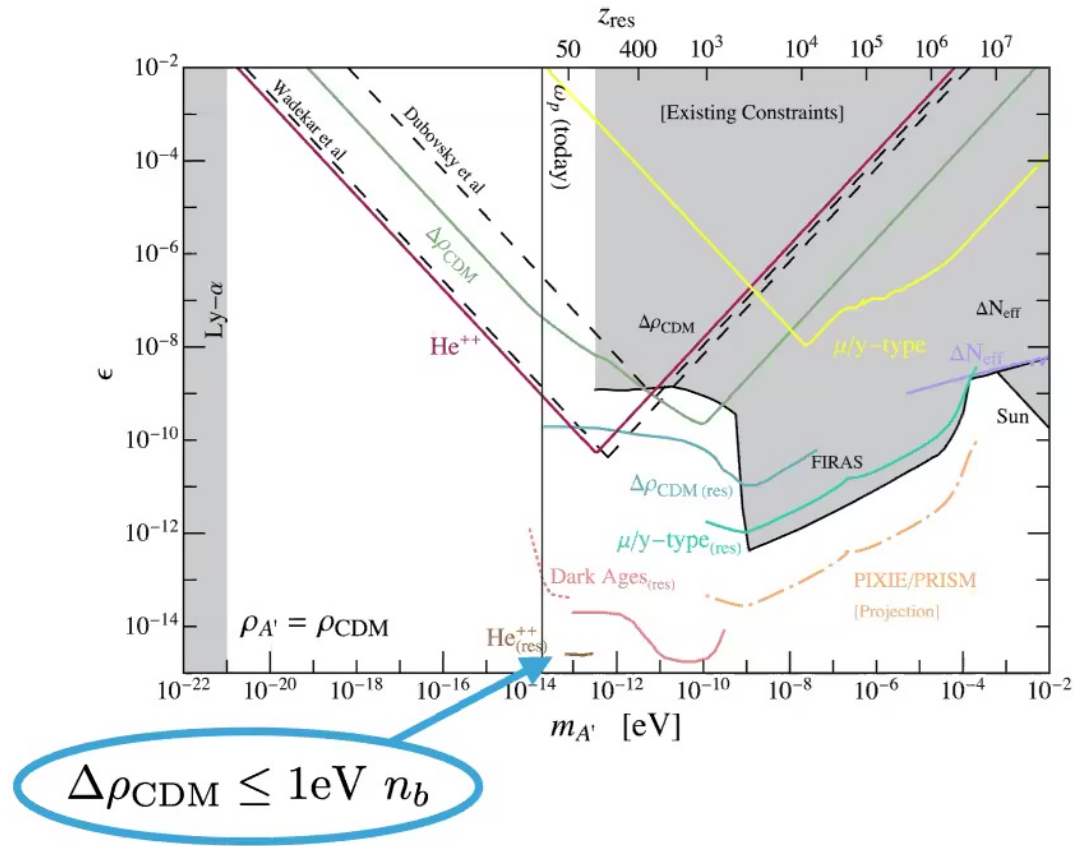




injecting too much energy after
H reionization completes







INHOMOGENEITIES

- ▶ so far all bounds are in the homogeneous limit: there is one special redshift at which $m_{A'} = \omega_p$

$$\frac{d}{dz} P_{A' \rightarrow \gamma}^{\text{homo}}(z) \simeq \frac{\pi \epsilon^2 m_{A'}}{3H(z)} \delta(z - z_{\text{res}}^{\text{homo}})$$

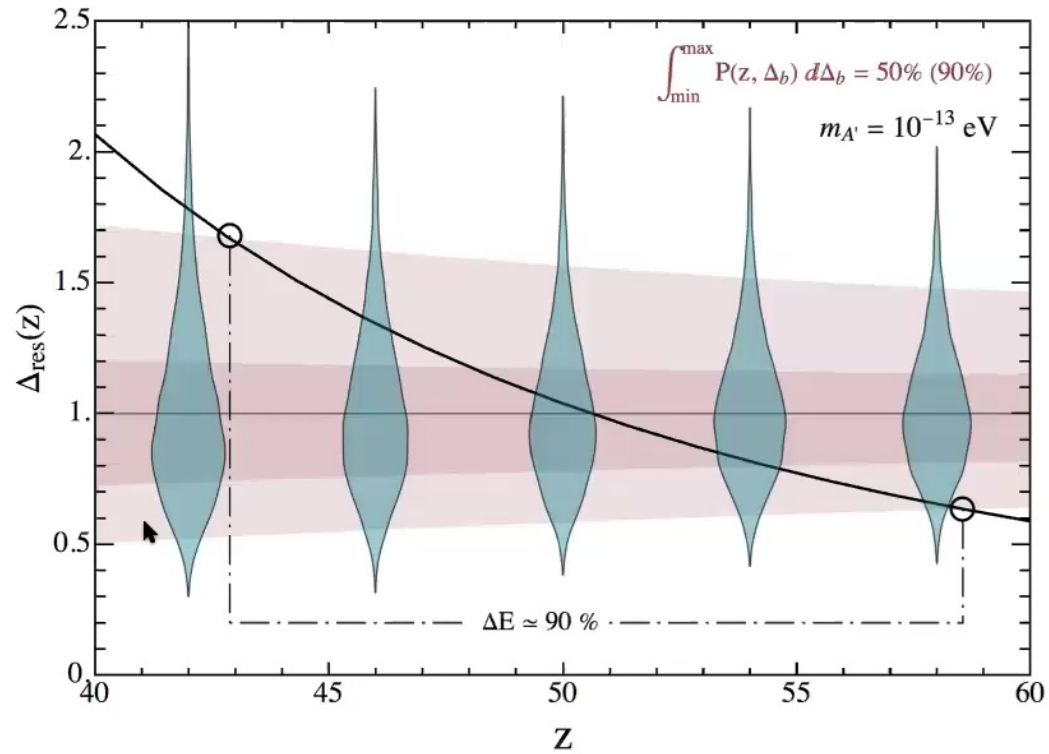
- ▶ including effects of inhomogeneous structure

$$\frac{d}{dz} P_{A' \rightarrow \gamma}(z, \Delta_b) \simeq \frac{\pi \epsilon^2 m_{A'}}{3H(z)} \delta(z - z_{\text{res}}(\Delta_b, m_{A'}))$$

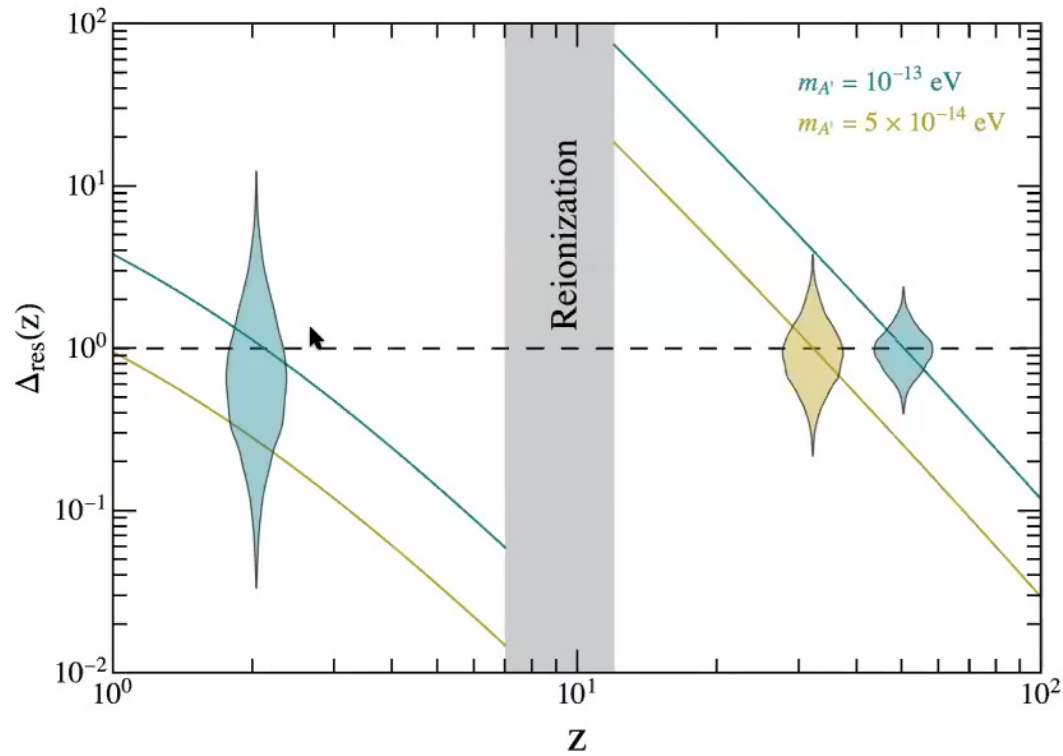
now depends on overdensity Δ_b

see also Caputo, Liu, Mishra-Sharma, Ruderman 2002.05165 and 2004.06733

INHOMOGENEITIES

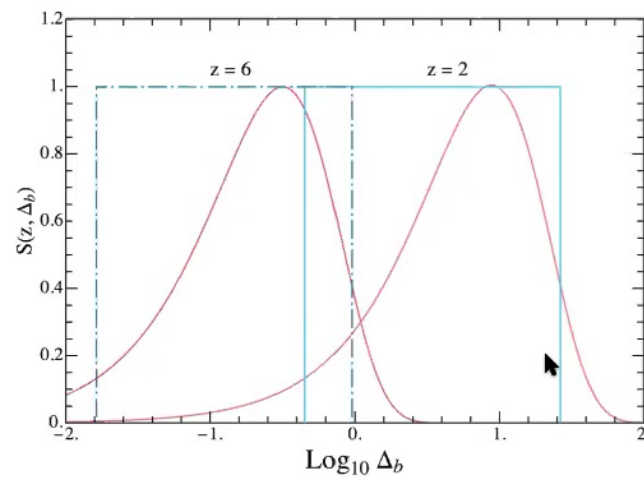


INHOMOGENEITIES



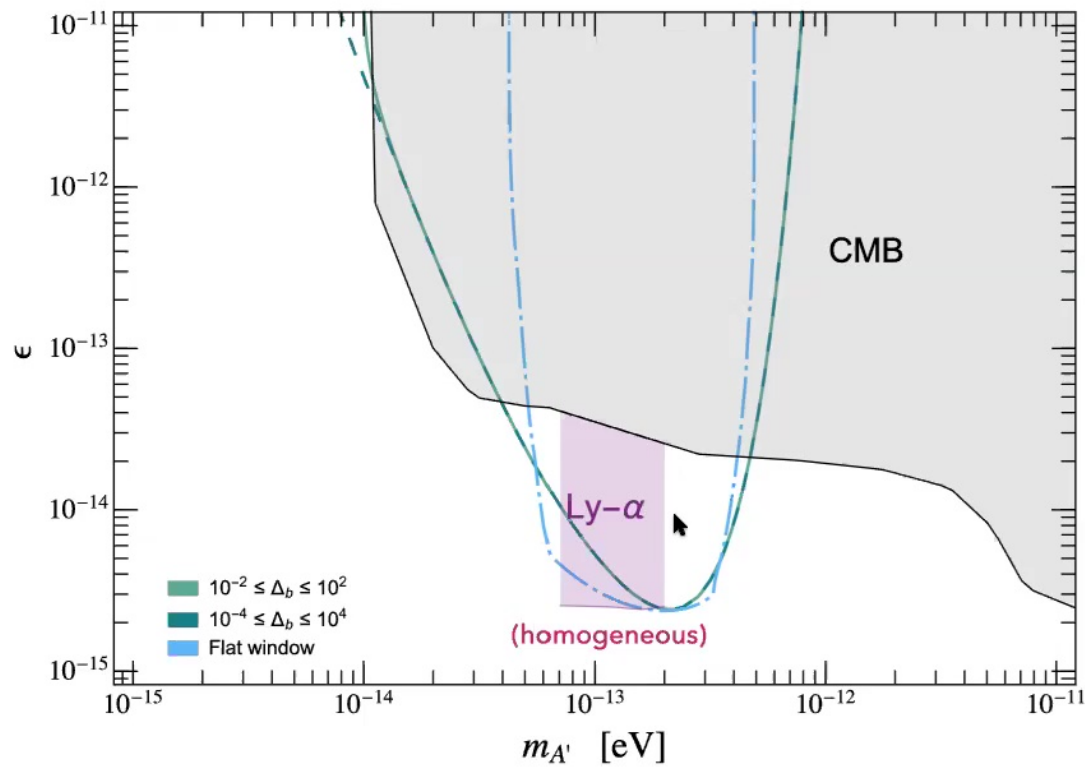
INHOMOGENEITIES

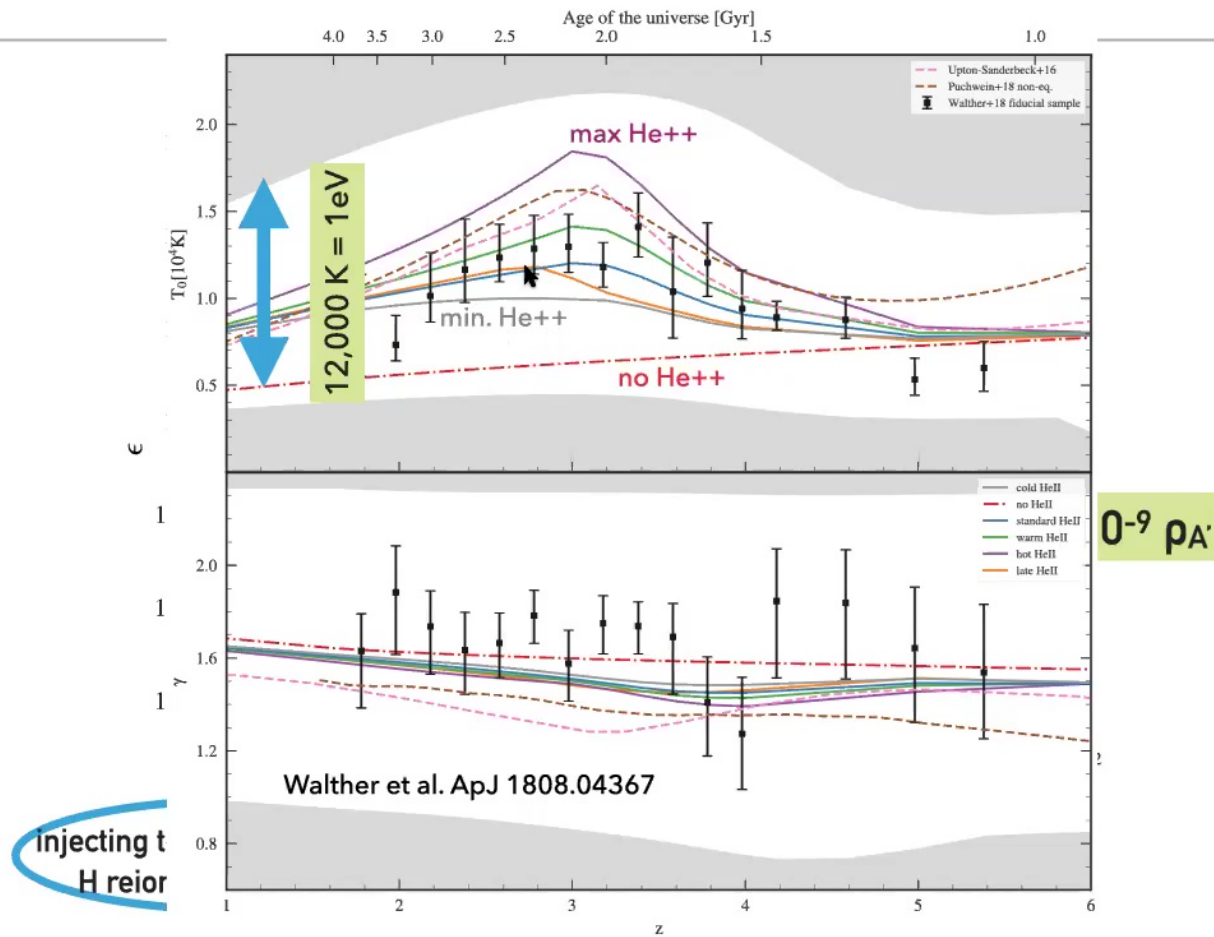
- ▶ constraints now depend on models of inhomogeneous structure formation and the observables of choice



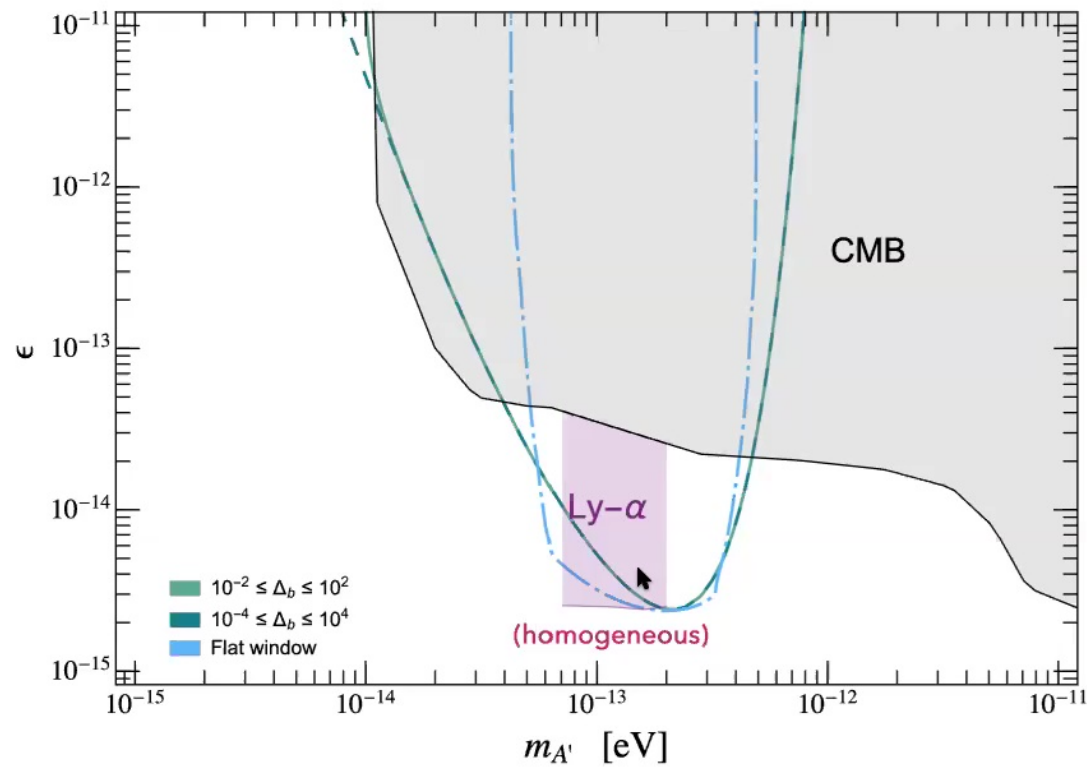
- ▶ e.g., observations during He++ epoch probed by Ly- α

INHOMOGENEITIES



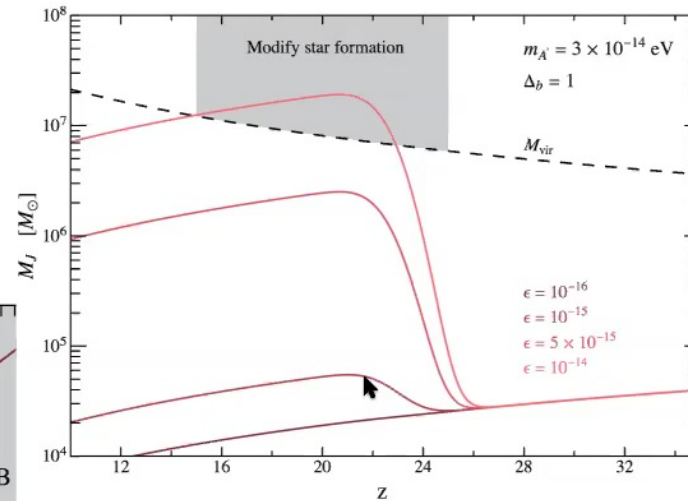
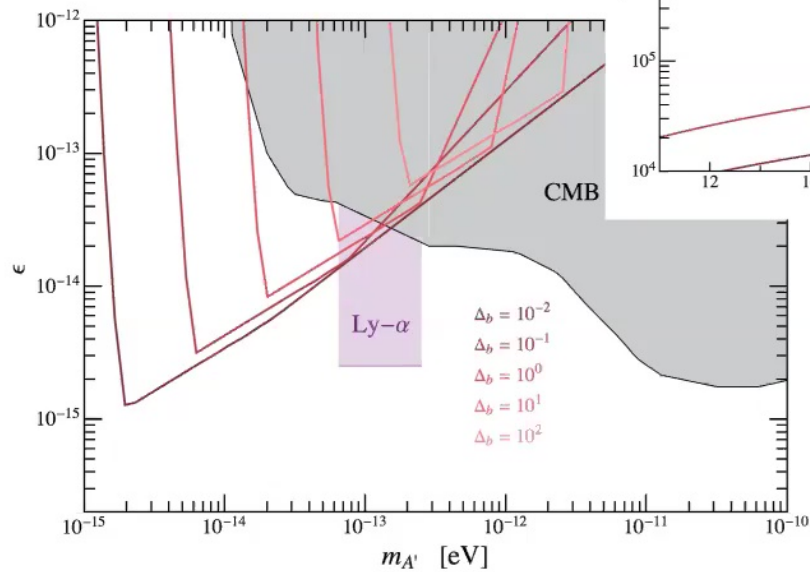


INHOMOGENEITIES



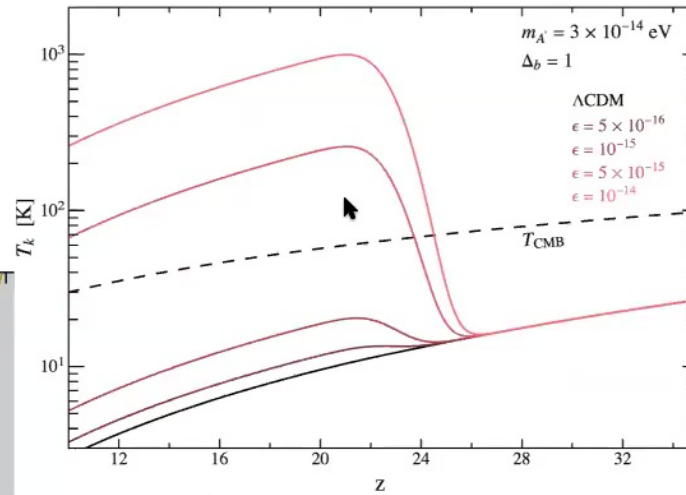
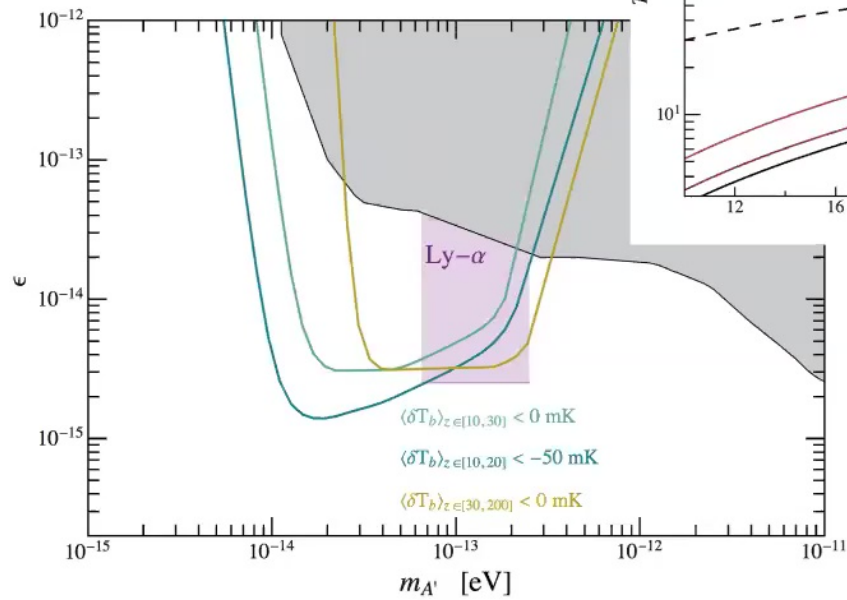
FUTURE POSSIBILITIES

modifying Jeans-mass-scale SM overdensities



FUTURE POSSIBILITIES

modifying 21cm signal



CONCLUSIONS

- ▶ Our understanding of dark photon dark matter has been changing qualitatively over the last few years
- ▶ Plasma mass leads to novel signatures unlike canonical dark matter scenarios
- ▶ Late-time conversion leads to exciting possibilities for future observations

